

New Trends in Fusion Research

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Credits and acknowledgments

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ILE Osaka, ESA, NASA, LLE, UCB, UKAEA, with

apologies to the many authors from whom I have 'stolen' viewgraphs



EUROPEAN FUSION DEVELOPMENT AGREEMENT



Lay-out of the course (*can still be changed!*)

- Mon Oct. 11 –1h
 - Fusion basics, Lawson criterion, plasmas, confinement schemes
 - Inertial confinement fusion
- Tue Oct 12 –2h
 - Magnetic confinement: principles and general challenges
 - The tokamak
 - Heating and current drive
 - Macroscopic equilibrium, stability, operational limits, disruptions
 - Plasma-wall interaction
- Wed Oct 13 –2h
 - Plasma diagnostics
 - Transport of energy and particles
 - The burning plasma regime
 - The future: ITER, the world burning plasma experiment

Web links

- Plasma physics lectures at EPFL
<http://crppwww.epfl.ch/lectures/>
- CRPP-EPFL <http://crppwww.epfl.ch>
- EFDA <http://www.efda.org>
- JET <http://www.jet.efda.org>
- ITER <http://www.iter.org>
- A useful US-based site with many fusion links
<http://www.fire.pppl.gov>

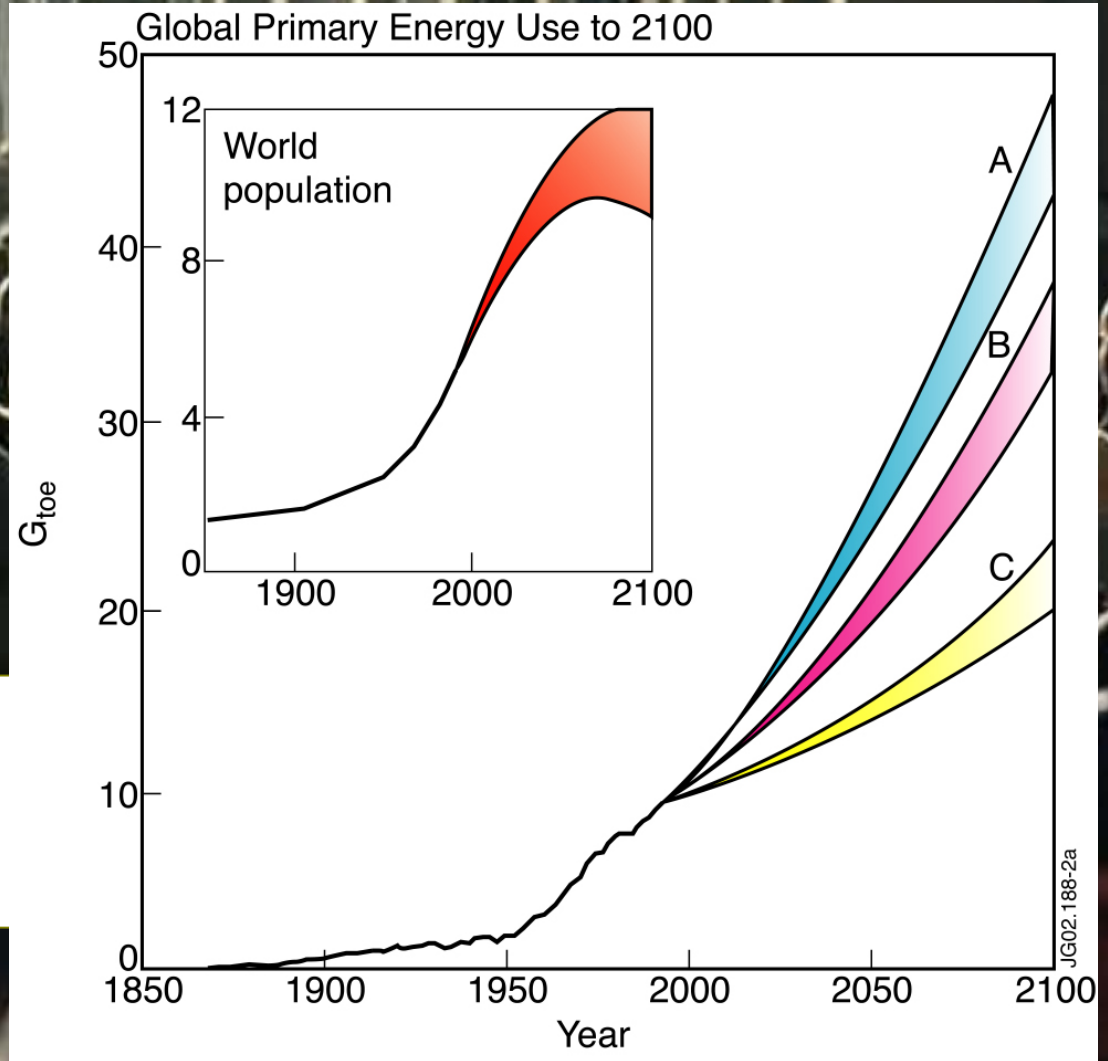
Lay-out of lecture 1 (today)

- Fusion basics
 - Energy needs
 - Urgency of alternates to fossil fuels: scarcity, environmental issues
- Plasmas
- Why *thermo*-nuclear fusion? Coulomb collisions in plasmas
- Fusion energy balance: Lawson criterion, breakeven and ignition
- Confinement schemes
- Inertial confinement fusion (*the view of a non-expert*)
 - Direct drive
 - Indirect drive
 - Drivers: lasers, ion beams, X-ray (Z-pinch)
 - Fast ignition

Growth of world population and energy demand

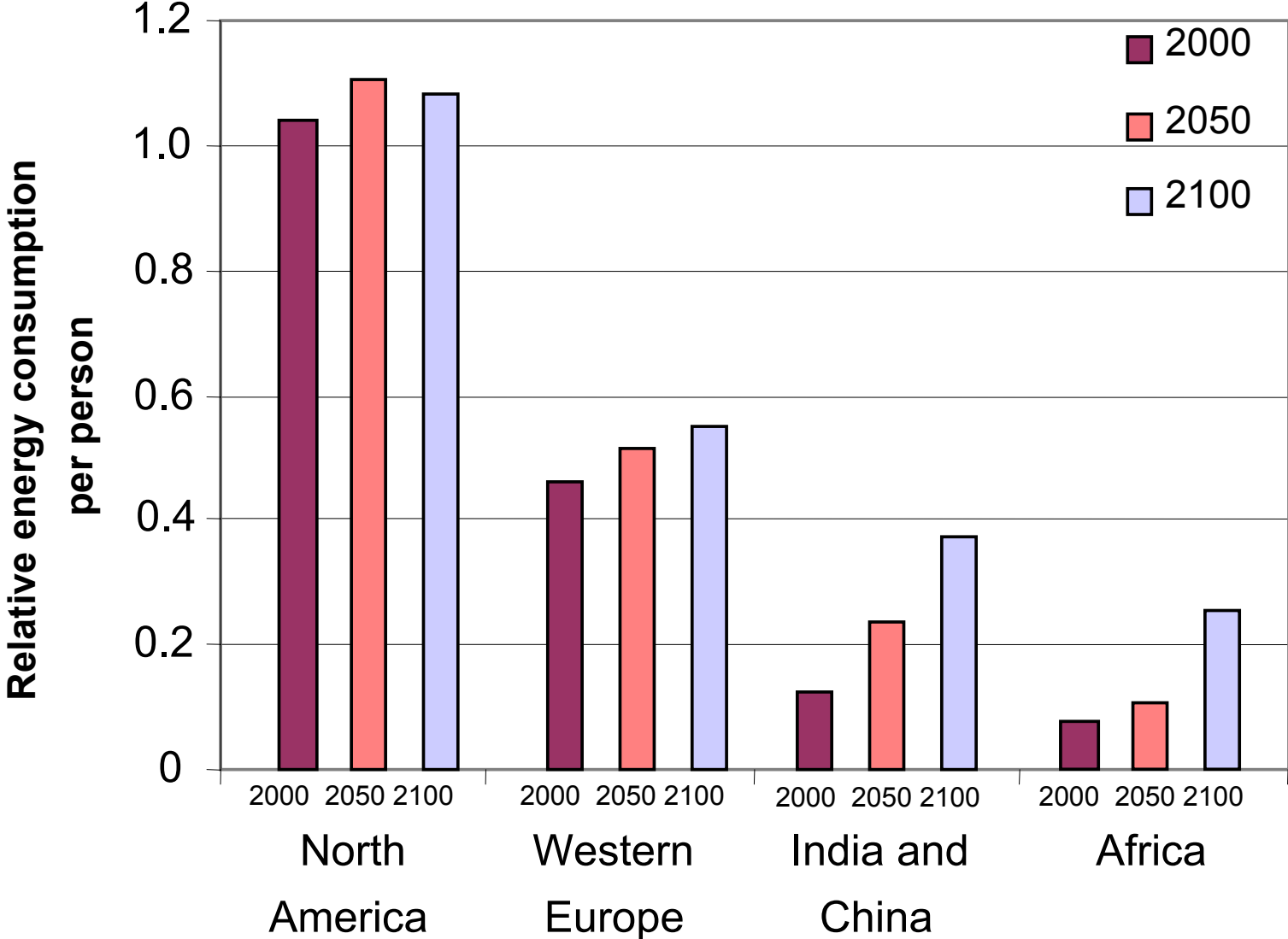


Energy demand grows even faster than world population



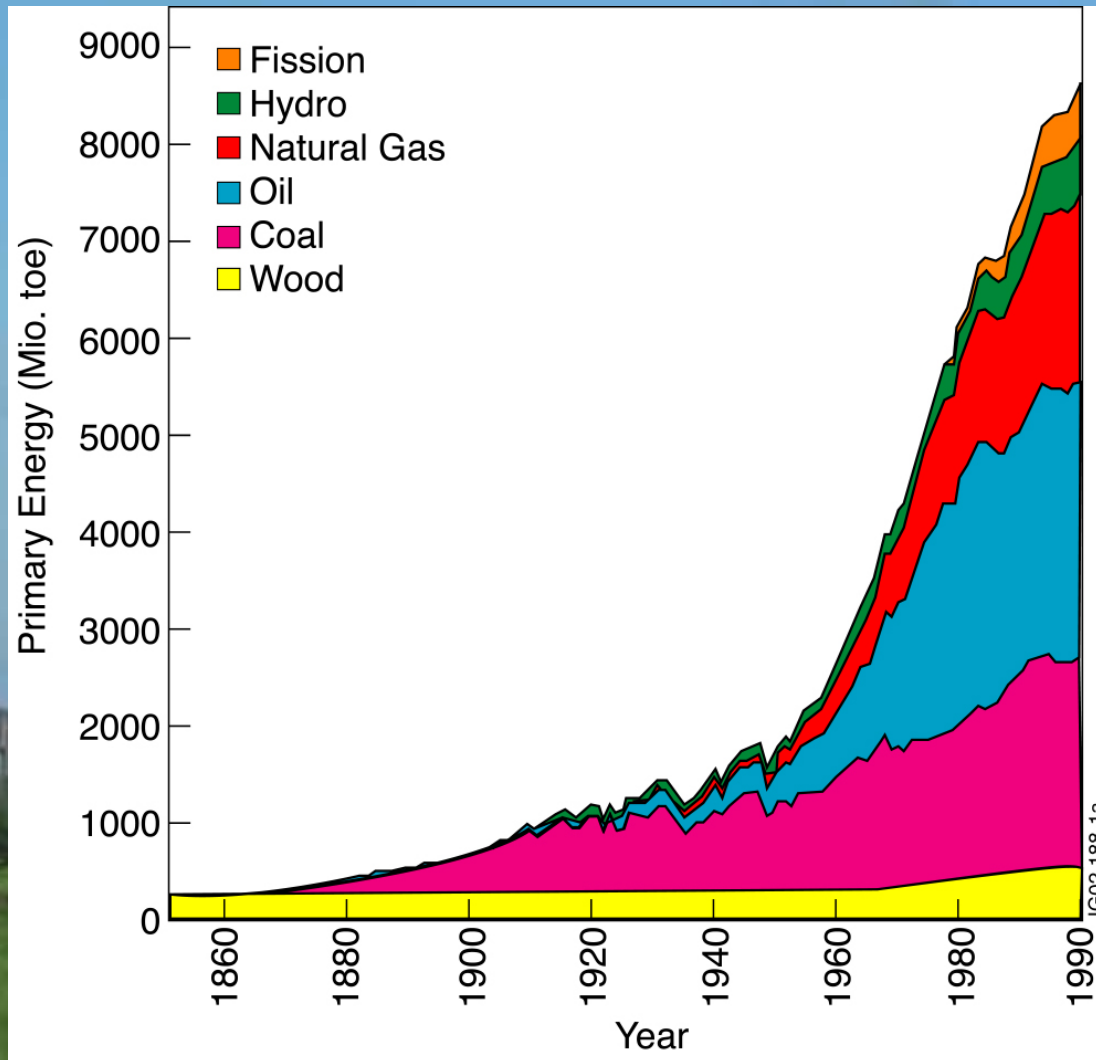
Unequal distribution of energy consumption

IIASA/WEC 1998: business as usual scenario



Reliance on fossil fuels

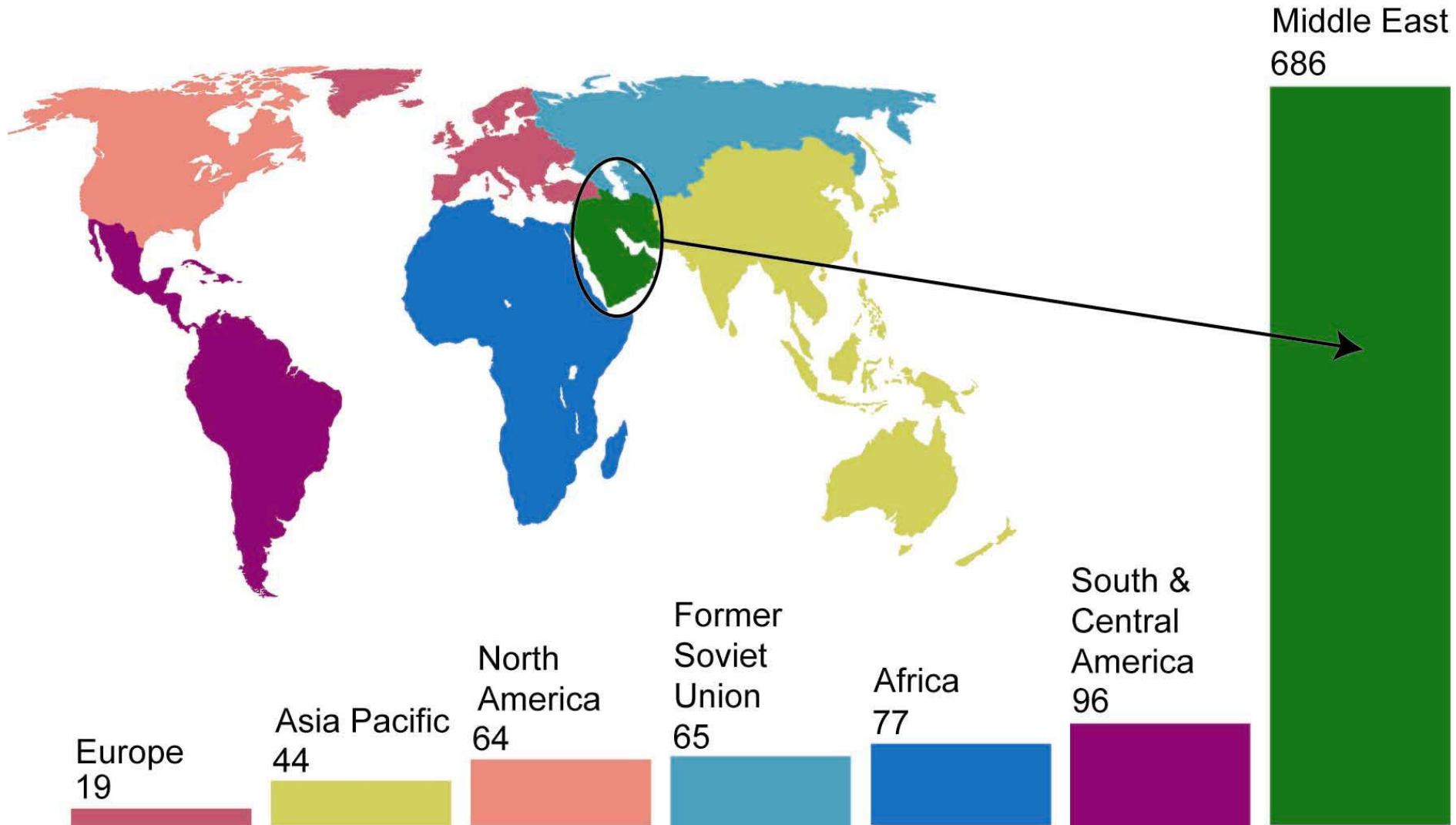
Fossil Fuels have been produced from decayed plant and animal matter over millions of years, cannot be re-formed in time



R.W.Bentley et al., Perspectives on the Future of Oil, Vol. 18, Nos. 2&3 2000, MULTI-SCIENCE PUBLISHING CO. LTD., 107 High Street, Brentwood, Essex, CM14 4RX, UK.

Urgency of alternates development

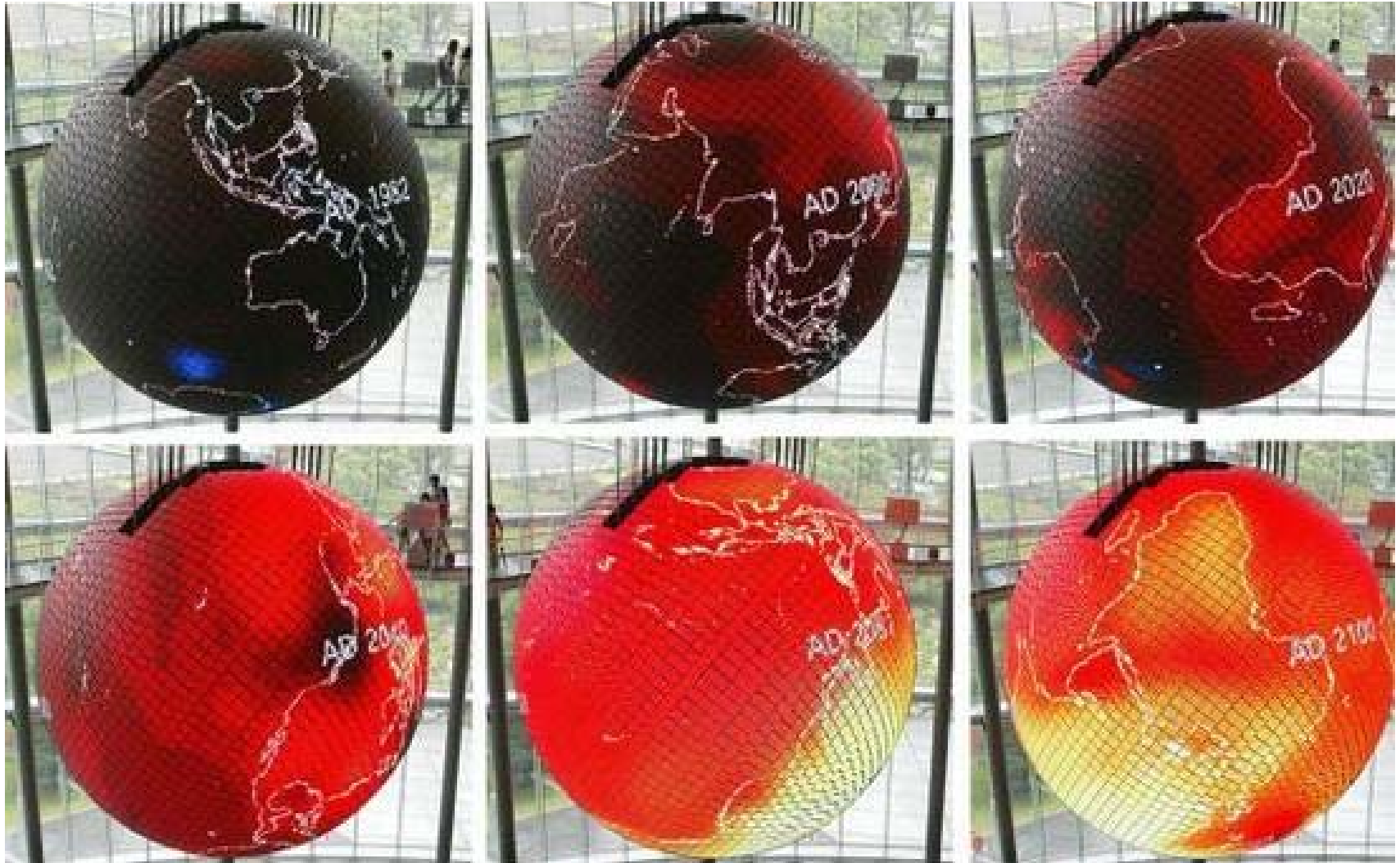
Oil geographical distribution



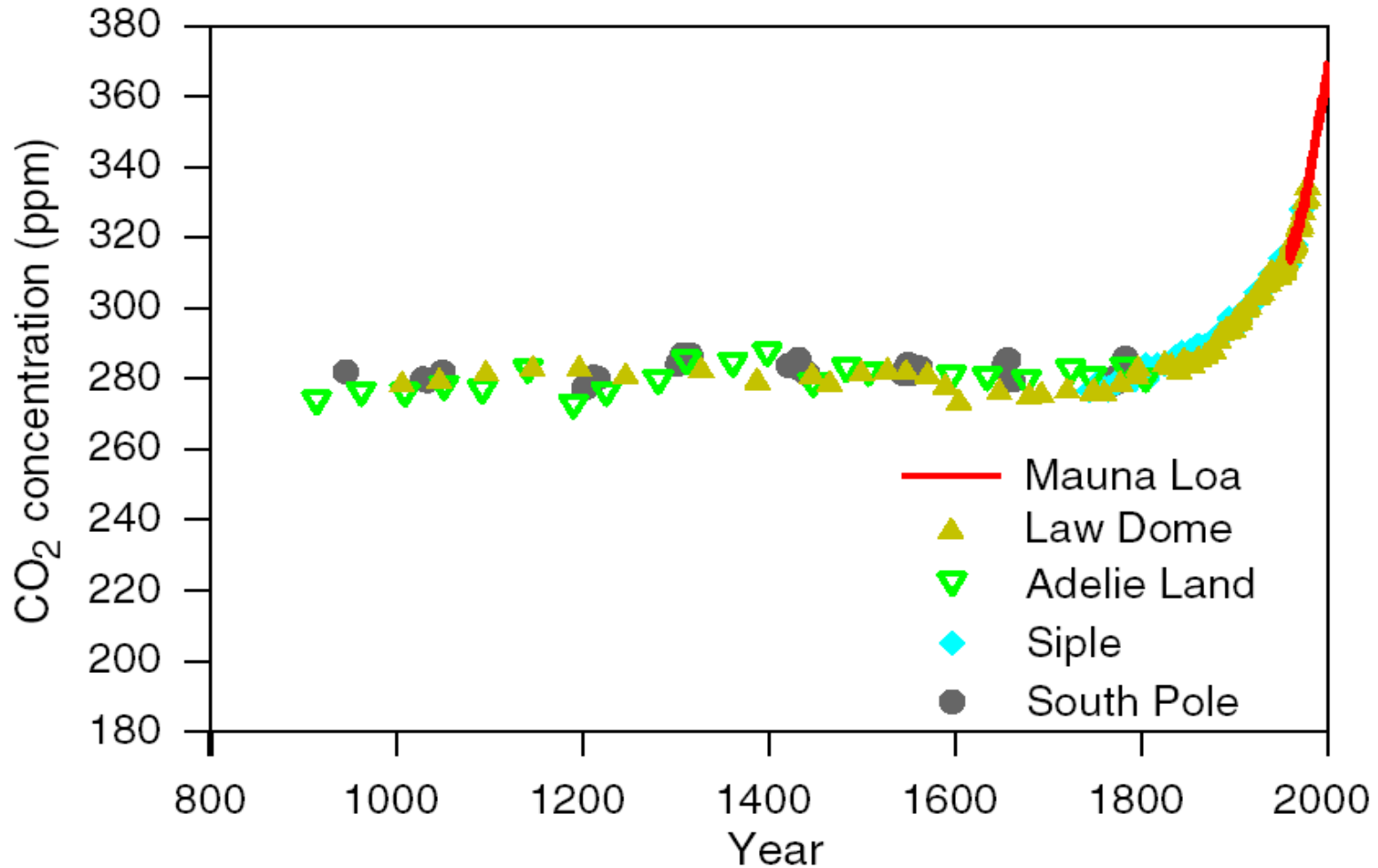
(source BP statistical review 2002)

Urgency of alternates development

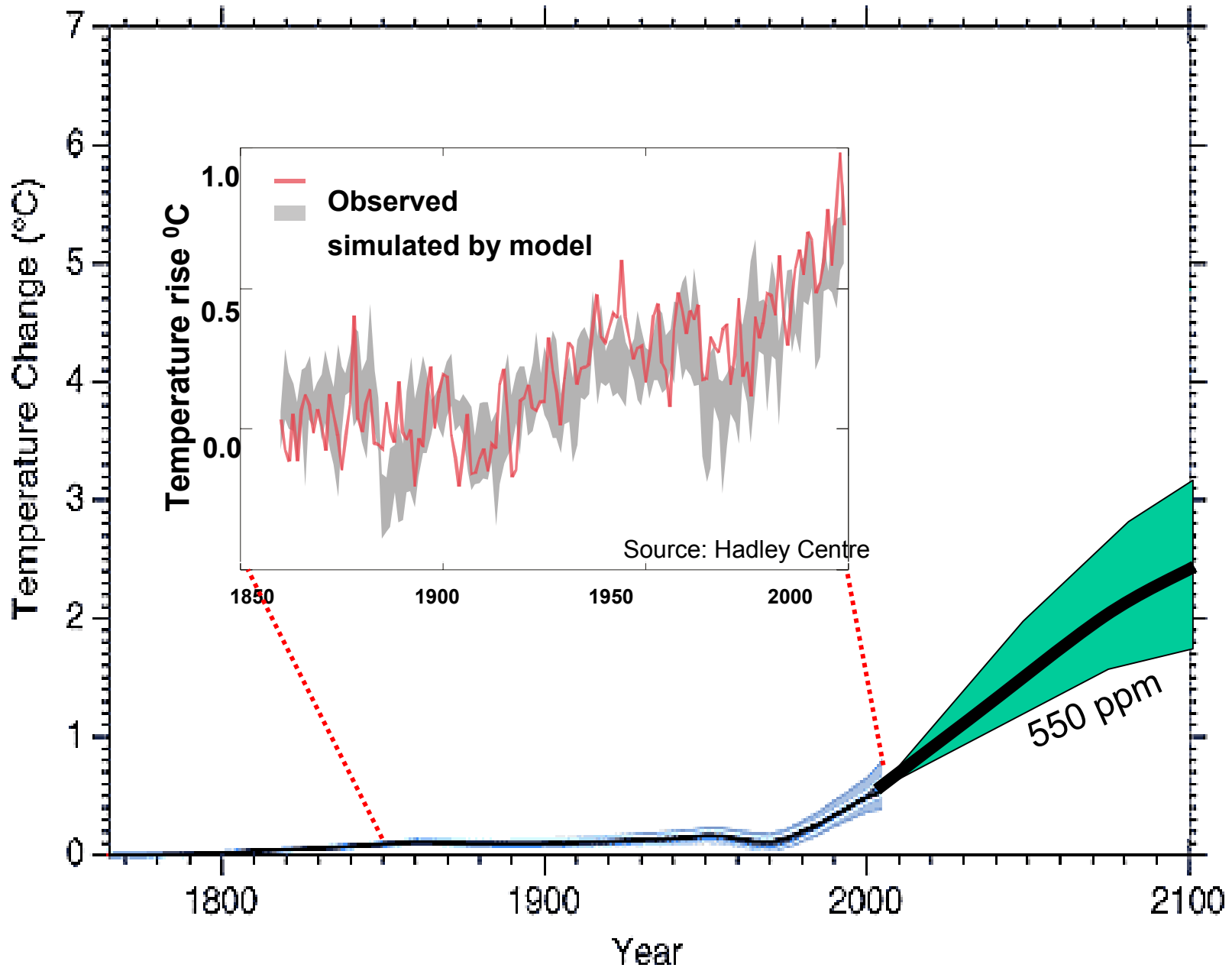
Environmental impact (CO₂ and global warming)



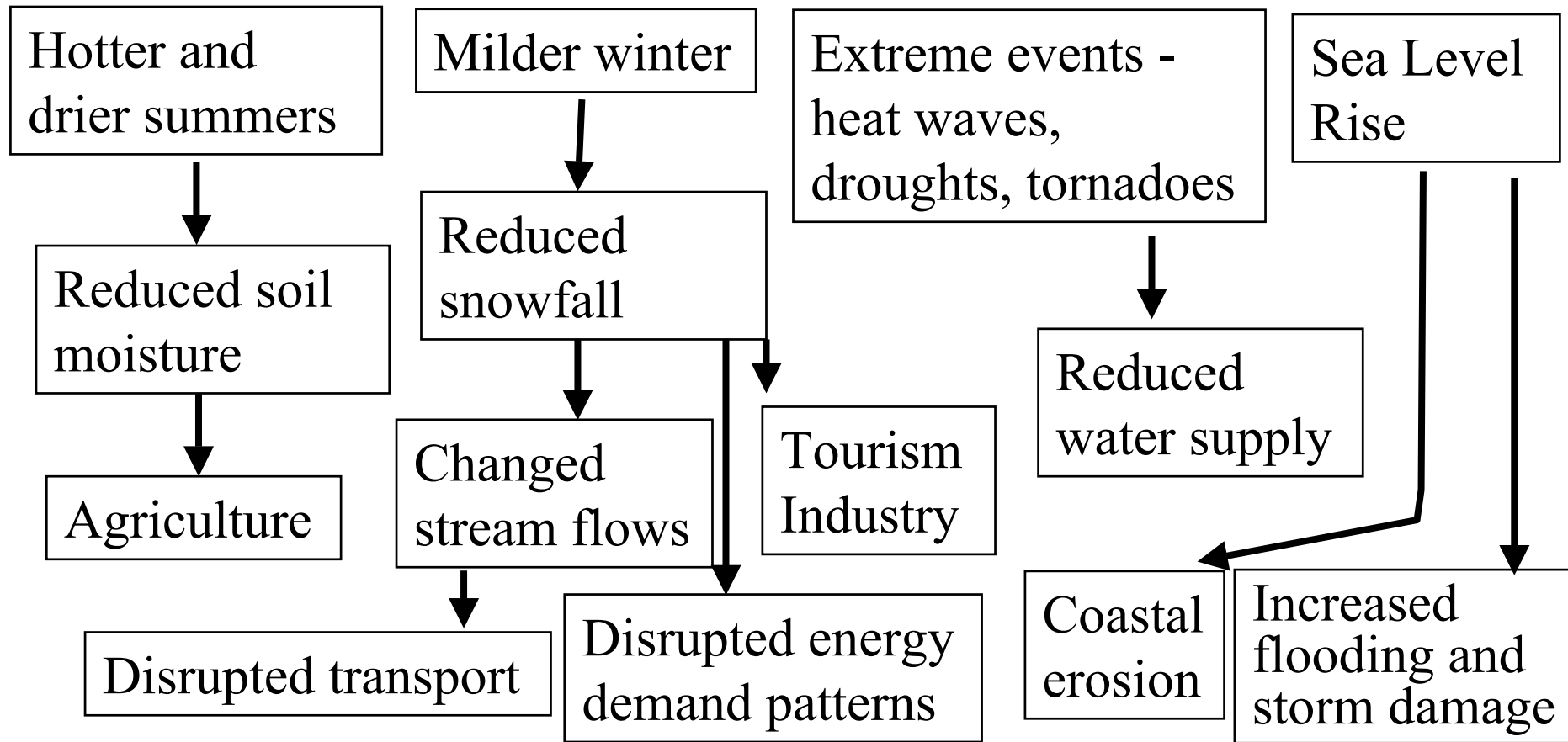
CO₂ is prime greenhouse gas



Modelling global warming



The effects of climate change



Ambitious goal for 2050 (when total world power market predicted to be 30TW)

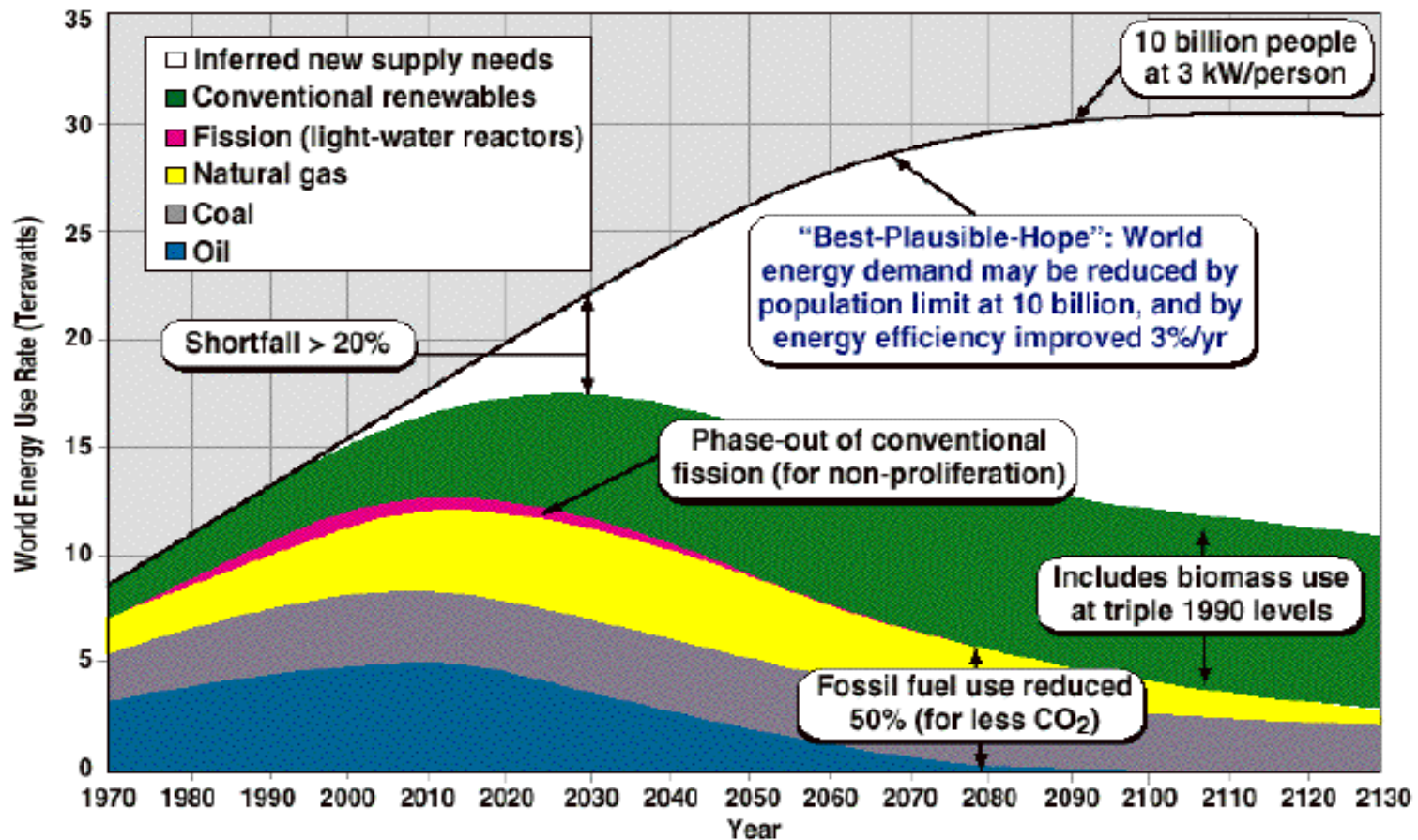
Limit CO₂ to twice pre-industrial level: will need 20 TW of CO₂-free power (today's world total power market is 13 TW)

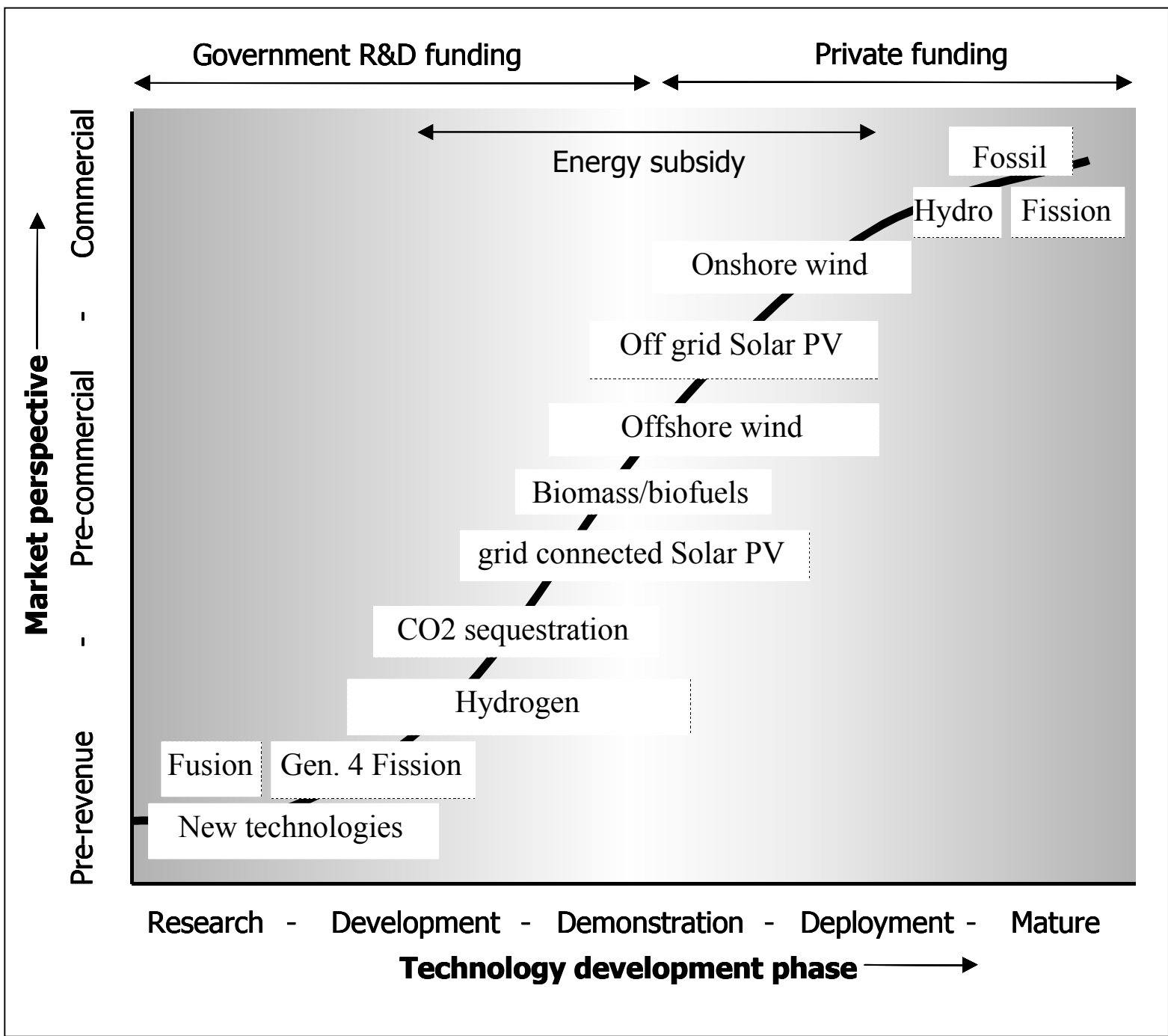
US DoE "The technology to generate this amount of emission-free power does'nt exist"

Urgency of alternates development

Scarcity of resources: related to CO₂ limitations

“BEST-PLAUSIBLE-HOPE” PROJECTION: 5 TW OF NEW NON-FOSSIL ENERGY SOURCES NEEDED BY 2030 (e.g., SOLAR, ADVANCED FISSION, FUSION)



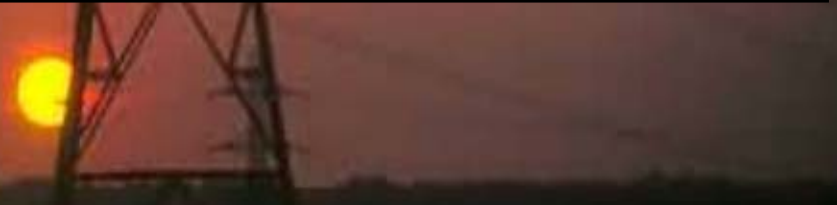


Alternative energy sources

Nuclear Fission

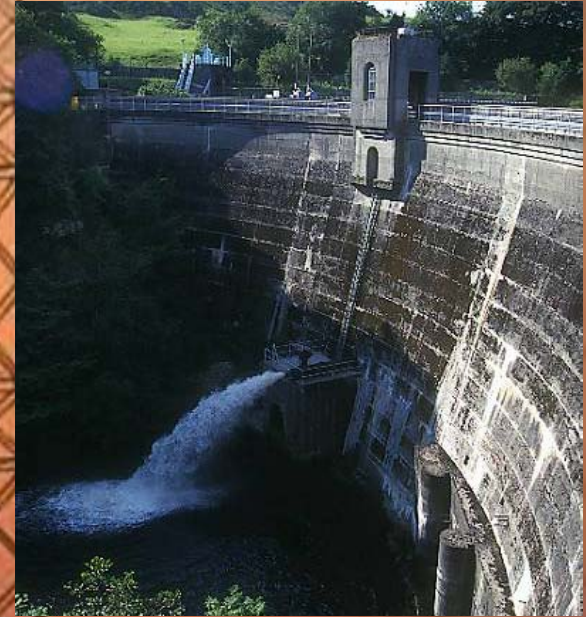


- **Long lived radioactive waste products (many thousands of years) that require transportation, re-processing and geological storage**
- **Public concerns on safety**



Alternative energy sources

Renewables



Renewables (wind, wave, solar, hydro) are the most attractive option at present and offer long term, clean energy reserves

However :

- Low energy density
- Fluctuations in time require storage systems


Ex.: Contribution to electricity of wind energy in 2002

Country or region	Cumul. installed (GW) 2002	share of electricity by wind
Germany	12	4%
Spain	5	5.5%
Denmark	2.8	18.5%
EU	24	2%
US	4.7	0.3%
Total World	32	0.45%

**Going to 20% of electricity in EU (240GW)
would mean 1500 big wind parks**

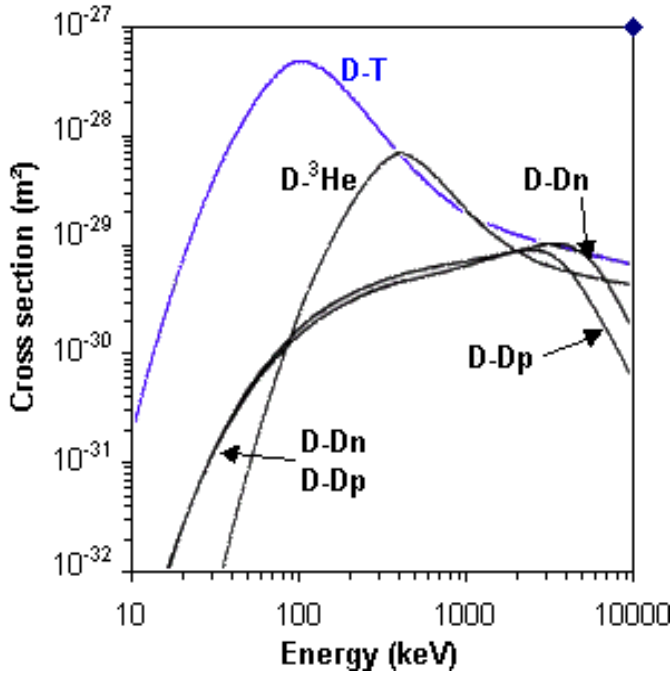


Horns Rev in Denmark: 160 MW (80 mills)

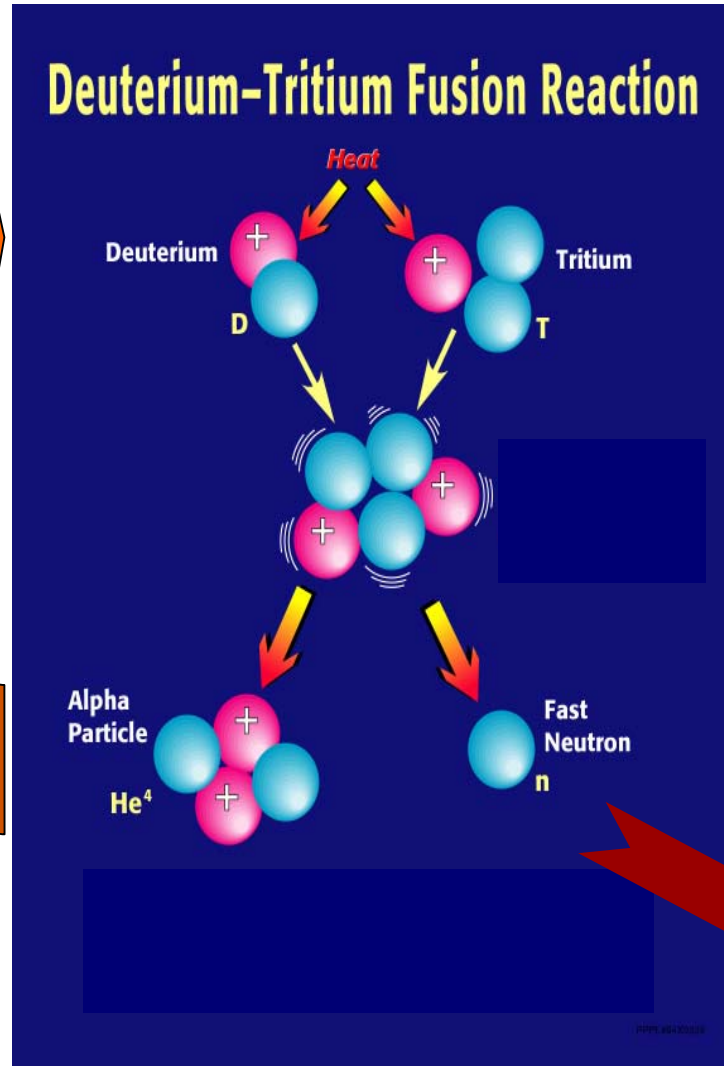


Thermonuclear Fusion

Fusion reactions



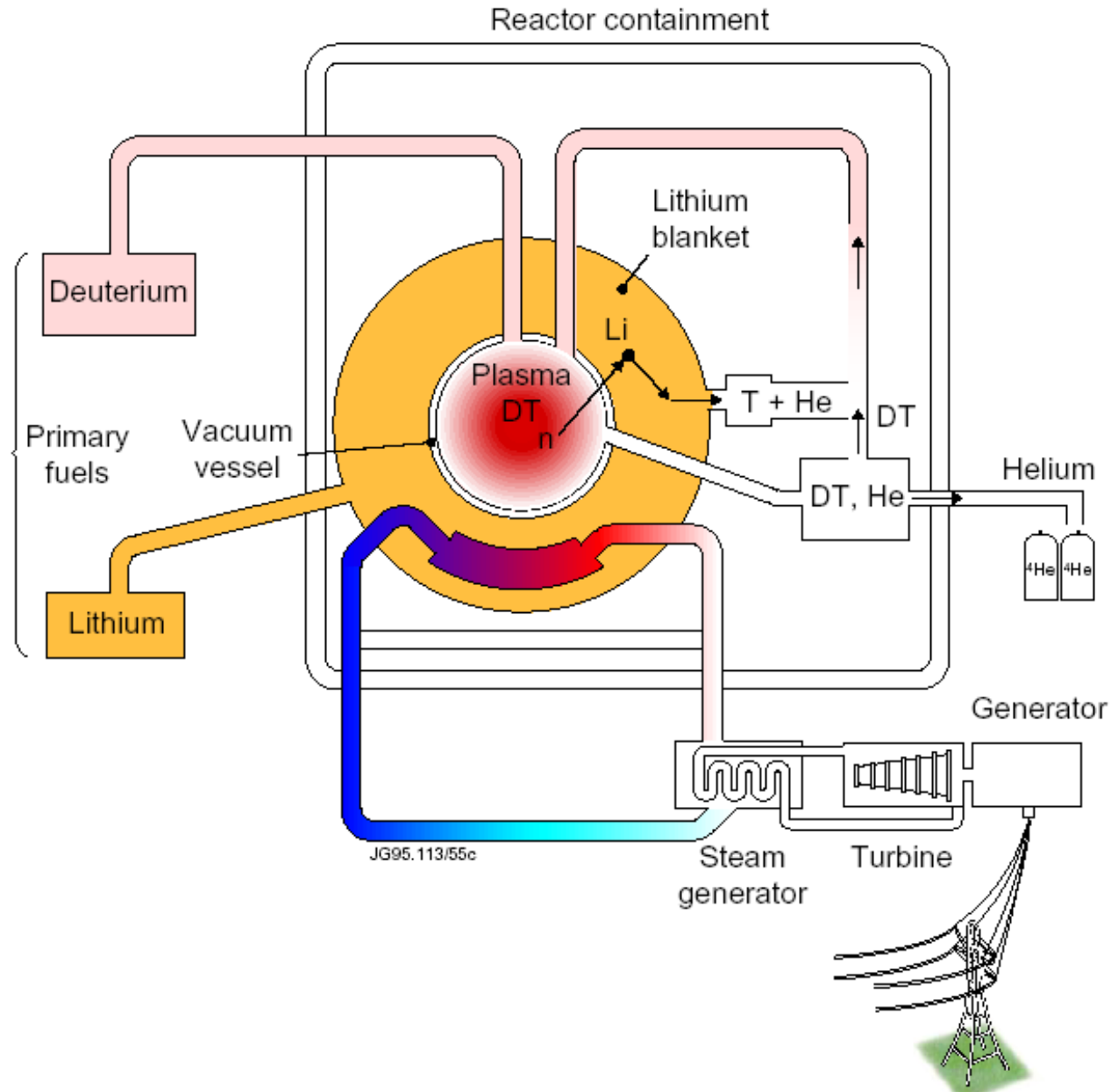
Plasma self-heating



Tritium replenishment

Energy for electricity

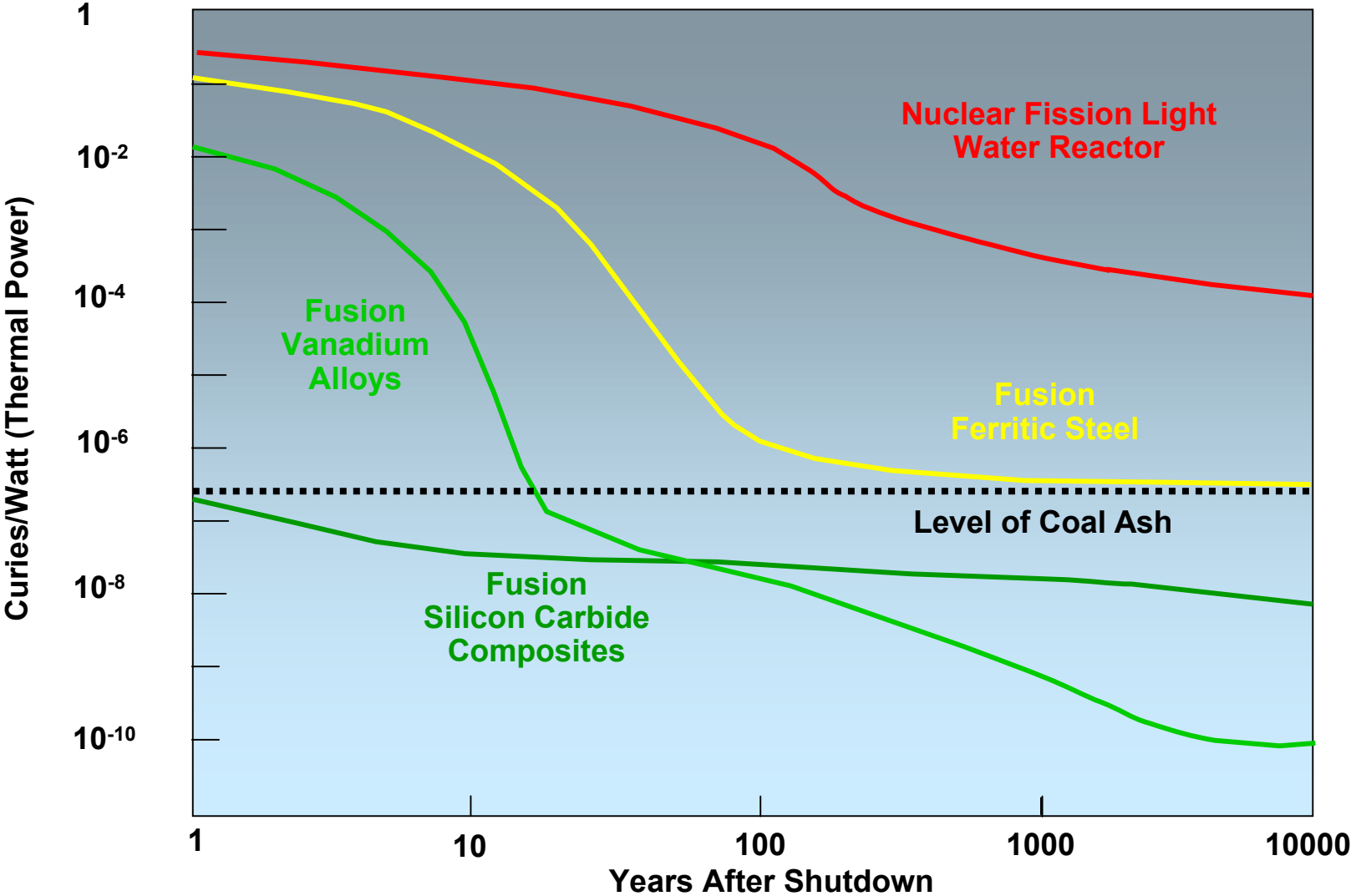
Schematic of a fusion power plant



Advantages of fusion energy

- High energy density fuel
 - 1g D-T \rightarrow 26000kW-hr (1g coal \rightarrow 0.003kW-hr)
- Abundant fuel, available everywhere
 - D is 1/6500 of H (OK for 10^{10} years)
 - Li is 17ppm of crustal rock (OK for 10^3 years)
- Environmental
 - no CO₂ emission
 - no high-level radioactive wastes
- No risk of nuclear accidents (<5min of fuel in reactor)
- No generation of weapons material
- Geographically concentrated, little use of land
- Not subject to local or seasonal variations

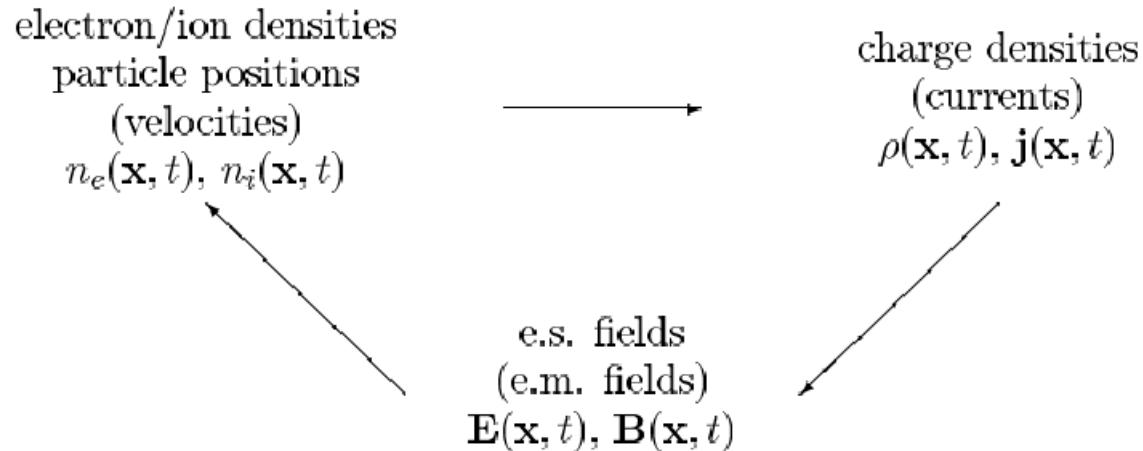
Comparison of activation in fusion and fission reactors



$T \sim 10\text{keV}$: matter is under the form of plasma

Definition of a plasma

- Ensemble of charged particles, globally neutral



- Exhibiting collective behaviour

- Local response to violation of neutrality

Static: screening of charge

Debye length λ_D

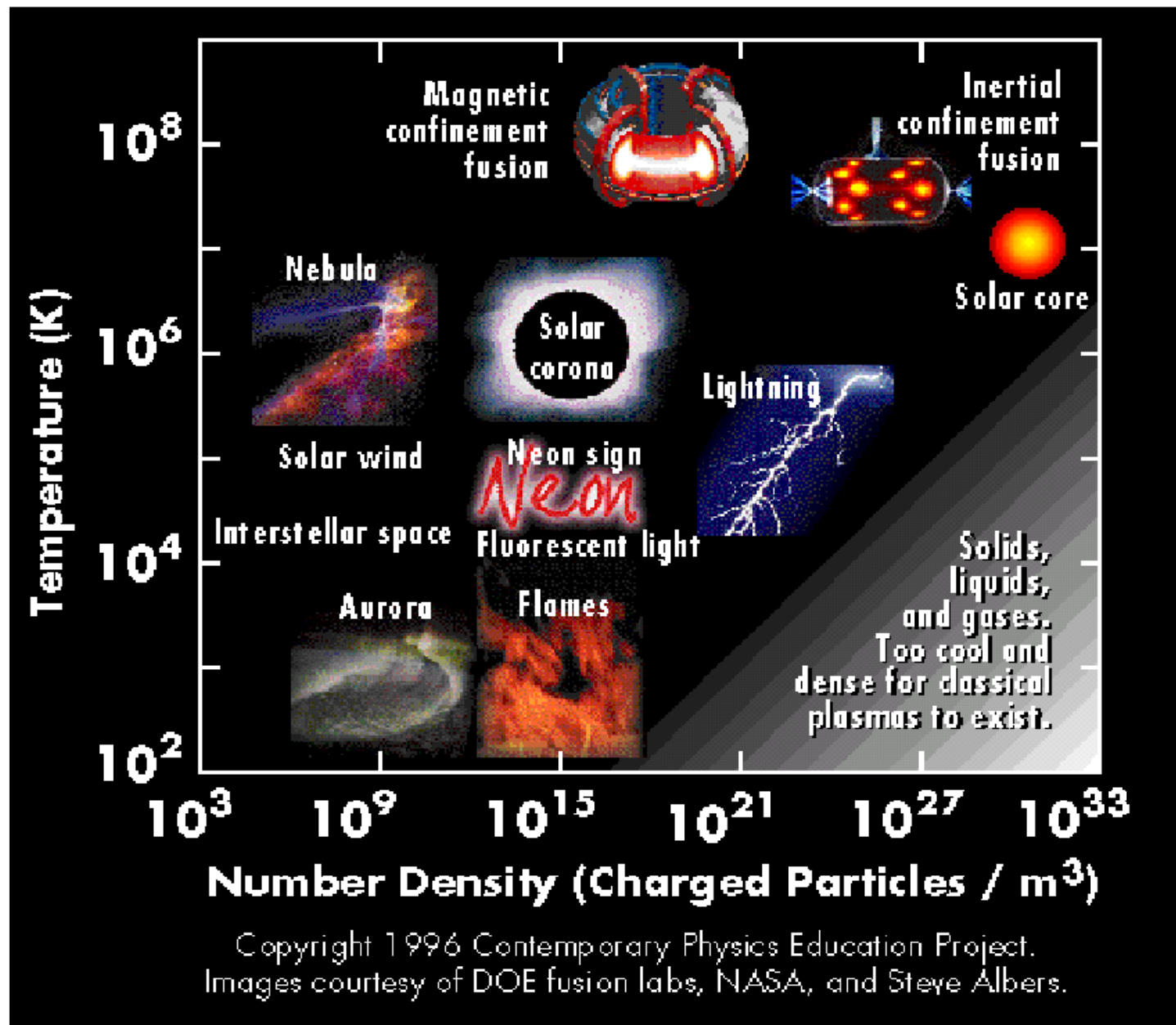
$$\phi(r) = \frac{e}{4\pi\epsilon_0} \frac{1}{r} \underbrace{\exp\left\{-\frac{r}{\lambda_D}\right\}}_{\text{term due to plasma collective interaction}}$$

$$\lambda_D = \sqrt{\frac{T}{n_0 e^2 / \epsilon_0}}$$

Dynamical: plasma oscillations

$$\omega_p = \frac{1}{\tau} = \sqrt{\frac{e^2 n}{\epsilon_0 m}} \quad \text{“plasma frequency”}$$

Examples of plasmas



Copyright 1996 Contemporary Physics Education Project.
Images courtesy of DOE fusion labs, NASA, and Steve Albers.

Why *thermo*-nuclear fusion ?

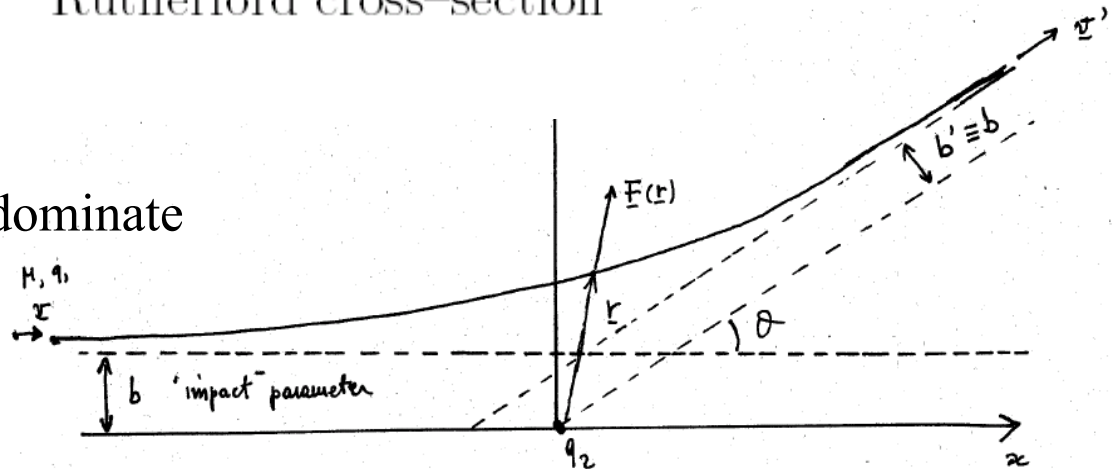
A note on collisions in plasmas

- Short range • Charged particles – neutrals (ionisation)
- inelastic • Nuclear (fusion reactions)
- Long range • Charged particles – charged particles (Coulomb)
 - ~ elastic – Approximate many interactions within the Debye sphere with binary interaction, consider all possible collisions and average

$$\frac{d\sigma}{d\Omega} = \frac{b_{90}^2}{4 \sin^4 \frac{\theta}{2}} \propto \frac{1}{v^4 \sin^4 \frac{\theta}{2}} \quad \text{“Rutherford cross-section”}$$

For small θ : $d\sigma/d\Omega \propto \theta^{-4}$;

→ Small angle collisions dominate



b_{90} : impact parameter for deflection of 90°

Coulomb collisions: cumulative effects and effective collision frequencies

- Consider all scatterers, integrate over b and particle distribution
- Ex. effective collision frequency for energy exchange

$$\bar{\nu}_{E_k}^{e/i} = \frac{1}{E_k} \int f_e(\mathbf{v}_e) d^3v \underbrace{\int_{b_{\min}}^{b_{\max}} \frac{m_e v^2}{2} \frac{m_e}{m_i} \left(\frac{2b_{90}}{b}\right)^2 n v 2\pi b db}_{\int_{b_{\min}}^{b_{\max}} \Delta E_k n v d\sigma = \frac{dE_k}{dt}} = K \frac{n}{T_e^{3/2}}$$

Hotter \rightarrow less collisional

$$K = \frac{2}{3} \sqrt{\frac{2}{\pi}} \frac{Z^2 e^4 m_e^{1/2}}{4\pi \epsilon_0^2 m_i} \ln \Lambda, \quad \Lambda \equiv \int_{b_{\min}}^{b_{\max}} \frac{db}{b}, \quad \text{“Coulomb logarithm”}$$

- We are considering only small angles

$$\implies b_{\min} \simeq \max\{b_{90}, \lambda_{\text{DeBroglie}}\} \quad (7)$$

- Due to the Debye screening effect, outside the Debye sphere the potential is screened, so the Coulomb collisions are no more effective

$$\implies b_{\max} \simeq \lambda_D. \quad (8)$$

Coulomb collisions: plasma resistivity

$$\eta = \frac{\sqrt{2}}{\pi^{3/2}} \frac{m_e^{1/2} Z e^2 \ln \Lambda}{12 \varepsilon_0^2 T_e^{3/2}} \propto T_e^{-3/2}$$

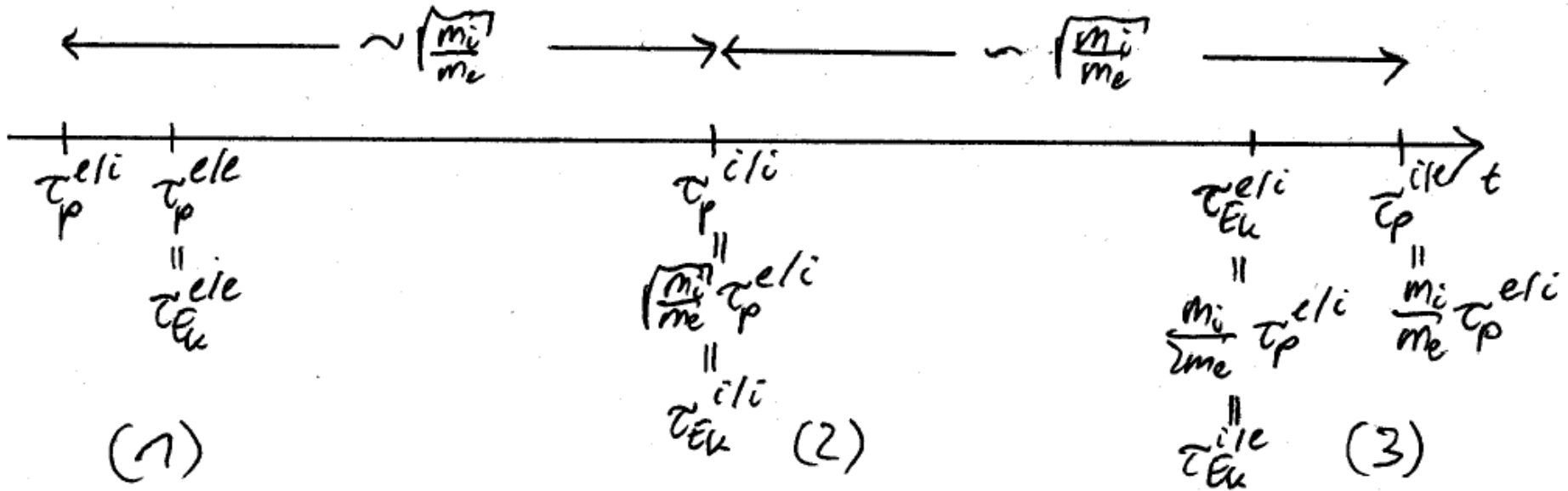
Resistivity - quantitative estimates

1. Plasma at 100 eV: $\eta \sim 6 \cdot 10^{-7} \Omega\text{m}$ [$\sim \eta$ of stainless steel]
2. Plasma at 1 keV: $\eta \sim 2 \cdot 10^{-8} \Omega\text{m}$ [$\sim \eta$ of copper]
3. For $T \gg 1$ keV plasma becomes almost superconducting

Observations The decrease of resistivity with temperature has two important consequences:

1. Magnetic flux is ‘frozen’ within plasma (e.g. solar wind carrying B-field with it)
2. Heating by current (‘ohmic heating’) becomes less and less effective at high T_e .

Coulomb collisions: characteristic time scales



- Ex. H plasma, $T_e = T_i = 10 \text{keV}$; $n = 10^{20} \text{m}^{-3}$

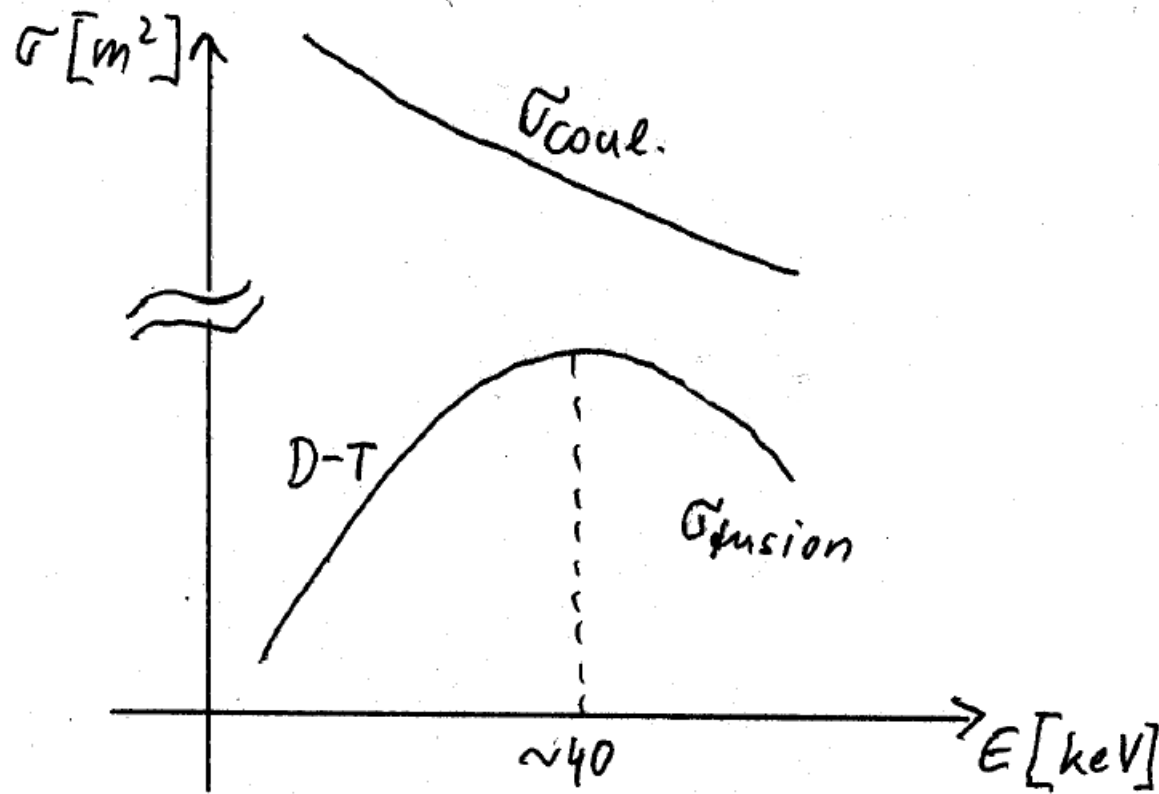
For momentum exchange and equilibrium within one species

$$\tau_{p,Ek}^{e/e} \sim 0.2 \text{ms}; \quad \tau_{p,Ek}^{i/I} \sim 10 \text{ms}$$

For thermal equilibrium between the two species

$$\tau_{Ek}^{e/i} \sim 0.5 \text{s}$$

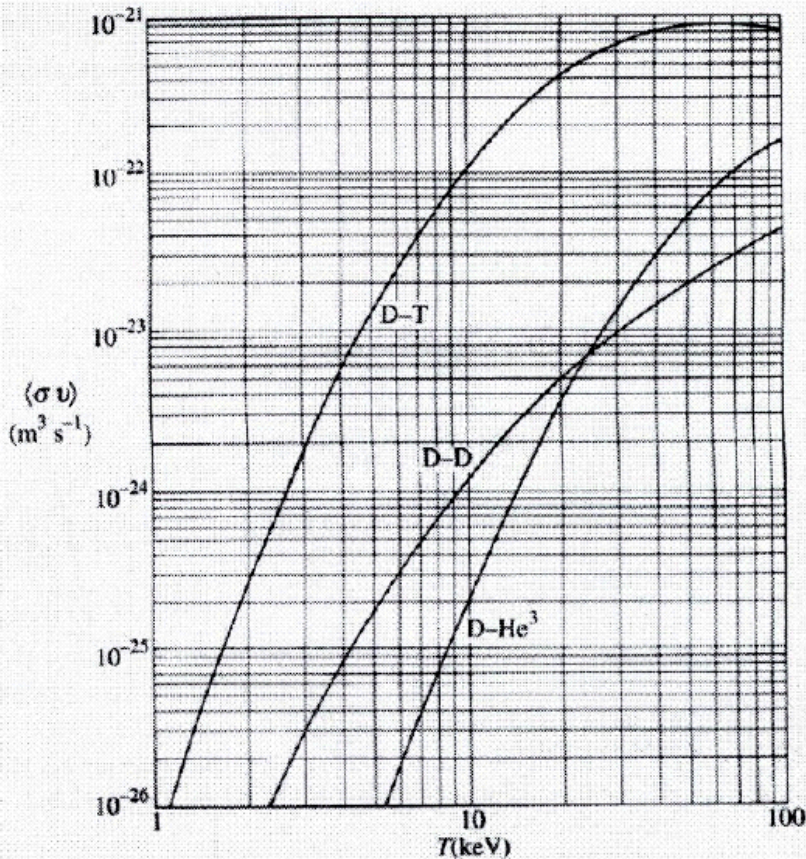
Coulomb collision σ is much larger than fusion σ for all energies



Fusion reactors must deal with 'thermal' plasmas

The Lawson criterion: energy production

- **Fusion power density** $\equiv P_f = n_D n_T \langle \sigma v \rangle E_f$; $E_f = 17.6 \text{ MeV}$
 $= \frac{1}{4} n^2 \langle \sigma v \rangle E_f$ ($n_D = n_T = n/2$)
- **Of this, 20% is in the α 's:** $P_\alpha = P_f / 5$



- $\langle \sigma v \rangle$ is the rate at which fusion reactions take place ('thermonuclear' fusion: we can average over Maxwellian distributions of D and T)
 - Ex.: $n = 5 \times 10^{20} \text{ m}^{-3}$, $\max(\langle \sigma v \rangle)$
 - To have $P_f \sim 1 \text{ GW}$ we need a volume $V \sim 6 \text{ m}^3$
 - If only it was so easy....

The Lawson criterion: losses

- **Losses**

- Thermal energy density

- $W \equiv 3nT$ is lost over characteristic time τ_E : $P_{\text{loss}} = W/\tau_E$

- Bremsstrahlung radiation

- $P_b = A Z^2 n^2 T^{1/2}$ (X-rays)

- Cyclotron emission

- $P_{\text{cycl}} = C n T B^2$ (micro-waves)

- but ~only for electrons, and mostly reabsorbed by plasma either directly or after reflection from metal walls \rightarrow negligible

- **Assumptions**

- Plasma is pure 50:50 D-T

- Efficiency of conversion of thermal energy into electricity $= \eta_1$

- Efficiency of conversion of electricity into plasma heating $= \eta_2$

- Total efficiency $\eta = \eta_1 \eta_2$

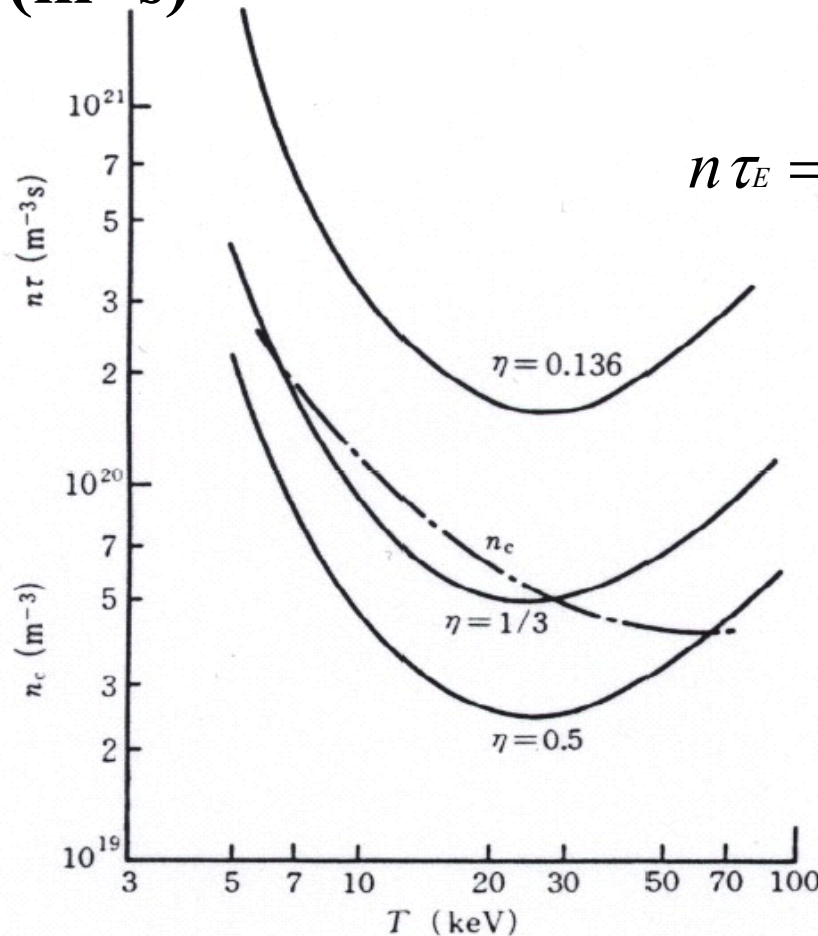
- The power density re-injected in plasma is $\eta(W/\tau_E + P_b + P_f)$

The Lawson criterion: breakeven

- Breakeven: Power reinjected = Losses

$$\eta(W/\tau_E + P_b + P_f) = W/\tau_E + P_b$$

$n\tau$ (m^{-3}s)



$$n\tau_E = \frac{3T}{\frac{\eta}{1-\eta} \frac{\langle\sigma v\rangle}{4} E_f - AT^{1/2}} = fct(T)$$

T (keV)

The Lawson criterion: ignition

- If α 's are confined, external heating is not needed (and bremsstrahlung can be neglected) if $P_\alpha = W/\tau_E$

Ignition condition

$$n \tau_E = \frac{3T}{\frac{\langle \sigma v \rangle}{4} E_\alpha} = fct(T)$$

– Ignition condition \leftrightarrow efficiency $\eta = 1/(1+E_\alpha/E_f^{\text{total}})=0.136$

– Fusion energy gain: $Q \equiv P_{\text{fusion}}/P_{\text{heat}} = 5 P_\alpha/P_{\text{heat}}$

– α heating fraction: $f_\alpha \equiv P_\alpha/(P_\alpha+P_{\text{heat}})=Q/(Q+5)$

$Q=1$

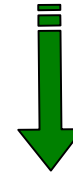
$f_\alpha=17\%$ *breakeven*

$Q=5$

$f_\alpha=50\%$

$Q=\infty$

$f_\alpha=100\%$ *ignition*

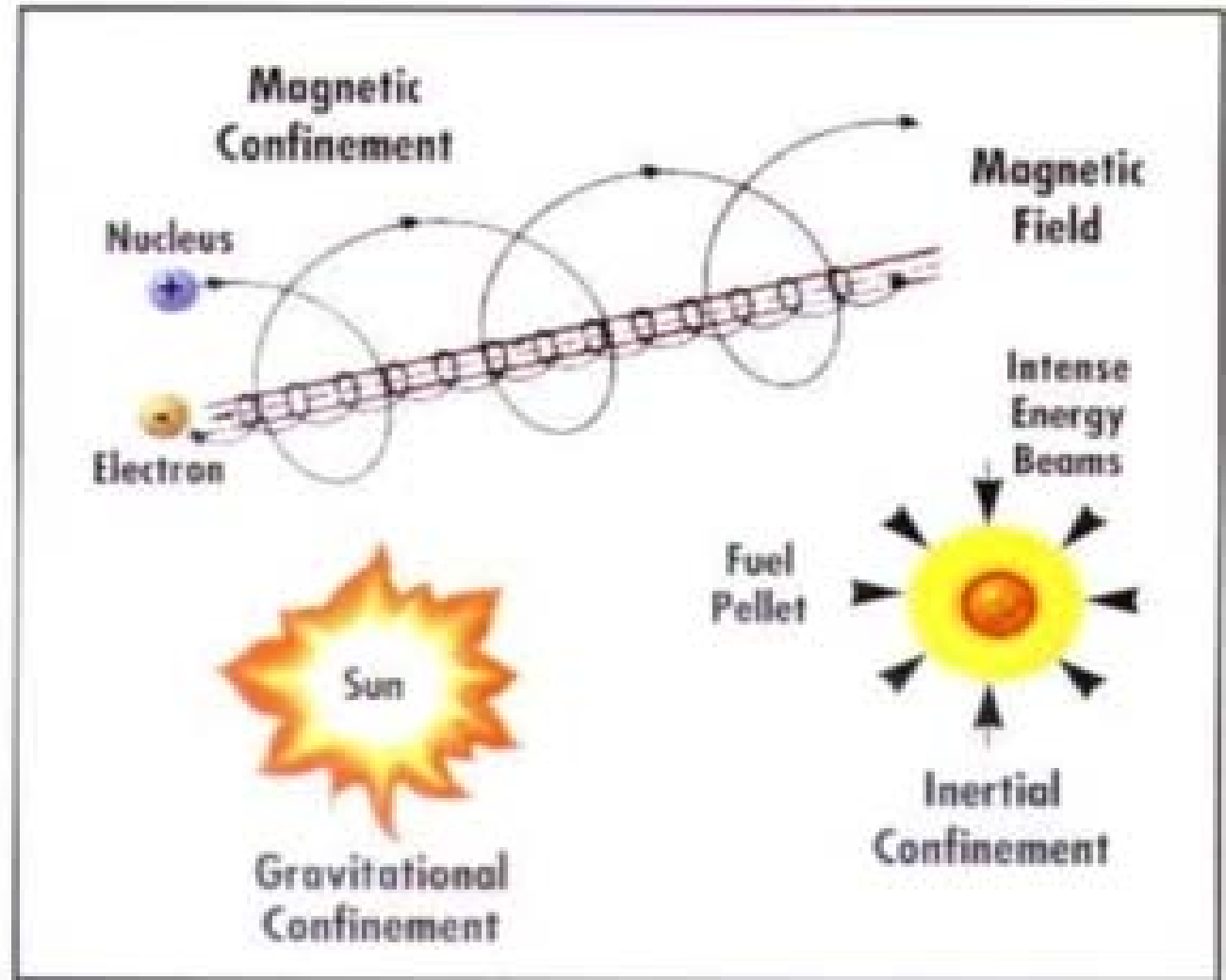


burning plasma

- *Need* $n \tau_E \sim 10^{21} m^{-3}s$ at $T \sim 10keV$
confinement *heating*

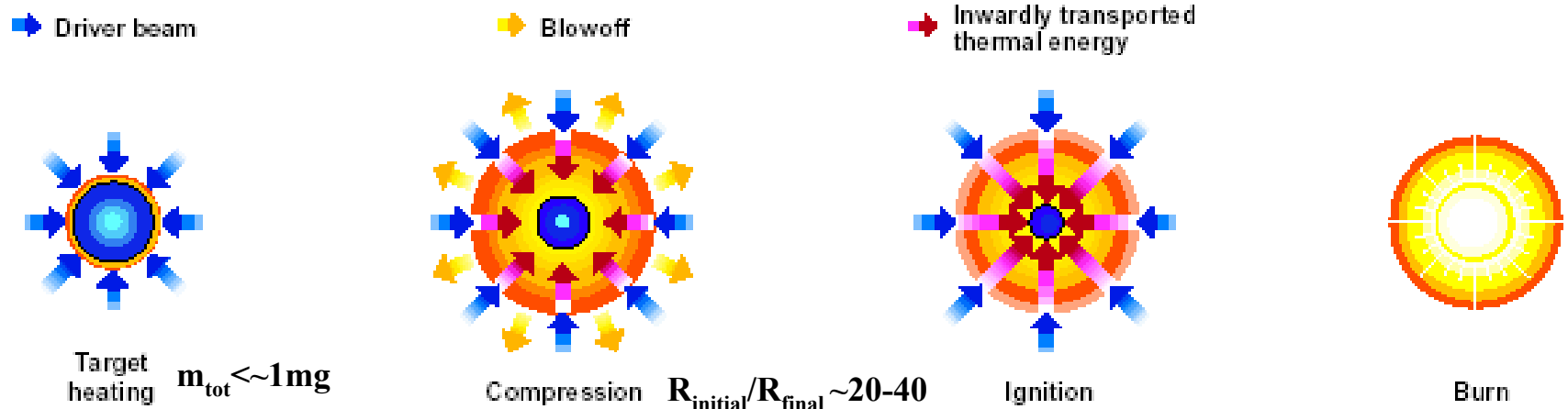
Plasma confinement

- *Magnetic*
 $n \sim 10^{20} \text{ m}^{-3}$
 $\tau_E \sim 10 \text{ s}$
- *Inertial*
 $n \sim 10^{31} \text{ m}^{-3}$
 $\tau_E \sim 10^{-10} \text{ s}$
- *Need to confine and heat the plasma*



Inertial Confinement Fusion

– A capsule with D-T is irradiated by lasers, X-rays, or particle beam



– Compression: need $\sim 10^{12}$ bar to reach 10^{31}m^{-3}

- Laser with $10^{16}\text{W/cm}^2 \rightarrow p_{\text{light}} \sim 10^6$ bar, largely insufficient
- Shock waves at the pellet surface, arriving at the center at the same time
- Once fusion starts, α heating sustains the reactions

– Heating to ignition must occur before ions fly away

- Energy flux F : $\tau_{\text{heat}} = U_{\text{th}} / (4\pi R^2 F)$; $U_{\text{th}} = 3nT(4/3\pi R^3)$

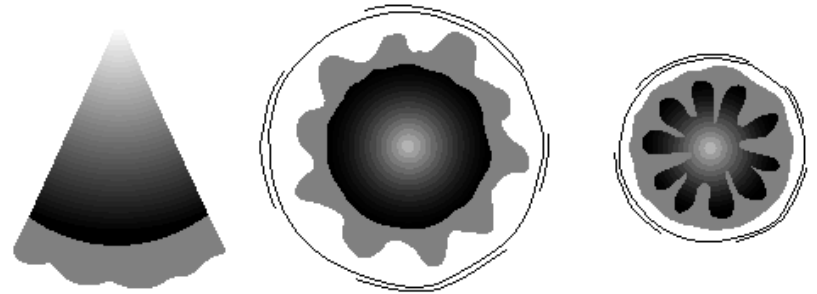
$$\tau_{\text{heat}} < \text{inertial time} = R/v_{\text{sound}} \quad (\sim 100\text{ps})$$

$$\rightarrow F > nT^{3/2}/m_i^{1/2} \sim 5 \times 10^{15} \text{ W/cm}^2$$



ICF: general issues

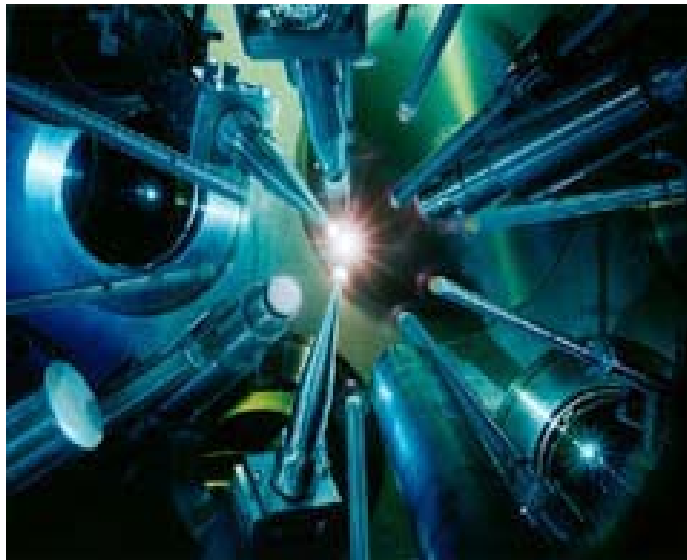
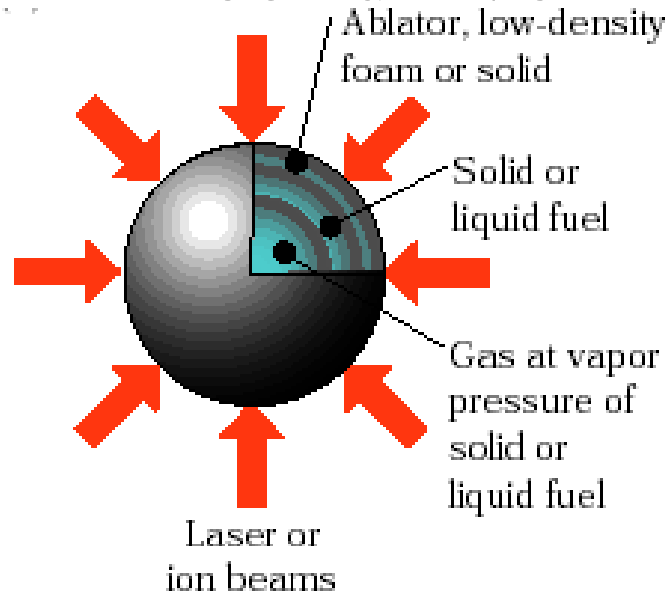
- Uniformity, stability of compression
 - Rayleigh-Taylor hydro-dynamic instability
 - Low density vaporised shell pushes high density D-T ice layer
 - Magnifies surface irregularities and may prevent ignition



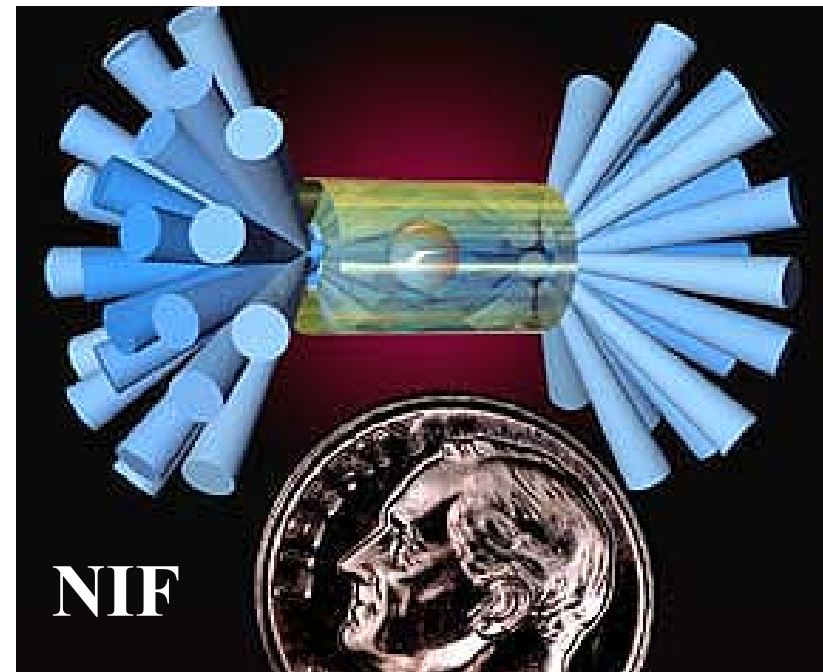
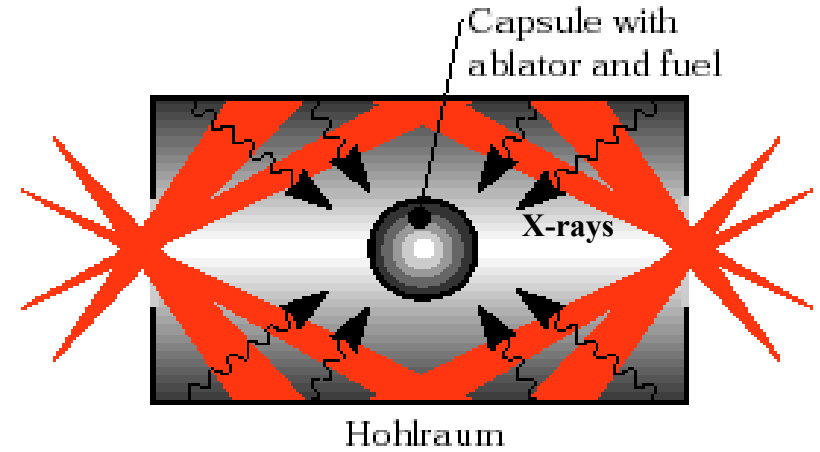
- Capsule design
- Efficiency of drivers
- Steady-state: extension of techniques from single pulse to many repetitive pulses for energy production
- Materials for first wall
 - Long lifetime, low induced radio-activity, ...
- Optimisation to reduce cost and increase efficiency

ICF: direct and indirect drive

Direct drive



Indirect drive

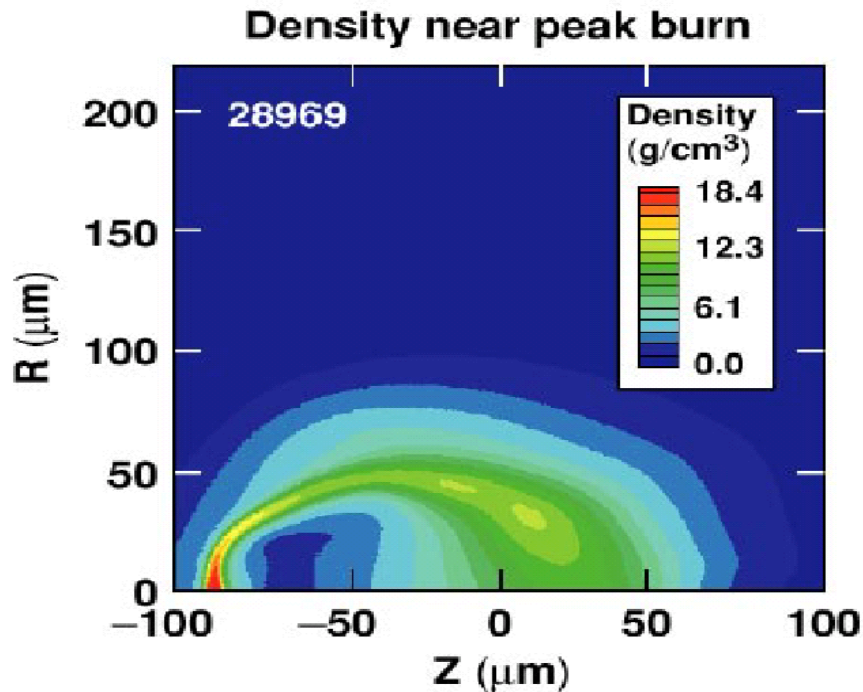


ICF – direct drive: ex. of results

Initial 2D hydrodynamic simulations show good agreement with experimental $\alpha=4$ cryogenic target results



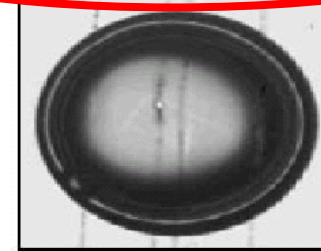
DRACO code simulated density contours



$\alpha = 4$ pulse, 17 kJ

100-μm thick ice layer

8-μm rms ice roughness

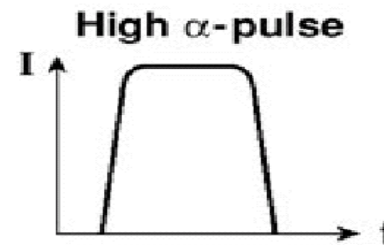


	Expt	1-D	2-D
Y_{1n}	5.95×10^9	5.60×10^{10}	5.32×10^9
Y_2	6.75×10^7	6.94×10^8	6.31×10^7
$\langle \rho R \rangle$	67	80.0	58
T_{ion}	2.5	1.7	2.0

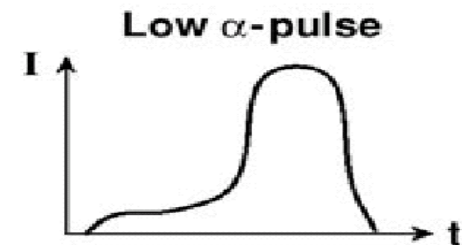
T.C.Sangster et al. Phys. Plasmas 10, 1937 (2003)

$$\alpha = P_{\text{fuel}}/P_{\text{Fermi}}$$

P_{Fermi} = Fermi degenerate pressure



Good stability



High gain

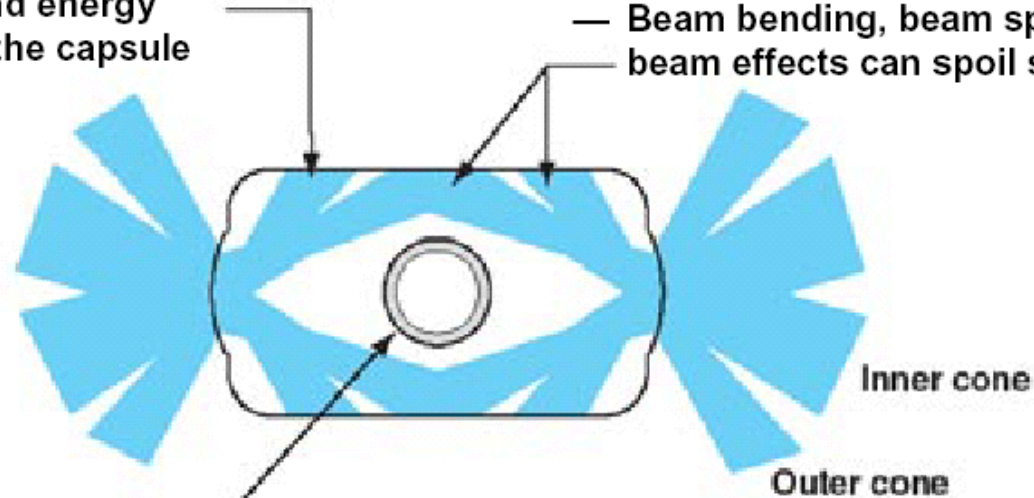
ICF – indirect drive: issues

Hohlraum Energetics

- Backscatter reduces laser absorption and energy coupled into the capsule

Indirect Drive Symmetry

- Unequal absorption between cones can spoil beam balance
- Beam bending, beam spray and cross beam effects can spoil symmetry

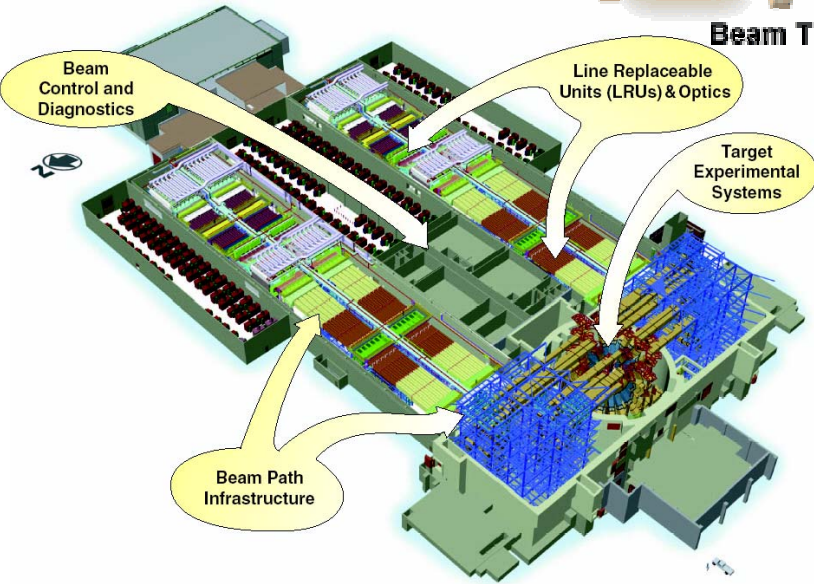
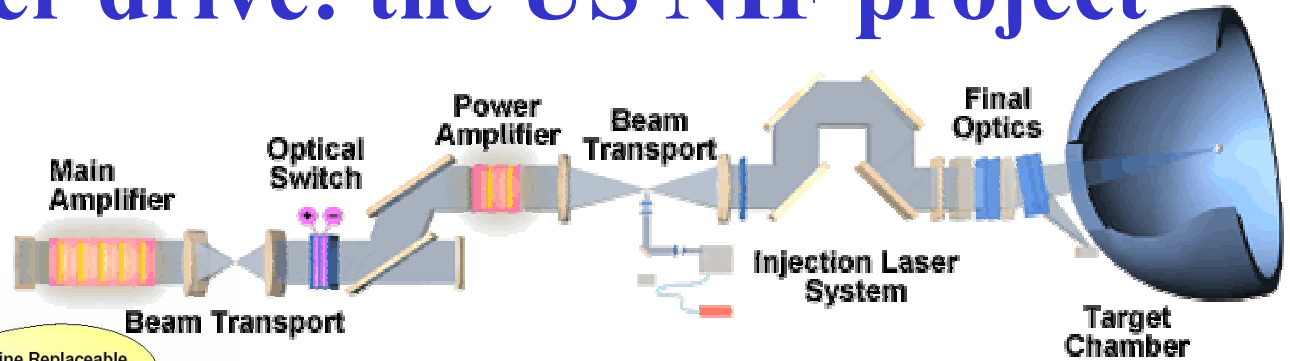


Implosion/compression to high density

- Time dependent scatter can spoil pulse shape
- Hot electrons from plasma waves can preheat the fuel making it less compressible

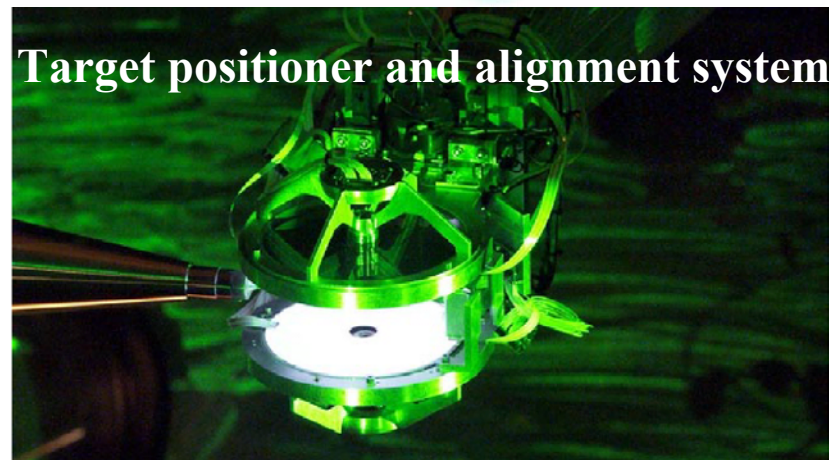
- Ex. of avenue for optimisation: higher hohlraum $T \rightarrow$ higher p , implosion velocity, compression \rightarrow ignition with less energy

ICF – laser drive: the US NIF project



~2009: 192 beams, ~1.8MJ, 500TW UV light ($0.35\mu\text{m}$); fusion yield~20MJ

Target chamber

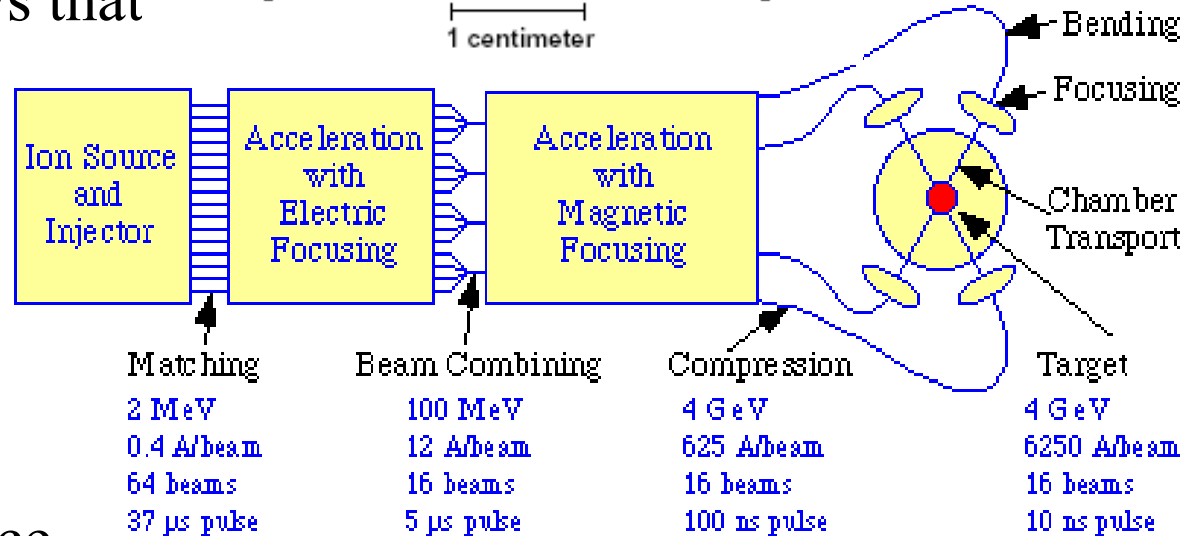
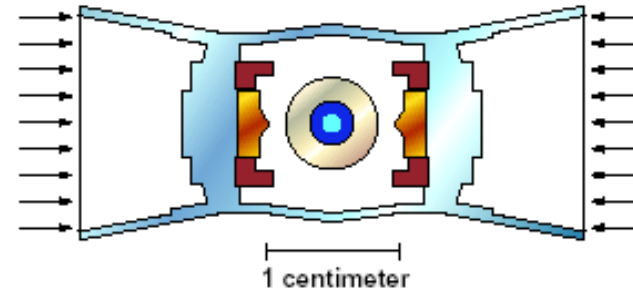


ICF – laser drive developments

- Today's laser drivers are limited for power plant use by
 - Efficiency (should be $>5\%$)
 - Repetition rate (should be $>0.1\text{Hz}$)
 - Damage to injection windows/mirrors by heat, n, debris
- New laser developments
 - Diode-Pumped Solid-State Laser Driver
 - Diodes instead of flashlamps to pump a solid-state laser could permit rapidly repeated firings, efficiency needed for power generation
 - Ex. Mercury laser at LLNL should reach: 100J, 10Hz, 10% efficiency, 3ns
 - Krypton-Fluoride Gas Laser Driver
 - With the KrF laser ($0.248\mu\text{m}$), the laser medium is a gas that can be circulated for heat removal to achieve high repetition rate
- Other driver methods
 - **Ion beams (indirect drive)**
 - **High rep rate z-pinch**

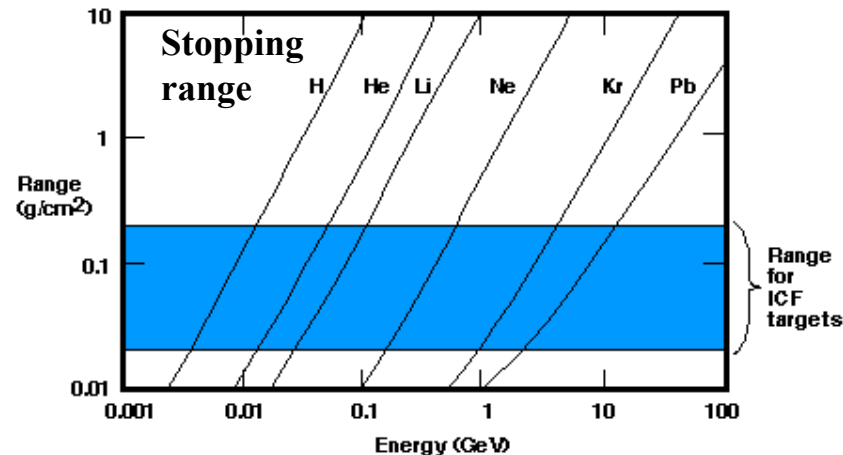
ICF – ion beam indirect drive

- Heavy ion beams
 - Ions hit target, energy gets converted into X-rays that compress pellet
 - Ex. Cs ions, 400 TW



Challenges

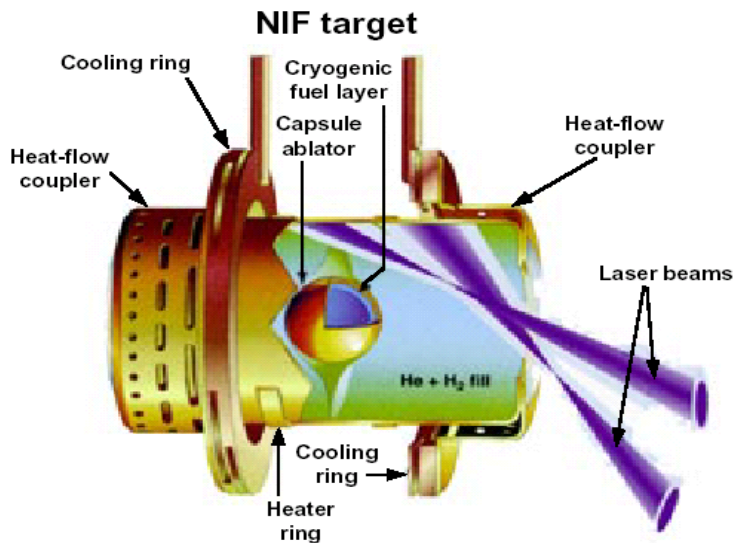
- Beam transport, space-charge, emittance
- Pulse compression
- Focus and deposition depth (light vs. heavy ions)
- Cost, but one accelerator could drive many target chambers



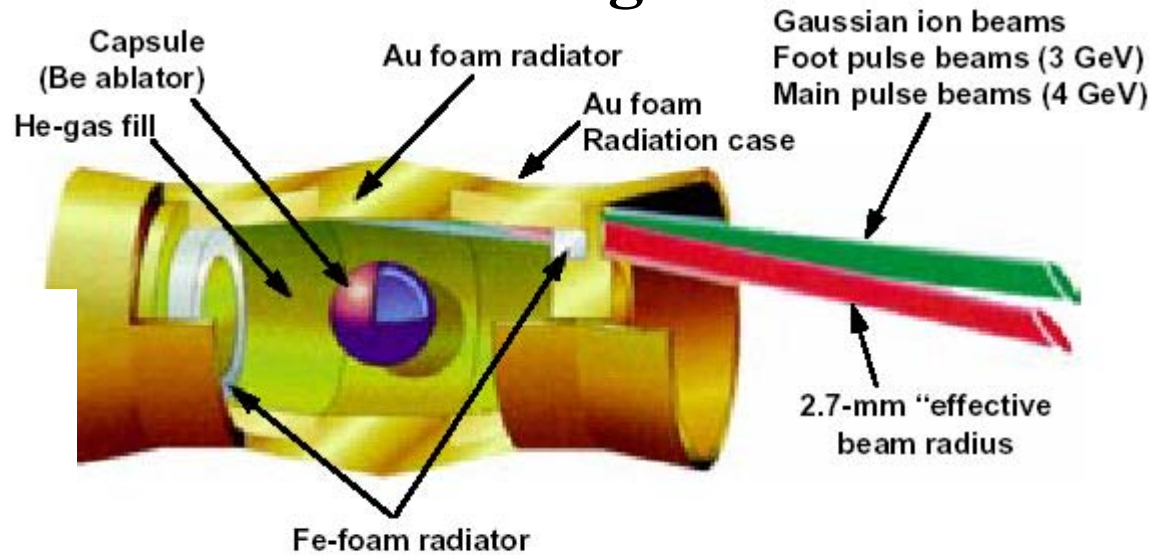
ICF – heavy ion beam target

- Similar issues to laser drive
 - Stability, ignition and burn propagation, symmetry control

Laser target



Ion beam target



Yield = 400 MJ

Driver (using NIF-like hohlraum to capsule radius ratio)

6 MJ of 4 GeV Pb ions \square gain 67

7.5 MJ of 8 GeV Pb ions \square gain 53

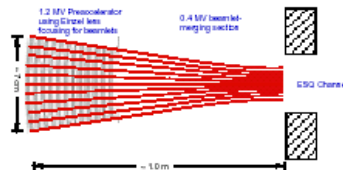
ICF – heavy ion driver: future developments

The heavy ion fusion program plans consists of distinct experiments on ion sources, beam transport, and focusing to be followed by an integrated beam experiment



NOW (next three years)

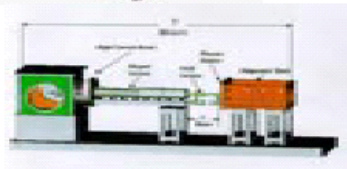
Brighter sources/
injector



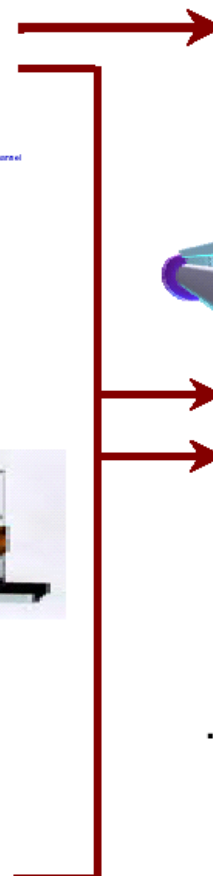
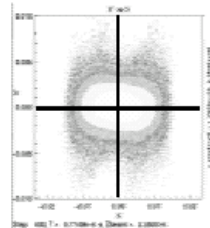
Maximum $\langle J \rangle$, B_n
Transport



Beam neutralization \rightarrow
min ϵ -limited focus r_f

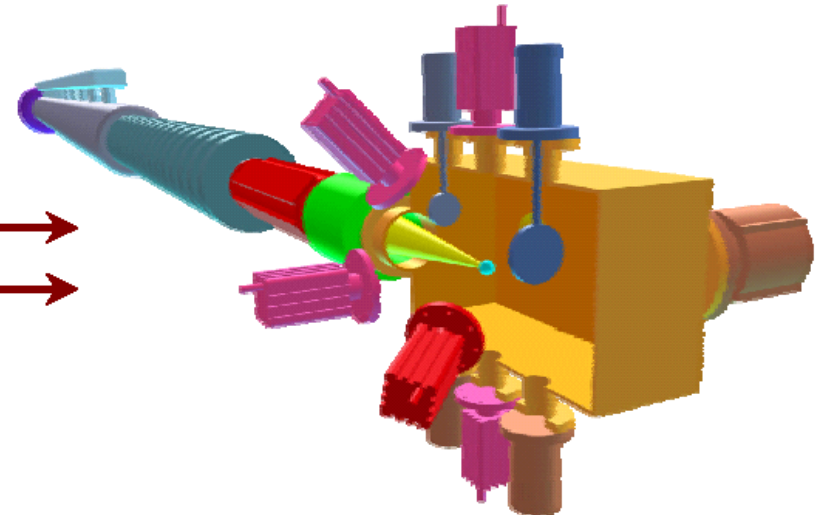


Theory/simulations



NEXT STEP

integrated beam experiments
(IBX)



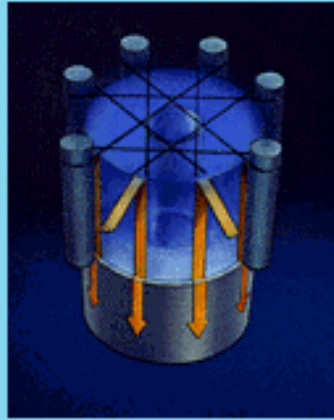
...to test source-to-target-integrated modeling

(Injection, acceleration, longitudinal
compression and final focus)

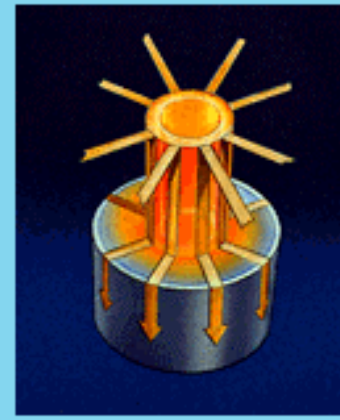
ICF – drive by X-rays from z-pinch

- Ex. Sandia NL
 - 360 tungsten wires (~1/10 of human hair) collapse, evaporate, form a plasma in high current pulse
 - Plasma emits X-rays:
 - $T \sim 150\text{eV}$
 - $E \sim 2.0\text{MJ}$
 - $P \sim 100\text{TW}$
 - Next generation: ~16MJ ?

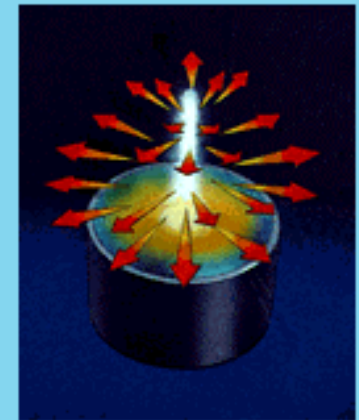
Z machine



Initiation

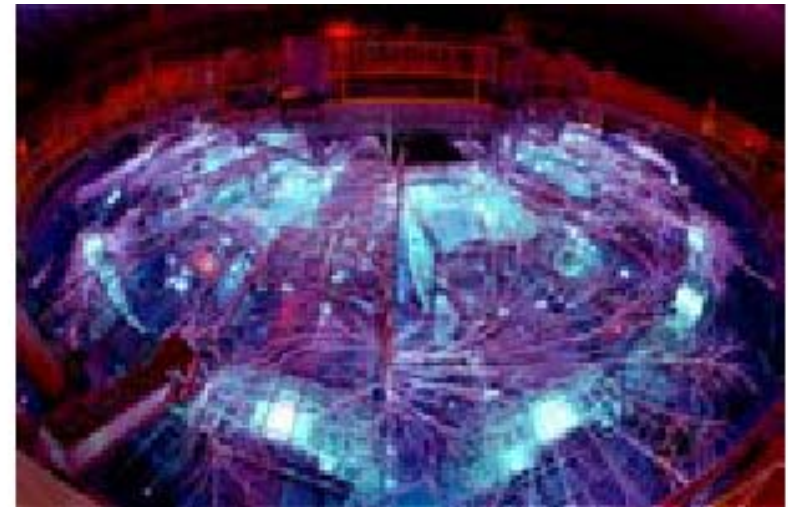


Implosion



Emission

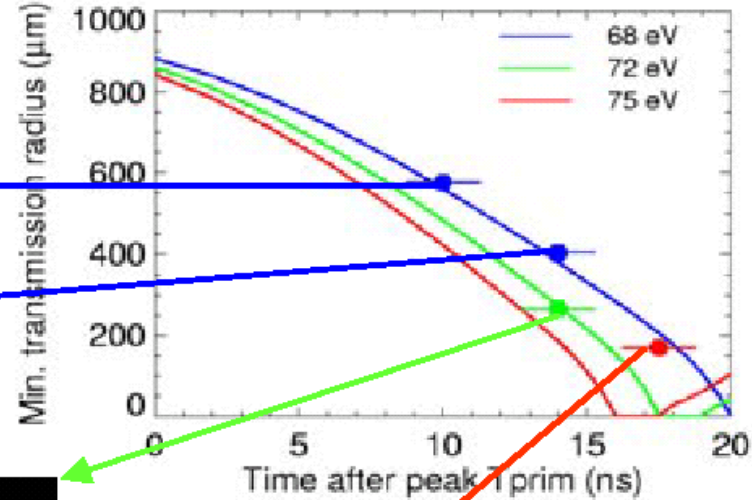
**X-rays-
Z –machine
(Sandia)**



ICF – drive by X-rays from z-pinch

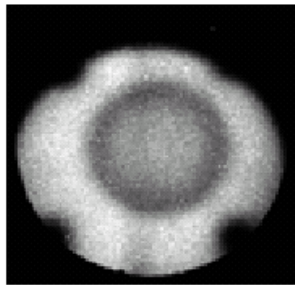
Progress in the symmetry of implosion of targets driven by indirect drive from double z-pinch

Radius vs. time

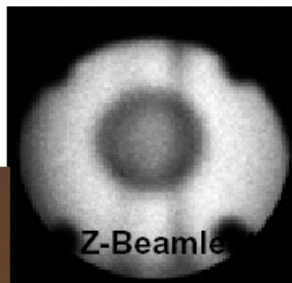


68-75 eV peak drive
~10 kJ absorbed

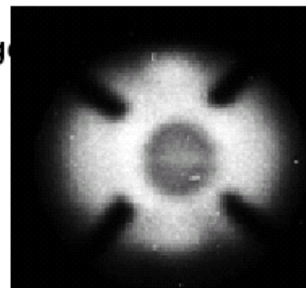
Z830



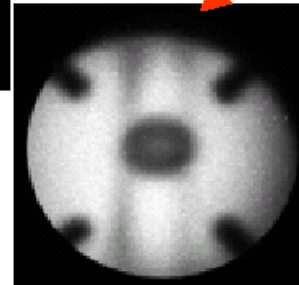
Z831



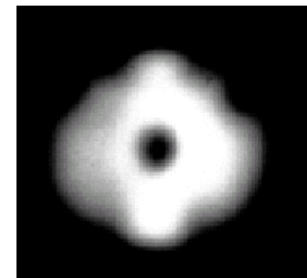
Z839



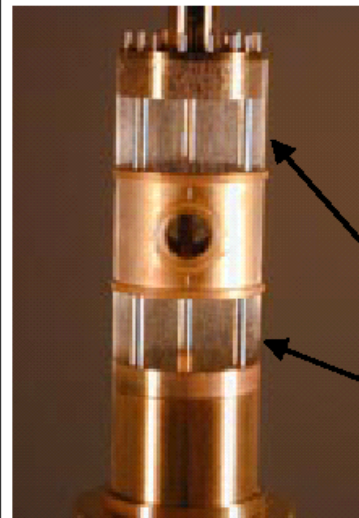
Z833



Peak density ~40 g/cc
CR >14



Z-pinch Wire arrays

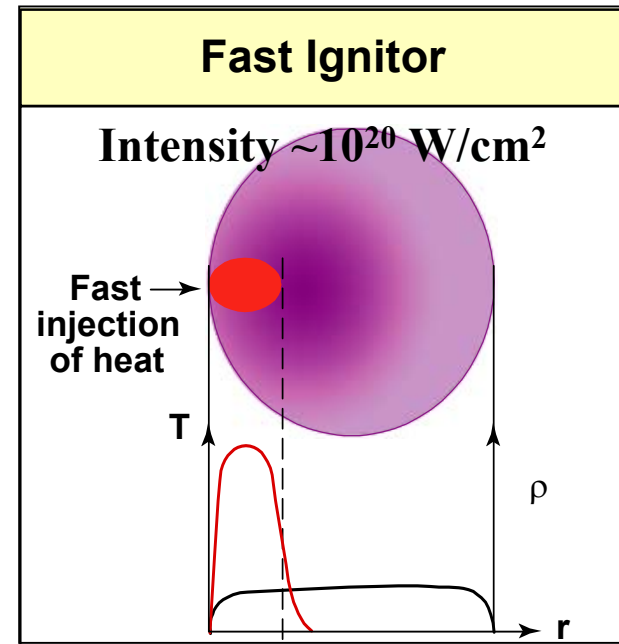
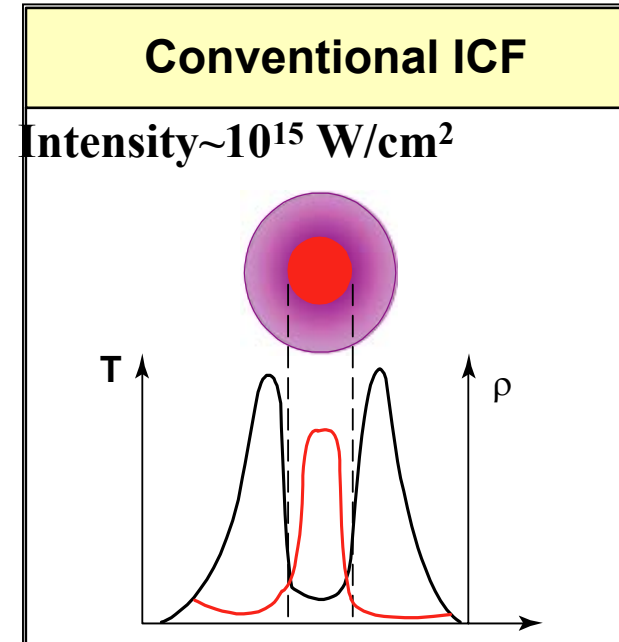


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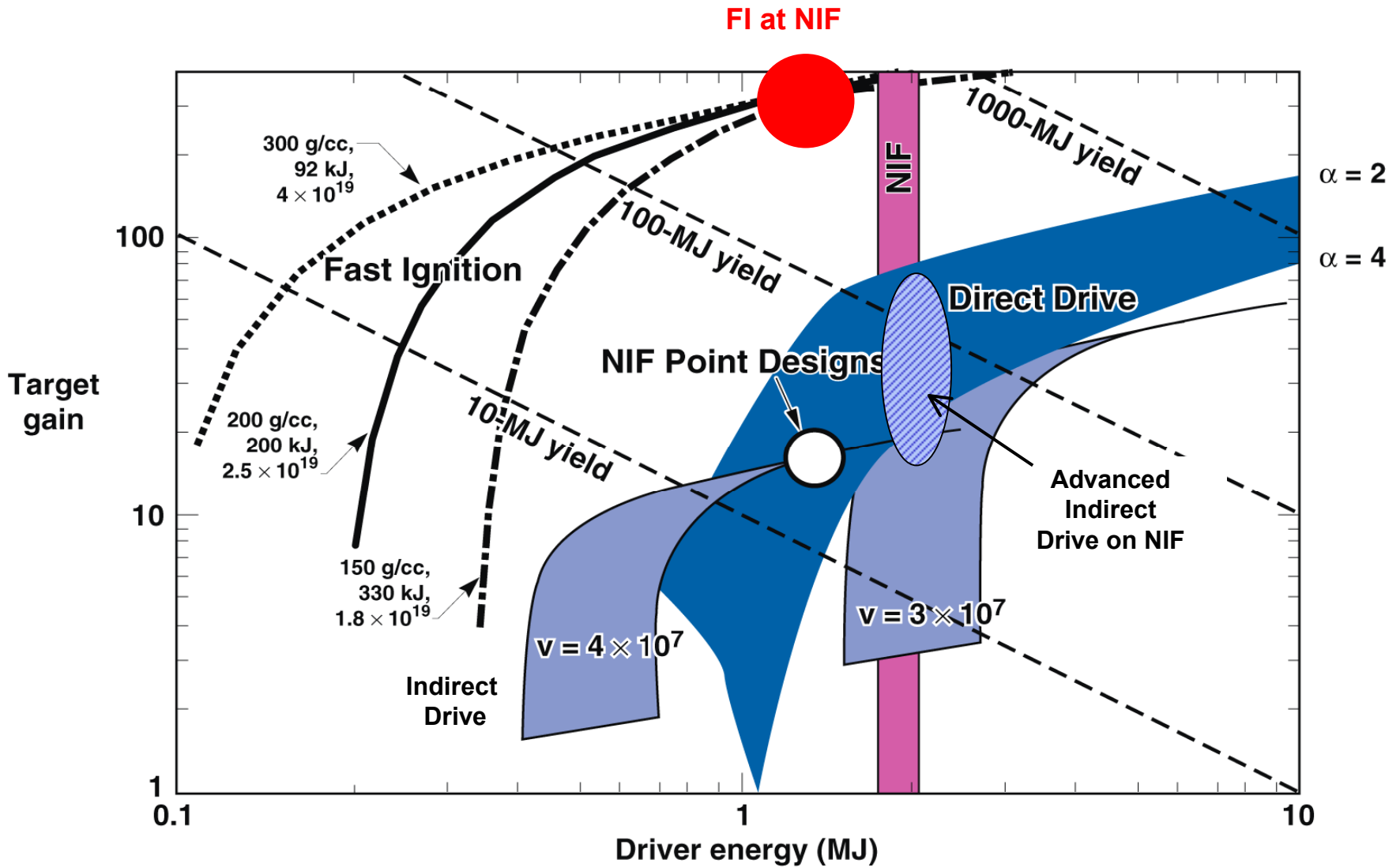
Time

ICF: fast ignition

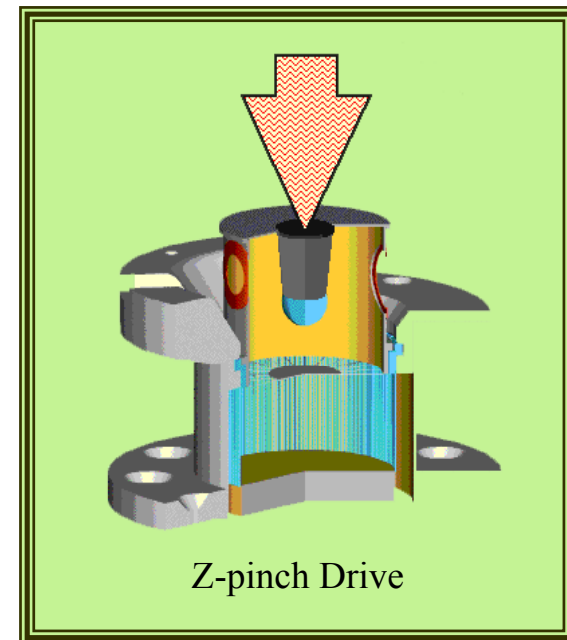
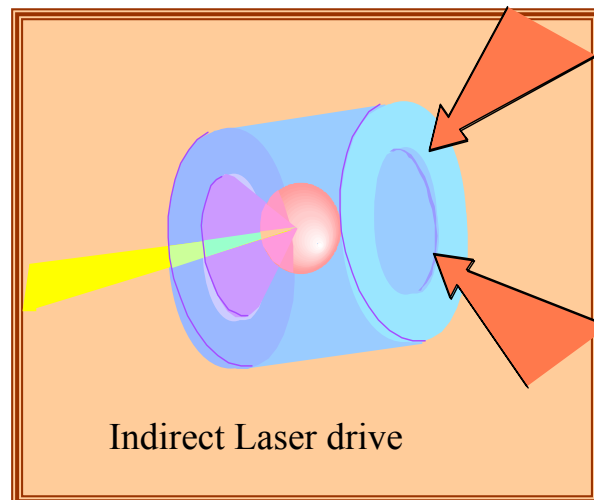
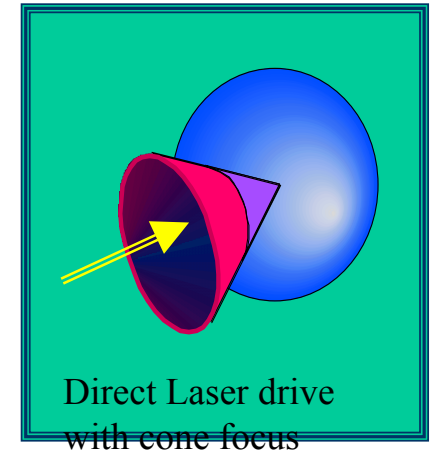
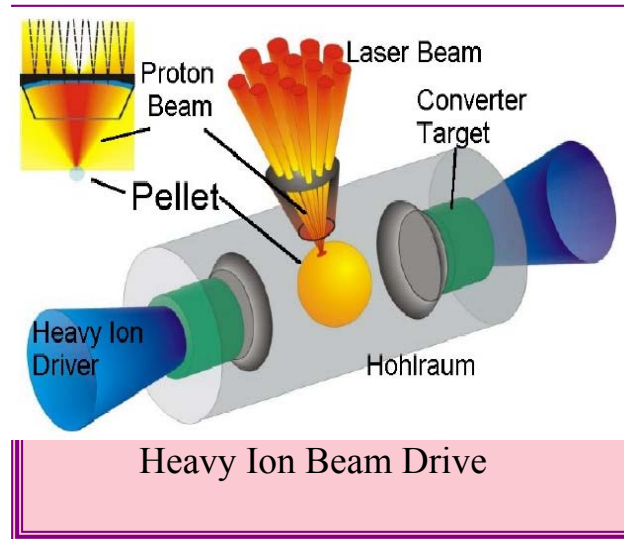
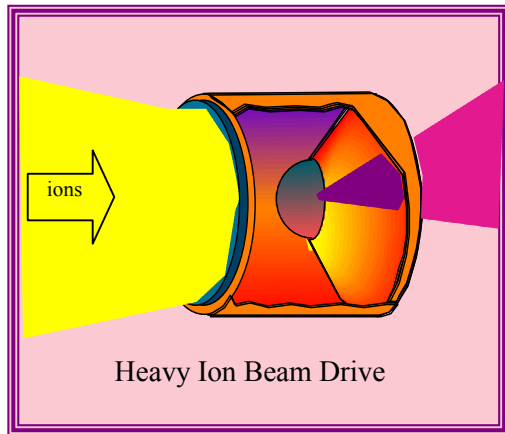
- Conventional ICF
 - D-T gas compressed by imploding solid D-T must form fusion ‘hot spot’, igniting and generating symmetrically propagating burn
 - Even if implosion is uniform, if hot spot is not symmetric, it squirts out, mixes with colder D-T and burn is prevented
- Fast ignition
 - Idea: decouple compression and ignition
 - No need for hot spot: at max compression, a very short ($<10^{-11}$ s) power pulse is injected on the side
 - *lower energy, inexpensive drivers could be used for simpler task of compressing fusion fuel if no need for a hot spot*



ICF: fast ignition may give higher gain



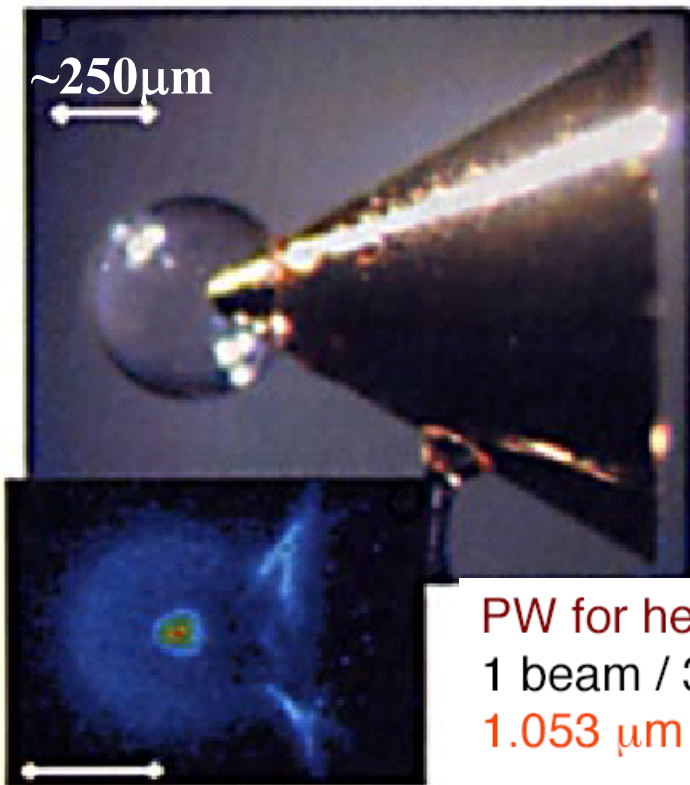
Fast ignition is compatible with all drivers



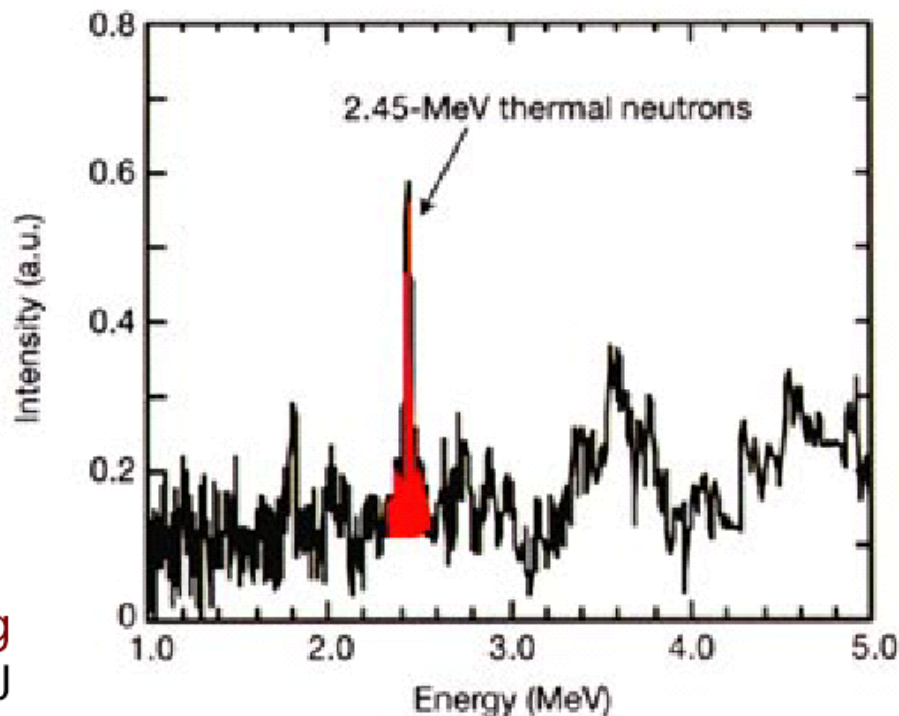
ICF – fast ignition: promising results?

Enhanced neutron output from fast heating of deuterated direct drive shell implosion on Gekko XIII laser (Japan, UK) R.

Kodama, et al., Nature 412, 798 (2001)



PW for heating
1 beam / 300 J
 $1.053 \mu\text{m}$ / 0.5ps



1.2 KJ compression pulse + 60 J, 100 tw fast heating pulse

Green lasers, $t = 1\text{-}2 \text{ nsec.}$
 $\sim 10^{14} \text{ W/cm}^2$

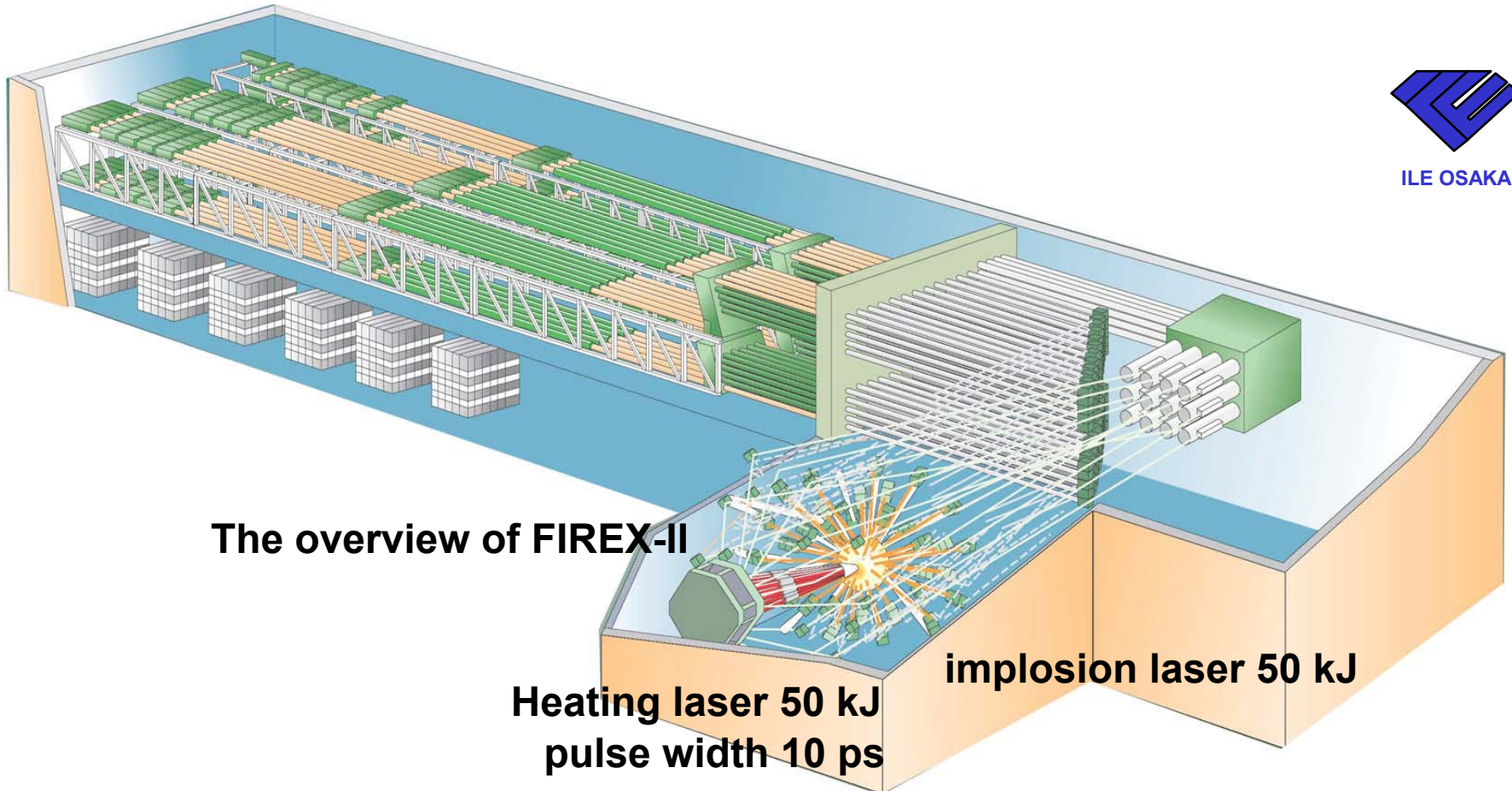
$1 \mu\text{m}$ laser, $t=0.5\text{ps}$
 $\sim 10^{19} \text{ W/cm}^2$

ICF – FI medium term development in Japan

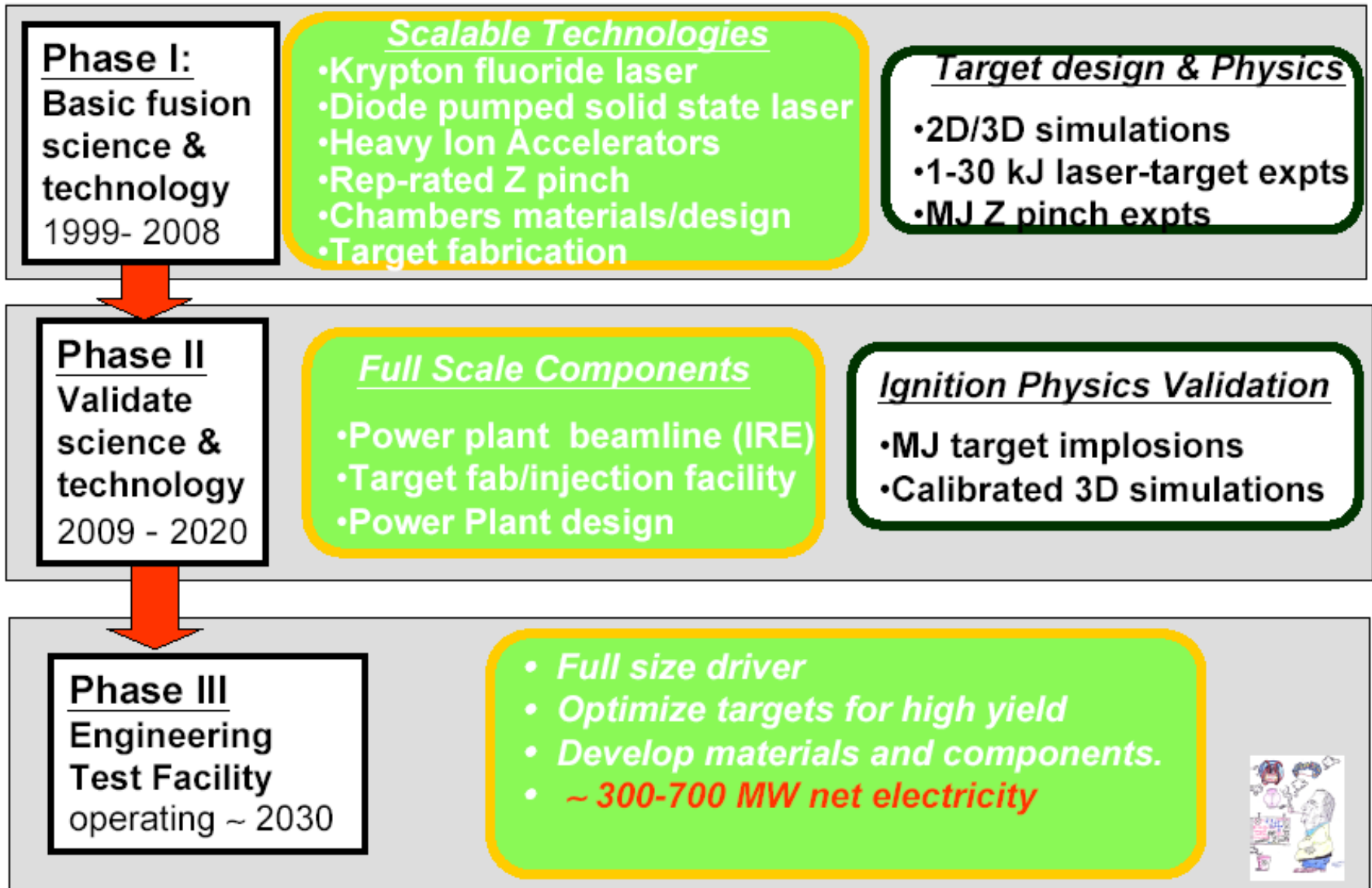
FIREX (Fast Ignition Realization Experiment)

Purpose: Establishment of fast ignition physics and ignition demonstration

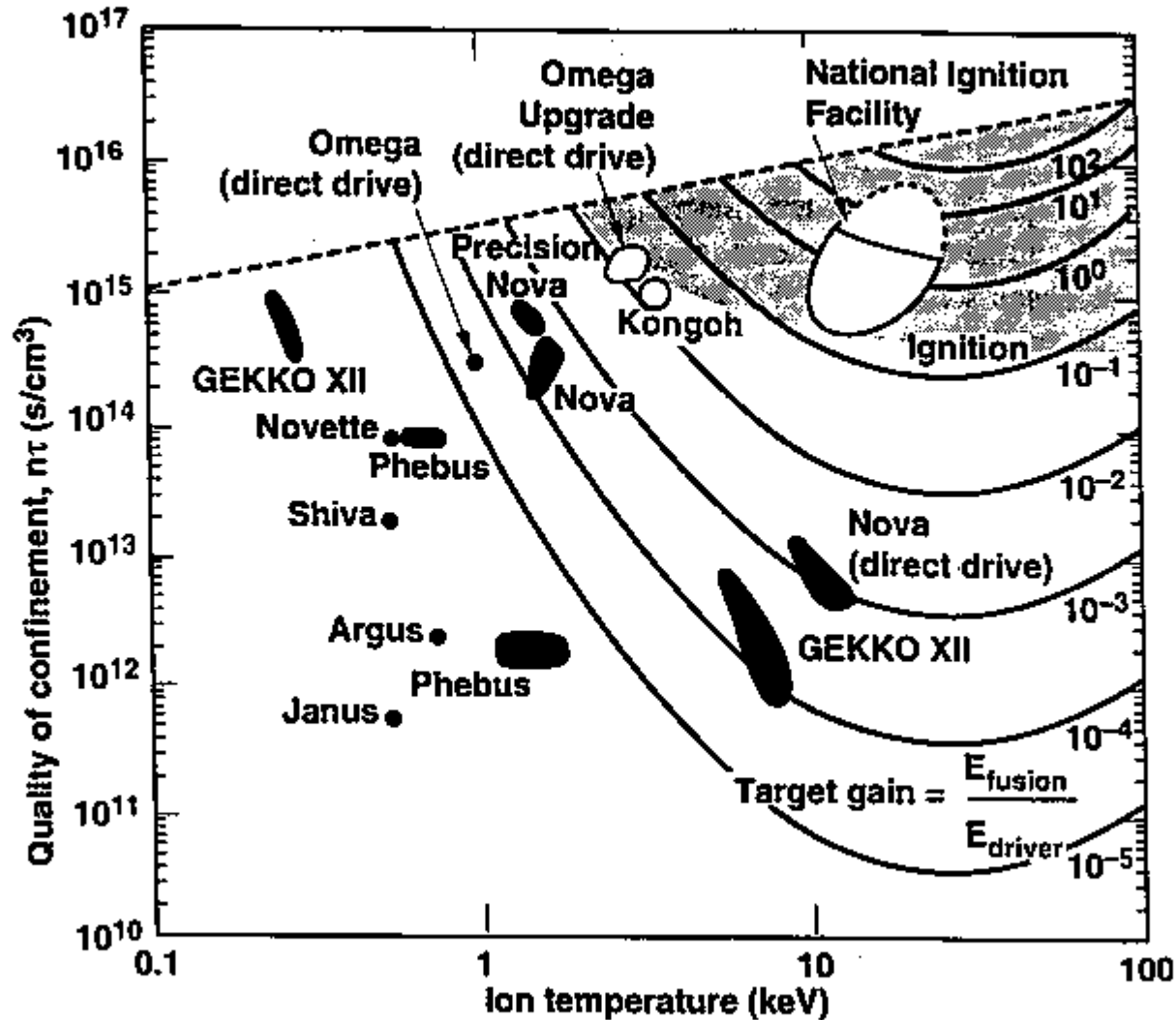
Starting Conditions: high density compression (already achieved)
heating by PW laser (1keV already achieved)



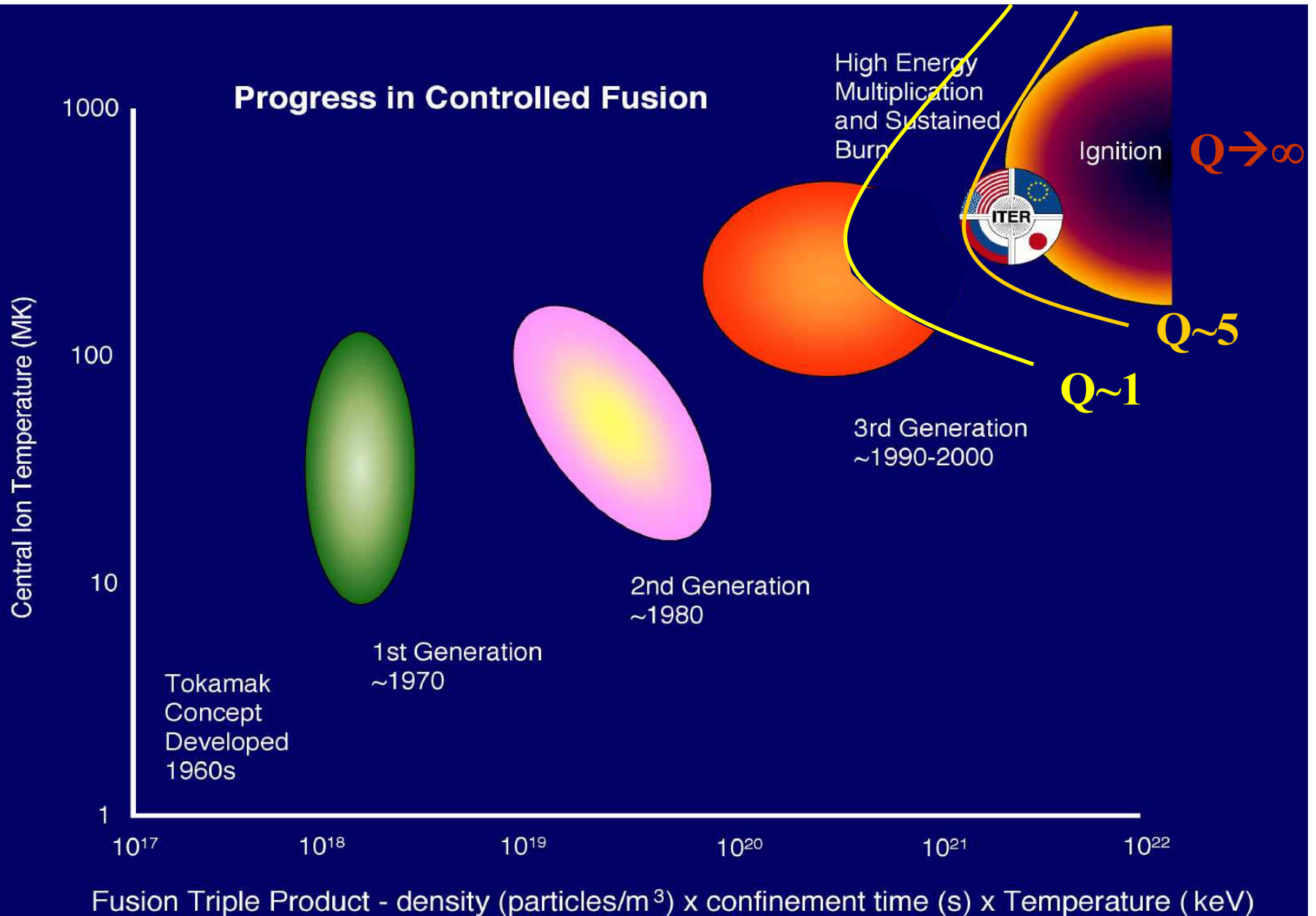
Foreseen ICF development steps (US view)



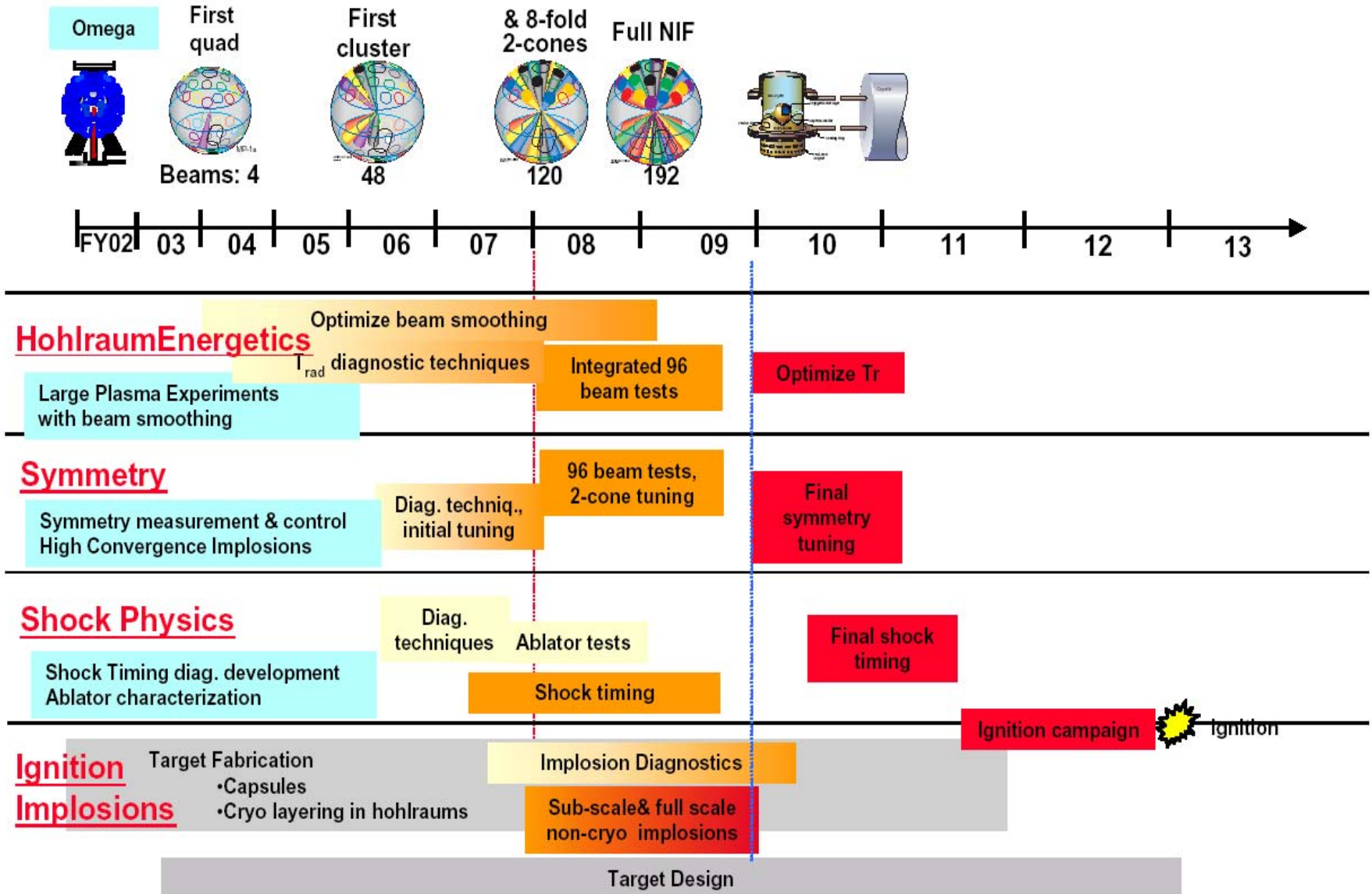
Summary of progress in inertial fusion



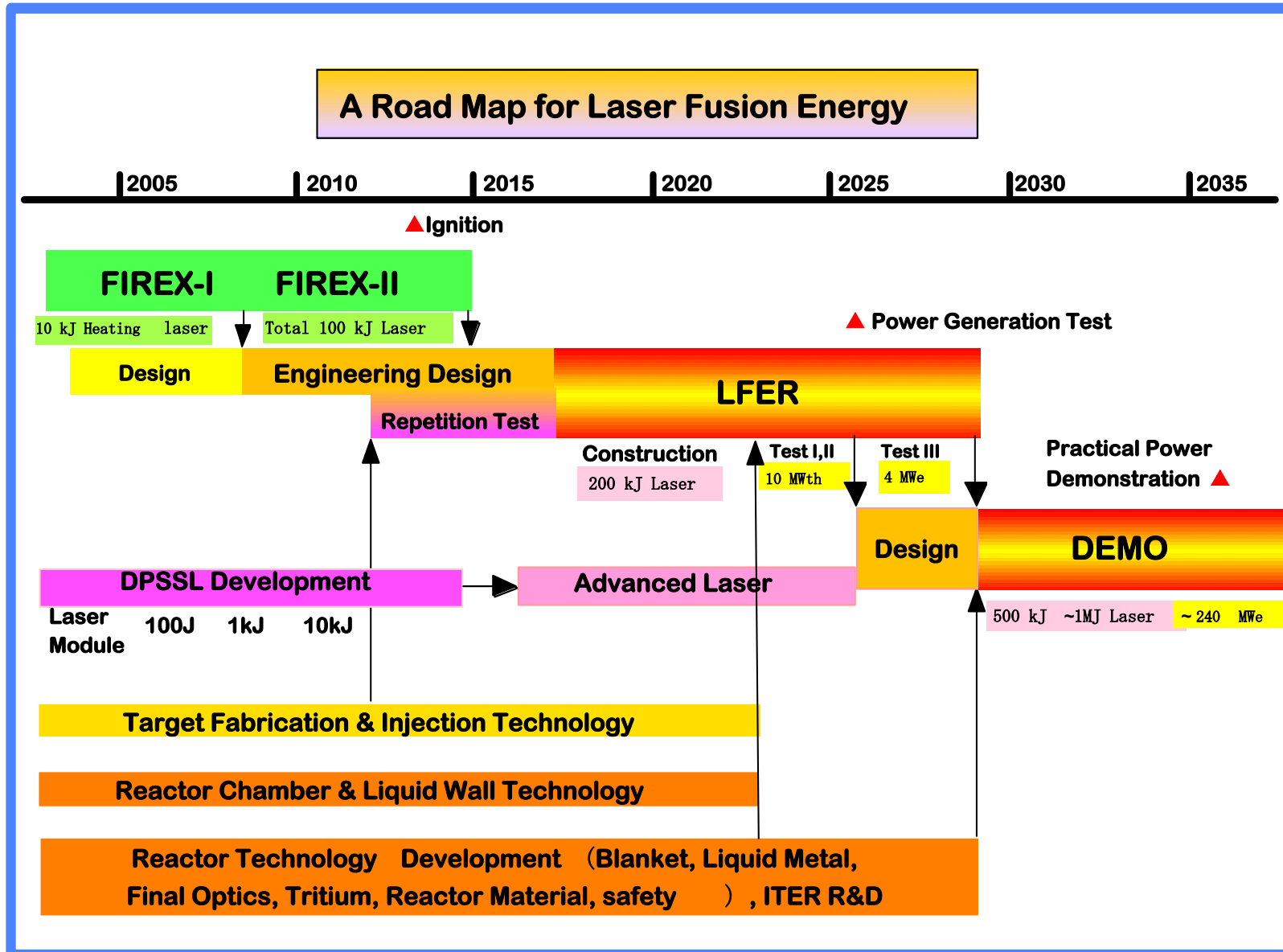
Progress in magnetic fusion



ICF – indirect drive: ignition plan on NIF



Hope from fast ignition roadmap towards a reactor in Japan



Taken from

Ken Tomabechi ¹⁾
Yasuji Kozaki ²⁾

1) IFE Forum
2) Institute of
Laser
Engineering,
Osaka University

ICF systems under development

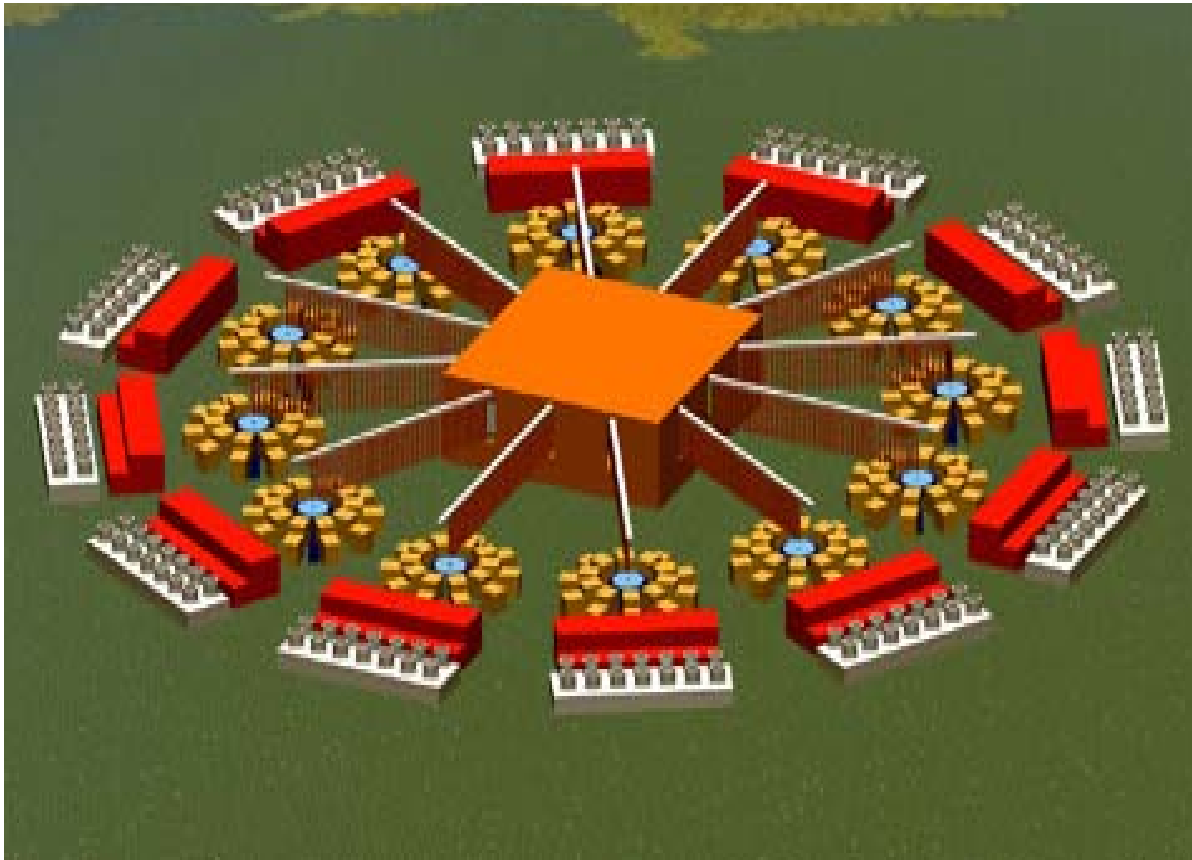
Table 1:
List of Major ICF Driver Facilities
and Their Operating Parameters
(Table Includes both Operating and Planned Facilities)

Location	Driver	Operating Parameters	Neutron Production per Shot
Sandia National Laboratory (USA)	PBFA-II (light ion beam)	36 Beams 100 TW (design) 10 TW (actual)	Unknown
Sandia National Laboratory (USA)	Z-pinch	2 MJ 290 TW 140 eV	D-T target not used yet.
Sandia National Laboratory (USA)	X-1 (successor to z-pinch) (Conceptual Design)	16 MJ 1000 TW	Projection Unknown
Europe	Heavy Ion Design for Ignition Facility (HIDIF) (Conceptual Design)	48 Beams 1 MJ 27 TW	Projection Unknown
Lawrence Livermore National Laboratory (USA)	NOVA laser	10 Beams ~40-70 kJ ~100 TW	10^8 - 3.6×10^{13}
Lawrence Livermore National Laboratory (USA)	National Ignition Facility (NIF)	192 Beams 1.8 MJ ~360 TW	10^{19} (projected at max 20 MJ yield scenario)
Osaka (Japan)	GEKKO-XII	12 Beams 15-30 kJ 0.1-10 ns	10^{13}
Osaka (Japan)	Kongoh (Under Design)	92 Beams 300 kJ 100 TW	?
Bordeaux (France)	Laser Mégajoule	1.8 MJs 120 TW	Same range as NIF
VNIIEP (Russia)	Iskra-5	12 Beams 15 kJ 0.25 ns	?

Sources: Schirmann and Tobin 1996; Gsponer and Hurni 1998; Velarde 1993; Livermore 1996b; Singer 1998.

Long term development of z-pinch

The long-range goal of Z-Pinch IFE is to produce an economically-attractive power plant using high-yield z-pinch-driven targets (~ 3 GJ) at low rep-rate (~ 0.1 Hz)



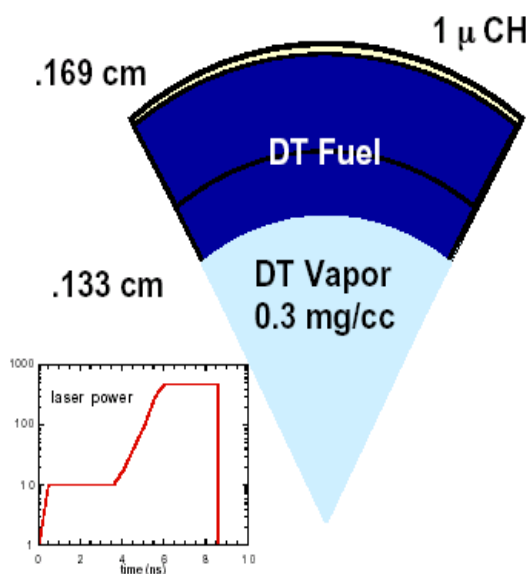
Z-Pinch IFE DEMO (ZP-3, the first study) used 12 chambers, each with 3 GJ at 0.1 Hz, to produce 1000 MWe

Control of hydrodynamic instabilities and laser imprint determine key features of laser direct drive targets



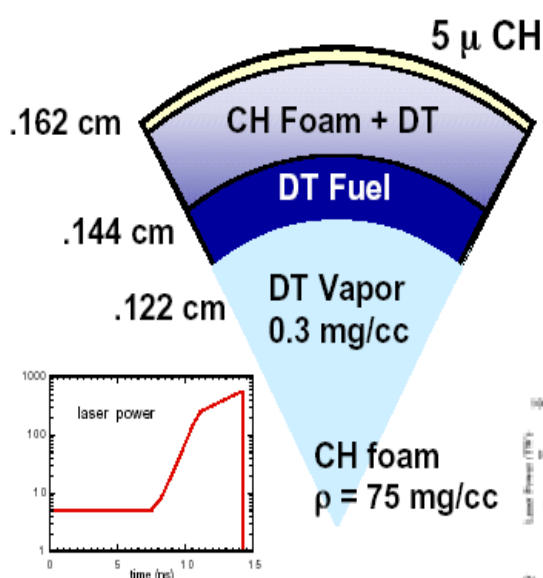
NIF BASELINE

PURE DT
(shock heated)

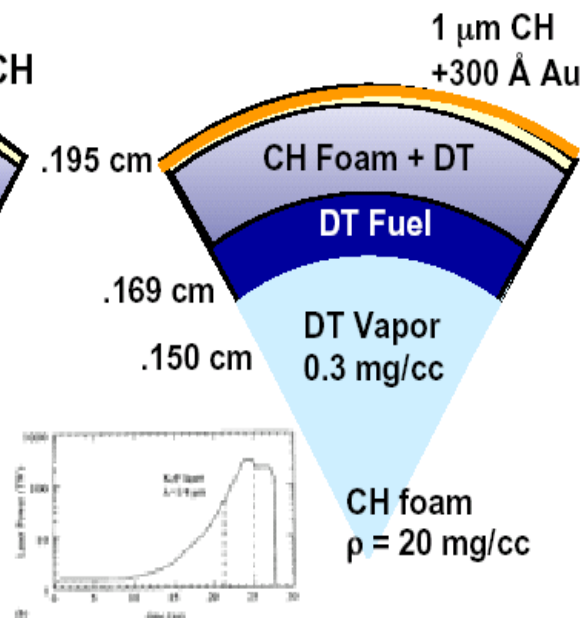


ADVANCED DESIGNS

**CH FOAM /DT+
DT**



**GOLD+
CH FOAM / DT +
DT**



Laser Type	Glass
Laser Energy	1.6 MJ (~60% absorbed)
1-D Gain	20-30

Glass	KrF
1.6 MJ (90% absorbed)	1.6 MJ
50-120	108

KrF (ISI)
1.3 MJ
127