

# Probing the hierarchy problem with the LHC

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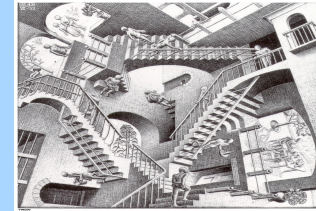
## SUSY

New particles at TeV scale  
stabilize  $m_H$



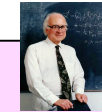
## Extra-dimensions

Additional dimensions  
 $\rightarrow M_{\text{gravity}} \sim M_{\text{EW}}$   
New states at TeV scale



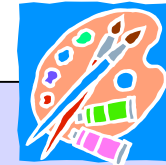
## Little Higgs

SM embedded in larger gauge group  
New particles at TeV scale, stable  $m_H$



## Technicolour

New strong interactions break EW symmetry  
 $\rightarrow$  Higgs (elementary scalar) removed  
New particles at TeV scale



## Split SUSY

Accept fine-tuning of  $m_H$   
(and of cosm. constant)  
by anthropic arguments  
Part of SUSY spectrum at TeV scale  
(for couplings unification and dark matter)



LHC potential for ~all these scenarios  
demonstrated since long time. Here:

- ① What can be done at the beginning ?
- ② Signal interpretation and constraints of underlying theory ?

# ① What can be done at the beginning ?

The first LHC data : from Summer 2007...

1 fb<sup>-1</sup> (10 fb<sup>-1</sup>) ≡ 6 months at 10<sup>32</sup> (10<sup>33</sup>) cm<sup>-2</sup>s<sup>-1</sup>  
at 50% efficiency → may collect  
several fb<sup>-1</sup> per experiment by end 2008

Channels ( <u>examples ...</u> )	Events to tape for 1 fb <sup>-1</sup> (per expt: ATLAS, CMS)	Total statistics from previous Colliders
$W \rightarrow \mu \nu$	$7 \times 10^6$	$\sim 10^4$ LEP, $\sim 10^6$ Tevatron
$Z \rightarrow \mu \mu$	$\sim 10^6$	$\sim 10^6$ LEP, $\sim 10^5$ Tevatron
$t\bar{t} \rightarrow W b W b \rightarrow \mu \nu + X$	$\sim 10^5$	$\sim 10^4$ Tevatron
$\tilde{g}\tilde{g} \quad m = 1 \text{ TeV}$	$10^2 - 10^3$	_____

With these data:

- Understand and calibrate detectors in situ using well-known physics samples  
e.g. -  $Z \rightarrow ee, \mu\mu$  tracker, ECAL, Muon chambers calibration and alignment, etc.  
-  $t\bar{t} \rightarrow b\bar{v} bjj$  jet scale from  $W \rightarrow jj$ , b-tag performance, etc.
- Measure SM physics at  $\sqrt{s} = 14 \text{ TeV}$  : W, Z, tt, QCD jets ... (omnipresent backgrounds to New Physics)

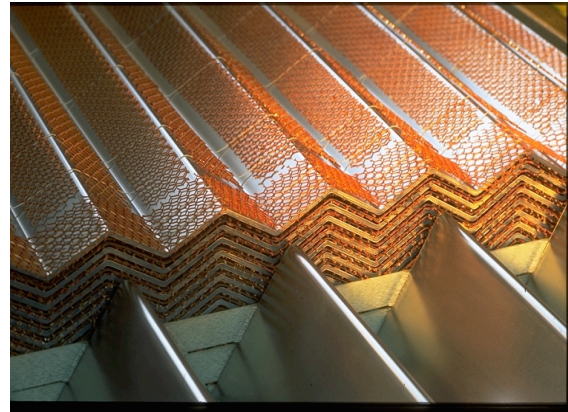
→ prepare the road to discovery ..... it will take a lot of time ...

## Preparing the detectors to explore the hierarchy problem ...

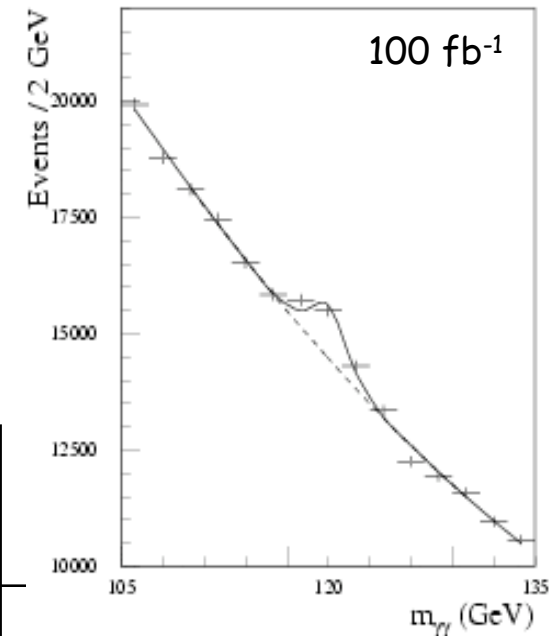
Example : the ATLAS electromagnetic calorimeter



Pb-liquid argon sampling calorimeter  
with Accordion shape, covering  $|\eta| < 2.5$

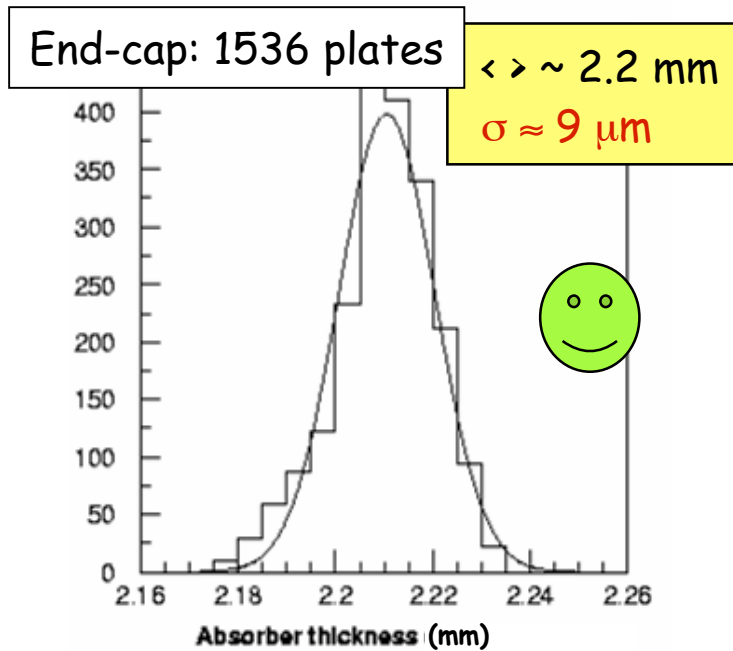


$H \rightarrow \gamma\gamma$  : to observe signal peak on top of huge  $\gamma\gamma$  background  
need mass resolution of  $\sim 1\%$   $\rightarrow$  response uniformity (i.e.  
total constant term of E-resolution)  $\leq 0.7\%$  over  $|\eta| < 2.5$



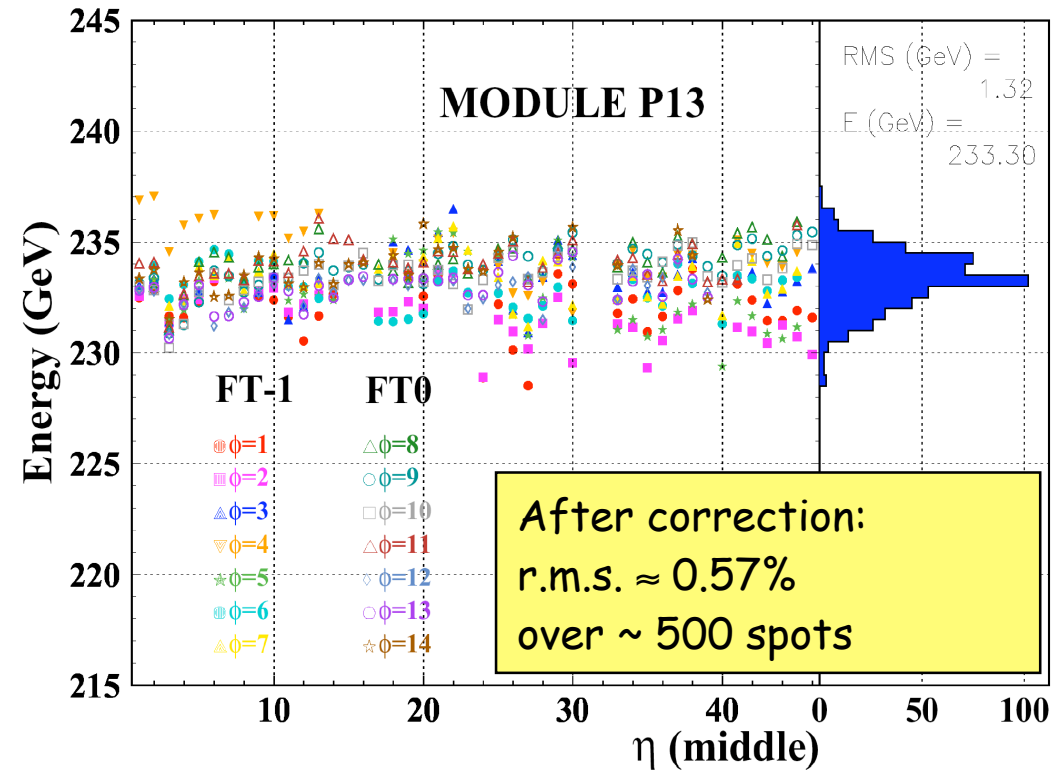
## ① Construction quality

Thickness of Pb plates must be uniform to 0.5% ( $\sim 10 \mu\text{m}$ )



## ② Test-beam measurements

Scan of a barrel module ( $\Delta\phi \times \Delta\eta = 0.4 \times 1.4$ ) with high-E electrons



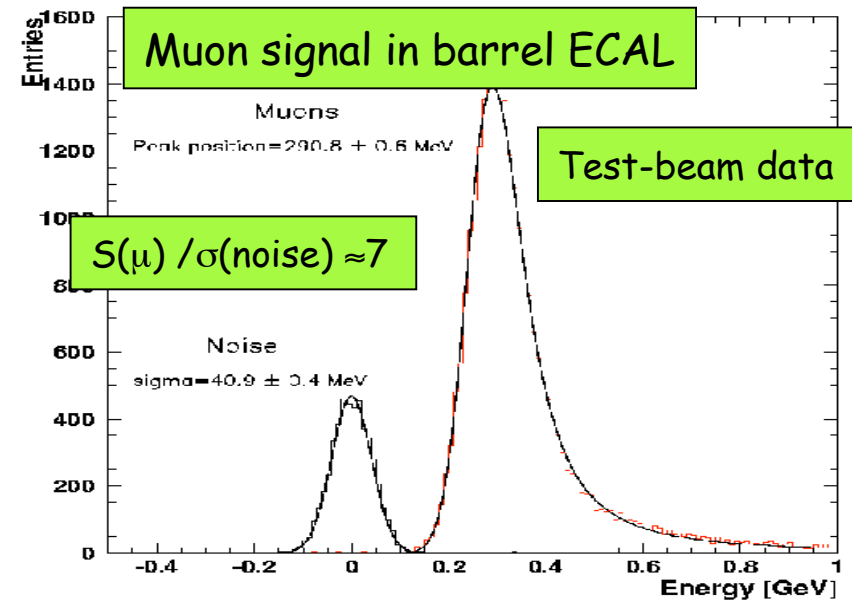
### ③ Cosmics runs:

Measured cosmic  $\mu$  rate in ATLAS pit : few Hz

→  $\sim 10^6$  events in  $\sim 3$  months of cosmics runs beginning 2007

→ enough for initial detector shake-down

→ ECAL : check calibration vs  $\eta$  to 0.5%



### ④ First collisions : calibration with $Z \rightarrow ee$ events (rate $\approx 1$ Hz at $10^{33}$ )

Use  $Z$ -mass constraint to correct long-range non-uniformities

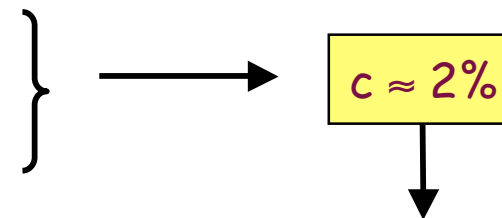
(module-to-module variations, effect of upstream material, etc.)

$\sim 10^5$   $Z \rightarrow ee$  events (few days data taking at  $10^{33}$ ) enough to achieve constant term  $c \leq 0.7\%$

Nevertheless, let's consider the worst (unrealistic ?) scenario : no corrections applied

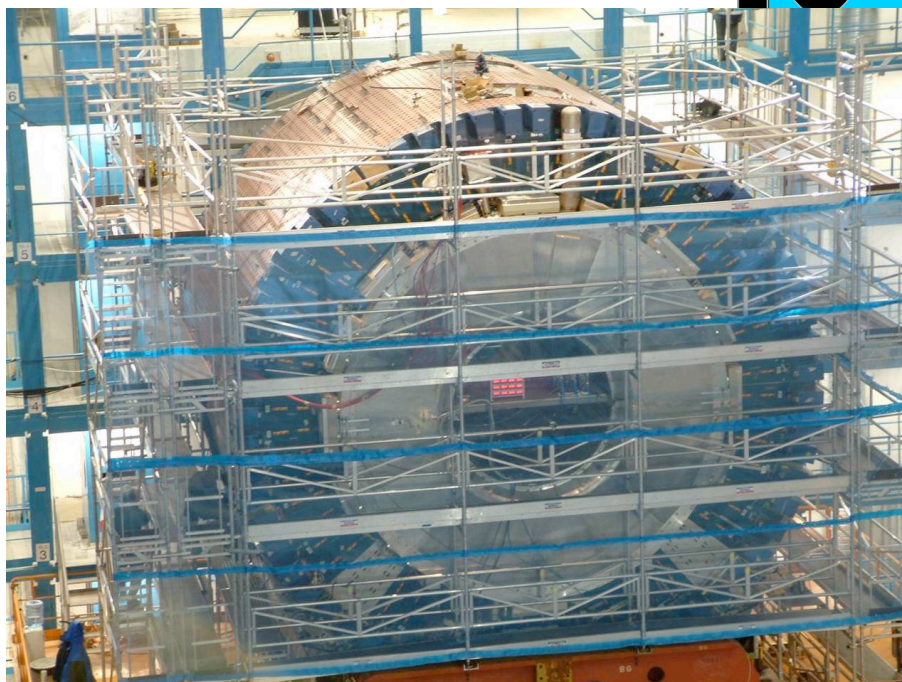
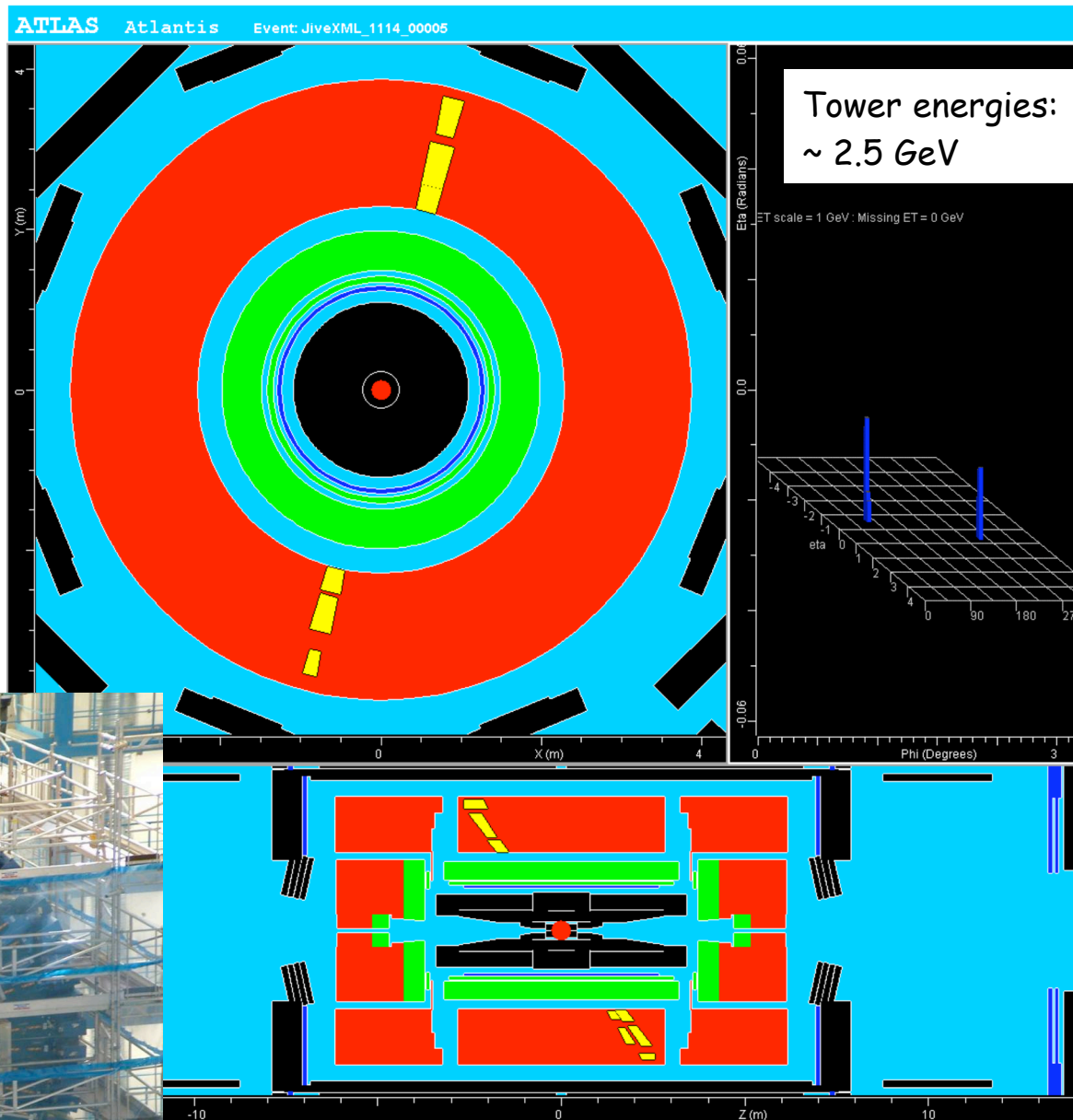
ECAL non-uniformity at construction level, i.e.:

- no test-beam corrections
- no calibration with  $Z \rightarrow ee$



$H \rightarrow \gamma\gamma$  significance  $m_H \sim 115$  GeV degraded by  $\sim 25\%$   
→ need 50% more L for discovery

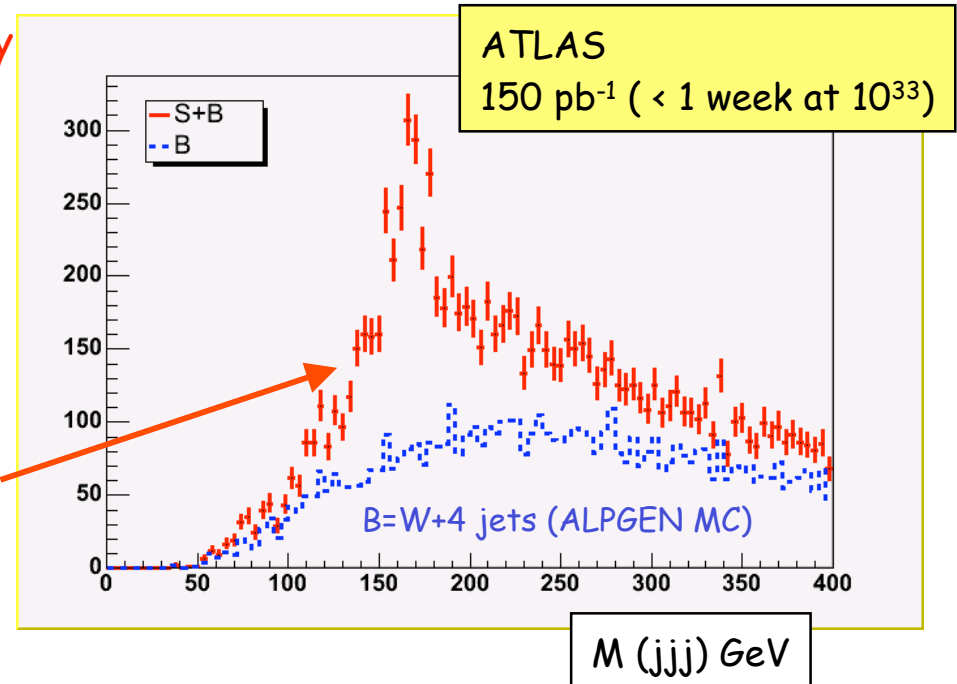
First cosmic muons  
observed by ATLAS  
in the pit on June 20th  
(recorded by hadron  
Tilecal calorimeter)



# Example of initial SM measurement : top signal and top mass (relevant to New Physics ....)

Bentvelsen et al.

- Use gold-plated  $t\bar{t} \rightarrow bW bW \rightarrow bl\nu bjj$  decay
- Very simple selection:
  - isolated lepton ( $e, \mu$ )  $p_T > 20$  GeV
  - exactly 4 jets  $p_T > 40$  GeV
  - no kinematic fit
  - no b-tagging required (pessimistic, assumes trackers not yet understood)
- Plot invariant mass of 3 jets with highest  $p_T$



Time	Events at $10^{33}$	Stat. error $\delta M_{top}$ (GeV)	Stat. error $\delta\sigma/\sigma$
1 year	$3 \times 10^5$	0.1	0.2%
1 month	$7 \times 10^4$	0.2	0.4%
1 week	$2 \times 10^3$	0.4	2.5%

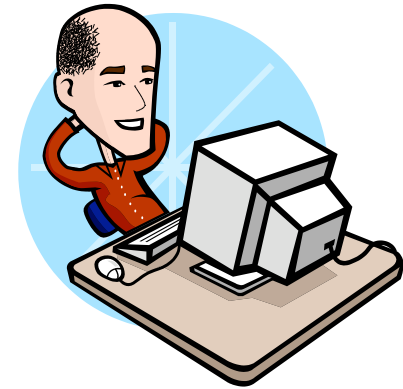
- top signal visible in few days also with simple selection and no b-tagging
- cross-section to  $\sim 20\%$
- top mass to  $\sim 7$  GeV (assuming b-jet scale to 10%)
- get feedback on detector performance :  $m_{top}$  wrong  $\rightarrow$  jet scale ?  
gold-plated sample to commission b-tagging
- $t\bar{t}$  is background to many searches



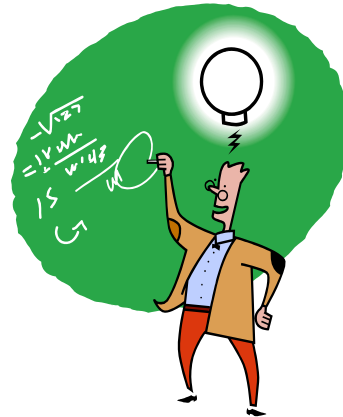
What about early discoveries ?

Three examples relevant to the hierarchy problem ...

An easy case : a new (narrow) resonance of mass  $\sim 1$  TeV decaying into  $e^+e^-$ ,  
e.g. a  $Z'$  or a Graviton  $\rightarrow e^+e^-$  of mass  $\sim 1$  TeV



An intermediate case : SUSY



A difficult case : a light Higgs ( $m_H \sim 115$  GeV)



# An "easy case" : $G \rightarrow e+e-$ resonance with $m \sim 1$ TeV

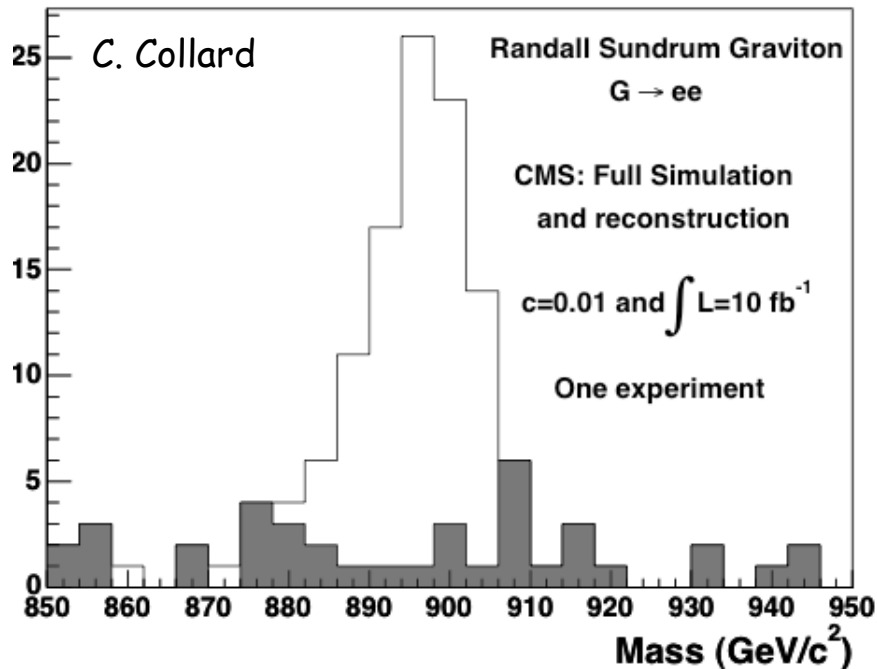
predicted in  
Randall-Sundrum  
Extra-dimensions

BR ( $G \rightarrow ee \approx 2\%$ ),  $c = 0.01$  (small/conservative coupling to SM particles)

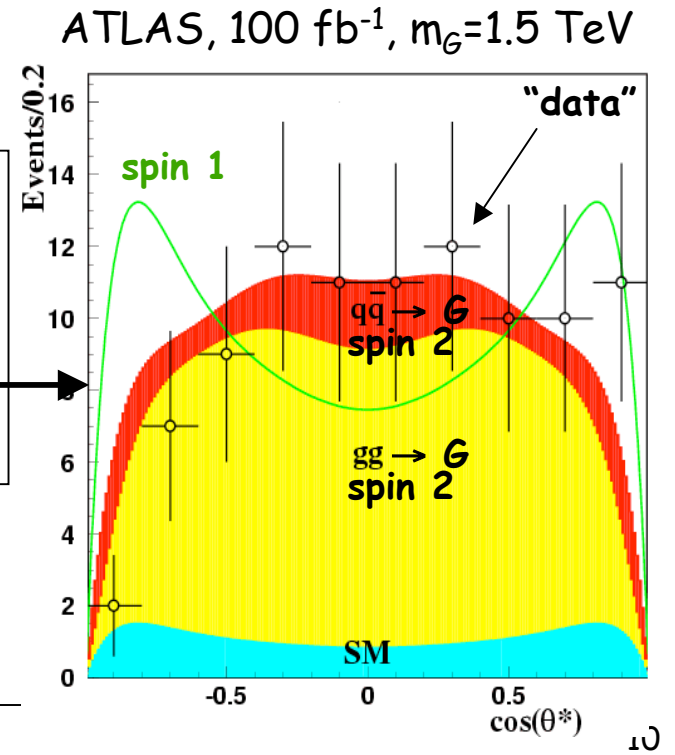
Mass (TeV)	Events for $10 \text{ fb}^{-1}$ (after all cuts)	$\int L dt$ for discovery ( $\geq 10$ observed events)
0.9	$\sim 80$	$\sim 1.2 \text{ fb}^{-1}$
1.1	$\sim 25$	$\sim 4 \text{ fb}^{-1}$
1.25	$\sim 13$	$\sim 8 \text{ fb}^{-1}$

**CMS**

- large enough signal for discovery with  $\int L dt < 10 \text{ fb}^{-1}$  for  $m < 1.3 \text{ TeV}$
- dominant Drell-Yan background small
- signal is mass peak above background



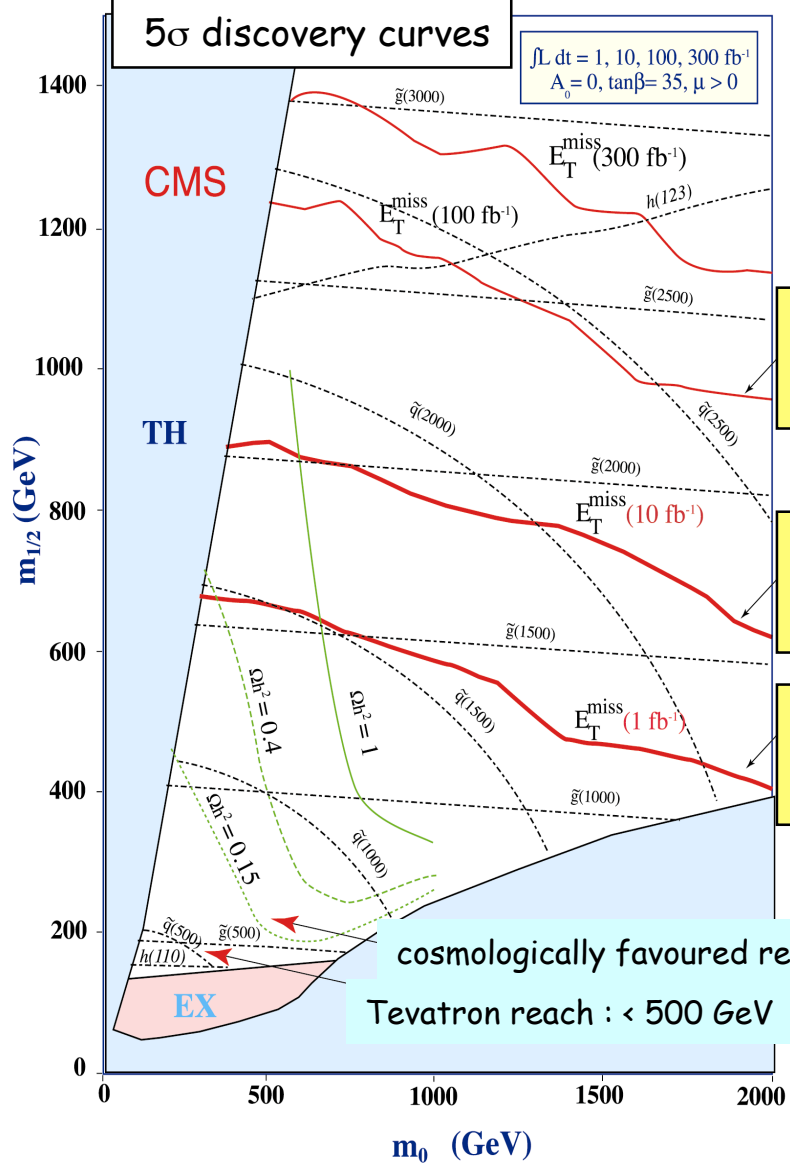
Graviton ( $s=2$ ) or  $Z'$  ( $s=1$ )?  
→ look at  $e^\pm$  angular distributions



# An "intermediate case" : SUPERSYMMETRY

If SUSY stabilizes  $m_H \rightarrow$  is at TeV scale  $\rightarrow$  could be found quickly ... thanks to:

- large  $\tilde{q}\tilde{q}, \tilde{q}\tilde{g}, \tilde{g}\tilde{g}$  cross-section  $\rightarrow \approx 100$  events/day at  $10^{33}$  for  $m(\tilde{q}, \tilde{g}) \sim 1$  TeV
- spectacular signatures

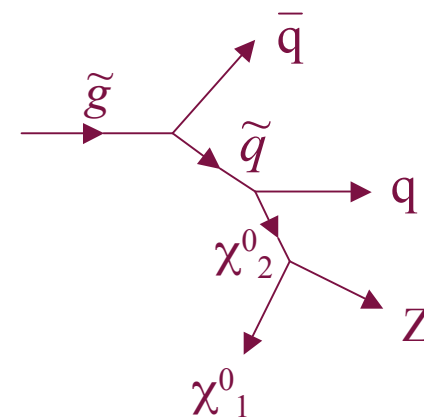


Using multijet +  $E_T^{miss}$  (most powerful and model-independent signature if R-parity conserved)

~ one year at  $10^{34}$ :  
up to ~2.5 TeV

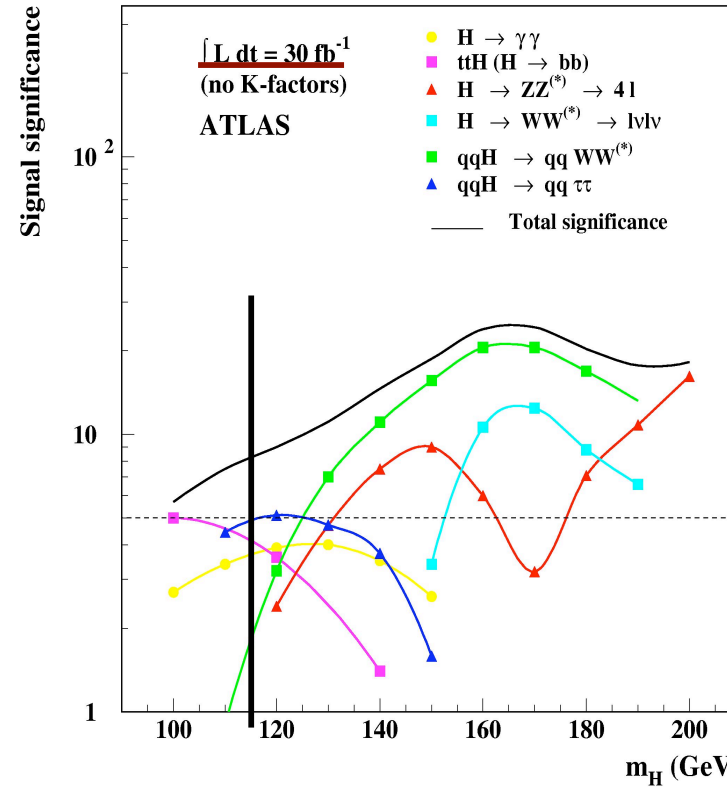
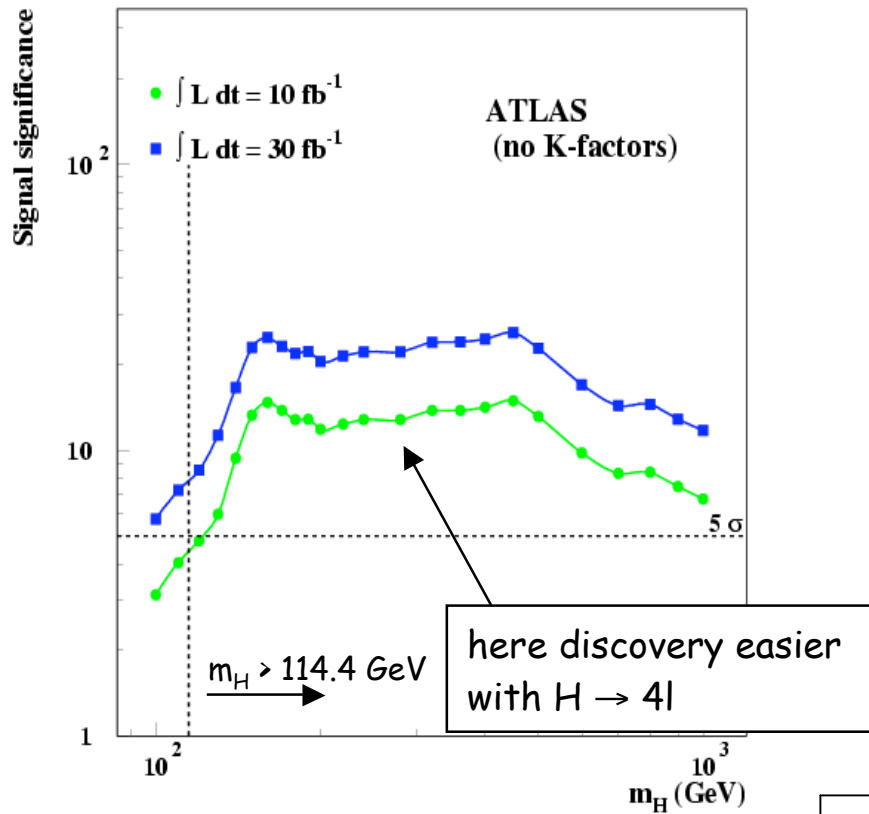
~ one year at  $10^{33}$ :  
up to ~2 TeV

~ one month at  $10^{33}$ :  
up to ~1.5 TeV



First/fast determination of SUSY (squark, gluino) mass scale from distribution of  $E_T^{miss} + \sum p_T(\text{jets})$

# A difficult case: a light Higgs ( $m_H \sim 115 \text{ GeV}$ ) ...



$m_H \sim 115 \text{ GeV}$      $10 \text{ fb}^{-1}$

total  $S/\sqrt{B} \approx 4^{+2.2}_{-1.3}$

ATLAS	$H \rightarrow \gamma\gamma$	$ttH \rightarrow ttbb$	$qqH \rightarrow qq\tau\tau$ ( $ll + l\text{-had}$ )
S	130	15	$\sim 10$
B	4300	45	$\sim 10$
$S/\sqrt{B}$	2.0	2.2	$\sim 2.7$

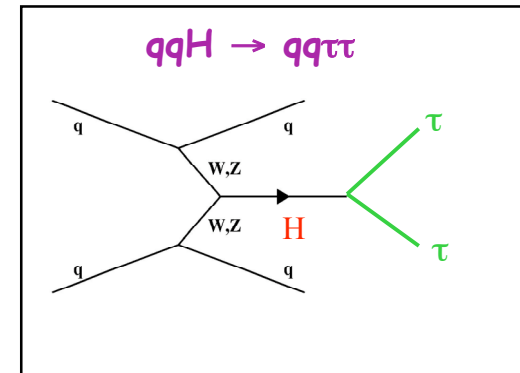
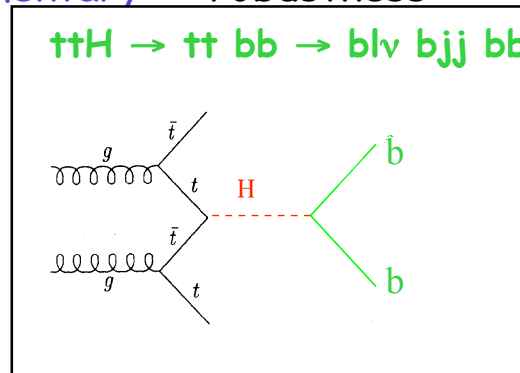
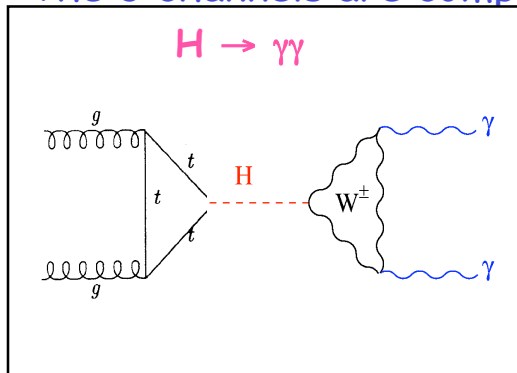
Full GEANT simulation, simple cut-based analyses

K-factors  $\equiv \sigma(\text{NLO})/\sigma(\text{LO}) \approx 2$  not included

## Remarks:

Each channel contributes  $\sim 2\sigma$  to total significance  $\rightarrow$  **observation of all channels important to extract convincing signal in first year(s)**

The 3 channels are complementary  $\rightarrow$  robustness:



- different production and decay modes
- different backgrounds
- different detector/performance requirements:
  - **ECAL crucial for  $H \rightarrow \gamma\gamma$**  (in particular response uniformity) :  $\sigma/m \sim 1\%$  needed
  - **b-tagging crucial for  $ttH$**  : 4 b-tagged jets needed to reduce combinatorics
  - **efficient jet reconstruction over  $|\eta| < 5$  crucial for  $qqH \rightarrow qq\tau\tau$**  : forward jet tag and central jet veto needed against background

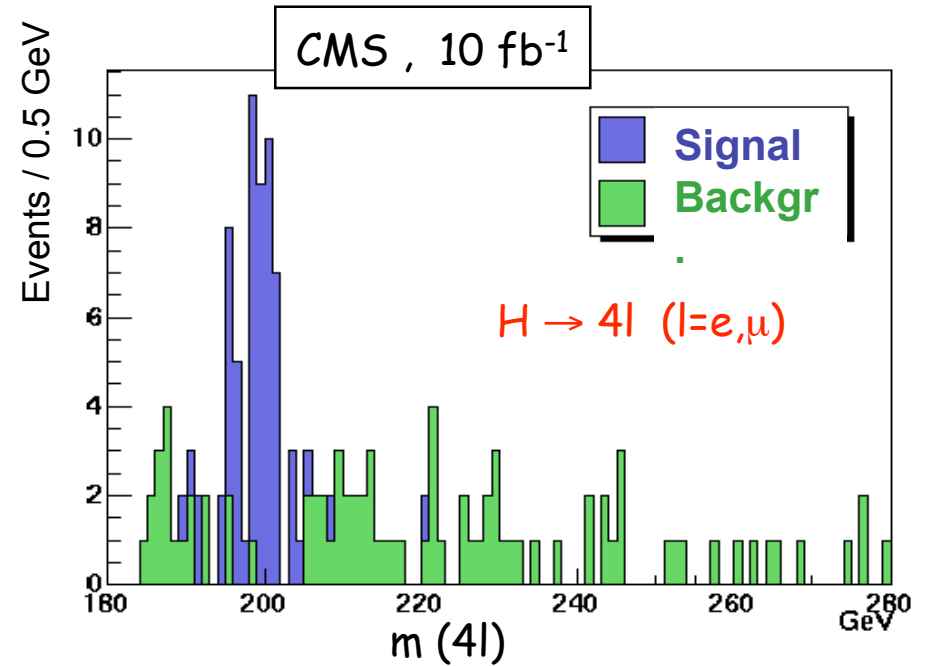
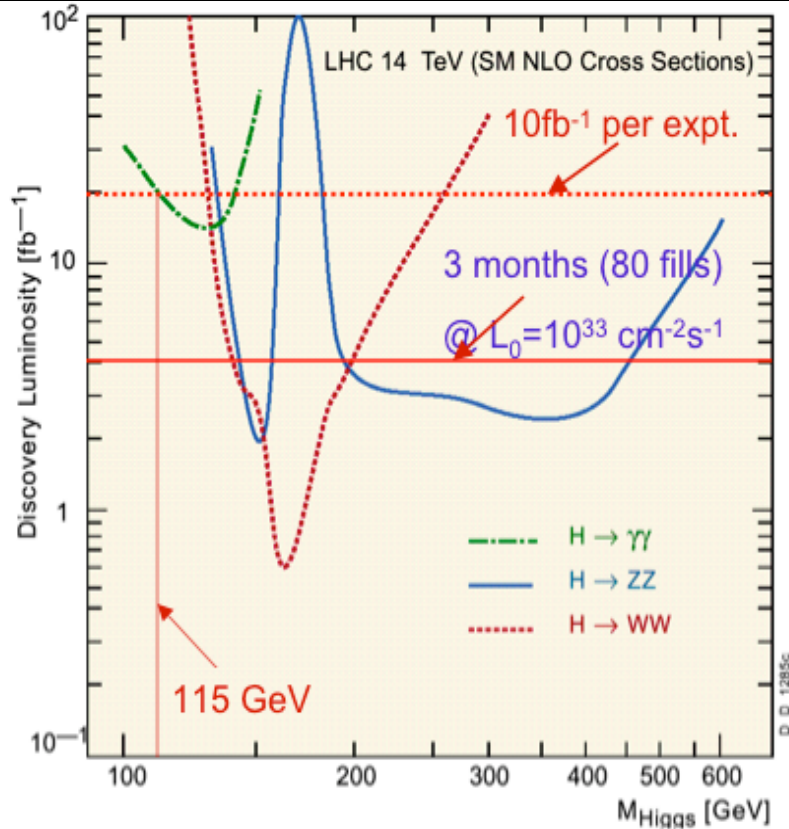
Note : -- **all require "low" trigger thresholds**

E.g.  $ttH$  analysis cuts :  $p_T(l) > 20 \text{ GeV}$ ,  $p_T(\text{jets}) > 15-30 \text{ GeV}$

-- **all require very good understanding (1-10%) of backgrounds**

If  $m_H > 180 \text{ GeV}$  : early discovery may be easier with  $H \rightarrow 4l$  channel

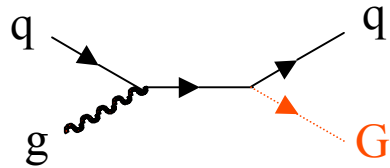
Luminosity needed for  $5\sigma$  discovery (ATLAS+CMS)



- $H \rightarrow WW \rightarrow l\nu l\nu$  : high rate ( $\sim 100$  evts/expt) but no mass peak  
→ not ideal for early discovery ...
- $H \rightarrow 4l$  : low-rate but very clean : narrow mass peak, small background

# Extra-dimensions (ADD models)

Look for a continuum of Graviton KK states :



→ topology is jet(s) + missing  $E_T$

$$\text{Cross-section} \approx \frac{1}{M_D^{\delta+2}}$$

$M_D$  = gravity scale

$\delta$  = number of extra-dimensions

ATLAS, 100 fb<sup>-1</sup>

	$\delta = 2$	$\delta = 3$	$\delta = 4$
$M_D^{\text{max}}$	9 TeV	7 TeV	6 TeV

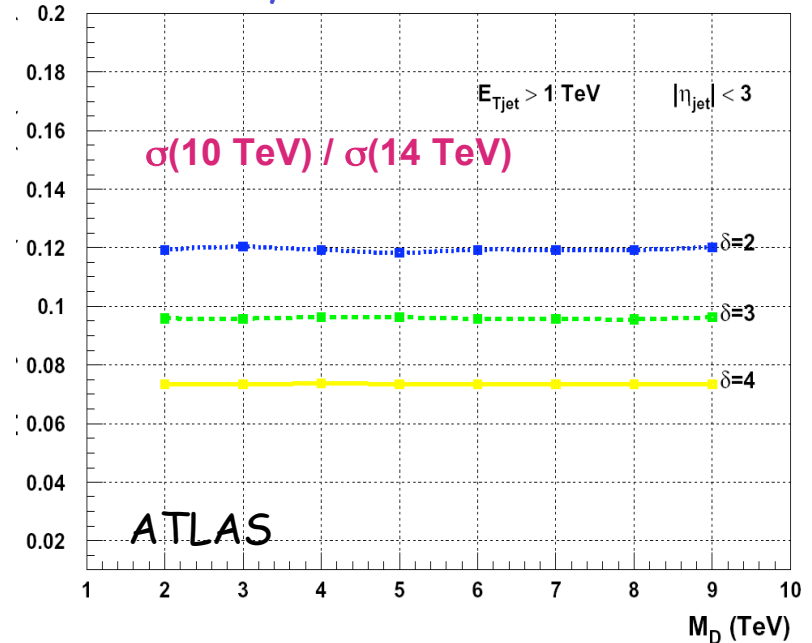
Discriminating between models:

- SUSY : multijets plus  $E_T^{\text{miss}}$  (+ leptons, ...)
- ADD : monojet plus  $E_T^{\text{miss}}$

To characterize the model need to measure  $M_D$  and  $\delta$

Measurement of cross-section gives ambiguous results: e.g.  $\delta=2, M_D=5$  TeV very similar to  $\delta=4, M_D=4$  TeV

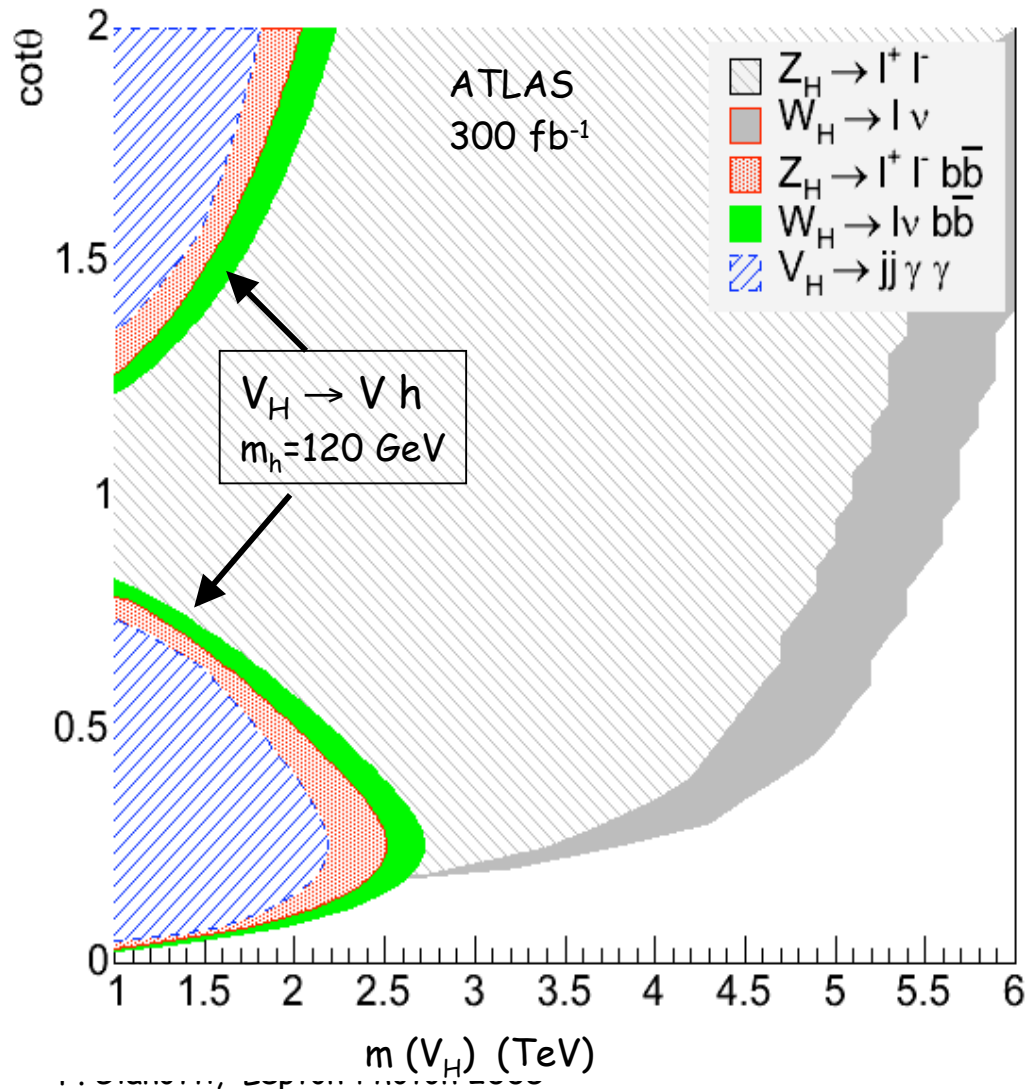
Solution may be to run at different  $\sqrt{s}$  :



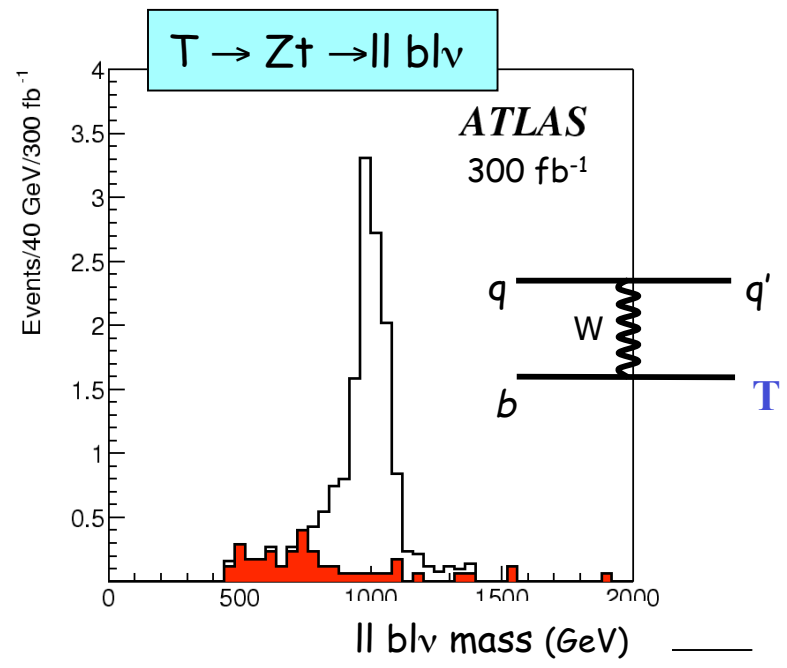
Good discrimination between various solutions possible with expected <5% accuracy on  $\sigma(10)/\sigma(14)$  for 50 fb<sup>-1</sup>

# Little Higgs models

Alternative approach to the hierarchy problem predicting heavy top  $T$  (EW singlet), new gauge bosons  $W_H, Z_H, A_H$  and Higgs triplet  $\Phi^0, \Phi^+, \Phi^{++}$

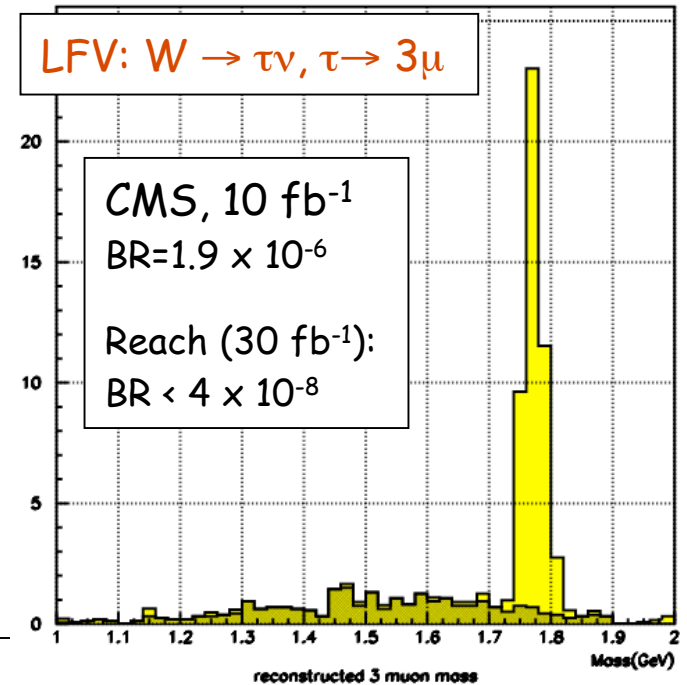
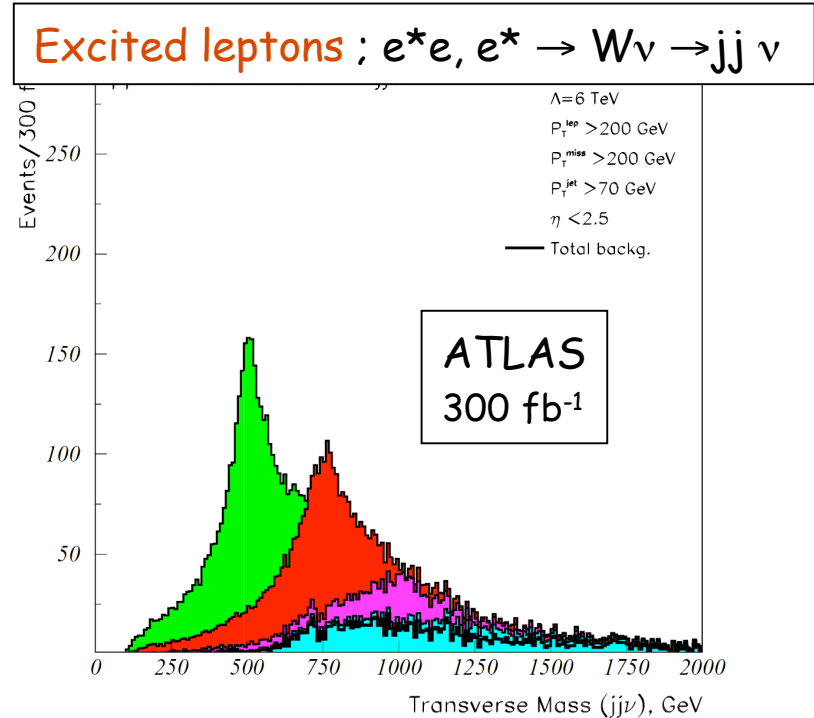
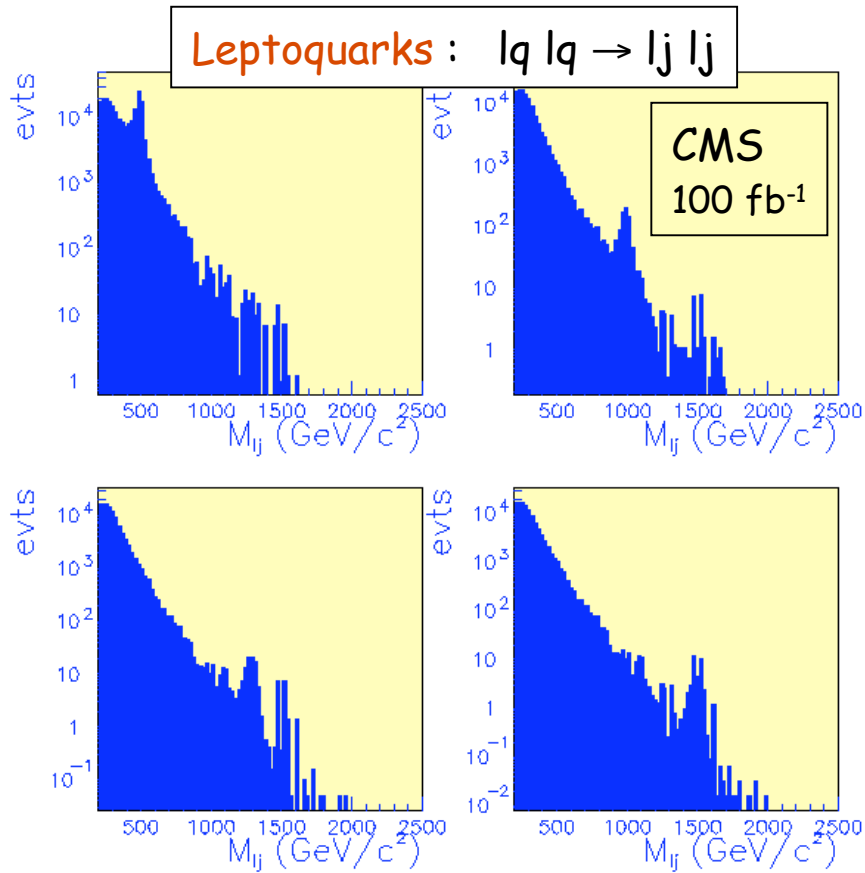


Observation of  $T \rightarrow Zt, Wb$  discriminates from 4<sup>th</sup> family quarks  
 Observation of  $V_H \rightarrow Vh$  discriminates from  $W', Z'$





# Other scenarios .....



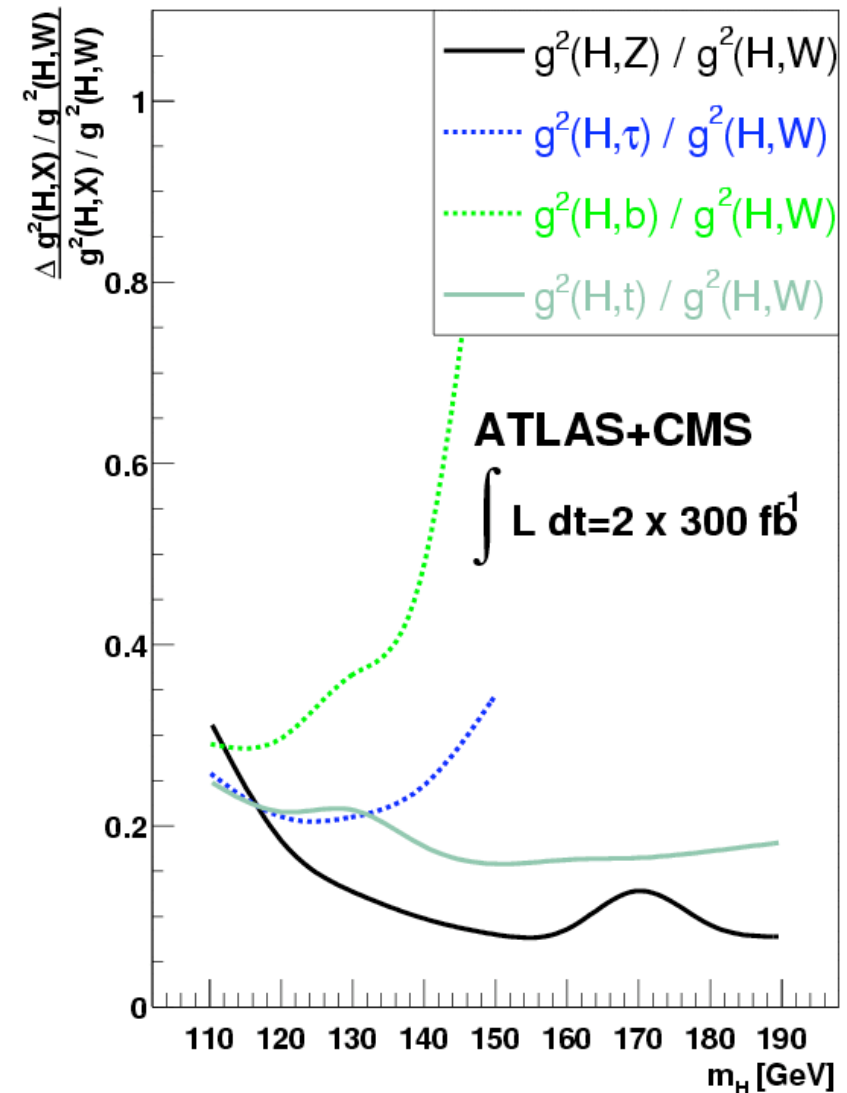
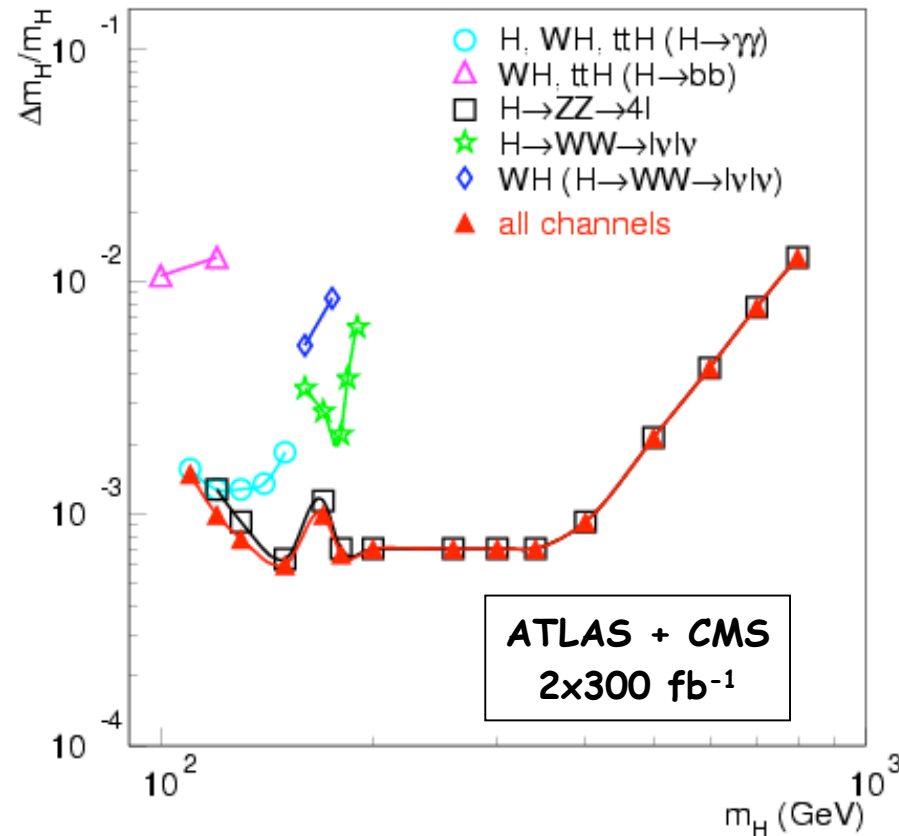
Large number of scenarios studied:

- ⇒ demonstrated detector sensitivity to many signatures
- robustness, ability to cope with unexpected scenarios
- ⇒ LHC direct discovery reach (hence exploration of hierarchy problem ... ) up to  $m \approx 5-6$  TeV

## ② Constraining the underlying theory ...

Courtesy M. Duehrssen

### Measurements of the SM Higgs parameters

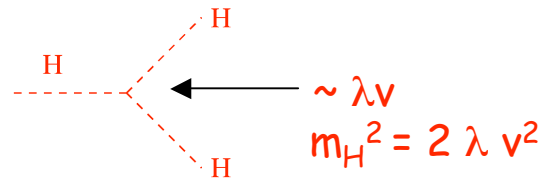


Lot of useful information to constrain the theory

(though not competitive with LC precision of e.g.  $\approx$  % on couplings)

## Higgs self-coupling $\lambda$

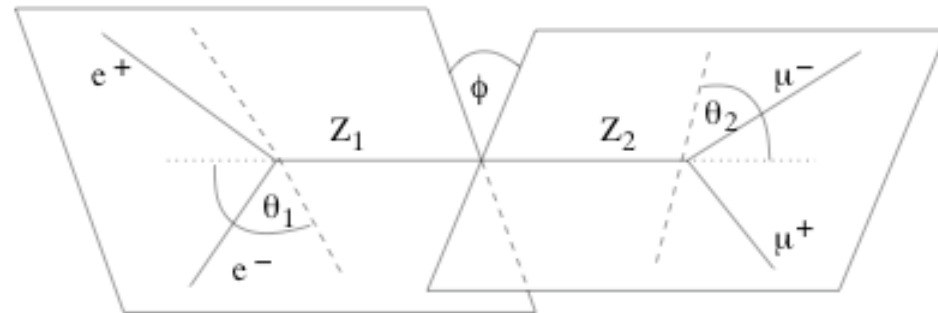
- not accessible at LHC
- may be constrained to  $\approx 20\%$  at Super-LHC ( $L=10^{35}$ )



## Higgs spin and CP

Promising for  $m_H > 180 \text{ GeV}$  ( $H \rightarrow ZZ \rightarrow 4l$ ),  
difficult at lower masses

Buszello et al. SN-ATLAS-2003-025

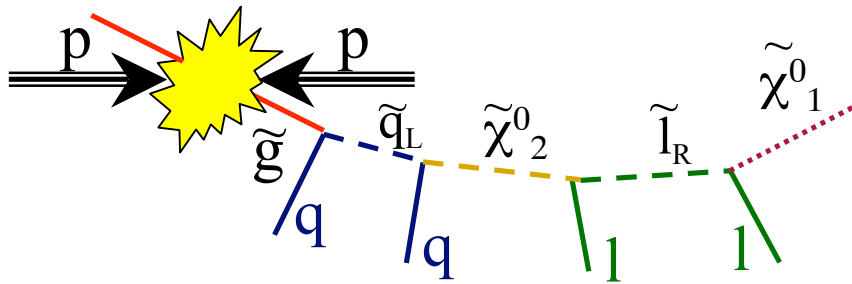


Significance for exclusion of  $J^{CP}=0^+$

ATLAS + CMS,  $2 \times 300 \text{ fb}^{-1}$

$m_H \text{ (GeV)}$	$J^{CP} = 1^+$	$J^{CP} = 1^-$	$J^{CP}=0^-$
200	$6.5 \sigma$	$4.8 \sigma$	$40 \sigma$
250	$20 \sigma$	$19 \sigma$	$80 \sigma$
300	$23 \sigma$	$22 \sigma$	$70 \sigma$

# Precise SUSY measurements

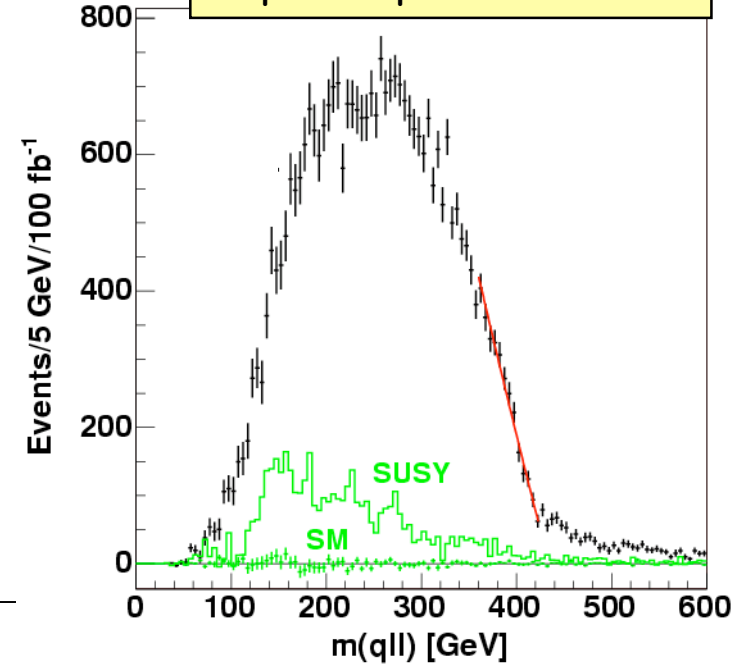
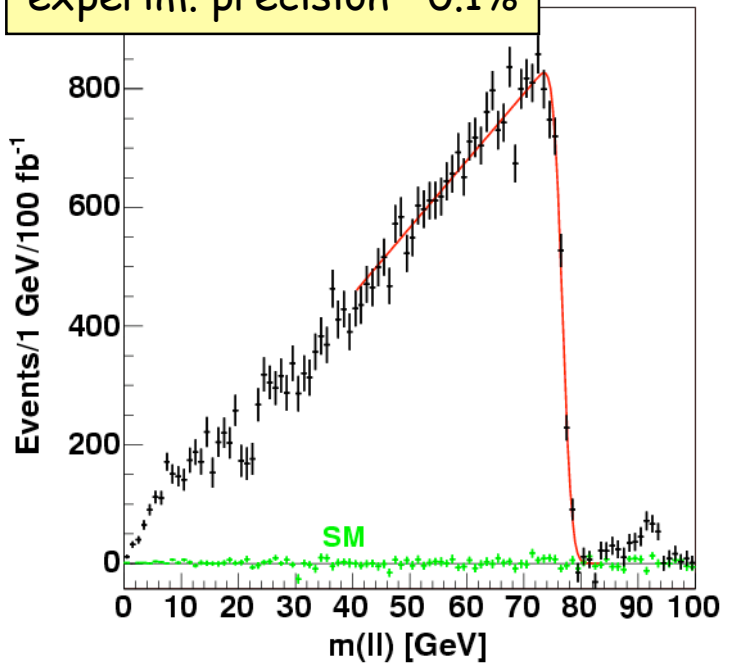


Mass peaks cannot be directly reconstructed ( $\chi_1^0$  undetectable) → measure invariant mass spectra (end-points, edges,..) of visible particles → deduce constraints on combinations of sparticle masses

$m(l+l^-)$  spectrum  
end-point : 77 GeV  
experim. precision ~0.1%

$$m(\tilde{q}_L \chi_2^0 \tilde{l}_R \chi_1^0) = 540, 177, 143, 96 \text{ GeV}$$

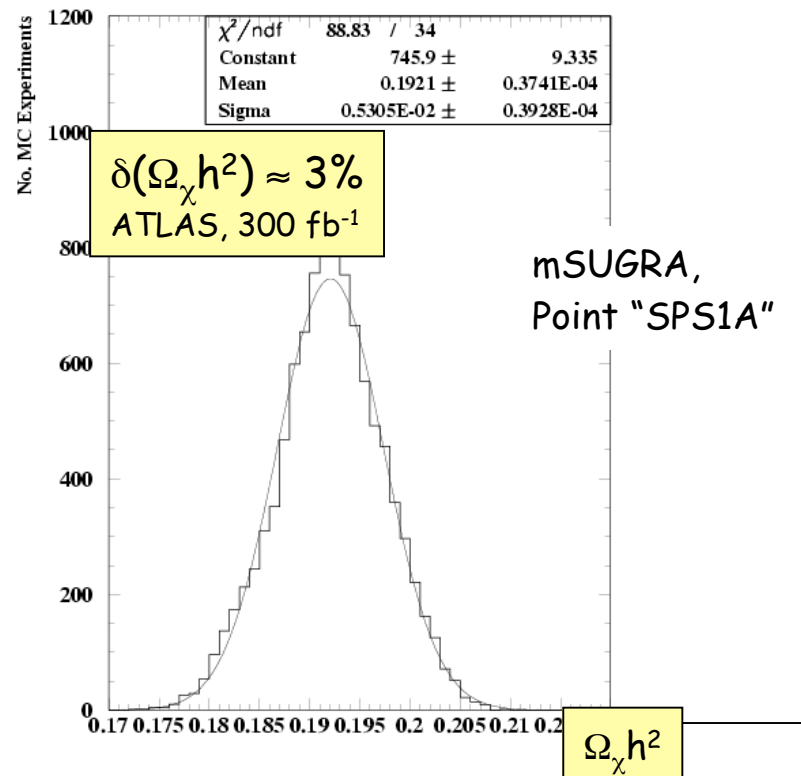
$m(l+l)^{\text{min}}$  spectrum  
end-point: 431 GeV  
experim. precision ~1 %



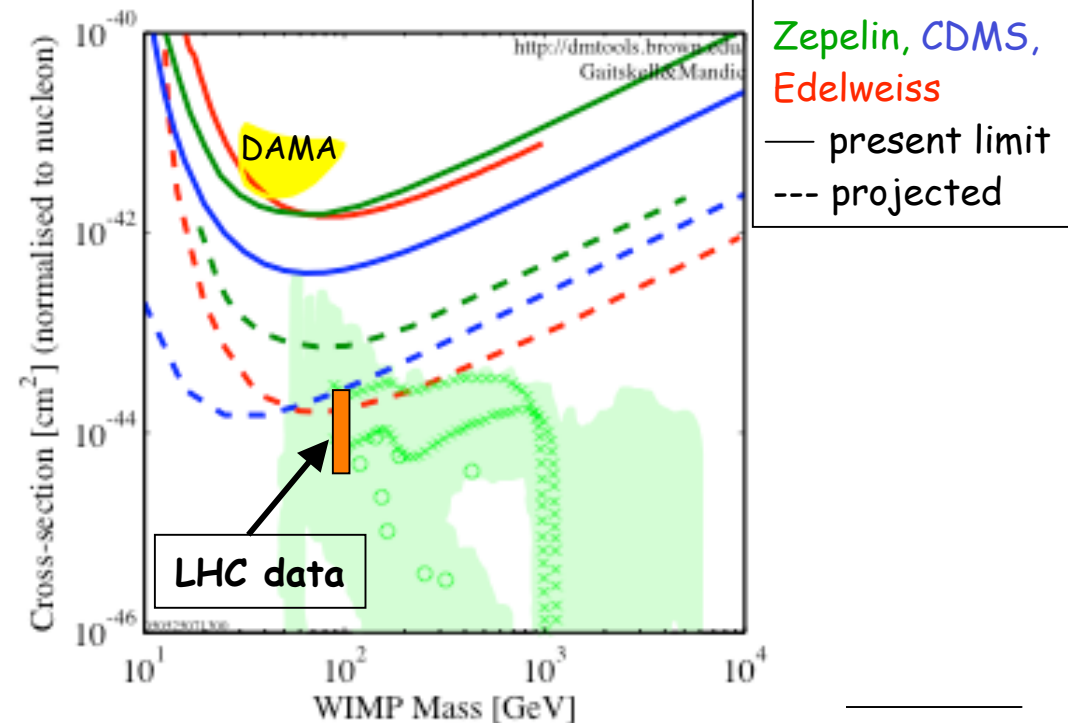
## Putting all measurements together:

- deduce several sparticle masses: typical precision 1%-20%  
Model-indep. (just kinematics), but interpretation is model-dep.
- from fit of model to all experimental measurements derive
  - sparticle masses with higher accuracy
  - fundamental parameters of theory to 1-30%
  - dark matter ( $\chi^0_1$ ) relic density and  $\sigma$  ( $\chi^0_1$  - nucleon)

demonstrated so far  
in mSUGRA (5 param.)  
and in more general  
MSSM (14 param.)



## Direct Dark Matter searches



## General strategy toward understanding the underlying theory

(SUSY as an example ...)

Discovery phase: inclusive searches ... as model-independent as possible

First characterization of model: from general features: Large  $E_T^{\text{miss}}$ ? Many leptons? Exotic signatures (heavy stable charged particles, many  $\gamma$ 's, etc.)? Excess of b-jets or  $\tau$ 's? ...

Interpretation phase:

- reconstruct/look for semi-inclusive topologies, eg.:
  - $h \rightarrow bb$  peaks (can be abundantly produced in sparticle decays)
  - di-lepton edges
  - Higgs sector: e.g.  $A/H \rightarrow \mu\mu, \tau\tau \Rightarrow$  indication about  $\tan\beta$ , measure masses
  - $tt$  pairs and their spectra  $\Rightarrow$  stop or sbottom production, gluino  $\rightarrow$  stop-top
- determine (combinations of) masses from kinematic measurements (e.g. edges ...)
- measure observables sensitive to parameters of theory (e.g. mass hierarchy)



At each step narrow landscape of possible models and get guidance to go on:

- lot of information from LHC data (masses, cross-sections, topologies, etc.)
- consistency with other data (astrophysics, rare decays, etc.)
- joint effort theorists/experimentalists will be crucial

## What the LHC can do and cannot do ....

SUSY as an example ...

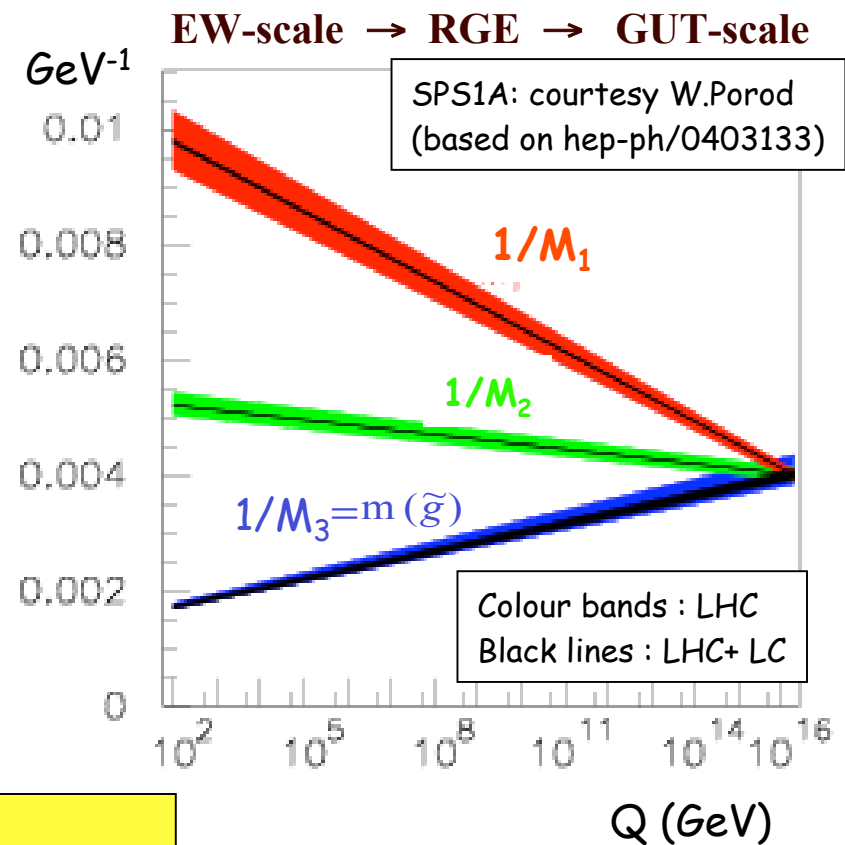
In general the LHC can (examples ...):

- discover SUSY up to  $m(\tilde{q}, \tilde{g}) \sim 2.5 \text{ TeV}$
- measure lightest Higgs  $h$  mass to  $\sim 0.1\%$
- derive sparticle masses (typically  $\tilde{q}, \tilde{g}, \chi^0_2$ ) from kinematic measurements
- constrain underlying theory by fitting a model to the data

More difficult or impossible (examples ...):

- disentangle squarks of first two generations
- observe / measure sleptons if  $m > 350 \text{ GeV}$
- measure full gaugino spectrum
- measure sparticle spin-parity and all couplings
- constrain underlying theory in model-indep. way

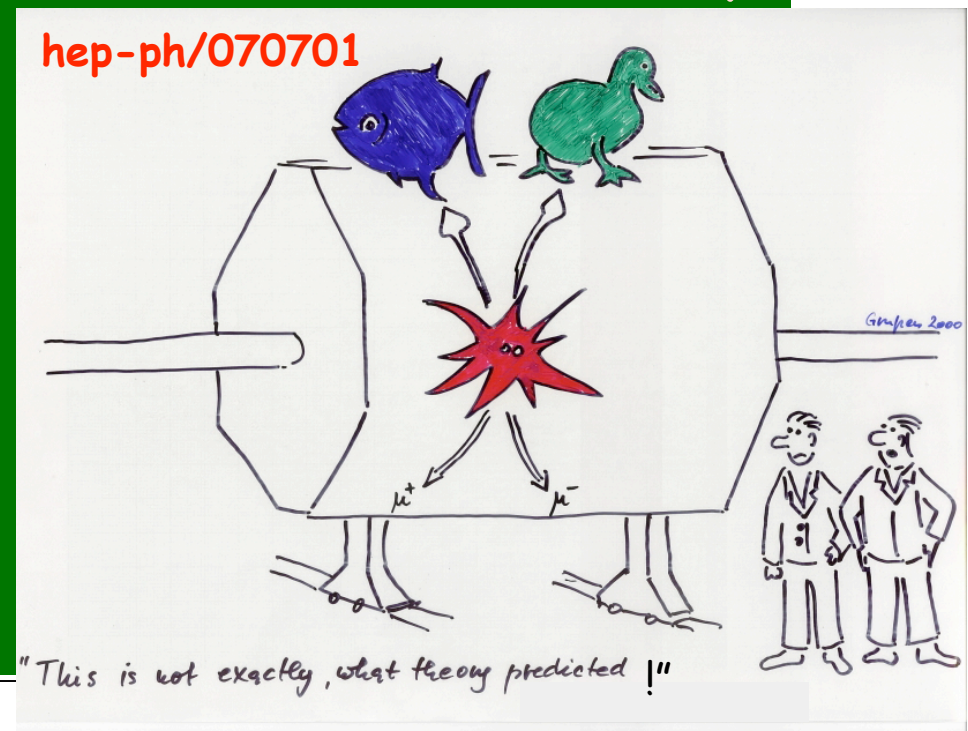
➔ complementarity with LC



Ultimate goal : from precise measurements of e.g. gaugino masses at the TeV scale reconstruct high-E theory

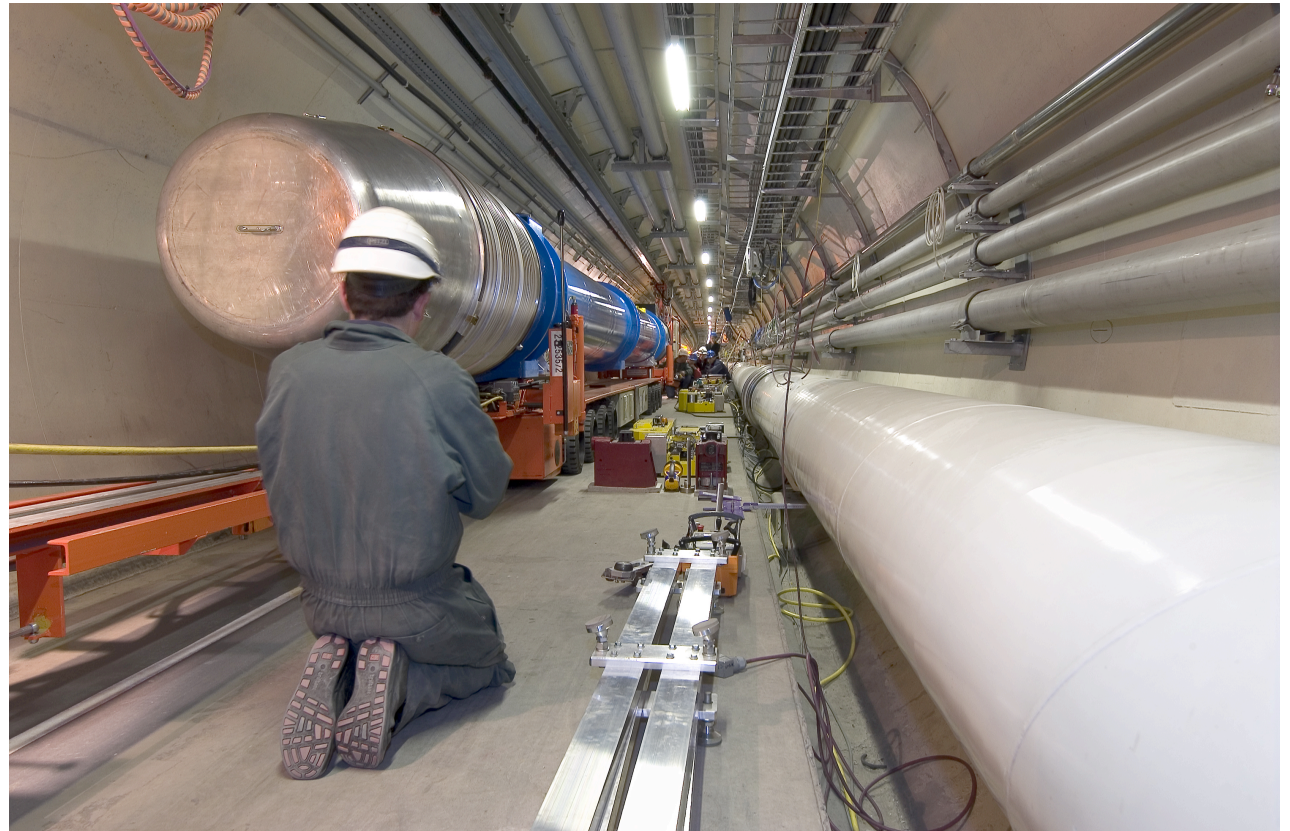
# Conclusions

- In 2 years from now, particle physics will enter a new epoch, hopefully the most glorious and fruitful of its history.
- Indeed, the hierarchy problem motivates strongly New Physics at the TeV scale
- The LHC will explore this scale in detail with direct discovery potential up to  $m \approx 5-6$  TeV
  - if New Physics is there, the LHC will find it
  - it will say final word about many TeV-scale predictions
  - it will tell us which are the right questions to ask, and how to go on





Has Nature prepared a "pleasant" welcome to the TeV-scale (striking signals with limited luminosity and non-ultimate detector performance) or shall we have to sweat through years of data taking and hard work before we can claim a discovery ?



Early determination of scale of New Physics would be crucial for planning of future facilities (ILC ? CLIC ? Underground Dark Matter searches ? .... )  
The future of our discipline will benefit from a quick feedback on SUSY and the rest .. !

Next challenge: efficient and as-fast-as-possible commissioning of machine and detectors of unprecedented complexity, technology and performance

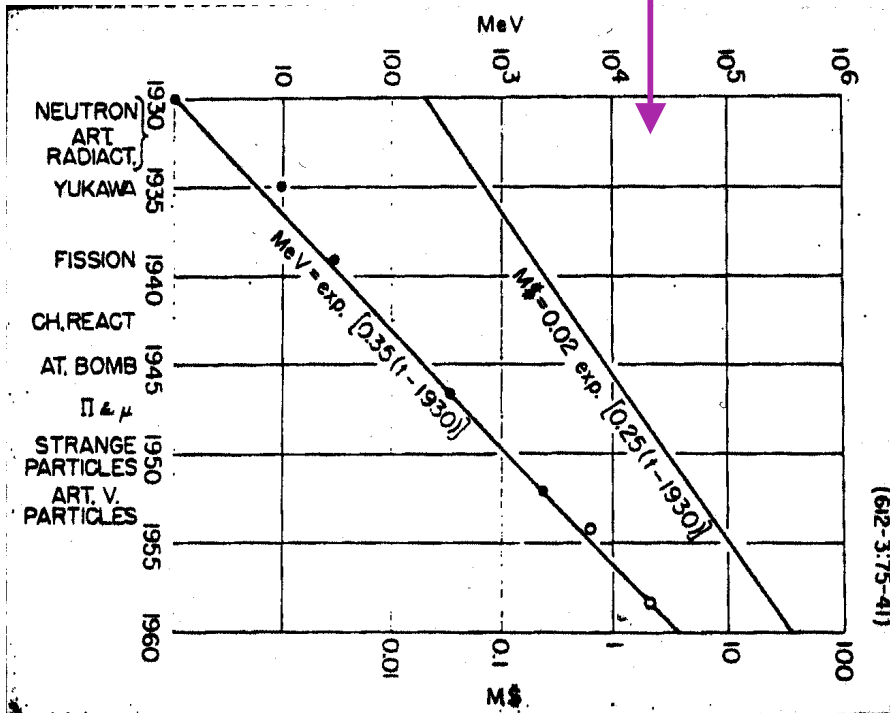
From E. Fermi, preparatory notes for a talk on  
 "What can we learn with High Energy Accelerators ?"  
 given to the American Physical Society, NY, Jan. 29th 1954

For these reasons... clamoring for higher and higher....

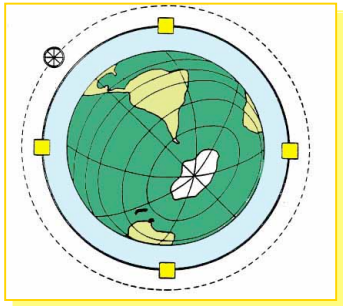
Slide 1 - MeV - M\$ versus time.  
 Extrapolating to 1994... 5 hi 9 Mev or hiest cosmic... 170 B\$.... preliminary design.... 8000 km, 20000 gauss

Slide 2 - 5 hi 15 ev machine.

What we can learn impossible to guess.... main element surprise.... some things look for but see others.... Experiments on pions.... sharpening knowledge.... ~~Spin and all had summary~~... certainly look for multiple production... *what experiments*



Fermi's extrapolation to year 1994:  
 2T magnets, R=8000 Km machine  
 $E_{beam} \sim 5 \times 10^3$  TeV , cost 170 B\$



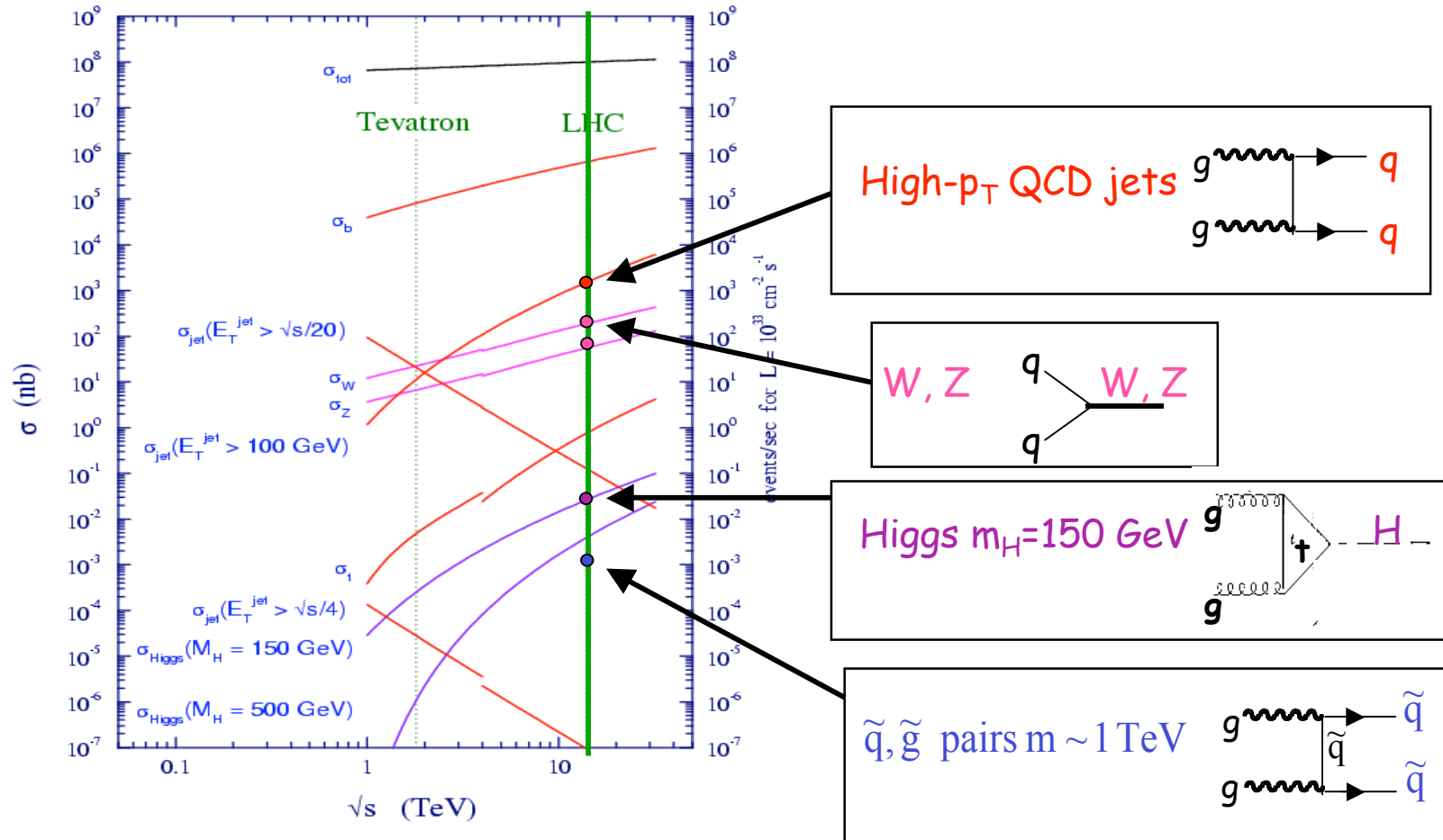
University of Chicago library  
 (thanks to M.Oreglia)

Many thanks to:

C. Collard, A. De Roeck, B. Gjelsten, K. Moening,  
L. Pape, G. Polesello, W. Porod, D. Tovey

# Back-up slides

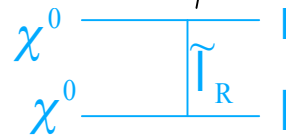
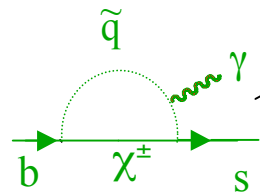
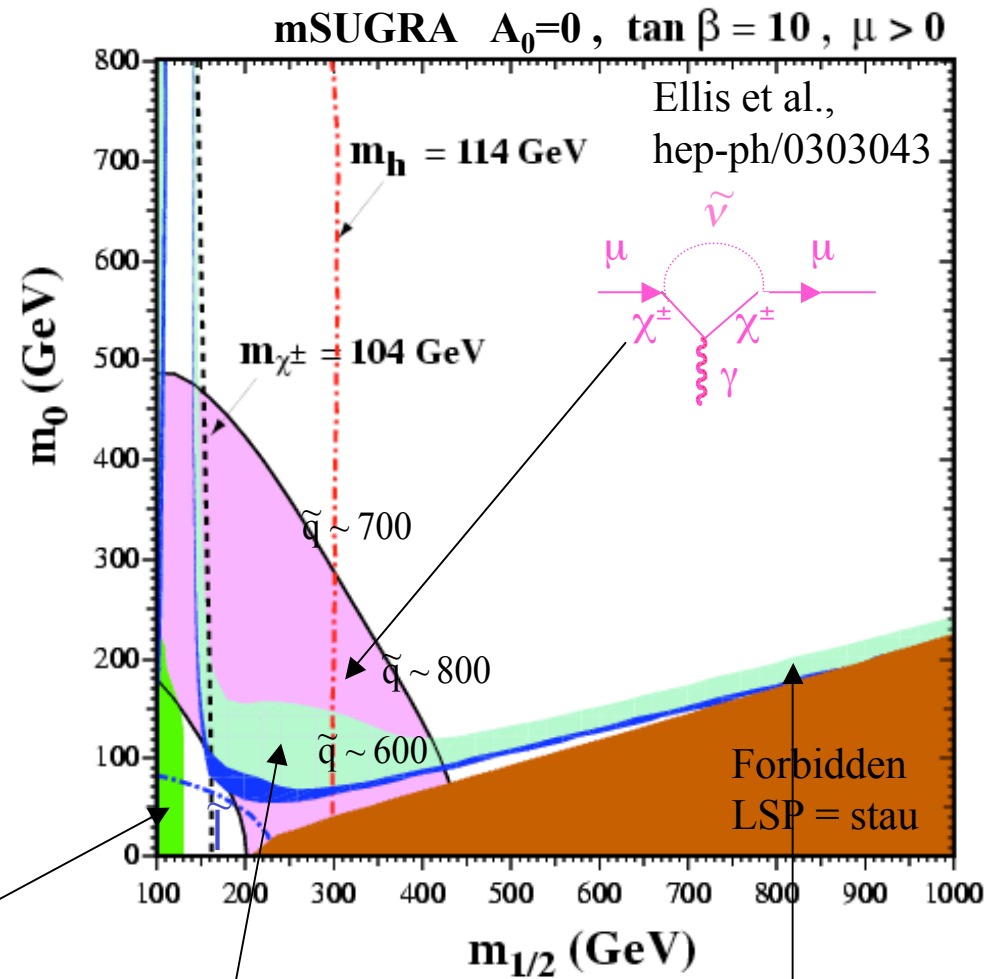
### ③ Huge (QCD) backgrounds (consequence of high energy ..)



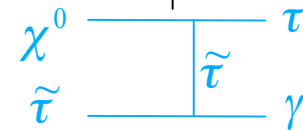
- No hope to observe light objects (W, Z, H?) in fully-hadronic final states  $\rightarrow$  rely on  $l, \gamma$
- Fully-hadronic final states (e.g.  $q^* \rightarrow qq$ ) can be extracted from backgrounds only with hard  $O(100 \text{ GeV})$   $p_T$  cuts  $\rightarrow$  works only for heavy objects
- Mass resolutions of  $\sim 1\%$  ( $10\%$ ) needed for  $l, \gamma$  (jets) to extract tiny signals from backgrounds
- Excellent particle identification: e.g.  $e/\text{jet}$  separation

# Combining Collider searches with other constraints (cosmology, ...)

- **Disfavoured by BR ( $b \rightarrow s\gamma$ )**  
 from CLEO, BELLE  
 $BR(b \rightarrow s\gamma) = (3.2 \pm 0.5) \cdot 10^{-4}$   
 used here
- **Favoured by  $g_\mu - 2$  (E821)**  
 assuming that  
 $\delta\alpha_\mu = (26 \pm 10) \cdot 10^{-10}$   
 is from SUSY ( $\pm 2 \sigma$  band)
- **Favoured by cosmology**  
 assuming  $0.1 \leq \Omega_\chi h^2 \leq 0.3$
- **Favoured by cosmology**  
 assuming  $0.094 \leq \Omega_\chi h^2 \leq 0.129$   
 i.e. new WMAP results



"bulk region"



"co-annihilation region"

Expected precision on mSUGRA parameters for six "LHC Points"

Point	$m_0$ (GeV)	$m_{1/2}$ (GeV)	$\text{tg}\beta$
1	$400 \pm 100$ (25%)	$400 \pm 8$ (2%)	$2 \pm 0.02$ (1%)
2	$400 \pm 100$ (25%)	$400 \pm 8$ (2%)	$10 \pm 1.2$ (12%)
3	$200 \pm 5$ (2.5%)	$100 \pm 1$ (1%)	$2 \pm 0.02$ (1%)
4	$800 \pm 35$ (4%)	$200 \pm 1.5$ (0.8%)	$10 \pm 0.6$ (6%)
5	$100 \pm 1.3$ (1.3%)	$300 \pm 1.5$ (0.5%)	$2 \pm 0.05$ (2.5%)
6 $\tan\beta = 45$	$218 \pm 30, 242 \pm 25$ (~ 10%)	$196 \pm 8, 194 \pm 6$ (3.5%)	$44 \pm 1.1, 45 \pm 1.7$ (~ 3%)

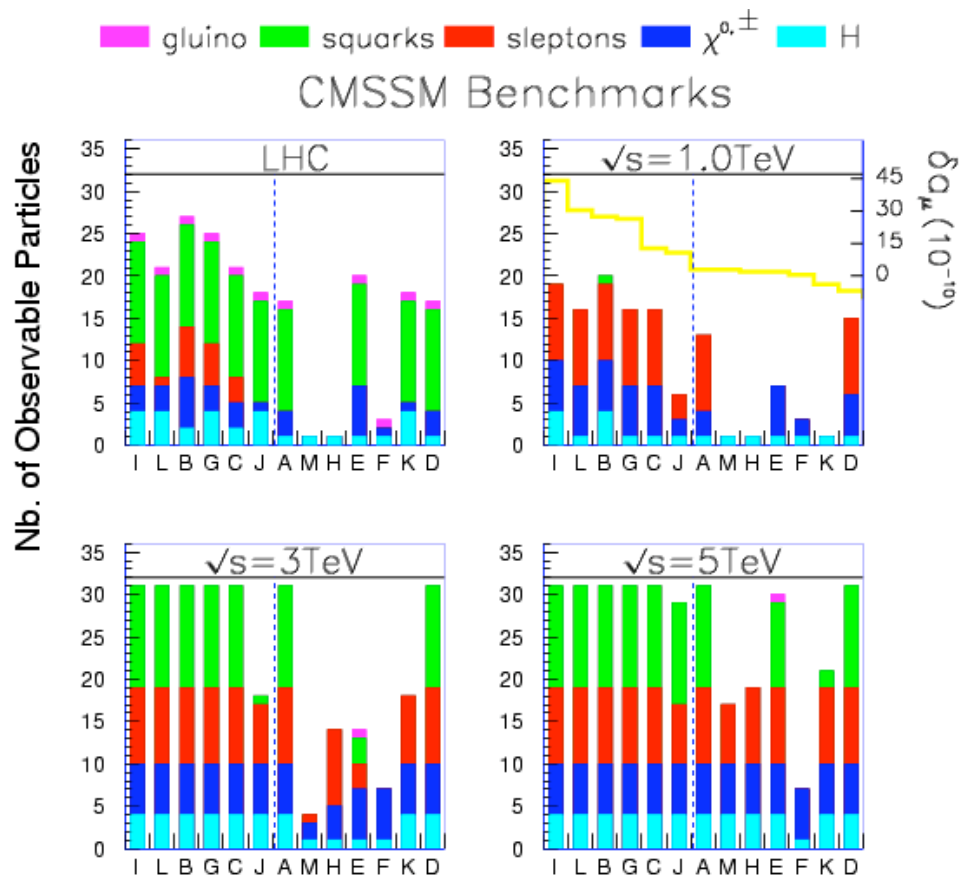
ATLAS  
300 fb<sup>-1</sup>

sign  $\mu$  determined  
except Point 6

$A_0 \sim$  unconstrained  
except Point 6

$\mu = +, -$

# Complementarity between LHC and future $e^+e^-$ Colliders



In general :

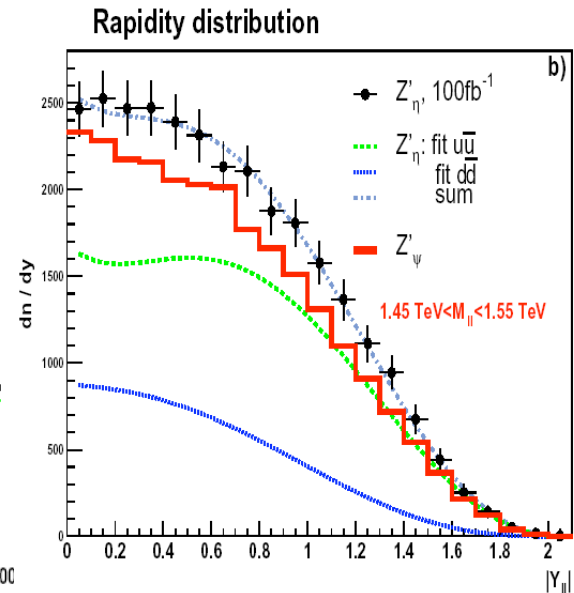
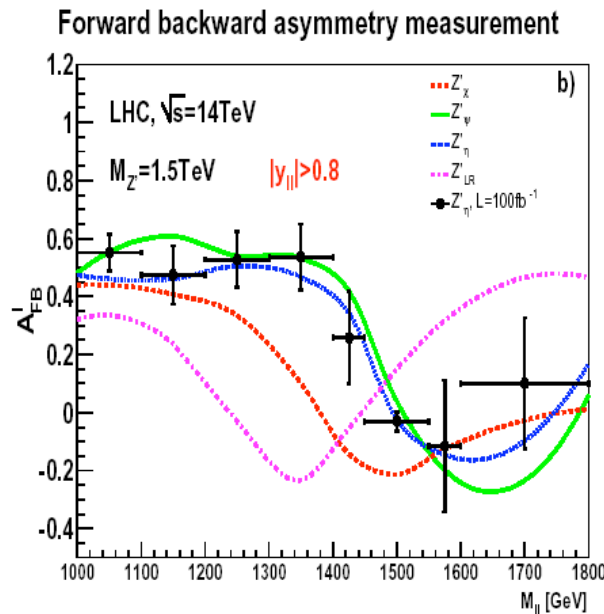
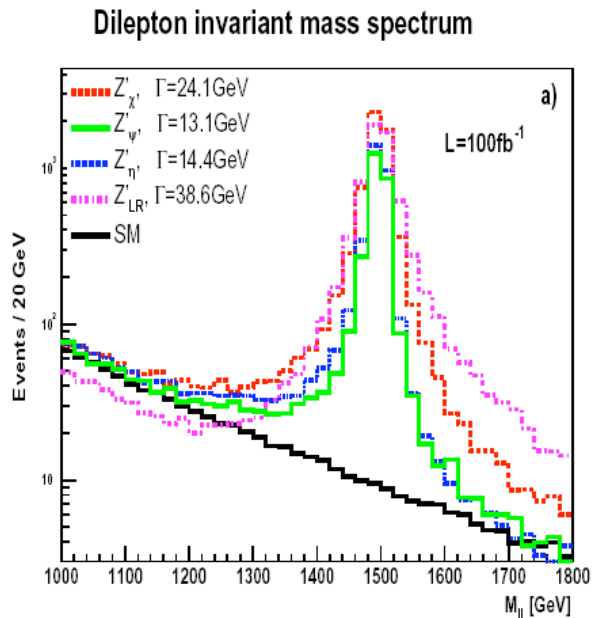
- LHC most powerful for  $\tilde{g}$  and  $\tilde{q}$  (strongly interacting) but can miss some EW sparticles (gauginos, sleptons) and Higgs bosons
- Depending on  $\sqrt{s}$ , LC should cover part/all EW spectrum (usually lighter than squarks/gluinos) → should fill holes in LHC spectrum. Squarks could also be accessible if  $\sqrt{s}$  large enough.

LC can perform precise measurements of masses (to  $\sim 0.1\%$ ), couplings, field content of sparticles with mass up to  $\sim \sqrt{s}/2$ , disentangle squark flavour, etc. (see lectures by M. Battaglia)



# Extended gauge groups : $Z' \rightarrow l+l-$

CMS



- Reach in 1 year at  $10^{34}$  : 4-5 TeV
- Discriminating between models possible up to  $m \sim 2.5$  TeV by measuring:
  - $\sigma \times \Gamma$  of resonance
  - lepton F-B asymmetry
  - $Z'$  rapidity

- HLT/DAQ deferrals limit available networking and computing for HLT → limit LVL1 output rate
- Large uncertainties on LVL1 affordable rate vs money (component cost, software performance, etc.)

Selections (examples ...)	LVL1 rate (kHz) $L = 1 \times 10^{33}$ no deferrals	LVL1 rate (kHz) $L = 2 \times 10^{33}$ no deferrals	LVL1 rate (kHz) $L = 2 \times 10^{33}$ with deferrals An example for illustration...
MU6,8,20	23	→ 19	→ 0.8
2MU6	---	0.2	0.2
EM20i,25,25	11	→ 12	→ 12
2EM15i,15,15	2	4	4
J180,200,200	0.2	0.2	0.2
3J75,90,90	0.2	0.2	0.2
4J55,65,65	0.2	0.2	0.2
J50+xE50,60,60	0.4	0.4	0.4
TAU20,25,25 +xE30	2	2	2
MU10+EM15i	---	0.1	0.1
Others (pre-scaled, etc.)	5	5	5
Total	~ 44	~ 43	~ 25

LVL1 designed for 75 kHz  
→ room for factor ~ 2 safety

Likely max affordable rate,  
no room for safety factor

### ③ Which data samples ?

Total trigger rate to storage at  $2 \times 10^{33}$   
reduced from  $\sim 540$  Hz (HLT/DAQ TP, 2000)  
to  $\sim 200$  Hz (now)

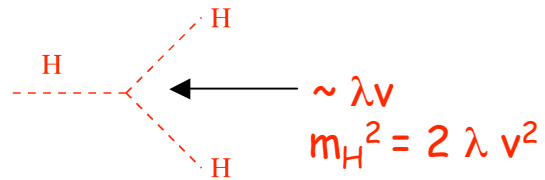
High-Level-Trigger output



Selection (examples ...)	Rate to storage at $2 \times 10^{33}$ (Hz)	Physics motivations (example)
$e25i, 2e15i$ $\mu20i, 2\mu10$	$\sim 40$ (55% W/b/c $\rightarrow eX$ ) $\sim 40$ (85% W/b/c $\rightarrow \mu X$ )	Low-mass Higgs ( $ttH, H \rightarrow 4\lambda, q\bar{q}$ ) W, Z, top, New Physics ?
$\gamma60i, 2\gamma20i$	$\sim 40$ (57% prompt $\gamma$ )	$H \rightarrow \gamma\gamma$ , New Physics (e.g. $X \rightarrow \gamma\gamma$ $m_X \sim 500$ GeV)
$j400, 3j165, 4j110$	$\sim 25$	Overlap with Tevatron for new $X \rightarrow jj$ in danger ...
$j70 + xE70$	$\sim 20$	SUSY : $\sim 400$ GeV squarks/gluinos
$\tau35 + xE45$	$\sim 5$	MSSM Higgs, New Physics (3 <sup>rd</sup> family !)? More difficult
$2\mu6 (+ m_B)$	$\sim 10$	Rare decays $B \rightarrow \mu\mu X$
Others (pre-scaled, exclusive, ...)	$\sim 20$	Only 10% of total !
<b>Total</b>	<b><math>\sim 200</math></b>	No safety factor included. "Signal" (W, $\gamma$ , etc.) : $\sim 100$ Hz

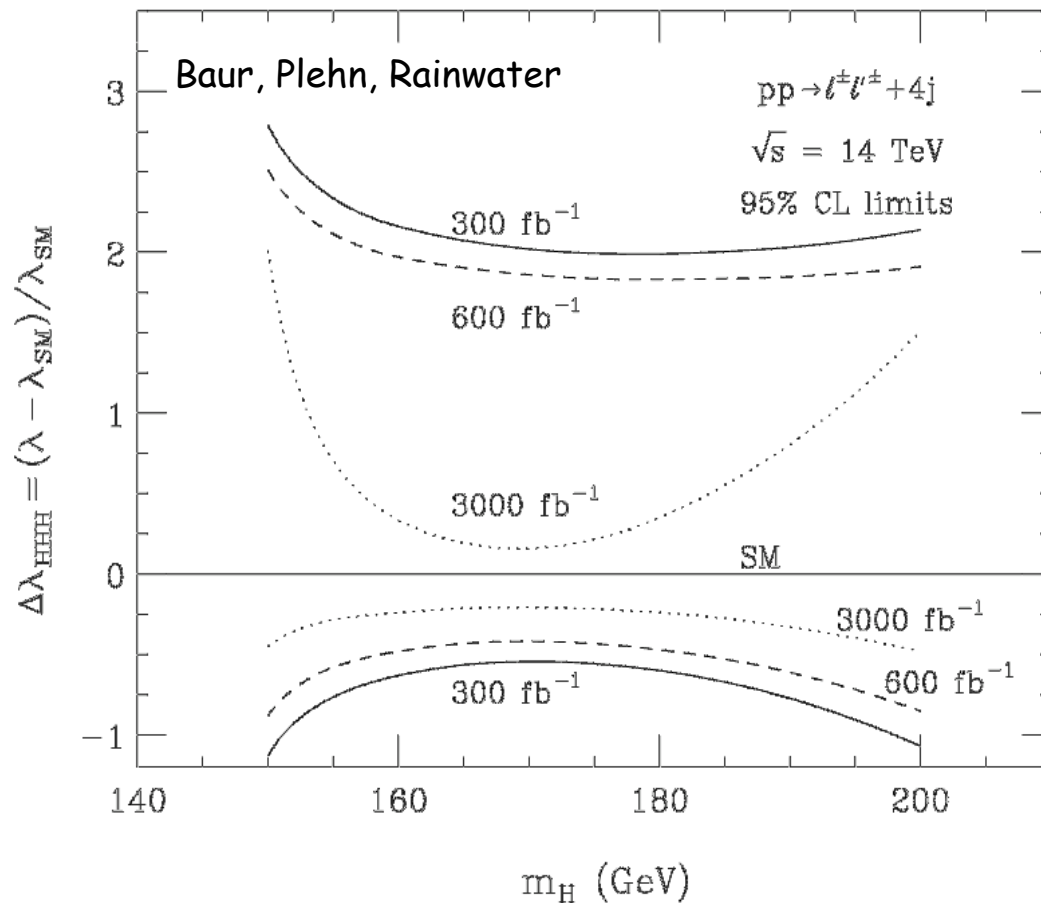
Best use of spare capacity when  $L < 2 \times 10^{33}$  being investigated

## Can the LHC measure the Higgs self-coupling $\lambda$ ?



$$HH \rightarrow W^+ W^- W^+ W^- \rightarrow l^\pm \nu jj l^\pm \nu jj$$

Higgs pair production is rate-limited at the LHC, but should be accessible with a luminosity upgrade to  $10^{35}$  (expected in  $\sim 2013$ )

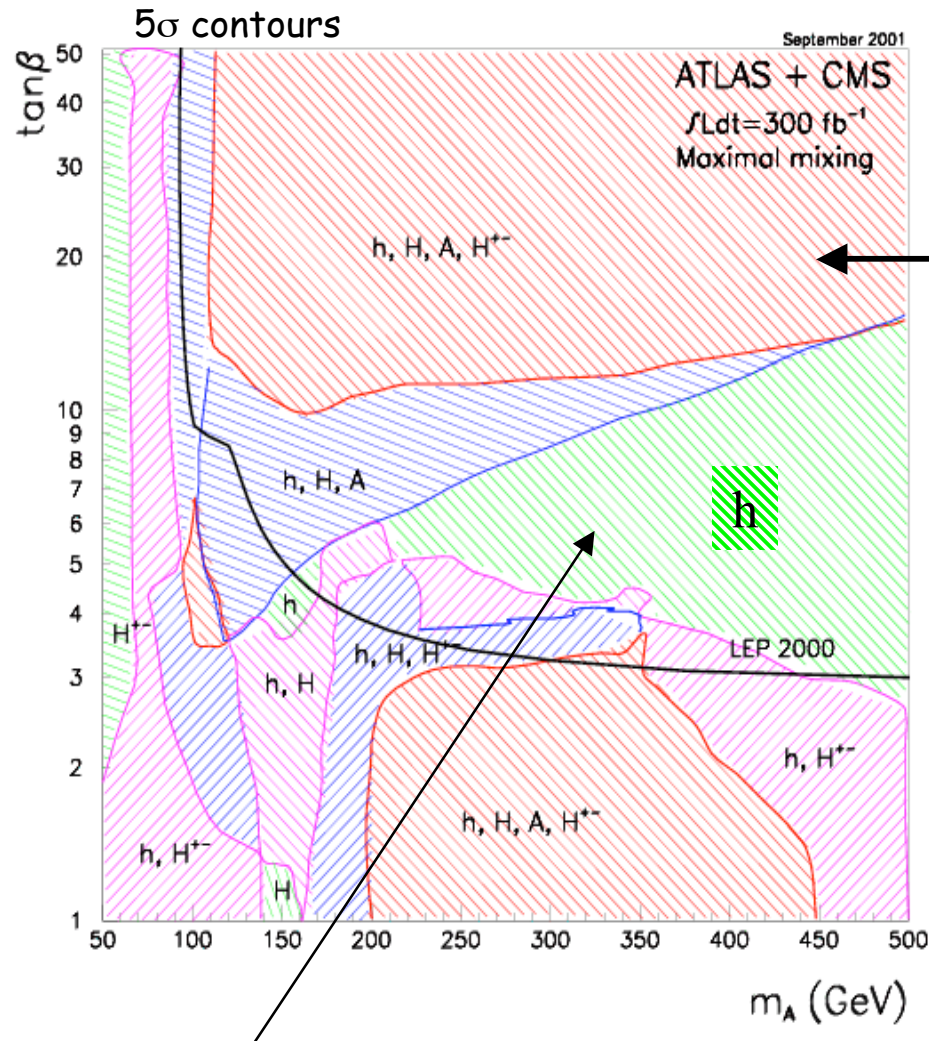


**LHC:  $\lambda = 0$  may be excluded at 95% CL.**

**SLHC:  $\lambda$  may be determined to 20-30% (95% CL)**

SUSY Higgs sector :  $h, H, A, H^\pm$

$m_h < 135 \text{ GeV}, \quad m_A \approx m_H \approx m_{H^\pm}$



$H, A \rightarrow \mu\mu, \tau$   
 $H^\pm \rightarrow \tau\nu, tb$

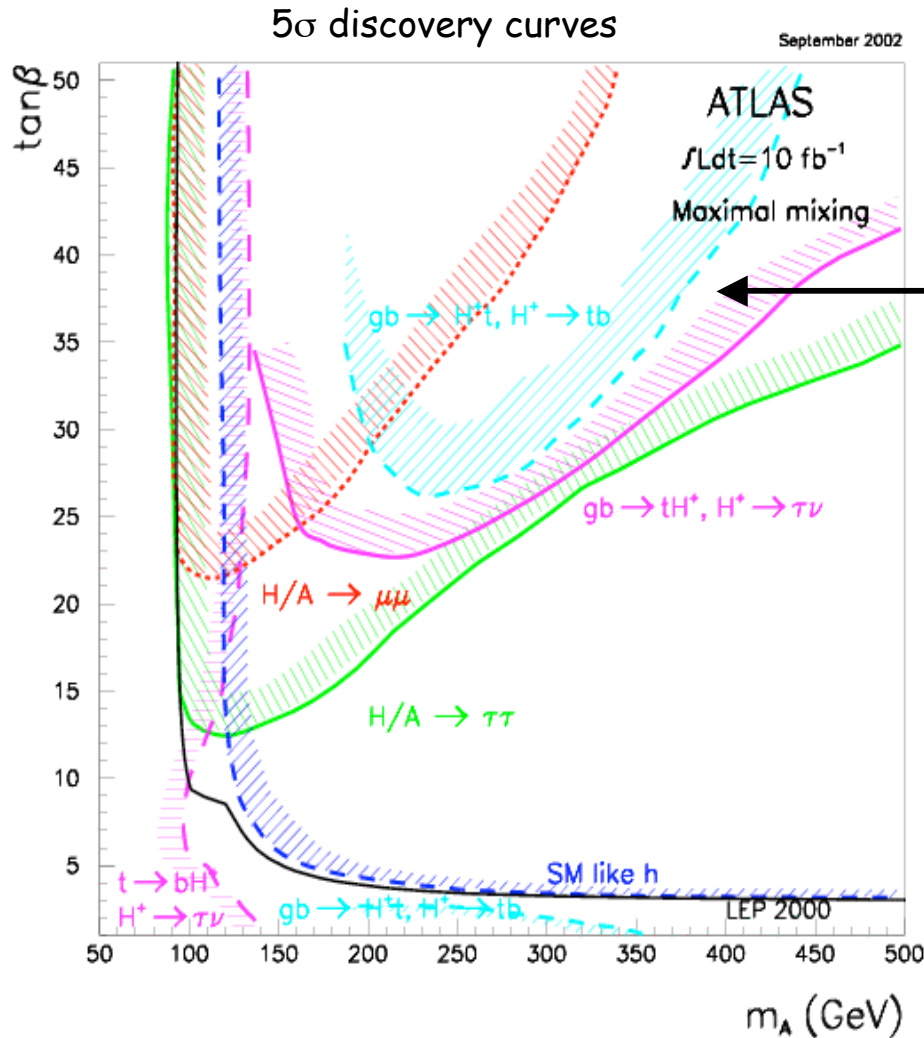
- 4 Higgs observable
- 3 Higgs observable
- 2 Higgs observable
- 1 Higgs observable

Assuming decays to  
 SM particles only

Here only  $h$  (SM - like) observable at LHC, unless  $A, H, H^\pm \rightarrow \text{SUSY}$   
 → LHC may miss part of the MSSM Higgs spectrum

Observation of full spectrum may require high-E ( $\sqrt{s} \approx 2 \text{ TeV}$ ) Lepton Collider

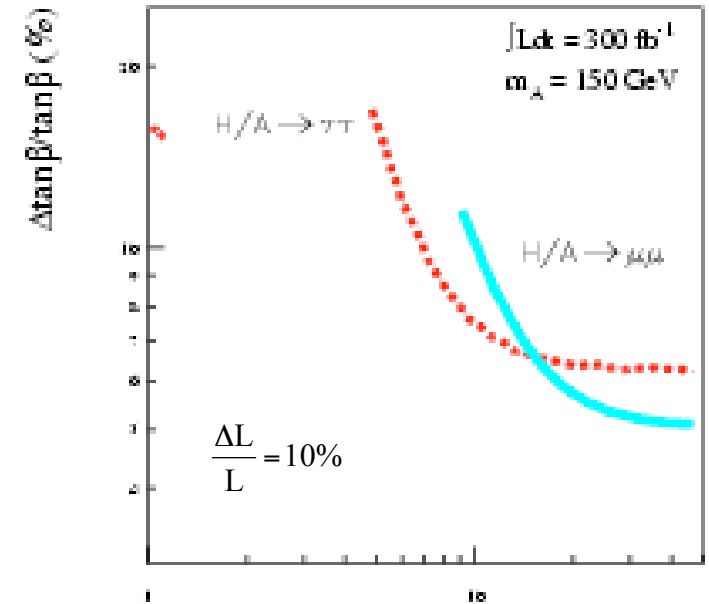
Most of MSSM Higgs plane already covered after 1 year of data taking ...



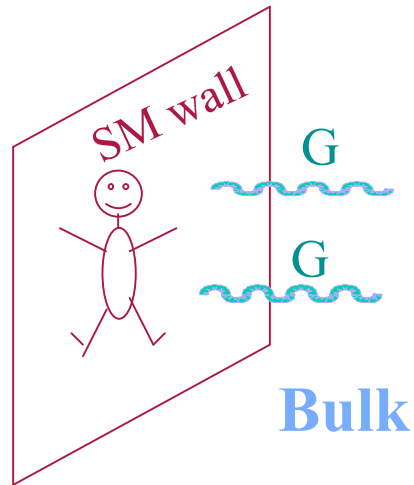
$A, H, H^\pm$  cross-section  $\sim tg^2\beta$



Measurement of  $tg\beta$



Large variety of channels and signatures accessible



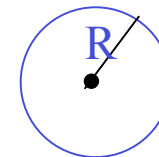
If gravity propagates  
in  $4 + \delta$  dimensions,  
a gravity scale  $M_D \approx 1 \text{ TeV}$  is possible

$$\left. \begin{aligned} V_4(r) &\sim \frac{1}{M_{\text{Pl}}^2} \frac{1}{r} \\ V_{4+\delta}(r) &\sim \frac{1}{M_D^{\delta+2}} \frac{1}{R^\delta} \frac{1}{r} \end{aligned} \right\} \text{at large distance}$$

$$M_{\text{Pl}}^2 \approx M_D^{\delta+2} R^\delta$$

- If  $M_D \approx 1 \text{ TeV}$  :
  - $\delta = 1$   $R \approx 10^{13} \text{ m}$   $\rightarrow$  excluded by macroscopic gravity
  - $\delta = 2$   $R \approx 0.7 \text{ mm}$   $\rightarrow$  limit of small- scale gravity experiments
  - ....
  - $\delta = 7$   $R \approx 1 \text{ Fm}$

$\rightarrow$  Extra-dimensions are compactified over  $R < \text{mm}$



- **Gravitons** in Extra-dimensions get **quantised mass**:

$$\left. \begin{aligned} m_k &\sim \frac{k}{R} & k = 1, \dots, \infty \\ \Delta m &\sim \frac{1}{R} & \text{e.g. } \Delta m \approx 400 \text{ eV } \delta = 3 \end{aligned} \right\} \rightarrow \begin{array}{l} \text{continuous tower} \\ \text{of massive gravitons} \\ \text{(Kaluza - Klein excitations)} \end{array}$$

$$\sigma \left[ \begin{array}{c} f \\ \swarrow \\ \text{G} \\ \nwarrow \\ f \end{array} \right] \approx \frac{1}{M_{\text{Pl}}^2} N_{\text{kk}} \approx \frac{1}{M_{\text{Pl}}^2} \left( \frac{\sqrt{s}}{\Delta m} \right)^\delta \approx \frac{1}{M_{\text{Pl}}^2} \sqrt{s}^\delta R^\delta \approx \frac{\sqrt{s}^\delta}{M_{\text{D}}^{\delta+2}}$$

Due to the large number of  $G_{\text{kk}}$ , the coupling SM particles - Gravitons becomes of EW strength



- Only one scale in particle physics : EW scale
- Can test geometry of universe and quantum gravity in the lab



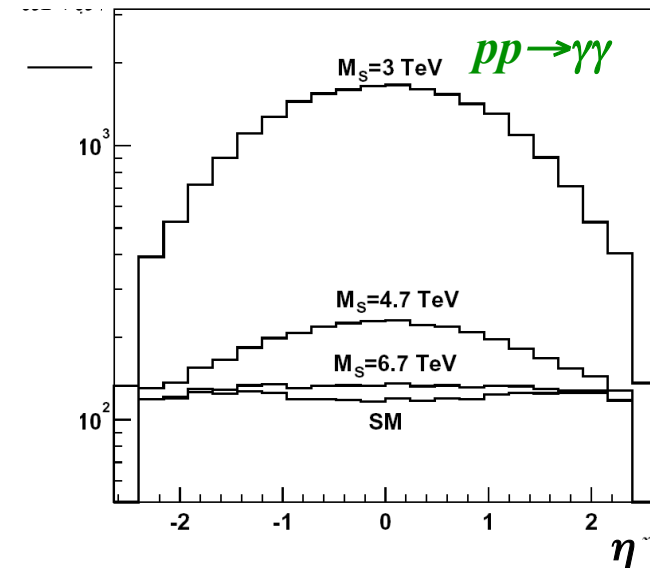
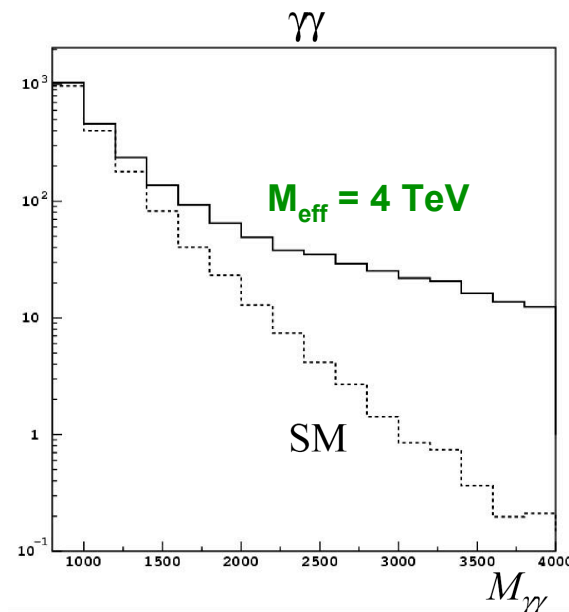
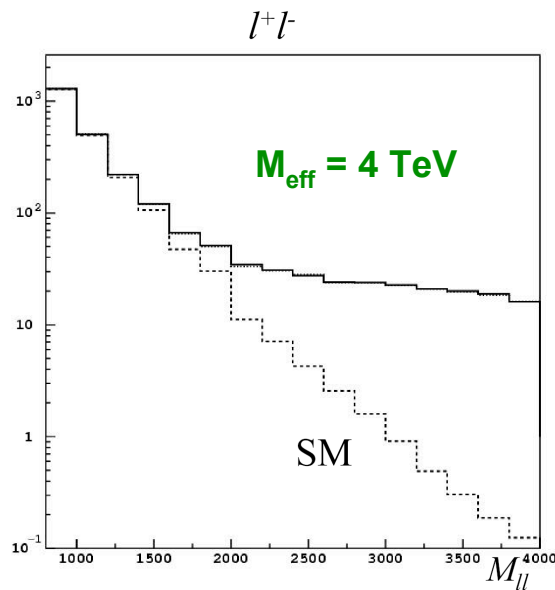


# Large Extra Dimensions

## *Virtual Exchange of Kaluza-Klein Gravitons*

Signatures:  $qq, gg \rightarrow \gamma\gamma, ll, (WW, tt, \dots)$

- excess over DY events in di-lepton, di-photon mass distributions
- some s-channel processes not present at tree-level in SM:  
 → more central production for  $\gamma\gamma$

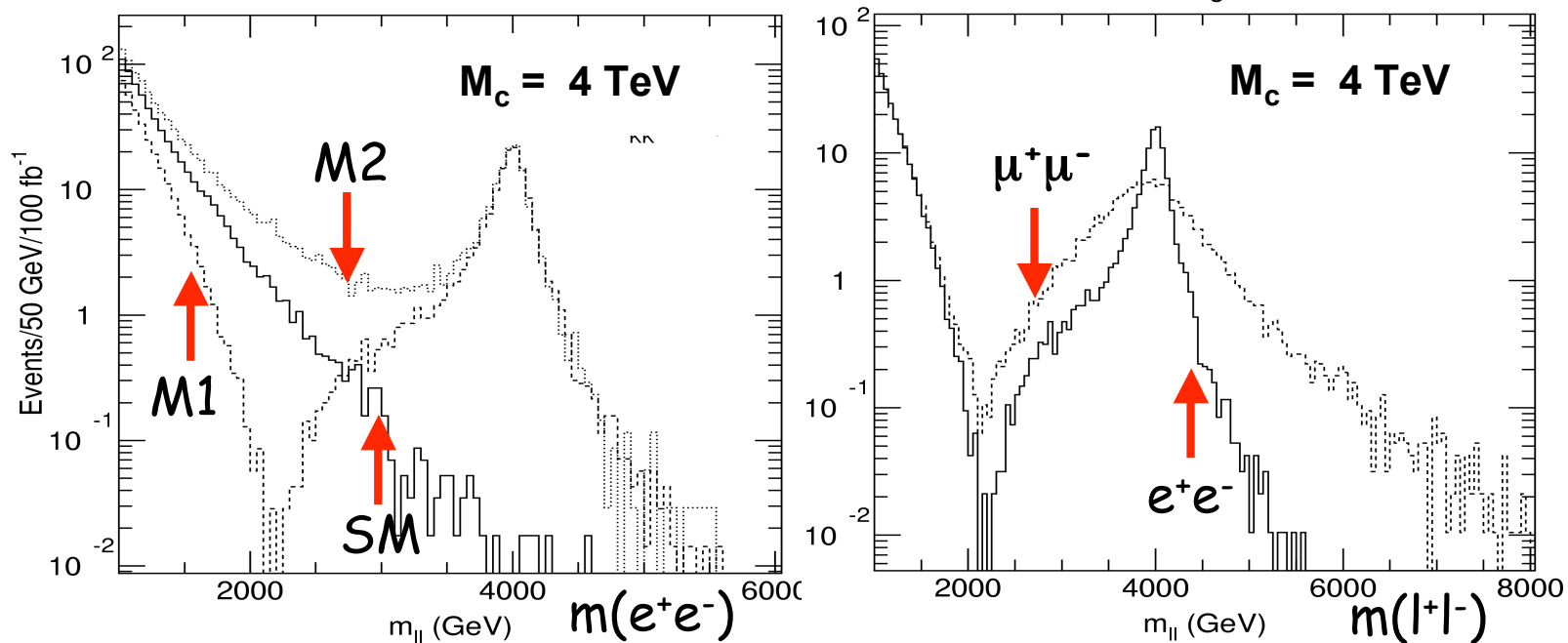


# TeV<sup>-1</sup>-sized Extra Dimensions

## $\gamma^{(1)}/Z^{(1)}$ Kaluza-Klein Gauge Bosons

Signatures:  $\gamma^{(1)}/Z^{(1)} \rightarrow e^+e^-, \mu^+\mu^-$

- 2 TeV electron in ATLAS:  $\Delta E/E \sim 0.7\%$  ( $\sim 20\%$  for a muon)
- acceptance for leptons:  $|\eta| < 2.5$

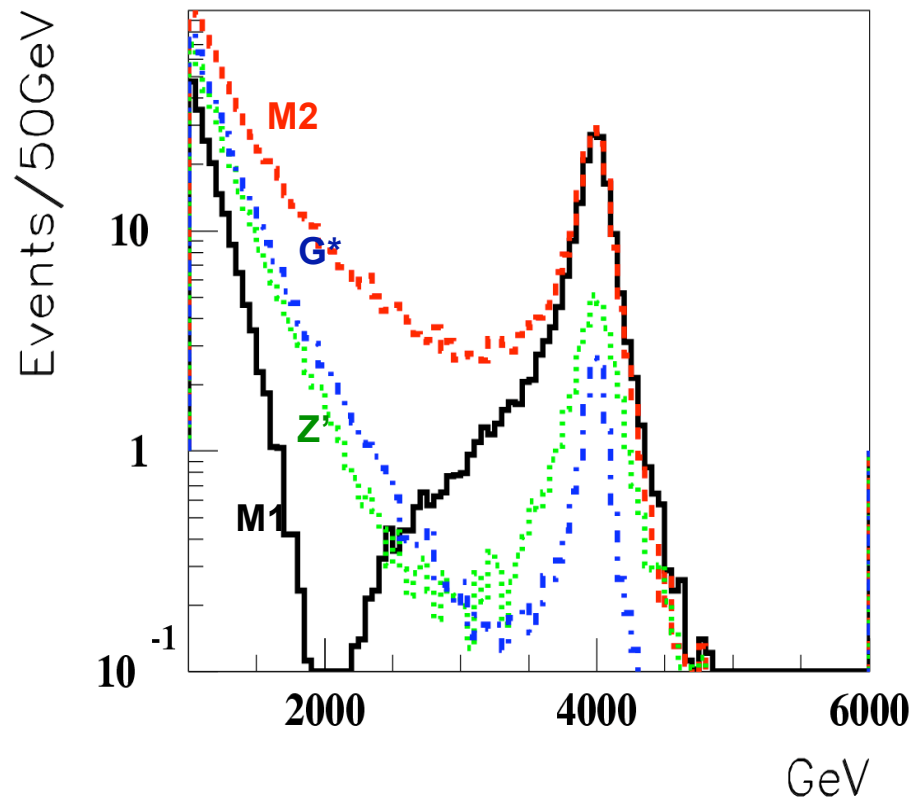


# TeV<sup>-1</sup>-sized Extra Dimensions

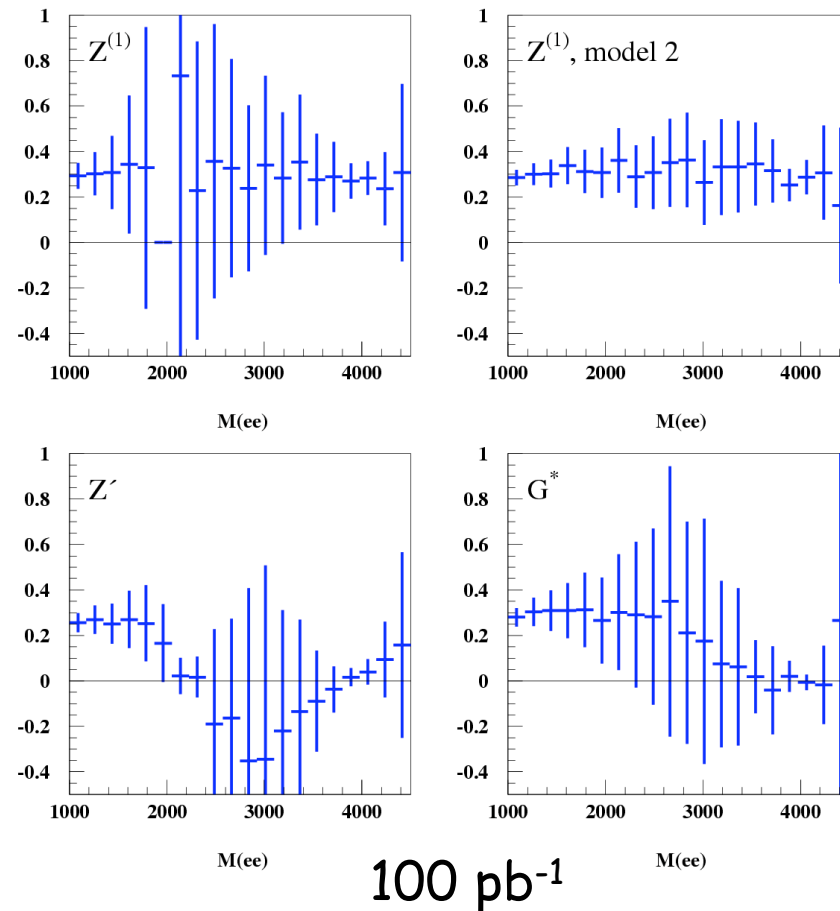
## $\gamma^{(1)}/Z^{(1)}$ Kaluza-Klein Gauge Bosons

Characterization of the model:

$Z^{(1)}$  or  $Z'$  or RS graviton ??

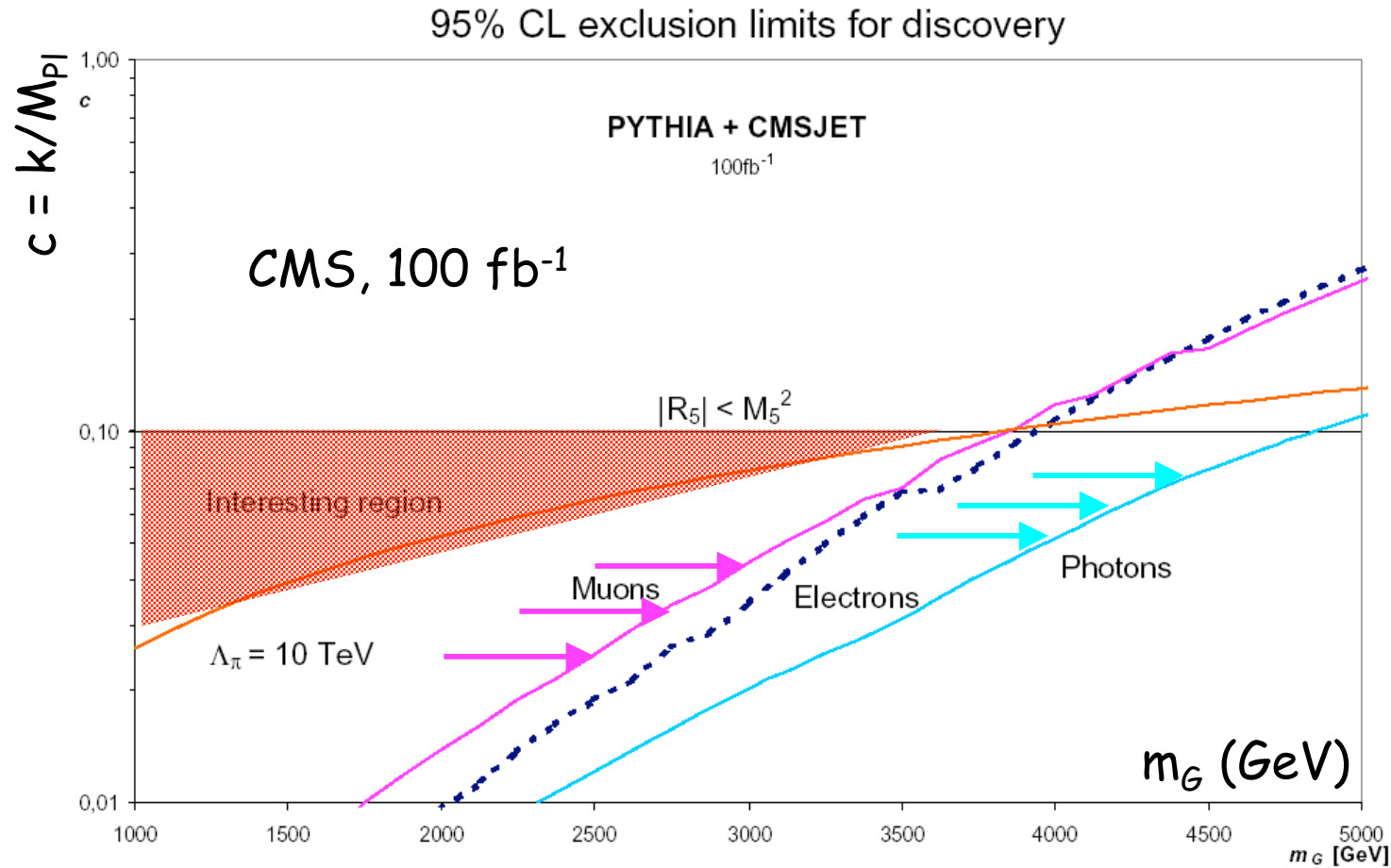


Forward-backward asymmetries:



# Warped Extra Dimension

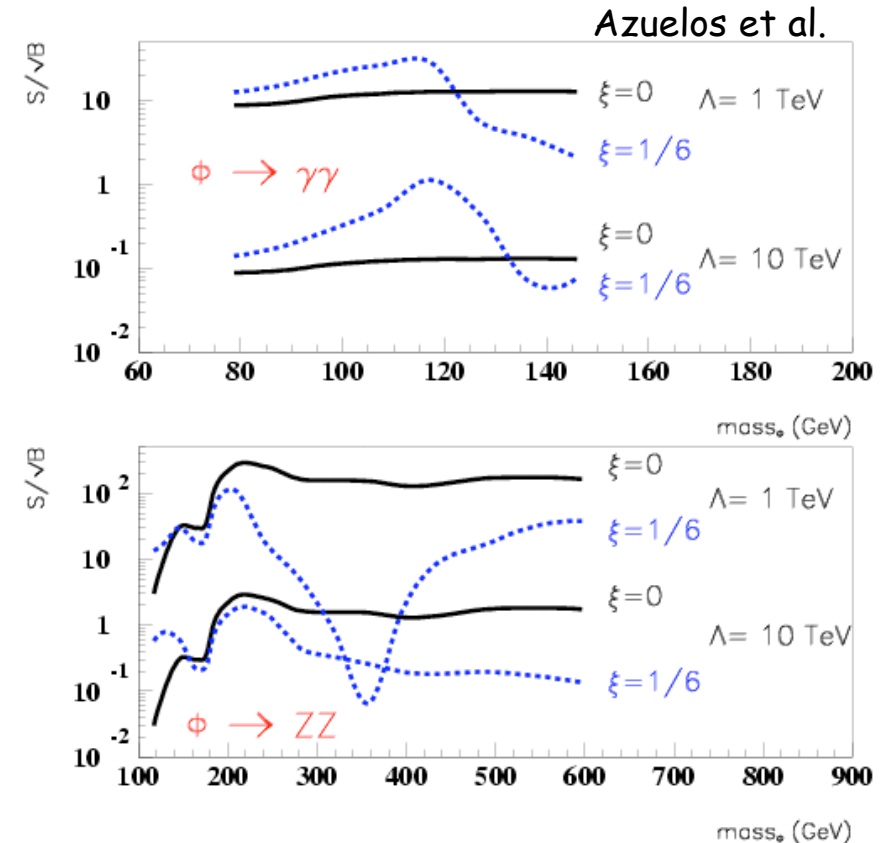
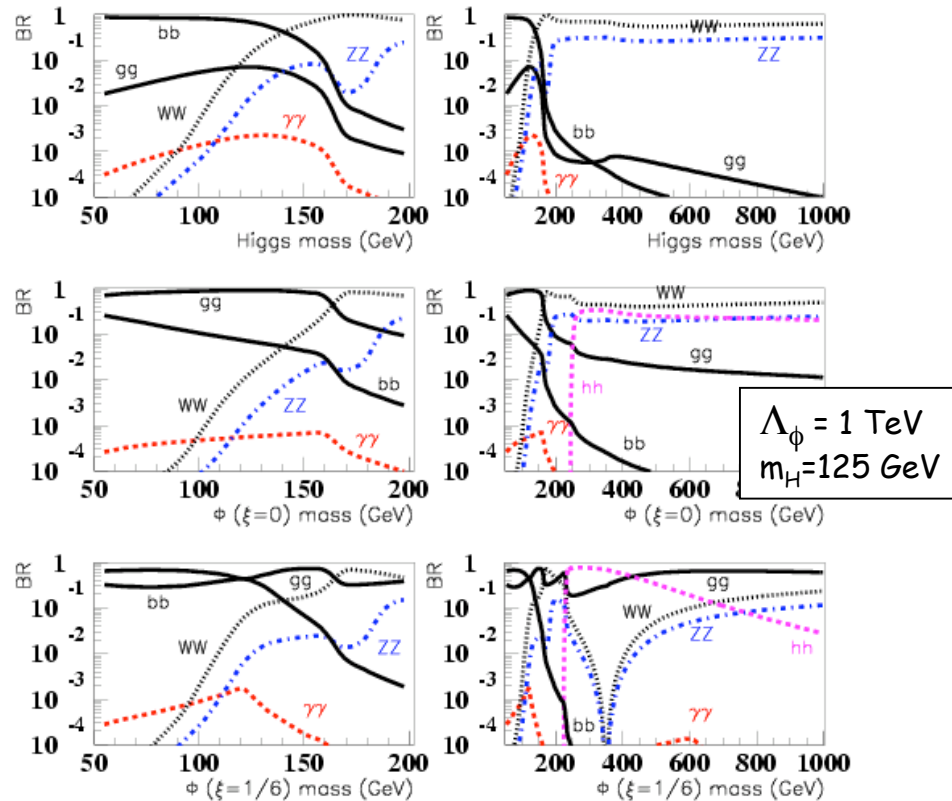
*Randall-Sundrum model: KK graviton narrow resonance*



→ LHC covers completely the interesting region

## Search for the Radion

- Scalar field which stabilises the size of the extra-dimension
- Parameters : radion mass ( $m_\phi$ ), radion vev ( $\Lambda_\phi$ ), h- $\phi$  mixing ( $\xi$ )



- Similar couplings as SM Higgs but with different strengths ( $\phi$ gg enhanced,  $\phi$ WW/ZZ suppressed in some cases);  $\phi \rightarrow hh$  important if open.  $\Gamma_\phi \ll \Gamma_H$
  - Precise measurements of couplings needed to disentangle  $\phi / H$ . LHC : 10-20%.
- More work needed here ...
- Impact on cosmological constant ?  $\rho_{\text{radion}} \sim M_D^8 / M_P^4 \approx \rho_\Lambda$

## Mini black holes production at LHC ?

... quite speculative for the time being ... many big theoretical uncertainties

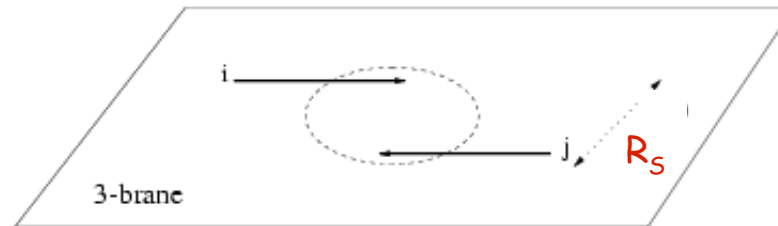
- Schwarzschild radius (i.e. within which nothing escapes gravitational force):

$$4\text{-dim.}, M_{\text{gravity}} = M_{\text{Planck}} \quad : \quad R_S \sim \frac{2}{M_{\text{Pl}}^2} \frac{M_{\text{BH}}}{c^2}$$

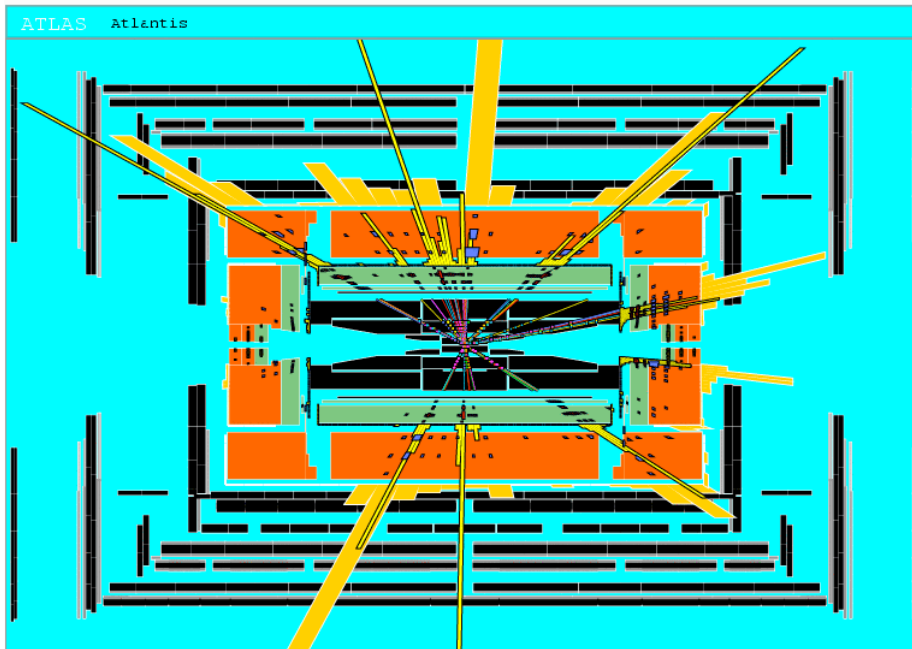
$$4 + \delta\text{-dim.}, M_{\text{gravity}} = M_D \sim \text{TeV} \quad : \quad R_S \sim \frac{1}{M_D} \left( \frac{M_{\text{BH}}}{M_D} \right)^{\frac{1}{\delta+1}}$$



Since  $M_D$  is low, tiny black holes of  $M_{\text{BH}} \sim \text{TeV}$  can be produced if partons  $ij$  with  $\sqrt{s_{ij}} = M_{\text{BH}}$  pass at a distance smaller than  $R_S$



- Large partonic cross-section :  $\sigma(ij \rightarrow \text{BH}) \sim \pi R_S^2$   
e.g. for  $M_D \sim 1 \text{ TeV}$  expected rate is  $\sim 1 \text{ Hz}$ !
- Black holes decay immediately ( $\tau \sim 10^{-26} \text{ s}$ ) by Hawking radiation (democratic evaporation) :
  - large multiplicity
  - small missing  $E$
  - jets/leptons  $\sim 5$
 } expected signature (quite spectacular ...)



A black hole event with  $M_{\text{BH}} \sim 8 \text{ TeV}$  in ATLAS

From preliminary studies : reach is  $M_{\text{D}} \sim 4 \text{ TeV}$  for any  $\delta$  in one year at high luminosity.

By testing Hawking formula  $\rightarrow$  proof that it is BH + measurement of  $\delta$

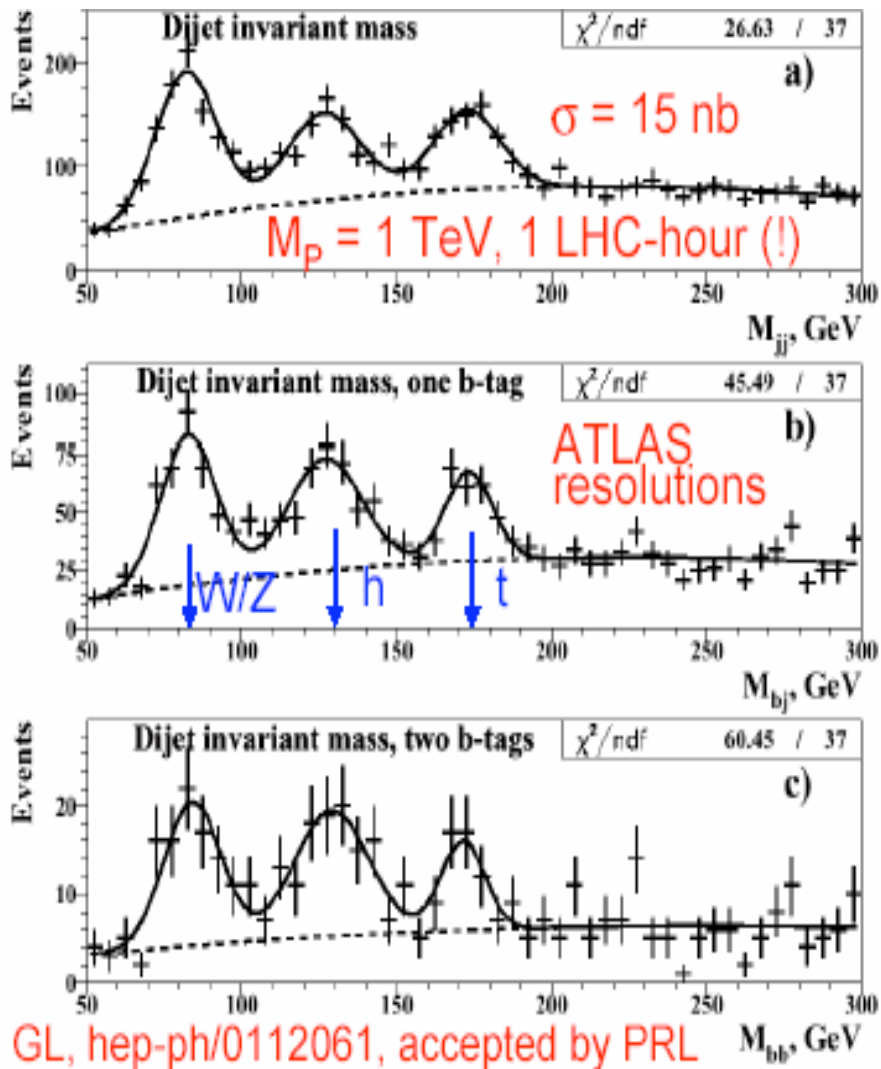
$$\log T_{\text{H}} = -\frac{1}{\delta + 1} \log M_{\text{BH}} + f(M_{\text{D}}, \delta)$$

precise measurements of  $M_{\text{BH}}$  and  $T_{\text{H}}$  needed  
( $T_{\text{H}}$  from lepton and photon spectra)

$M_{\text{D}}$  from cross-section

Note: mini-BH should also be produced by ultra-high-energy cosmic neutrinos and observed by Auger (Feng and Shapere, hep-ph/0109106)

# Higgs production in black hole decays



Higgs of mass 130 GeV may be discovered in  
 one hour ( $M_p = 1 \text{ TeV}$ )  
 one day ( $M_p = 2 \text{ TeV}$ )  
 one month ( $M_p = 3 \text{ TeV}$ )  
 one year ( $M_p = 4 \text{ TeV}$ )  
 at  $10^{34}$



$$pp \rightarrow Z_R \rightarrow N_i N_i \rightarrow l j j l j j$$

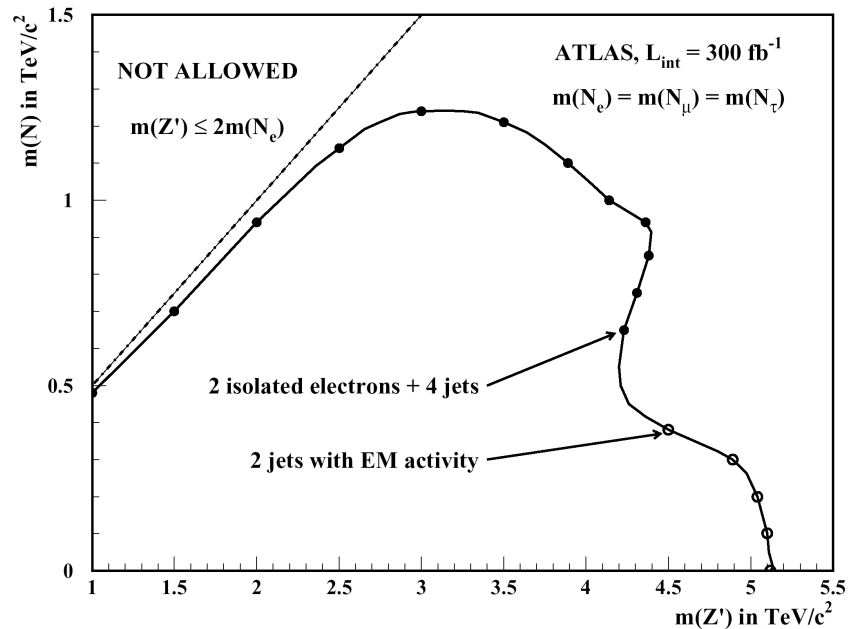
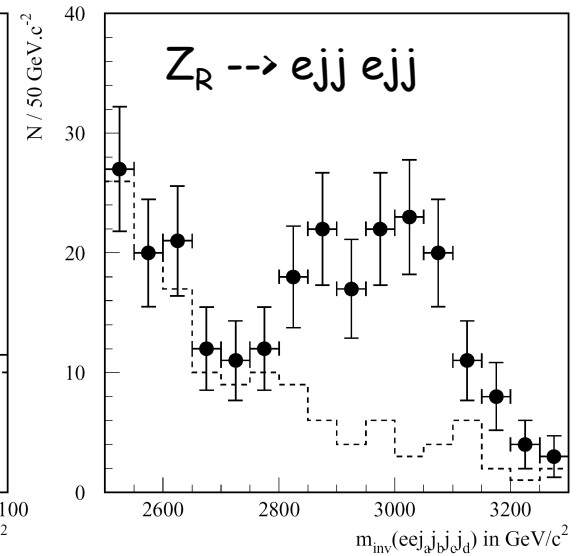
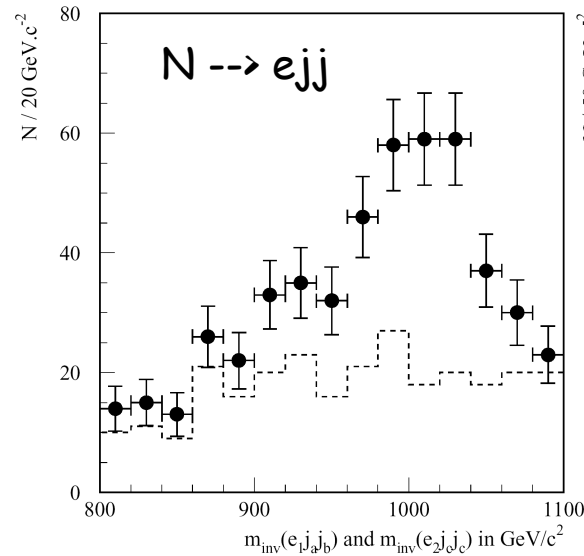
backgrounds:

*t tbar*

*DY, WW, ZW, ZZ*

*LRSM bckg: W<sub>R</sub>...*

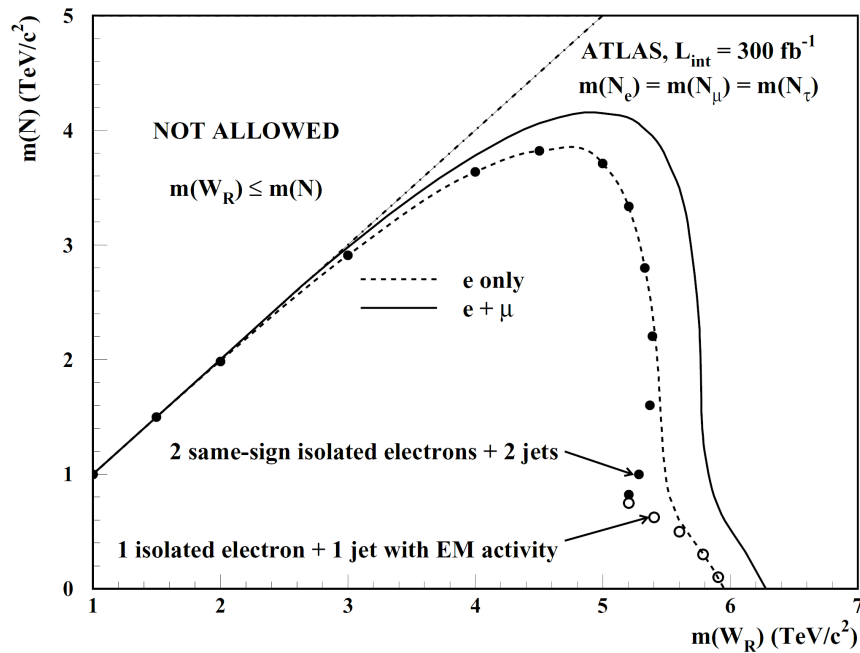
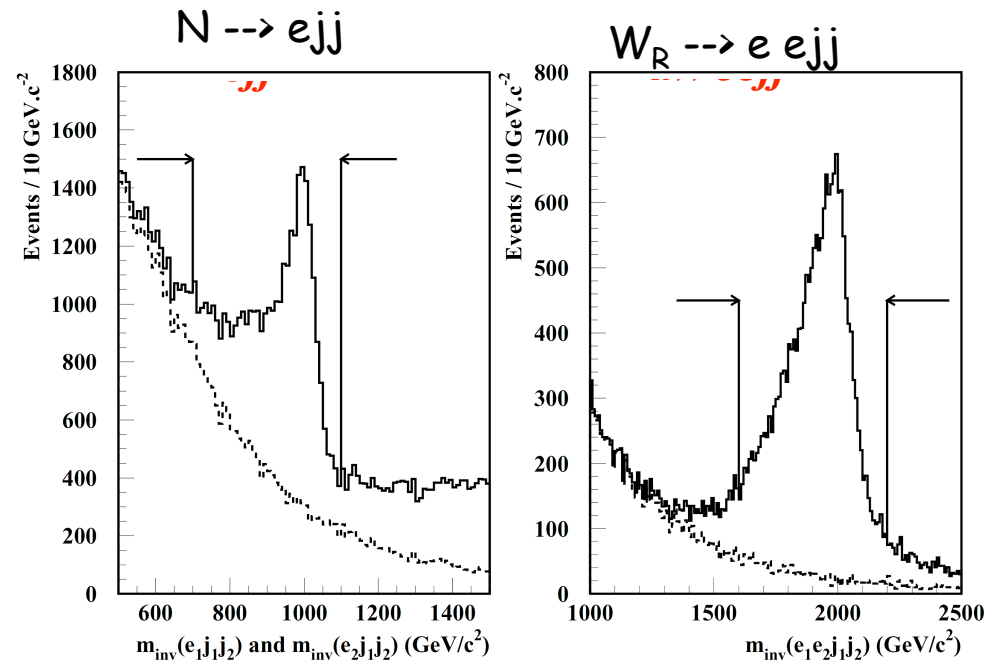
⇒ cuts on  
 $m_{ee}$ ,  $pT(\text{jets})$



FB asymmetry gives a  
measure of  $\kappa = g_R/g_L$

$$pp \rightarrow W_R \rightarrow | N_l \rightarrow | l j j$$

backgrounds:  
*t tbar*  
*DY, WW, ZW, ZZ*



# Heavy leptons

**backgrounds:**  
 $t\bar{t}$ ,  $WW$ ,  $WZ$ ,  $ZZ$

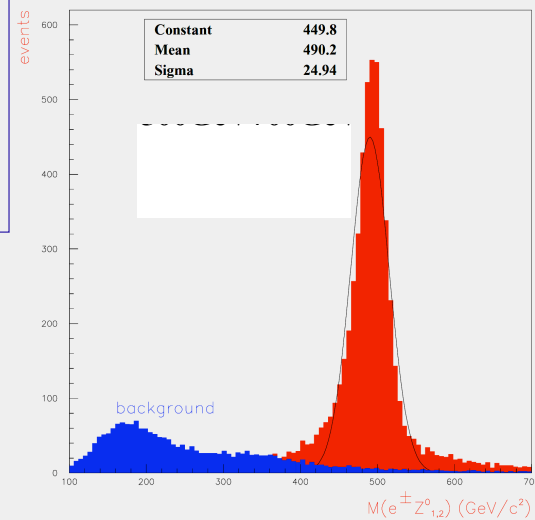
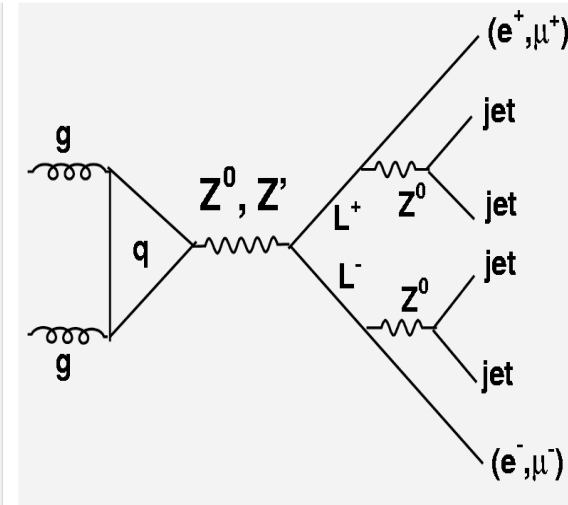
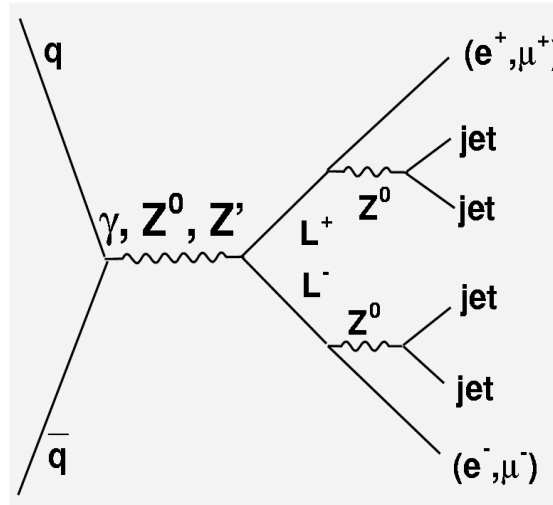
**also:**  
 6-lepton channel

**Experimental considerations:**  
 - high energy leptons, jets

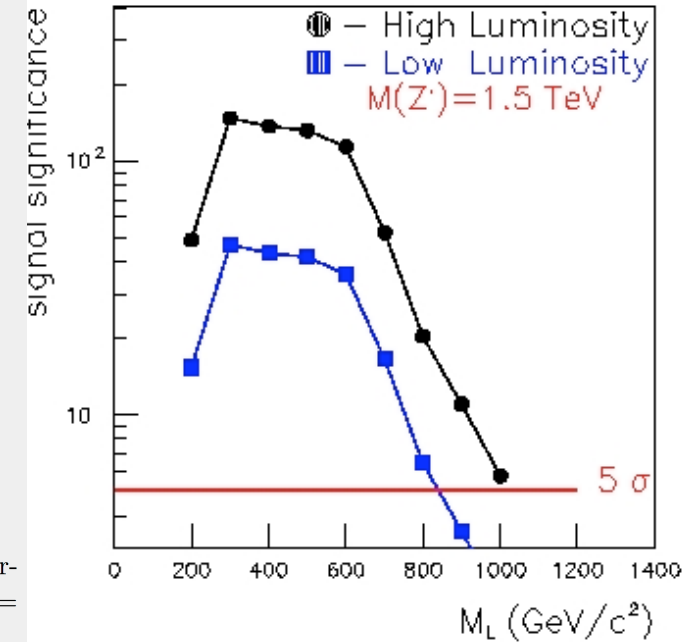
**Systematics:**  
 - large NLO corrections

## conclusion:

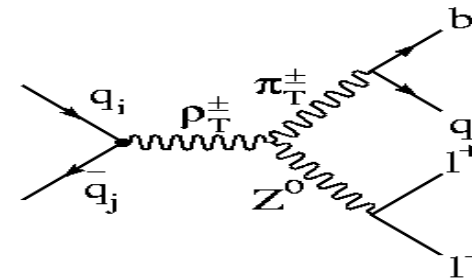
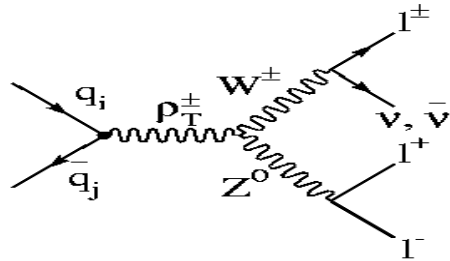
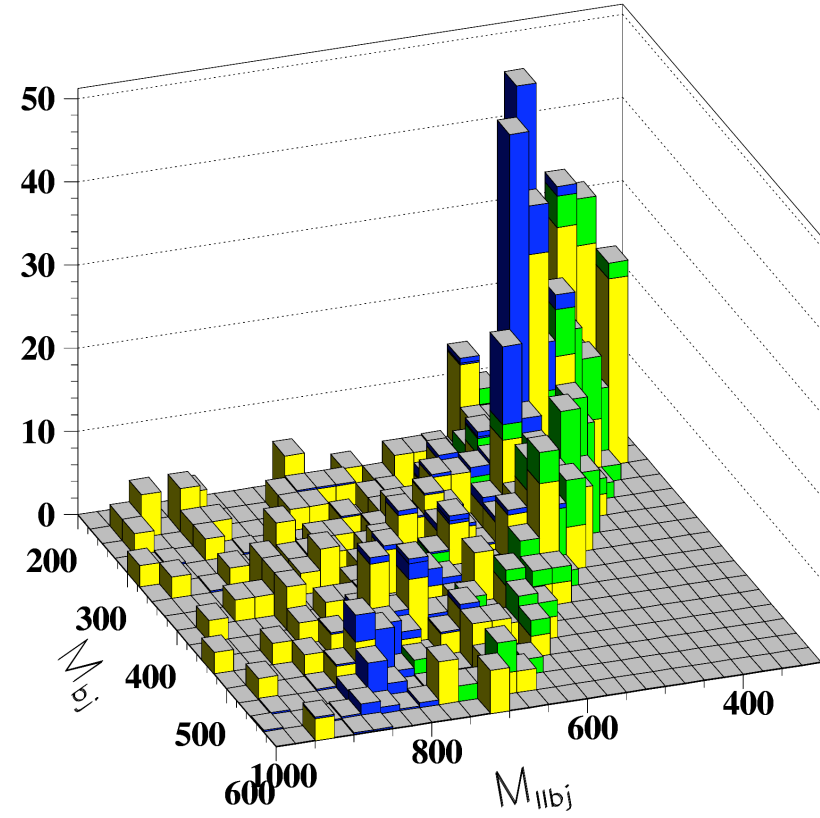
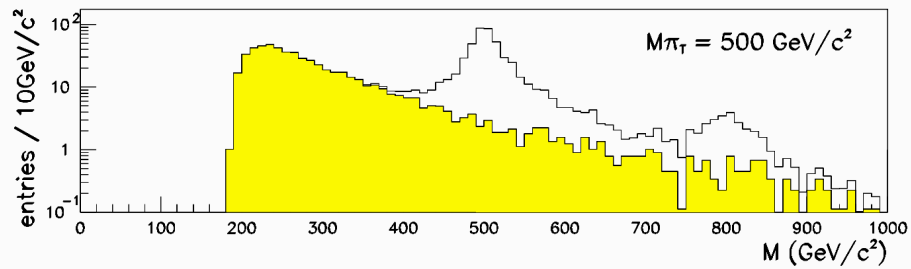
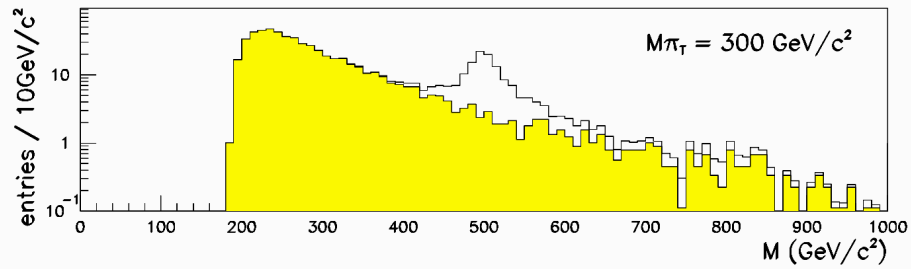
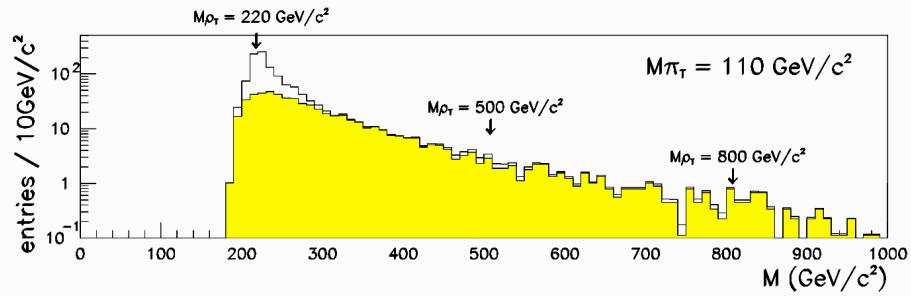
ATLAS can discover this sequential charged heavy leptons up to  $M_L = 0.9 / 1.0$  TeV (low/high luminosity)



**Figure 13:** Signal to background comparison for  $M_L = 0.5$  TeV/ $c^2$  and  $M_{Z'} = 0.7$  TeV/ $c^2$ , for  $L \rightarrow e + Z^0$  channel.

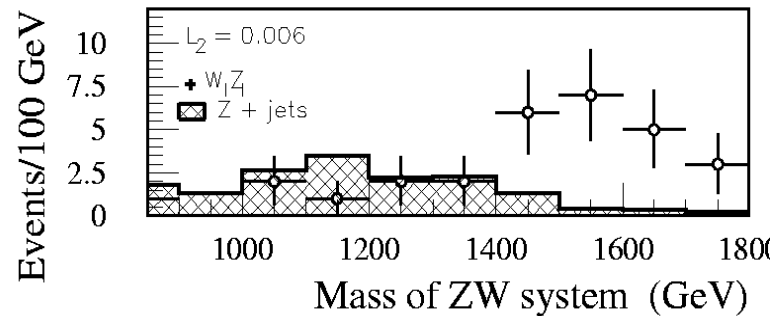
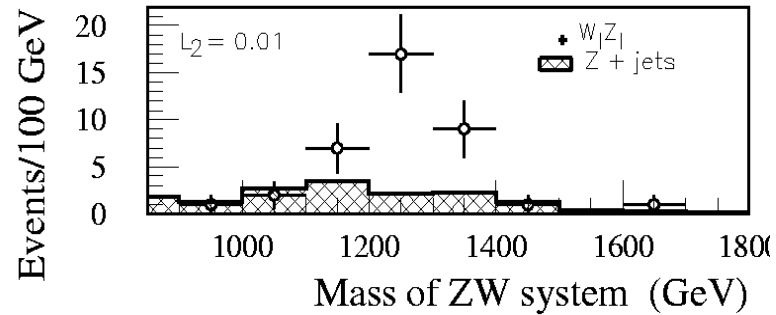


# Technicolour

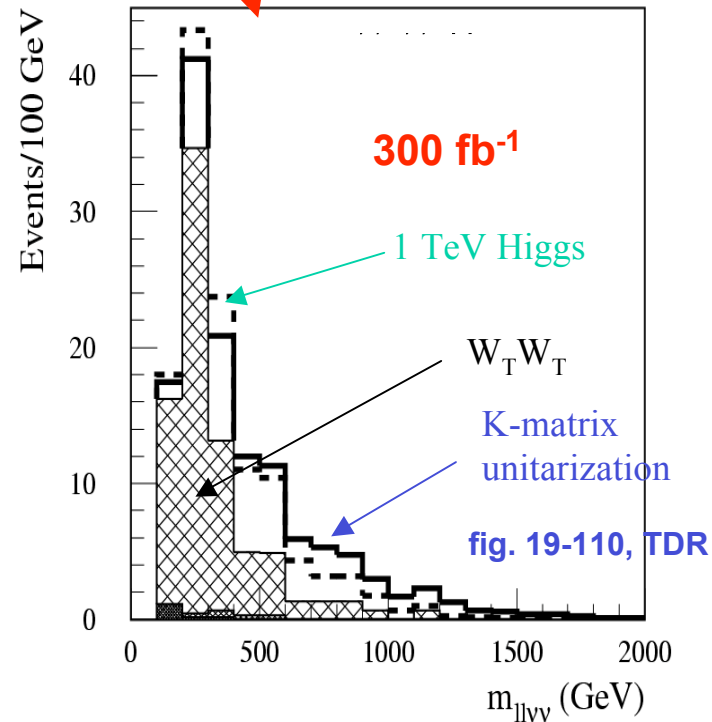


# Strong symmetry breaking : $V_L V_L$ scattering

Resonant scattering  
 $qq \rightarrow qq$   $WZ \rightarrow qq$   $jj$   $ll$



*Non-resonant  $VV$  scattering:  
 needs high luminosity, good  
 understanding of  
 backgrounds*

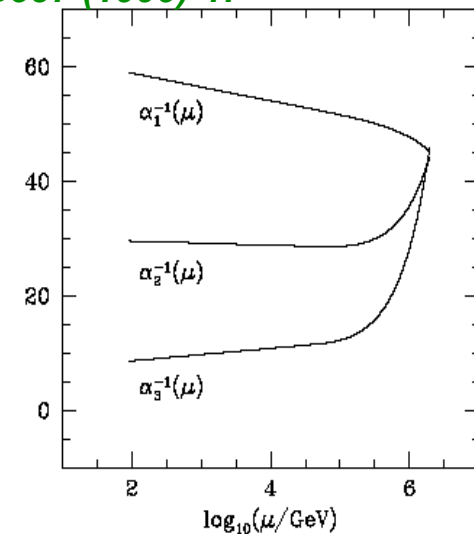
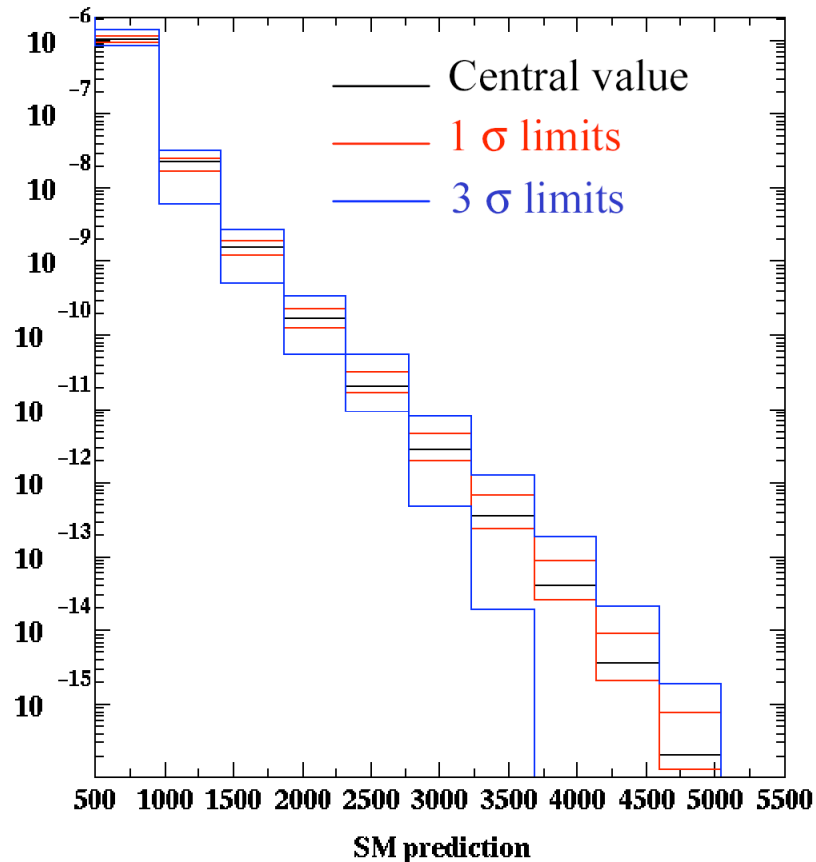


# Coupling Unification at TeV

- KK states affect running of gauge couplings
- Dijet cross-section

K.R. Dienes, E. Dudas and T. Gherghetta,  
*Nucl.Phys. B537 (1999) 47*

SM prediction zones



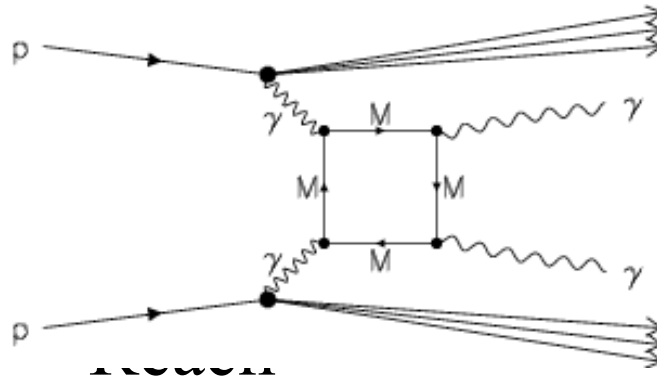
-Sensitivity of deficit in jet cross section,  $\sim 10$  TeV, at parton level

-PDF uncertainties limit reach to 1 TeV

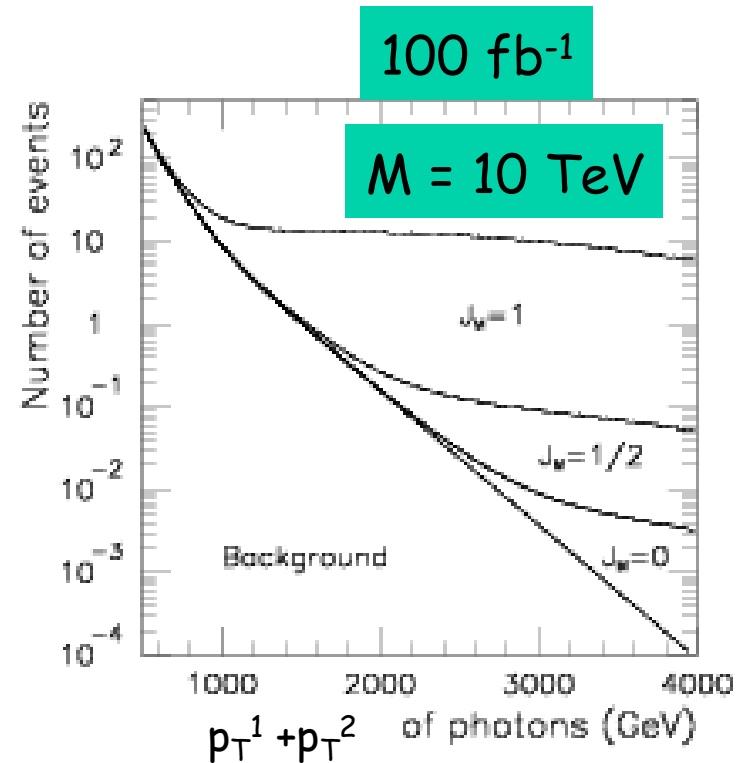
# Monopoles

- Motivation

- Restores Maxwell's equ'ns symmetry & explains charge quantization



- 10-20 TeV  
(spin dependent)



	LHC	LC	LHC+LC	SPS1a
$\tan \beta$	10.22±9.1	10.26±0.3	10.06±0.2	10
$M_1$	102.45±5.3	102.32±0.1	102.23±0.1	102.2
$M_2$	191.8±7.3	192.52±0.7	191.79±0.2	191.8
$M_3$	578.67±15	fixed 500	588.05±11	589.4
$M_{\tilde{\tau}_L}$	fixed 500	197.68±1.2	199.25±1.1	197.8
$M_{\tilde{\tau}_R}$	129.03±6.9	135.66±0.3	133.35±0.6	135.5
$M_{\tilde{\mu}_L}$	198.7±5.1	198.7±0.5	198.7±0.5	198.7
$M_{\tilde{\mu}_R}$	138.2±5.0	138.2±0.2	138.2±0.2	138.2
$M_{\tilde{e}_L}$	198.7±5.1	198.7±0.2	198.7±0.2	198.7
$M_{\tilde{e}_R}$	138.2±5.0	138.2±0.05	138.2±0.05	138.2
$M_{\tilde{q}^3_L}$	498.3±110	497.6±4.4	521.9±39	501.3
$M_{\tilde{t}_R}$	fixed 500	420±2.1	411.73±12	420.2
$M_{\tilde{b}_R}$	522.26±113	fixed 500	504.35±61	525.6
$M_{\tilde{q}^2_L}$	550.72±13	fixed 500	553.31±5.5	553.7
$M_{\tilde{c}_R}$	529.02±20	fixed 500	531.70±15	532.1
$M_{\tilde{s}_R}$	526.21±20	fixed 500	528.90±15	529.3
$M_{\tilde{q}^1_L}$	550.72±13	fixed 500	553.32±6.5	553.7
$M_{\tilde{u}_R}$	528.91±20	fixed 500	531.70±15	532.1
$M_{\tilde{d}_R}$	526.2±20	fixed 500	528.90±15	529.3
$A_\tau$	fixed 0	-202.4±89.5	352.1±171	-253.5
$A_t$	-507.8±91	-501.95±2.7	-505.24±3.3	-504.9
$A_b$	-784.7±35603	fixed 0	-977±12467	-799.4
$m_A$	fixed 500	399.1±0.9	399.1±0.8	399.1
$\mu$	345.21±7.3	344.34±2.3	344.36±1.0	344.3

Table 5.29: Result for the general MSSM parameter determination in SPS1a using the mass measurements given in Tab. 5.25. Shown are the nominal parameter values and the result after fits to the different data sets. All masses are given in GeV.