

Photoproduction of Heavy Vector Mesons at the LHC

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Physics at LHC

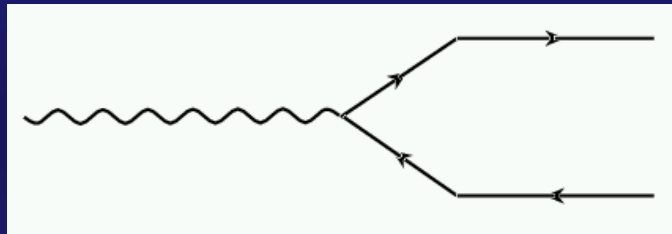
13-17 July 2004 . Vienna . Austria



Photoproduction of heavy Vector Mesons - A probe of the nucleon/nucleus gluon density

The photon wave function can be written as a Fock decomposition:

$$|\gamma\rangle = C_{\text{bare}} |\gamma_{\text{bare}}\rangle + C_{\rho} |\rho\rangle + C_{\omega} |\omega\rangle + C_{\phi} |\phi\rangle + \dots + C_q |qq\rangle$$

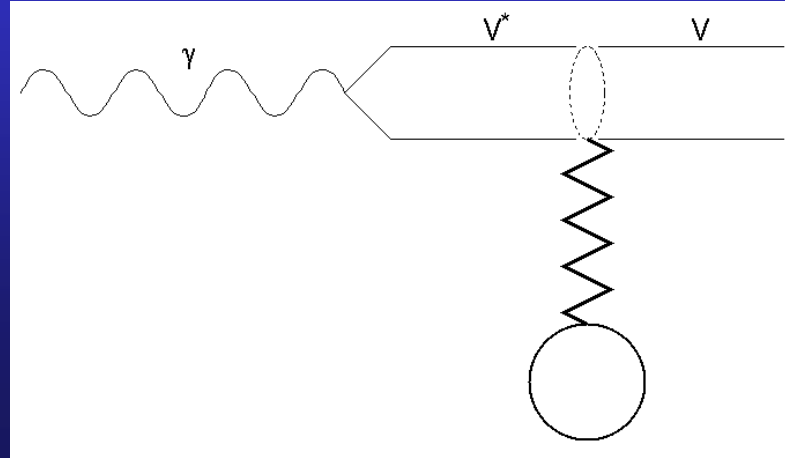


A photon is not always a photon...

With a certain probability it will appear as a $q\bar{q}$ fluctuation

Conservation of quantum numbers ($\gamma: J^P = 1^-$) \Rightarrow The photon tend to fluctuate to a vector meson ($\rho, \omega, \phi, J/\Psi$). Vector Meson Dominance.

A photon in the vector meson state will interact strongly (hadronically).
 The hadronic component can materialize if the virtual qq-pair is
 knocked on mass shell.



The probability to find the photon in the vector meson state V :

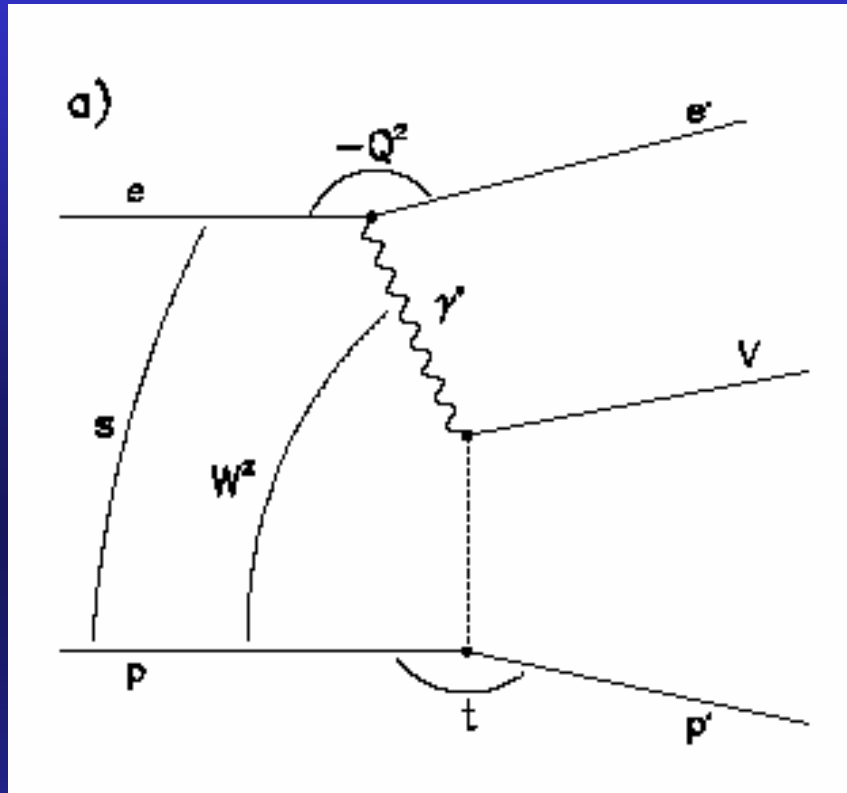
$$C_V = \frac{\sqrt{4\pi\alpha_{em}}}{f_V}$$

f_V – photon-vector-meson coupling

Vector Meson Dominance:

$$\frac{d\sigma}{dt}(\gamma A) = C_V^2 \frac{d\sigma}{dt}(VA)$$

Studied at HERA in ep collisions:

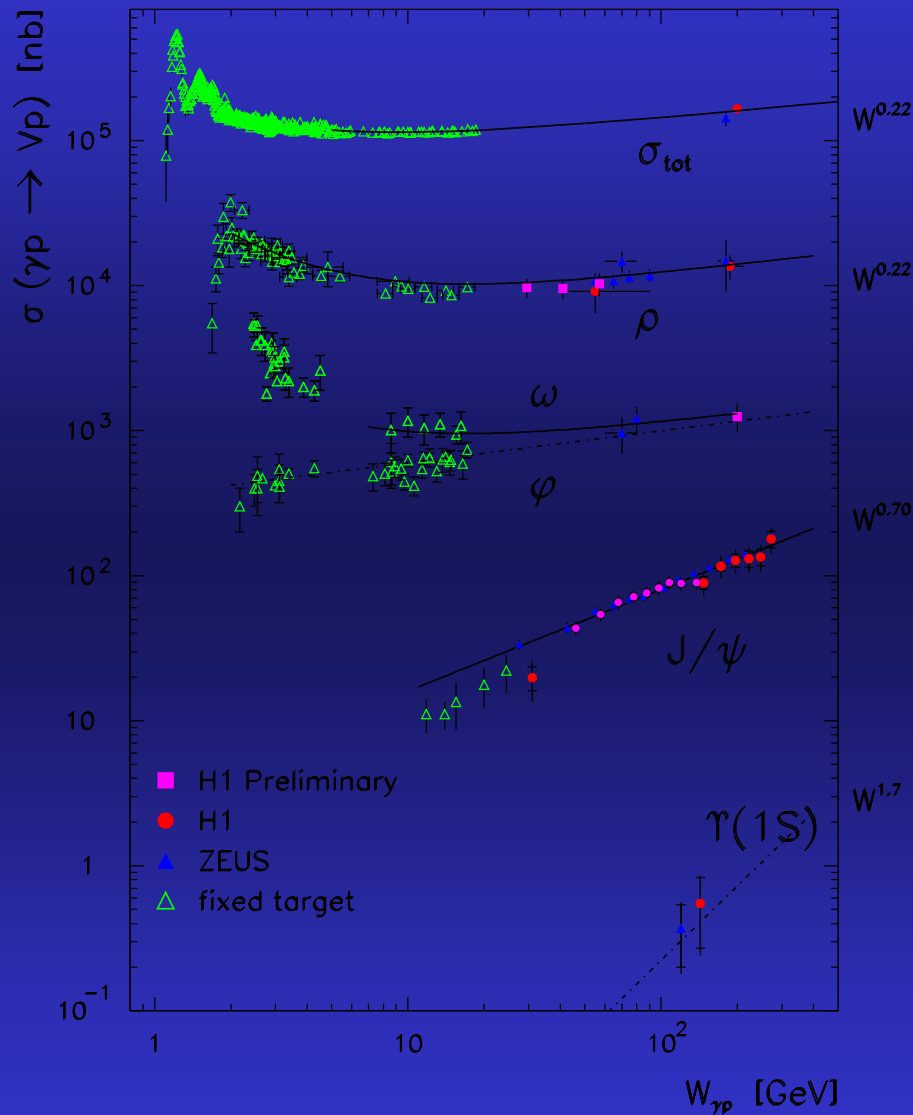


$W_{\gamma p}$: photon-proton
CM energy

t : (momentum transfer
from proton) ²

$-Q^2$: virtuality of the
photon; for protons or
nuclei, $Q^2 \approx 0$.

Summary of HERA Results



- $W_{\gamma p}$ up to 200 GeV

- $\rho^0 \leftrightarrow 10\%$ of σ_{TOT}

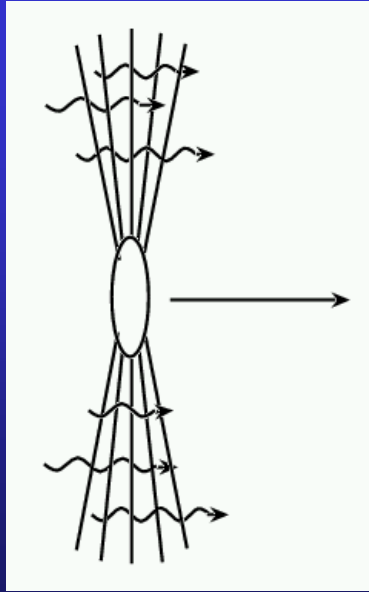
- Light mesons:

$$\sigma \propto W_{\gamma p}^{0.22}$$

- J/ψ :

$$\sigma \propto W_{\gamma p}^{0.80}$$

- a few tens of Υ seen



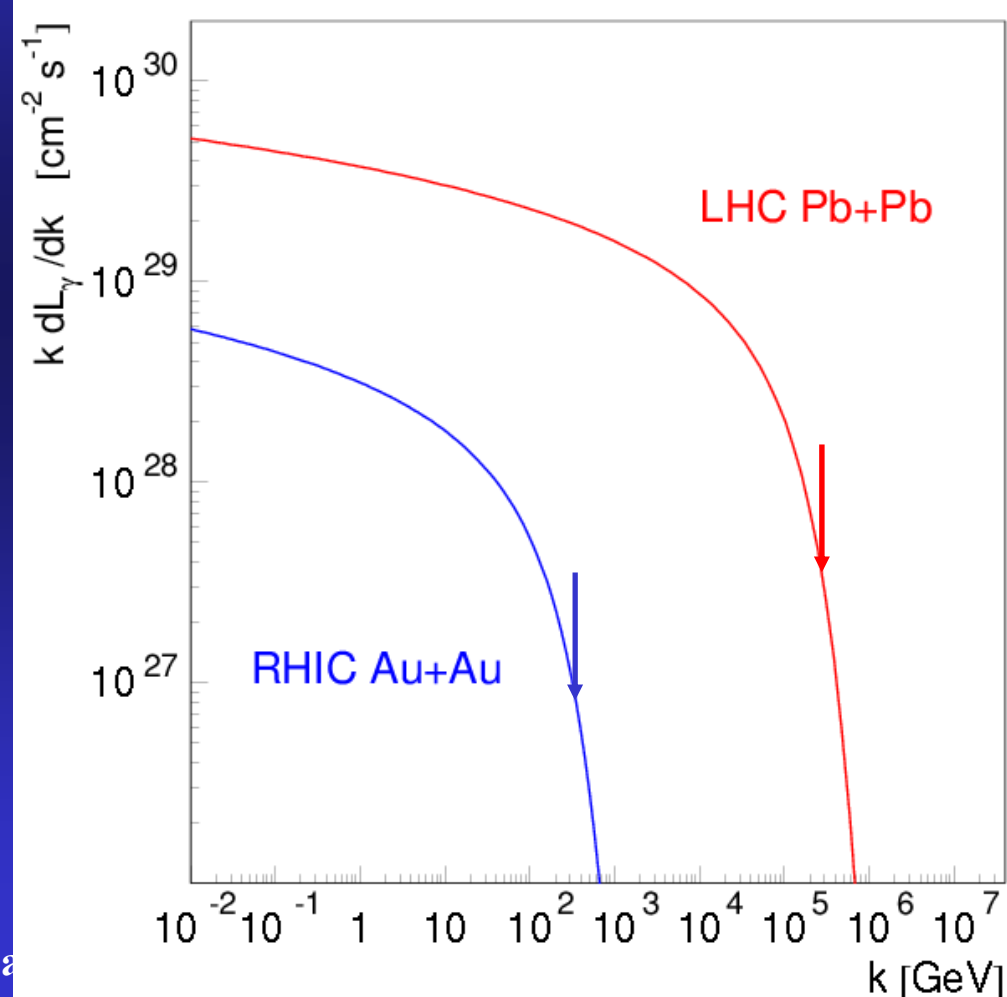
Electromagnetic field \leftrightarrow An equivalent flux of photons.

(Fermi 1924, Weizsäcker-Williams 1935)

Equivalent photon luminosity

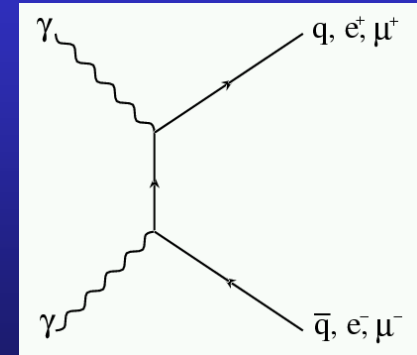
The photon spectrum extends to $\sim \gamma/R$

$\Leftrightarrow W_{\gamma p} \approx 1000 \text{ GeV}$
in Pb+Pb collisions
($Q^2 \approx 0$)



Ultra-peripheral Interactions Menu

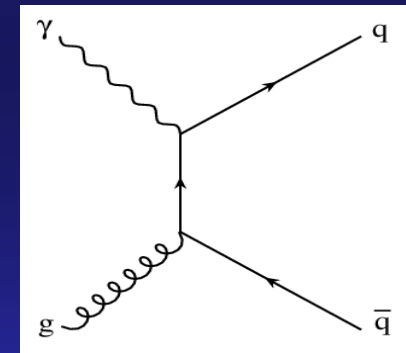
1. Purely electromagnetic, two-photon



2. Photonuclear

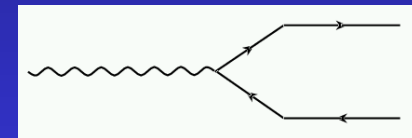
a) direct

Example: gamma+gluon



b) resolved

Vector Meson Dominance



Calculation of Vector Meson cross sections in Heavy-Ion Interactions

Phenomenological model based on scaling data of γp to γA .

J. Nystrand, S. Klein PRC 60(1999)014903

The ingredients:

Photon spectrum:

Weizsäcker-Williams

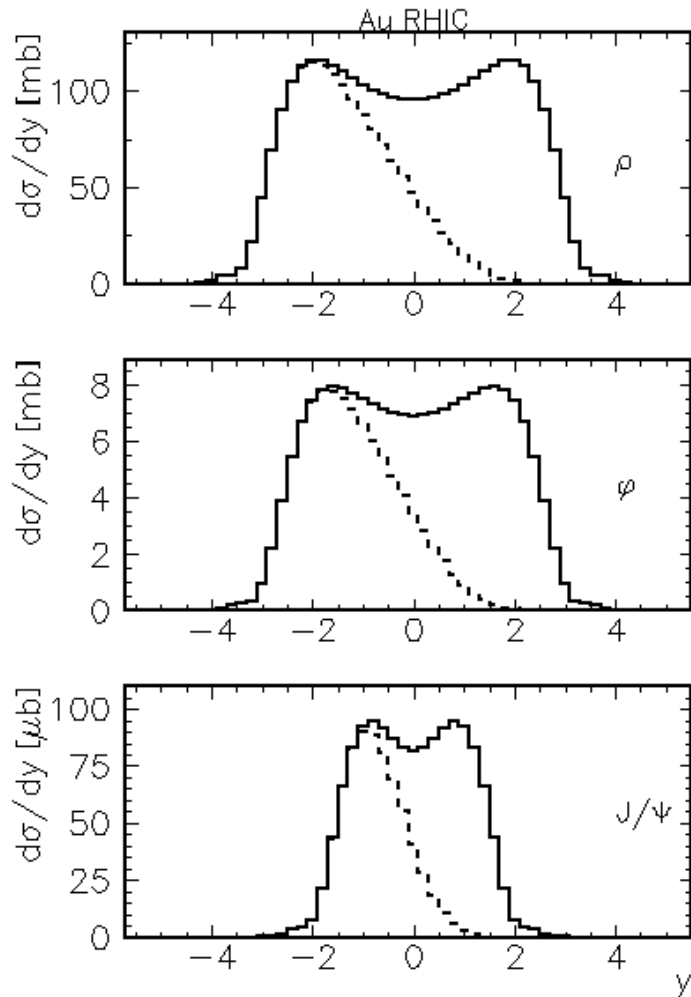
Input photon-nucleon data:

parameterized from results at HERA and fixed target

Scaling $\gamma p \rightarrow \gamma A$:

- 1) Neglecting cross terms - γ fluctuates into V which scatters elastically
- 2) Shadowing through a Glauber model
- 3) nuclear momentum transfer from Form factor

The model predicts cross sections, rapidity and p_T distributions of vector mesons at RHIC and LHC. For Au+Au 200 GeV at RHIC:



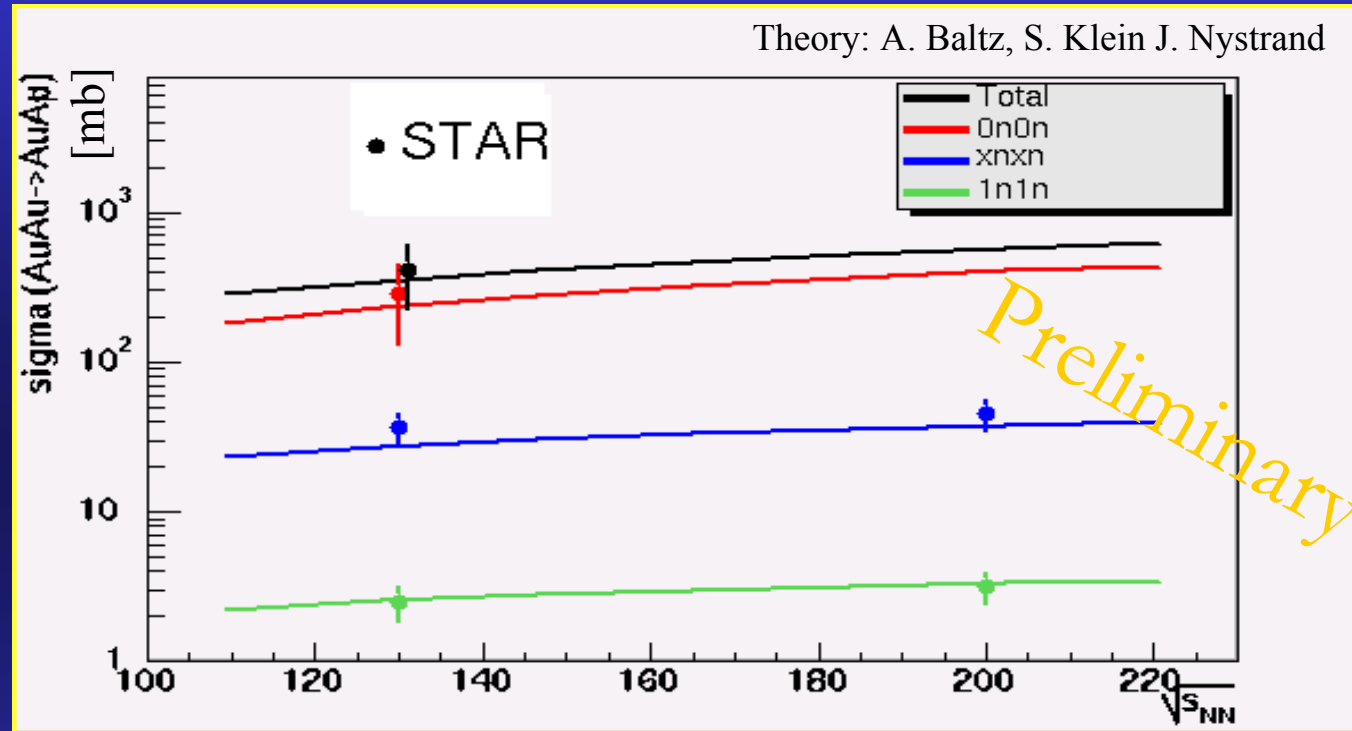
σ [mb] (prod. rate)

ρ	590	(120 Hz)
ω	59	(12 Hz)
ϕ	39	(7.9 Hz)
J/ψ	0.29	(0.058 Hz)

Cross sections in the 1-600 mb range!

The p_T distribution determined by the nuclear Form Factor,
 $p_T \sim 1/R \sim 50 \text{ MeV}/c$

Data on $\text{Au}+\text{Au}\rightarrow\text{Au}+\text{Au}+\rho^0$ at RHIC in good agreement with calculations



STAR Collaboration, Quark Matter 2002

⇒ Triggering + Analysis techniques work; photon spectrum and basic photonuclear cross sections well understood.

γp data from HERA

$$J/\Psi: \sigma = 1.5 W_{\gamma p}^{0.80} \quad [\text{nb}, W \text{ in GeV}]$$

\Rightarrow Extrapolate this to LHC energies

Too little data on Υ , QCD predicts

$$\left. \frac{d\sigma}{dt} \right|_{t=0} = \frac{\alpha_s^2 \Gamma_{ee}}{3\alpha M_V^5} 16\pi^3 \left[xg\left(x, \frac{M_V^2}{4}\right) \right]^2 \quad \text{Ryskin 1993}$$

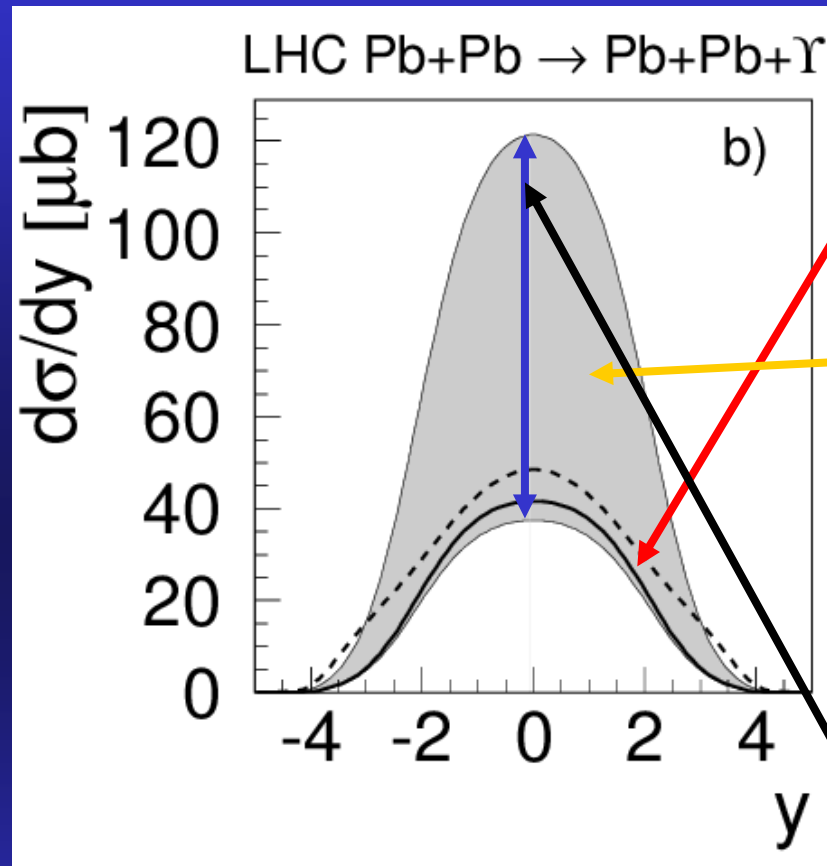
Further developments: Martin, Ryskin, Tubner Phys. Lett B

454 (1999)339 and Frankfurt, McDermott, Strikman JHEP 02(1999)002.

\Rightarrow Can be parameterized as

$$\Upsilon: \sigma = 0.06 W_{\gamma p}^{1.7} \quad [\text{pb}, W \text{ in GeV}]$$

Υ in ultra-peripheral Pb+Pb collisions at the LHC



A² scaling of QCD prediction

A² scaling of exp. data from HERA.

Uncertainty in measured cross section (mainly poor statistics).

Mid-rapidity $y=0 \Leftrightarrow$
 γp CM energy $W_{\gamma p} = 230 \text{ GeV}$,
 $x=2 \cdot 10^{-3}$

What can one expect in pp/ $\bar{p}p$?

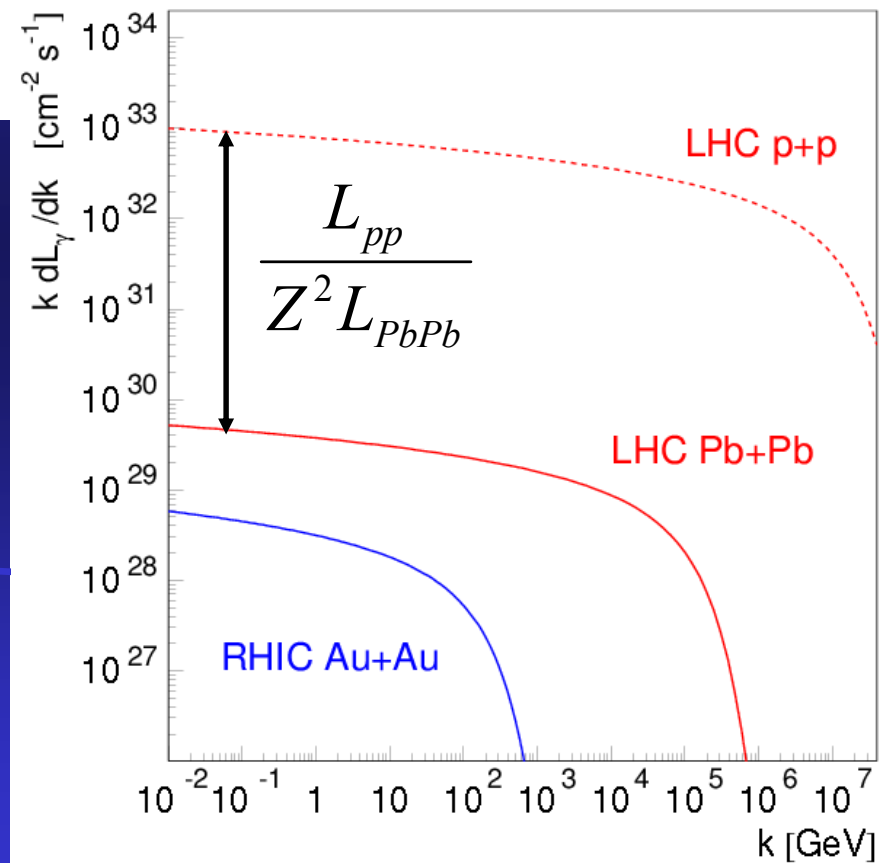
Equivalent photon spectrum of a proton:

$$\frac{dN}{dk} = \frac{\alpha}{2\pi k} \left[1 + \left(1 - \frac{2k}{\sqrt{s}} \right)^2 \right] \left(\ln(A) - \frac{11}{6} + \frac{3}{A} - \frac{3}{2A^2} + \frac{1}{3A^3} \right)$$

$$A = \frac{(1 + 0.71 \text{ GeV}^2)}{Q_{\min}^2}$$

Drees and Zeppenfeld (1989)

Point charge, $b > 0.7 \text{ fm}$



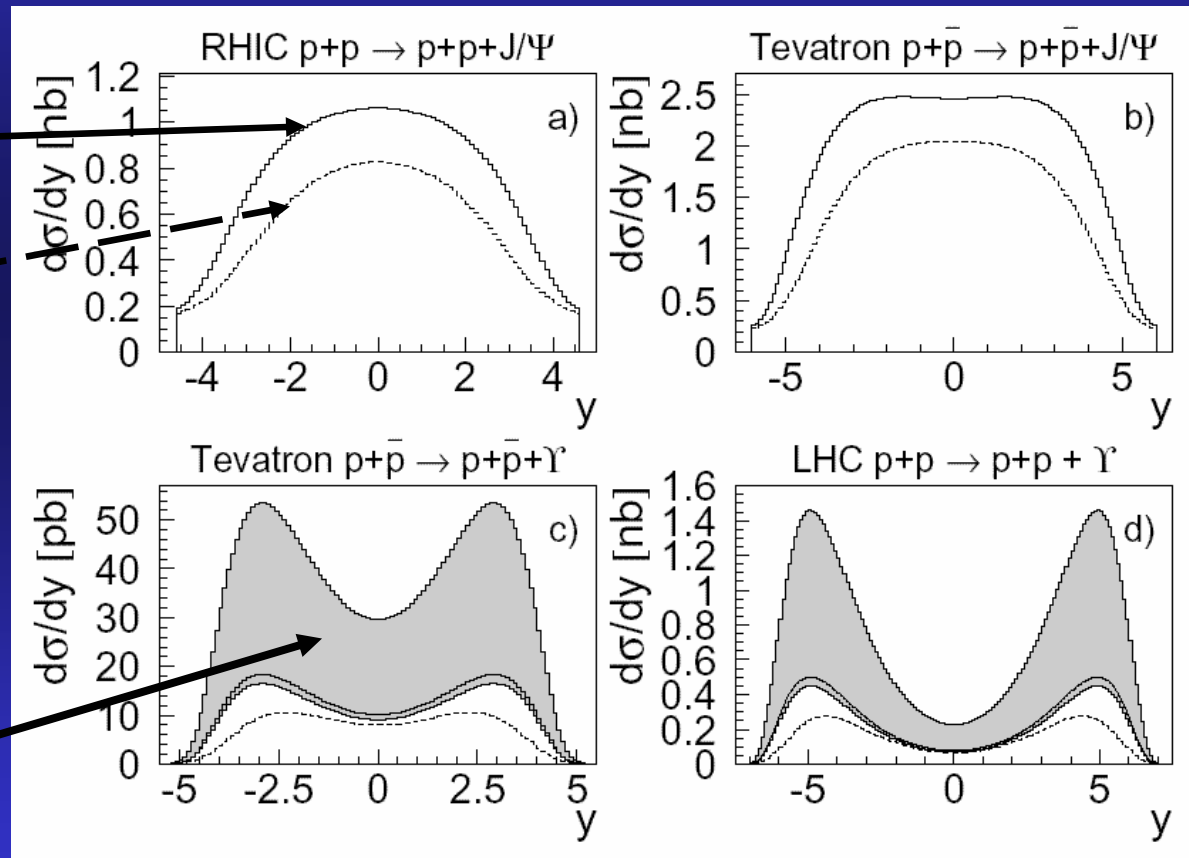
$\sigma(pp \rightarrow pp + V)$

$$\frac{d\sigma}{dy} = k \frac{dn_\gamma}{dk} \sigma(\gamma p \rightarrow V p)$$

Drees & Zeppenfeld
photon spectrum

point charge, $b > 1 \text{ fm}$

Uncertainty in experimental
cross section (mainly poor
statistics).



$$\sigma(pp \rightarrow pp + V)$$

$$\frac{d\sigma}{dy} = k \frac{dn_\gamma}{dk} \sigma(\gamma p \rightarrow Vp)$$

Kinematics at mid-rapidity ($y=0$)

J/ ψ

Υ

Tevatron

$$W_{\gamma p} = 80 \text{ GeV} \quad x \approx 1 \cdot 10^{-3}$$

$$W_{\gamma p} = 130 \text{ GeV} \quad x \approx 5 \cdot 10^{-3}$$

LHC

$$W_{\gamma p} = 210 \text{ GeV} \quad x \approx 2 \cdot 10^{-4}$$

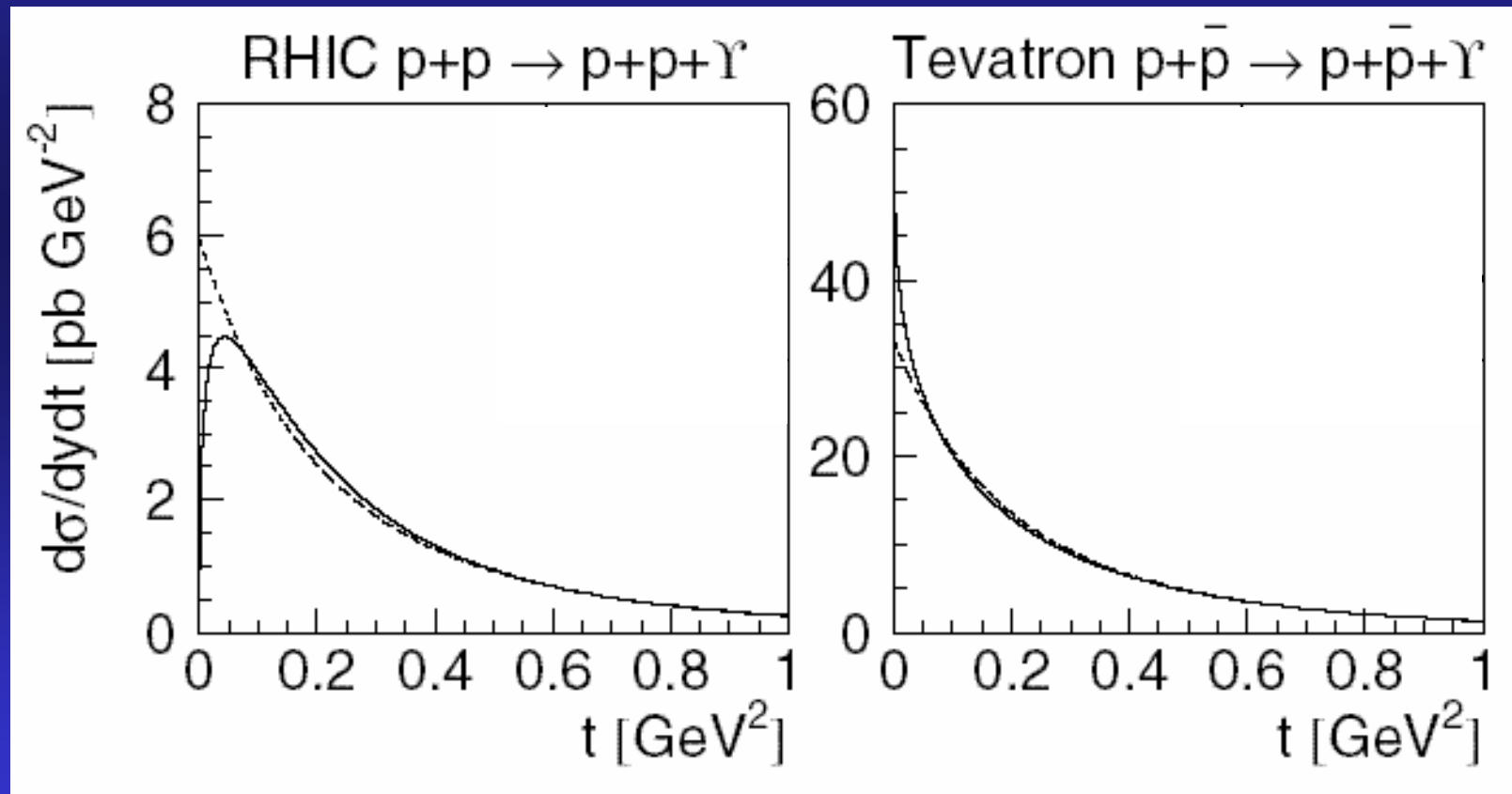
$$W_{\gamma p} = 350 \text{ GeV} \quad x \approx 6 \cdot 10^{-4}$$

Lower x can be reached away from $y=0$, but separation of photon-emitter and photon-target non-trivial.

Transverse momentum spectrum

Dominated by proton form factor

Low $p_T < \sim 1 \text{ GeV}/c$

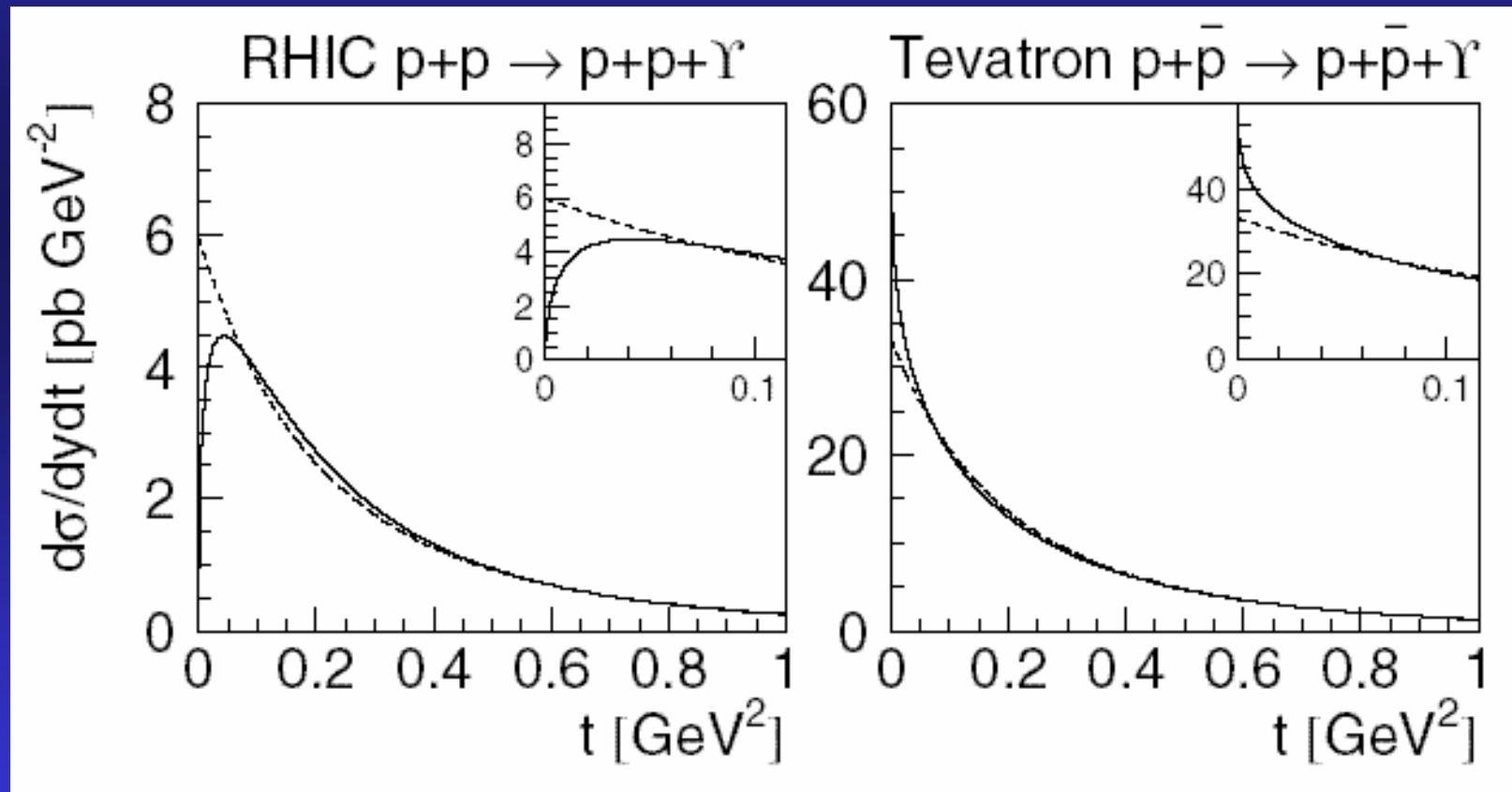


Transverse momentum spectrum

Dominated by proton form factor

Low $p_T < \sim 1 \text{ GeV}/c$

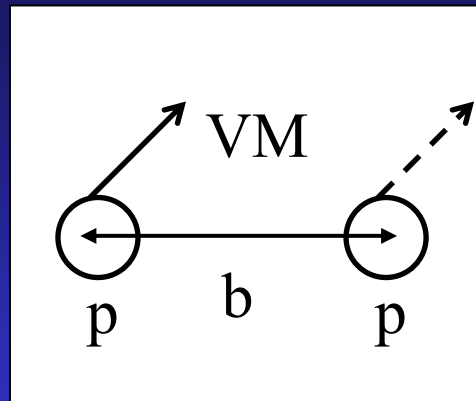
— with interference
- - - without interference



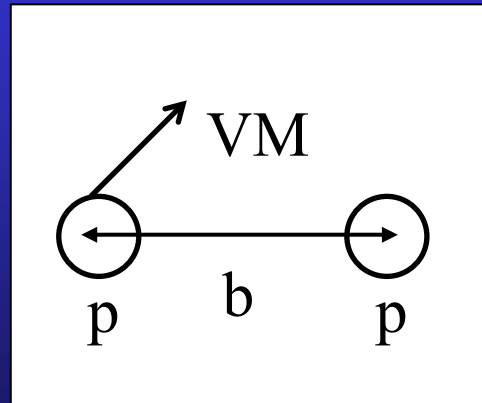
Transverse momentum spectrum

At very low p_T ($p_T \ll 1/\langle b \rangle$), not possible to distinguish photon-emitter and photon-target.

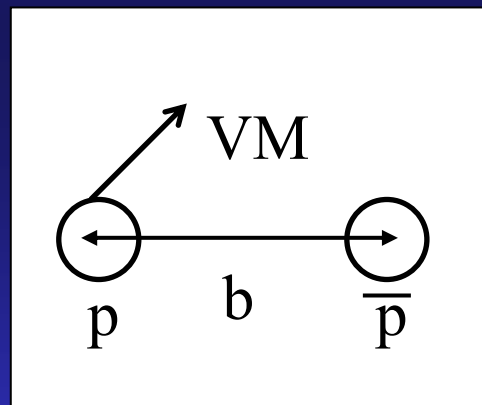
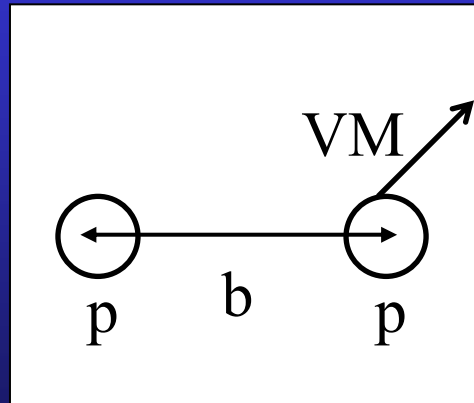
\Rightarrow Add amplitudes (not cross sections) with correct sign.



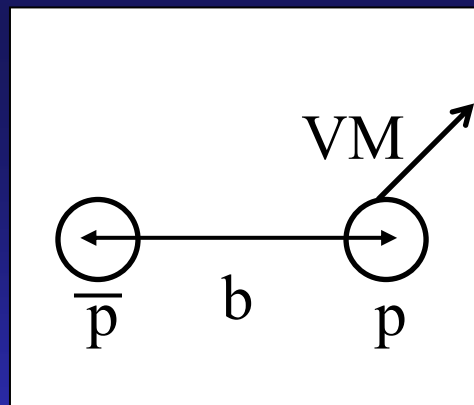
Exchanging photon-emitter and photon-target



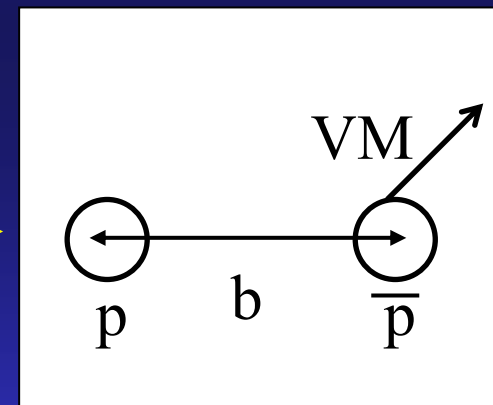
\rightarrow
P



\rightarrow
P



\rightarrow
C



$J^{PC} = 1^{--} \Rightarrow$ Destructive interference in pp (AA),
constructive interference in $p\bar{p}$.

Comparisons with hadronic production mode (pp)

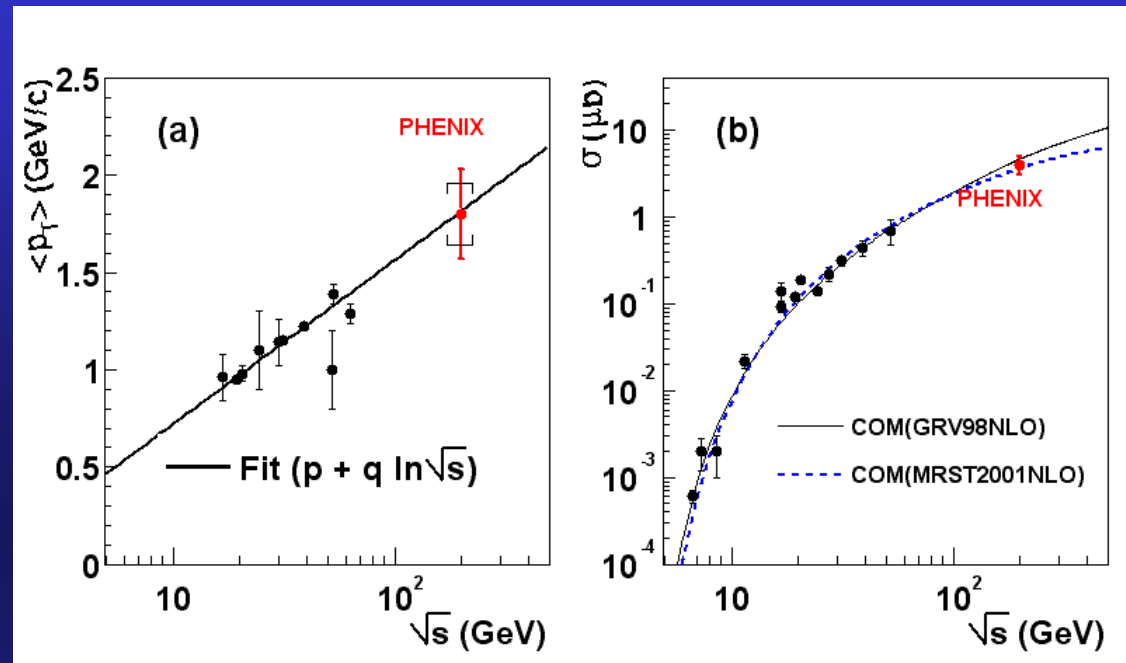
J/ψ

RHIC, $\sqrt{s}=200$ GeV:

$$\sigma = 2.70 \pm 0.40 \mu\text{b}$$

Photoproduction:

$$\sigma \approx 7 \text{ nb}$$



S.S. Adler et al. (PHENIX Collaboration)
 PRL 92(2004)051802, $p+p \rightarrow J/\psi + X$

γ

Tevatron, $\sqrt{s}=1.96$ TeV:

$$d\sigma/dy|_{y=0} = 30.4 \pm 3.1 \text{ nb}$$

Photoproduction:

$$d\sigma/dy|_{y=0} = 10\text{-}25 \text{ pb}$$

Comparisons with hadronic production mode at LHC

$pp \rightarrow \Upsilon :$

hadronic: $0.15\text{--}0.28 \mu\text{b}^*$ photo: 3.5 nb

ratio: $\sim 2 \cdot 10^{-2}$

*ALICE PPR / R. Vogt

$PbPb \rightarrow \Upsilon :$

Total hadronic (all centralities) $\sigma_{AA} \approx A^2 \sigma_{pp}$

hadronic: $6.5\text{--}12.1 \text{ mb}$ photo: $170 \mu\text{b}$

ratio: $\sim 2 \cdot 10^{-2}$

Much better background rejection in PbPb:

a) Higher multiplicities (hadronic interactions)

b) Coherence \Rightarrow photoproduced Υ w/ very low p_T

In pp, hadronic 'background' must be suppressed by a factor 10^3 !

Rapidity gaps:

$$\langle dn_{\text{ch}}/dy \rangle \approx 4 - 6 \text{ in pp at Tevatron/LHC}$$

Probability of having a gap of width Δy :

$$\exp(- \langle dn_{\text{ch}}/dy \rangle \cdot \Delta y)$$

$\Rightarrow \Delta y \approx 2$ will be sufficient

Further rejection from p_T distribution, $p_T < 1 \text{ GeV}/c$

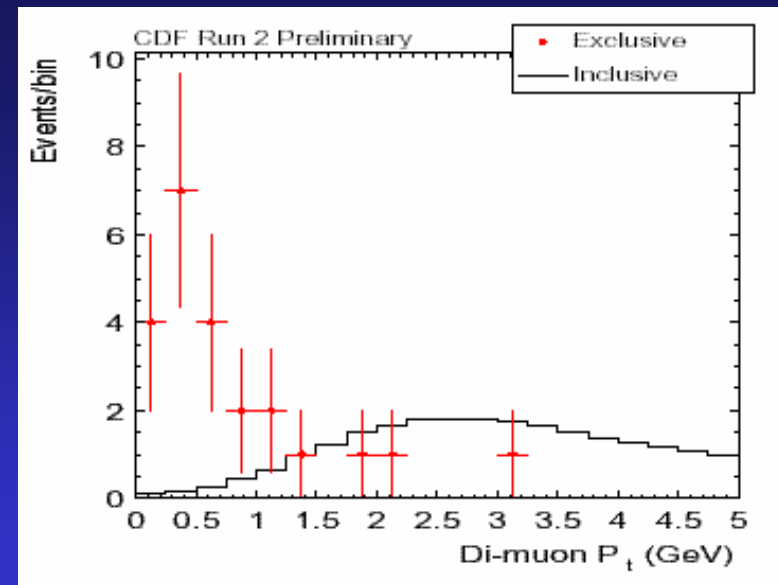
Some indication of J/ψ at the Tevatron

Presented at "Small-x and Diffraction", Fermilab,
September 2003 (CDF Collaboration, Angela Wyatt)

Searched for $p\bar{p} \rightarrow p\bar{p} + \chi_c$ via Pomeron-Pomeron,
 $\chi_c \rightarrow J/\psi + \gamma$

Found also a sample of J/ψ's without γ's !

p_T distribution of
exclusive J/ψ's:



Kinematics at mid-rapidity - pp vs. PbPb

$W_{\gamma p}$ – photon-proton CM energy

x – Bjorken- x of gluon

$y=0$

J/ ψ

Υ

Tevatron

$$W_{\gamma p} = 80 \text{ GeV} \quad x \approx 1 \cdot 10^{-3}$$

$$W_{\gamma p} = 130 \text{ GeV} \quad x \approx 5 \cdot 10^{-3}$$

LHC pp

$$W_{\gamma p} = 210 \text{ GeV} \quad x \approx 2 \cdot 10^{-4}$$

$$W_{\gamma p} = 350 \text{ GeV} \quad x \approx 6 \cdot 10^{-4}$$

LHC PbPb

$$W_{\gamma p} = 130 \text{ GeV} \quad x \approx 6 \cdot 10^{-4}$$

$$W_{\gamma p} = 230 \text{ GeV} \quad x \approx 2 \cdot 10^{-3}$$

Note: AA and pp competitive for selected reaction channels/final states.

Not suitable for measuring $\sigma_{\text{tot}}(\gamma\gamma)$ or $\sigma_{\text{tot}}(\gamma p)$, for example.

- Tagging of beam-nuclei not possible.
- $Q^2 \approx \hbar/R \approx 0$.
- The corresponding QCD processes must be suppressed.

Summary

- Some more details in PRL 92(2004)142003 (hep-ph/0311164).
- Interesting possibility to study exclusive production of heavy vector mesons in pp and $p\bar{p}$ collisions, and in Ultra-peripheral AA collisions.
- A probe of the nucleon/nucleus gluon density at low Bjorken- x .
- Extends the energy range studied at HERA.