The $a_0 - a_2$ pion scattering length from $K^+ \to \pi^+ \pi^0 \pi^0$ decay

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A "commercial" for $K_+ \to \pi^+ \nu \bar{\nu}$ and $K_L \to \pi^0 \nu \bar{\nu}$

 $K_+ \to \pi^+ \nu \bar{\nu} \text{ and } K_L \to \pi^0 \nu \bar{\nu} :$ a unique physics opportunity

Insight on CKM and CP violation equivalent to BELLE + BABAR

Short distance reach comparable to LEP200 but

Independent of either programs!

A "commercial" for Neutral Hyperons

We can put the final word on Ξ^0 and Λ beta decay

And put the final word on V_{us} ... and flavor symmetry breaking

Summary

- Pion scattering lengths.
- The method.
- Theoretical precision.
- Further theoretical work.

N. Cabibbo, arXiv: hep-ph0405001

Pion scattering lengths and chiral dynamics

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Weinberg 1966

$$a_0 m_{\pi^+} = \frac{7m_{\pi^+}^2}{16\pi f_{\pi}^2} = 0.159$$
$$a_2 m_{\pi^+} = \frac{-m_{\pi^+}^2}{8\pi f_{\pi}^2} = -0.045$$

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Colangelo et al. 2001 :
$$a_0 m_{\pi^+} = 0.220 \pm 0.005$$

 $a_2 m_{\pi^+} = -0.0444 \pm 0.0010$

$$(a_0 - a_2)m_{\pi^+} = 0.265 \pm 0.004$$

Can we match this precision?



$\pi^+\pi^- \to \pi^0\pi^0$ re-scattering from $K^+ \to \pi^+\pi^+\pi^-$.

$$K^+ \to \pi^+ \pi^0 \pi^0$$
 and $K^+ \to \pi^+ \pi^+ \pi^-$

7

The unperturbed amplitudes are well known — see PDG, but with the best data we would like to fit our own. A simple parametrization:

$$K^{+} \to \pi^{+} \pi^{0} \pi^{0} \qquad : \qquad \mathcal{M}_{0}(s_{1}, s_{2}, s_{3}) = A_{\mathrm{av}}^{0}(1 + g^{0}(s_{\pi\pi} - s_{0})/2m_{\pi^{+}}^{2})$$
$$K^{+} \to \pi^{+} \pi^{+} \pi^{-} \qquad : \qquad \mathcal{M}_{+}(s_{1}, s_{2}, s_{3}) = A_{\mathrm{av}}^{+}(1 + g^{+}(s_{\pi\pi} - s_{0})/2m_{\pi^{+}}^{2}),$$

with g^0, g^+ the linear slope parameters.

Let us write:

$$\mathcal{M}(K^+ \to \pi^+ \pi^0 \pi^0) = \mathcal{M} = \mathcal{M}_0 + \mathcal{M}_1$$

where \mathcal{M}_1 is the contribution of the re-scattering graph.

We must consider two cases, according to whether $M_{\pi\pi}$ is above or below the $\pi^+\pi^-$ threshold.

$$s_{\pi\pi} > 4m_{\pi^+}^2 : \qquad \mathcal{M}_1 = i2\frac{(a_0 - a_2)m_{\pi^+}}{3}\mathcal{M}_{+,\text{thr}}\sqrt{(s_{\pi\pi} - 4m_{\pi^+}^2)/s_{\pi\pi}}$$
$$s_{\pi\pi} < 4m_{\pi^+}^2 : \qquad \mathcal{M}_1 = -2\frac{(a_0 - a_2)m_{\pi^+}}{3}\mathcal{M}_{+,\text{thr}}\sqrt{(4m_{\pi^+}^2 - s_{\pi\pi})/s_{\pi\pi}}$$

where $\mathcal{M}_{+,\text{thr}}$ is the value of \mathcal{M}_{+} at the $\pi^{+}\pi^{-}$ threshold.

- Above threshold: \mathcal{M}_1 imaginary
- Below threshold: \mathcal{M}_1 real and negative

A technical clarification

I-spin is badly broken between the $\pi^0 \pi^0$ and the $\pi^+ \pi^-$ threshold which is just the region we are interested in! How to define I=0 and I=2 scattering lengths?

> Solution we have adopted: Define $(a_0 - a_2)/3$ through:

$$M(\pi^{+}\pi^{-} \to \pi^{0}\pi^{0})|_{\text{S-wave}} = imv_{\pi^{+}}(a_{0} - a_{2})/3 + O(v_{\pi^{+}}^{2})$$

where $v_{\pi^{+}} = \sqrt{(s_{\pi\pi} - 4m_{\pi^{+}}^{2})/s_{\pi\pi}}$

The key to the measurement

Below threshold there is an interference term which is absent above the threshold:

$$|\mathcal{M}|^{2} = \begin{cases} (\mathcal{M}_{0})^{2} + (\mathcal{M}_{1})^{2} + 2\mathcal{M}_{0}\mathcal{M}_{1} & : \quad s_{\pi\pi} < 4m_{\pi^{+}}^{2} \\ (\mathcal{M}_{0})^{2} + |\mathcal{M}_{1}|^{2} & : \quad s_{\pi\pi} > 4m_{\pi^{+}}^{2} \end{cases}$$

The interference term is proportional to $a_0 - a_2$



The $s_{\pi\pi}$ invariant mass distribution with/without the re-scattering correction, in arbitrary units.



The $s_{\pi\pi}$ invariant mass distribution with/without the re-scattering correction, in arbitrary units.



The $s_{\pi\pi}$ invariant mass distribution (1/4 of existing NA48 data) with fit above the $\pi^+\pi^-$ threshold. As analyzed by Italo Mannelli.







Precision

• Below threshold define
$$\delta = \sqrt{\frac{4m_{\pi^+}^2 - s_{\pi\pi}}{4m_{\pi^+}^2}}$$

- The differential rate is predicted with errors $\sim \delta^3$
- The theoretical error on $a_0 a_2$ is $\sim \delta_{max}^2$

One can strike a balance between theoretical uncertainty and statistical error.



Theoretical precision can be improved by restricting the range of $s_{\pi\pi}$.

Space for improvement

- Experimental fit to terms $O(\delta^3)$ in differential rate.
- Compute these terms in Chiral Perturbation Theory
- Compute radiative corrections

The theoretical error can be made very small.