



LUND UNIVERSITY



Academic Training Lectures

CERN

4, 5, 6, 7 April 2005

Monte Carlo Generators for the LHC

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CERN and Lund University

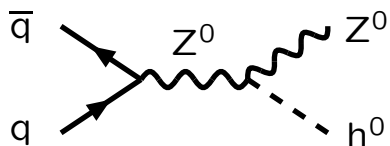
1. (Monday) Introduction and Overview; Matrix Elements
2. (Tuesday) Parton Showers; Matching Issues
3. (Wednesday) Multiple Interactions and Beam Remnants
4. (today) Hadronization and Decays; Summary and Outlook

Event Physics Overview

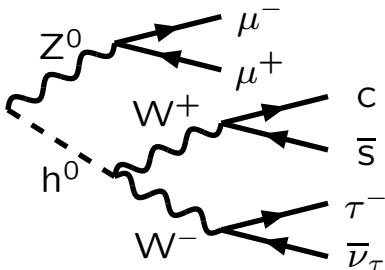
Repetition: from the “simple” to the “complex”,
or from “calculable” at large virtualities to “modelled” at small

Matrix elements (ME):

- 1) Hard subprocess:
 $|\mathcal{M}|^2$, Breit-Wigners,
parton densities.

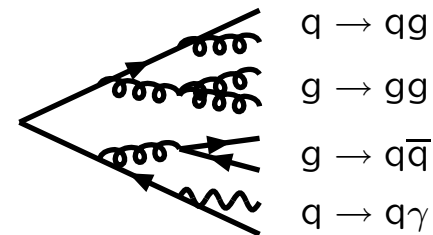


- 2) Resonance decays:
includes correlations.

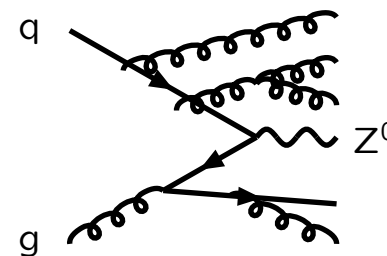


Parton Showers (PS):

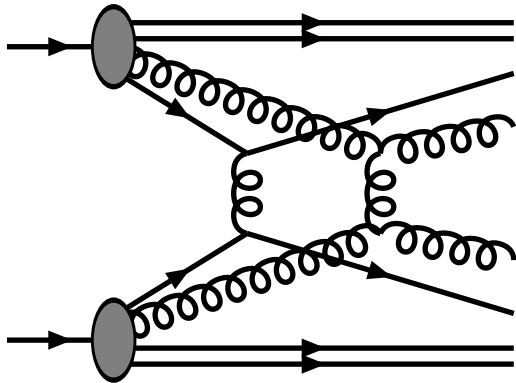
- 3) Final-state parton showers.



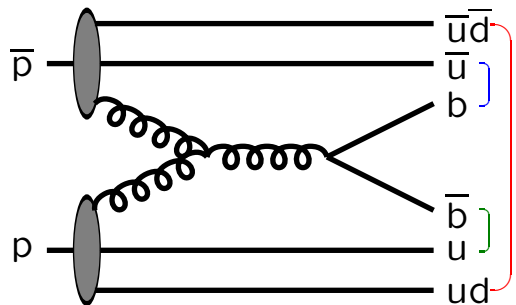
- 4) Initial-state parton showers.



5) Multiple parton-parton interactions.

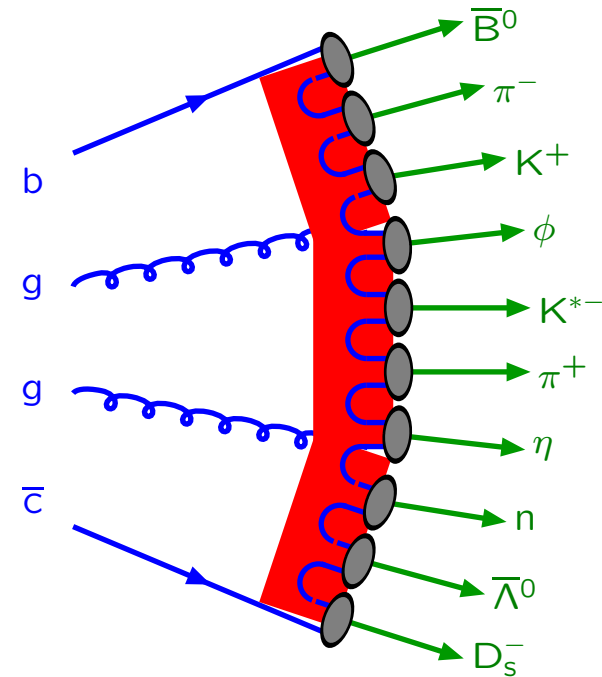


6) Beam remnants, with colour connections.

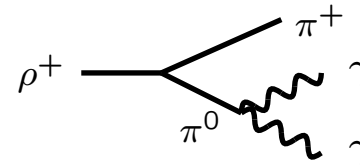


5) + 6) = Underlying Event

7) Hadronization



8) Ordinary decays: hadronic, τ , charm, ...



Hadronization/Fragmentation models

Perturbative \rightarrow nonperturbative \implies not calculable from first principles!

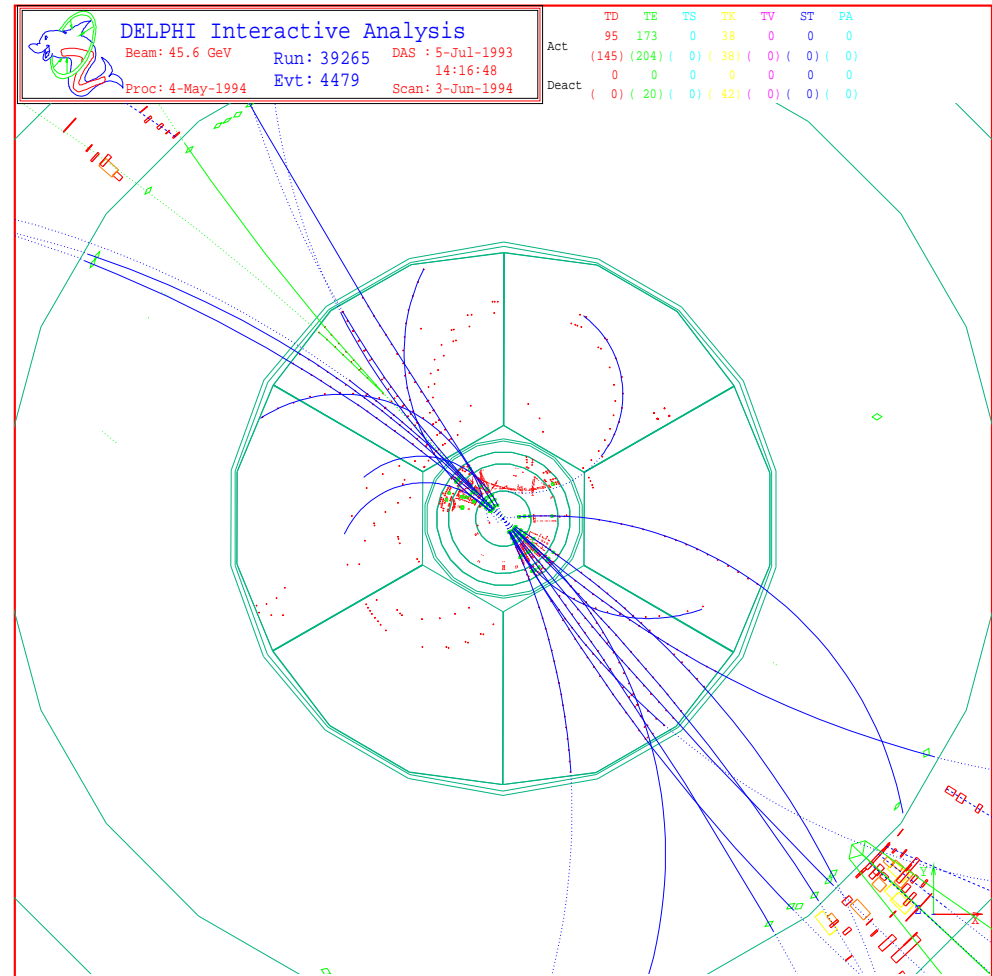
Model building = ideology + “cookbook”

Common approaches:

- 1) **String** Fragmentation
(most ideological)
- 2) **Cluster** Fragmentation
(simplest?)
- 3) **Independent** Fragmentation
(most cookbook)
- 4) **Local Parton–Hadron Duality**
(limited applicability)

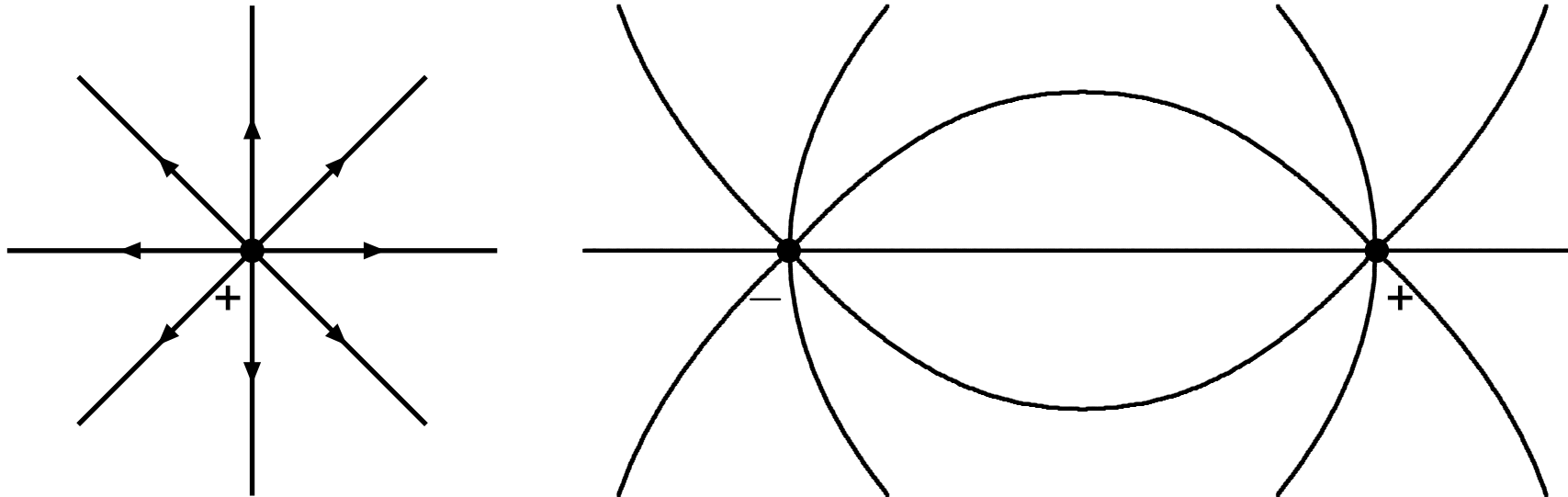
Best studied in

$$e^+e^- \rightarrow \gamma^*/Z^0 \rightarrow q\bar{q}$$



The Lund String Model

In QED, field lines go all the way to infinity

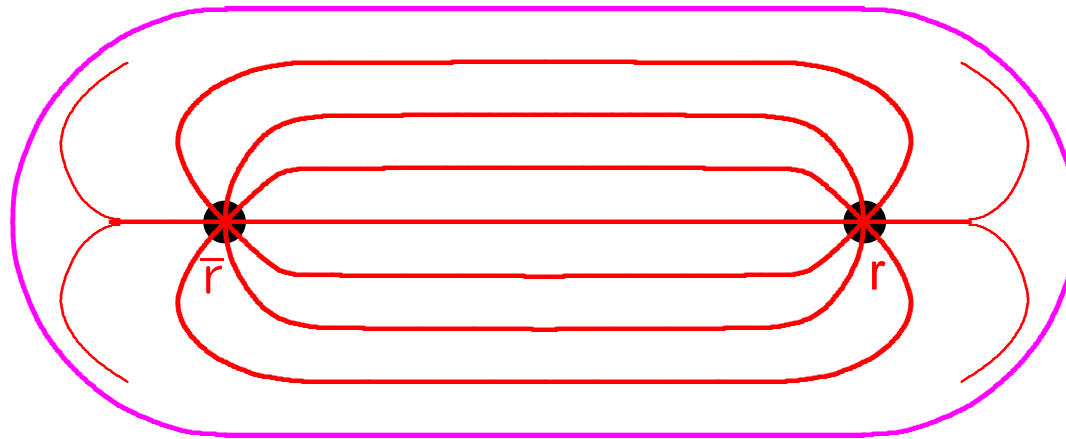


since photons cannot interact with each other.

Potential is simply additive:

$$V(\mathbf{x}) \propto \sum_i \frac{1}{|\mathbf{x} - \mathbf{x}_i|}$$

In QCD, for large charge separation, field lines seem to be compressed to tubelike region(s) \Rightarrow **string(s)**



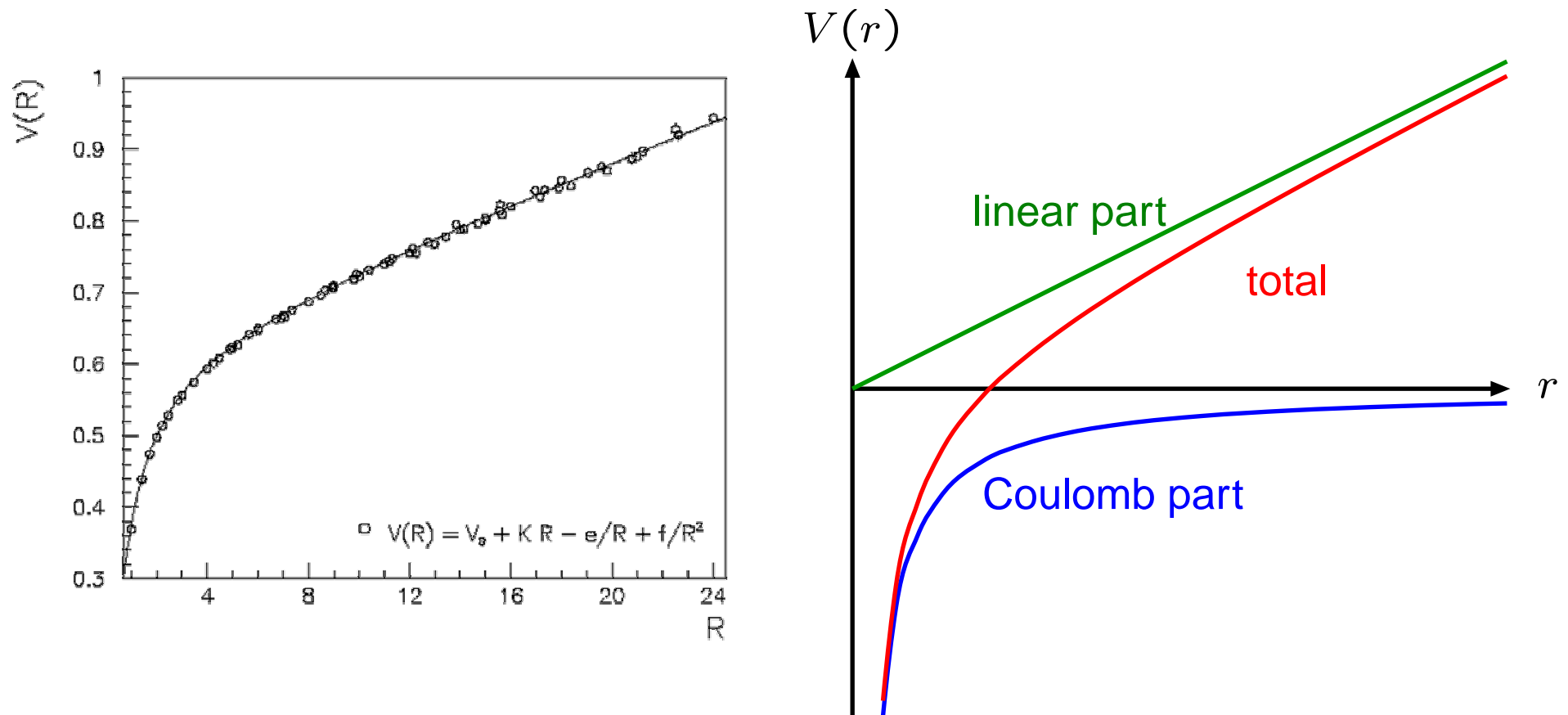
by self-interactions among soft gluons in the “vacuum”.
(Non-trivial ground state with quark and gluon “condensates”).
Analogy: vortex lines in type II superconductor)

Gives linear confinement with string tension:

$$F(r) \approx \text{const} = \kappa \approx 1 \text{ GeV/fm} \iff V(r) \approx \kappa r$$

Separation of transverse and longitudinal degrees of freedom
 \Rightarrow simple description as 1+1-dimensional object – string –
with Lorentz invariant formalism

Linear confinement confirmed e.g. by quenched lattice QCD



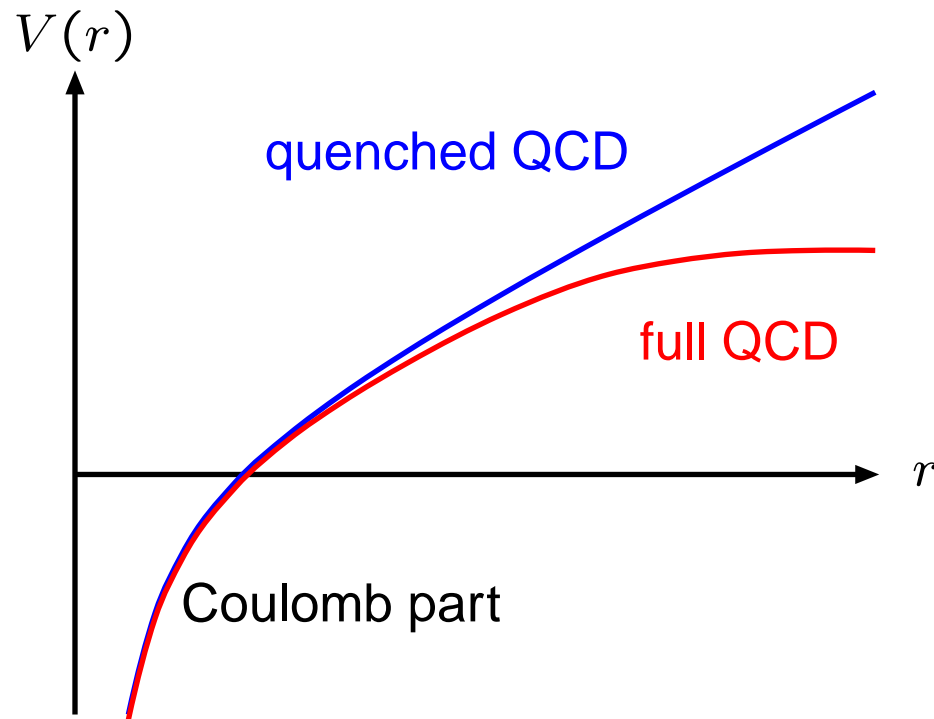
$$V(r) \approx -\frac{4\alpha_s}{3r} + \kappa r \approx -\frac{0.13}{r} + r$$

(for $\alpha_s \approx 0.5$, r in fm and V in GeV)

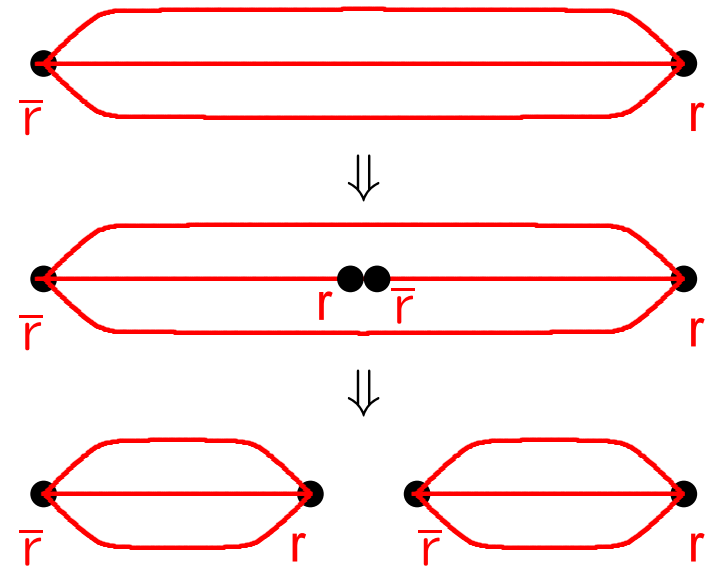
$V(0.4 \text{ fm}) \approx 0$: Coulomb important for internal structure of hadrons,
not for particle production (?)

Real world (??, or at least unquenched lattice QCD)

\implies nonperturbative string breakings $gg \dots \rightarrow q\bar{q}$



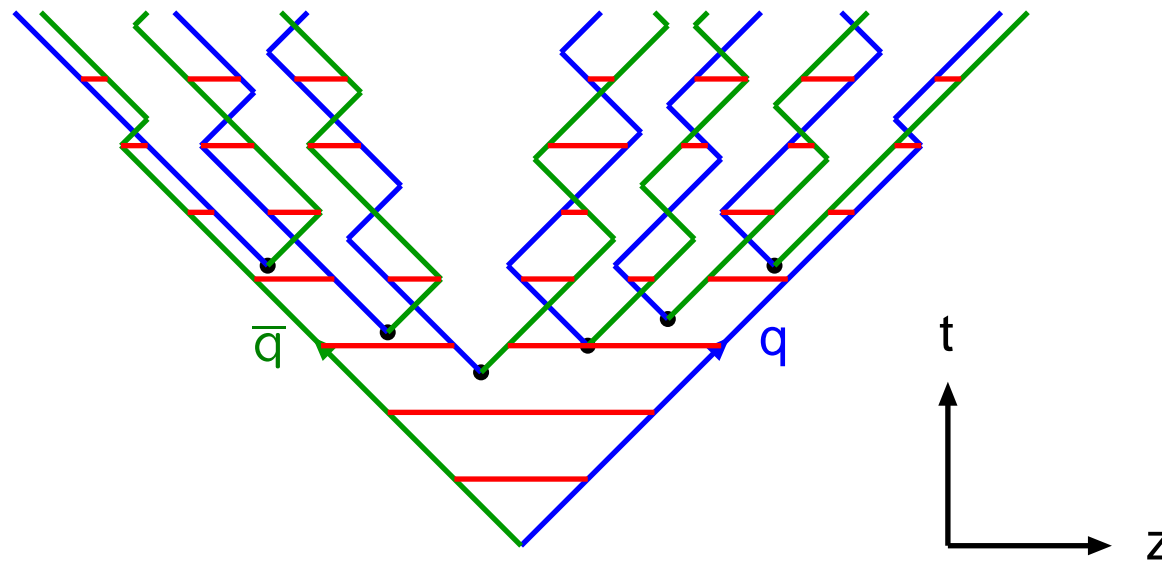
simplified colour representation:



Repeat for large system \Rightarrow *Lund model*
which neglects Coulomb part:

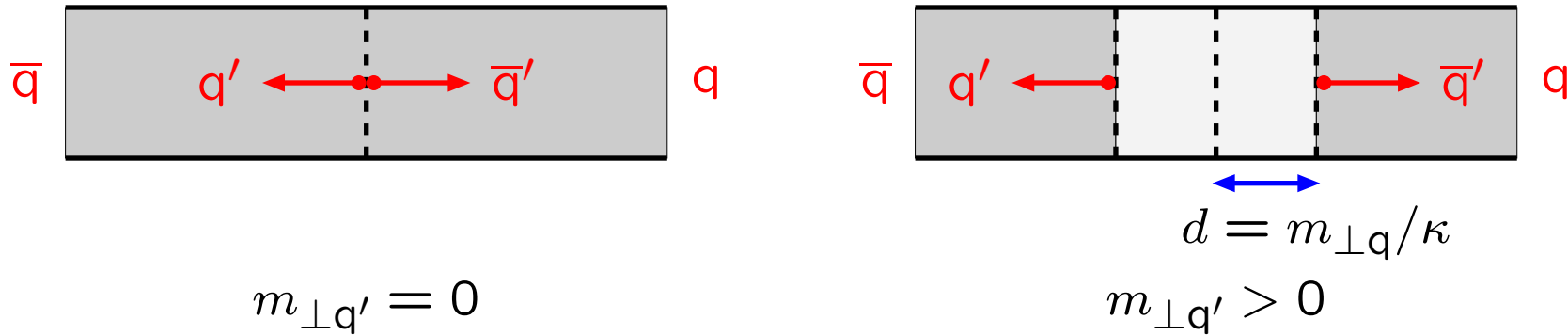
$$\left| \frac{dE}{dz} \right| = \left| \frac{dp_z}{dz} \right| = \left| \frac{dE}{dt} \right| = \left| \frac{dp_z}{dt} \right| = \kappa$$

Motion of quarks and antiquarks in a $q\bar{q}$ system:



gives simple but powerful picture of hadron production
(with extensions to massive quarks, baryons, ...)

How does the string break?



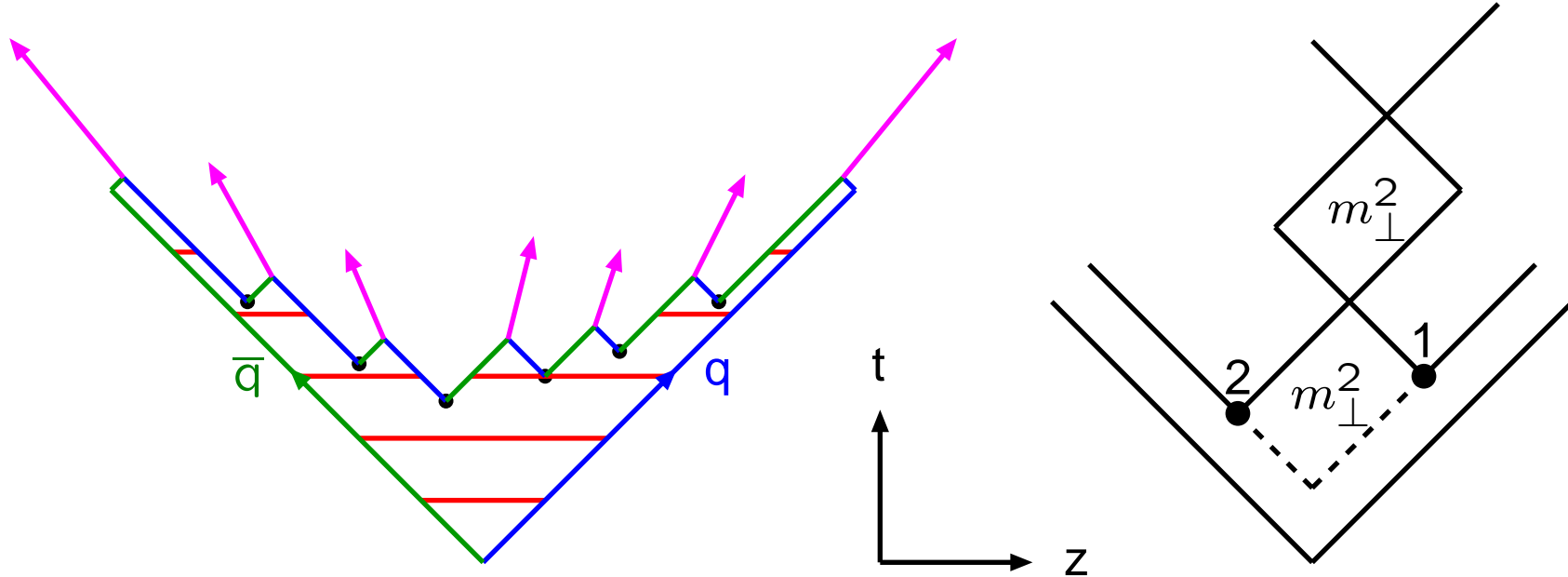
String breaking modelled by tunneling:

$$\mathcal{P} \propto \exp\left(-\frac{\pi m_{\perp q}^2}{\kappa}\right) = \exp\left(-\frac{\pi p_{\perp q}^2}{\kappa}\right) \exp\left(-\frac{\pi m_q^2}{\kappa}\right)$$

- 1) common Gaussian p_{\perp} spectrum
- 2) suppression of heavy quarks $u\bar{u} : d\bar{d} : s\bar{s} : c\bar{c} \approx 1 : 1 : 0.3 : 10^{-11}$
- 3) diquark \sim antiquark \Rightarrow simple model for baryon production

Hadron composition also depends on spin probabilities, hadronic wave functions, phase space, more complicated baryon production, ...
 \Rightarrow "moderate" predictivity (many parameters!)

Fragmentation starts in the middle and spreads outwards:

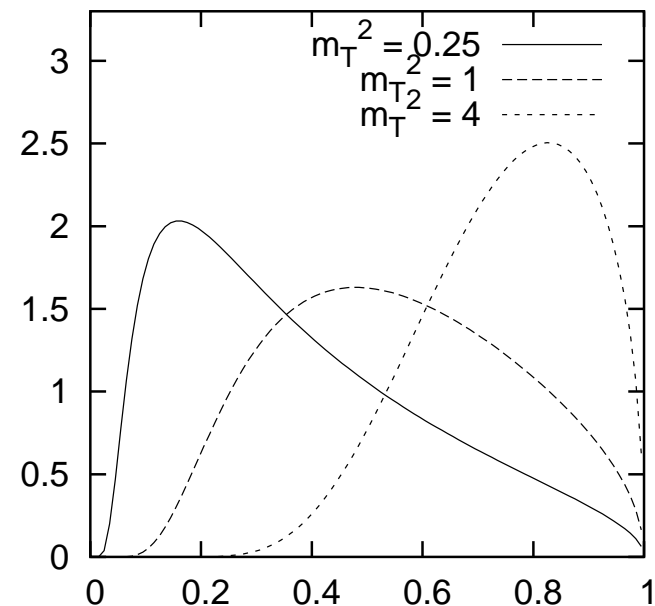


$f(z)$, $a = 0.5$, $b = 0.7$

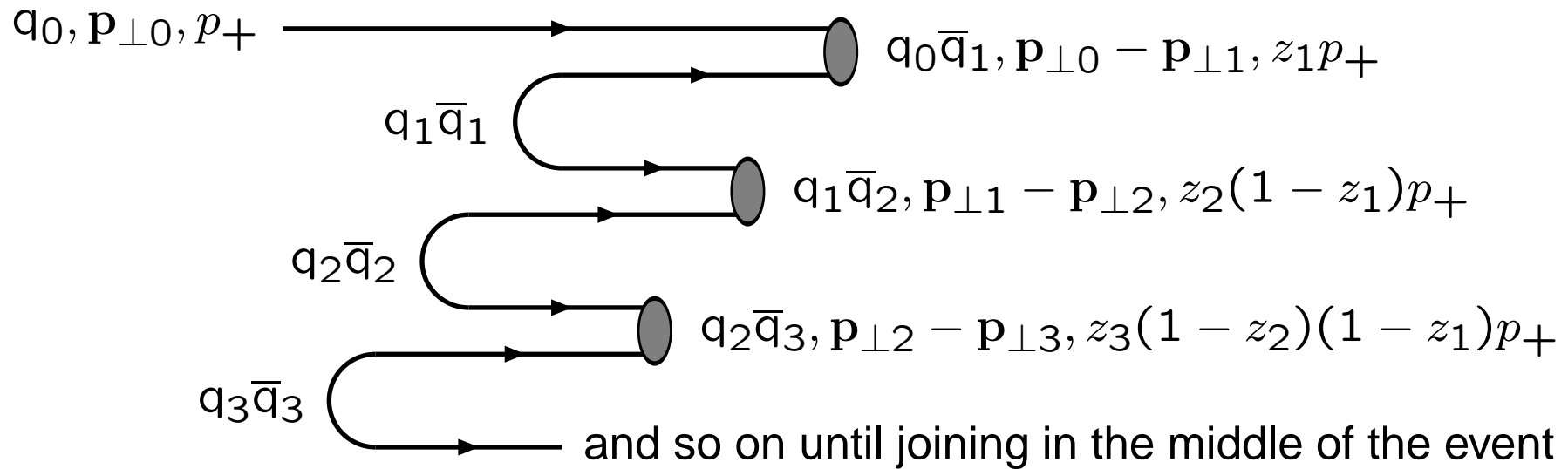
but breakup vertices causally disconnected
 \Rightarrow can proceed in arbitrary order
 \Rightarrow *left-right symmetry*

$$\begin{aligned} \mathcal{P}(1, 2) &= \mathcal{P}(1) \times \mathcal{P}(1 \rightarrow 2) \\ &= \mathcal{P}(2) \times \mathcal{P}(2 \rightarrow 1) \end{aligned}$$

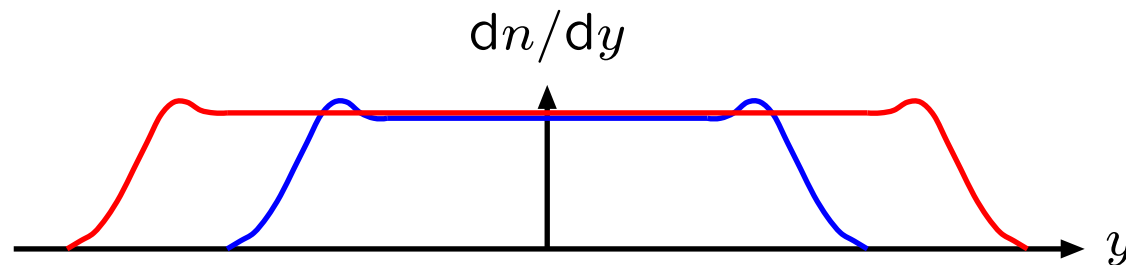
\Rightarrow Lund symmetric fragmentation function
 $f(z) \propto (1 - z)^a \exp(-bm_{\perp}^2/z)/z$



The iterative ansatz

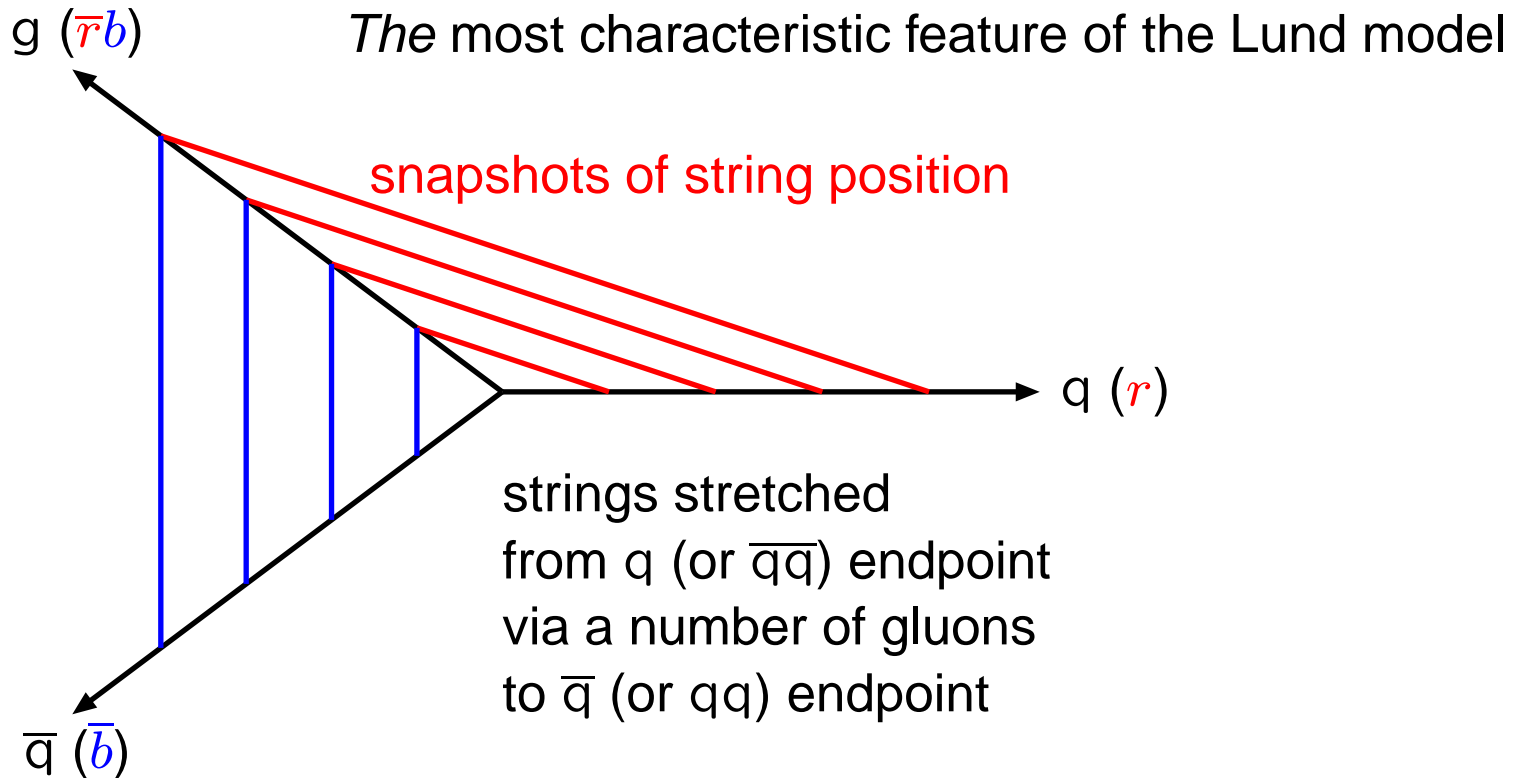


Scaling in lightcone $p_{\pm} = E \pm p_z$ (for $q\bar{q}$ system along z axis) implies flat central rapidity plateau + some endpoint effects:



$\langle n_{\text{ch}} \rangle \approx c_0 + c_1 \ln E_{\text{cm}}, \sim$ Poissonian multiplicity distribution

The Lund gluon picture



Gluon = kink on string, carrying energy and momentum

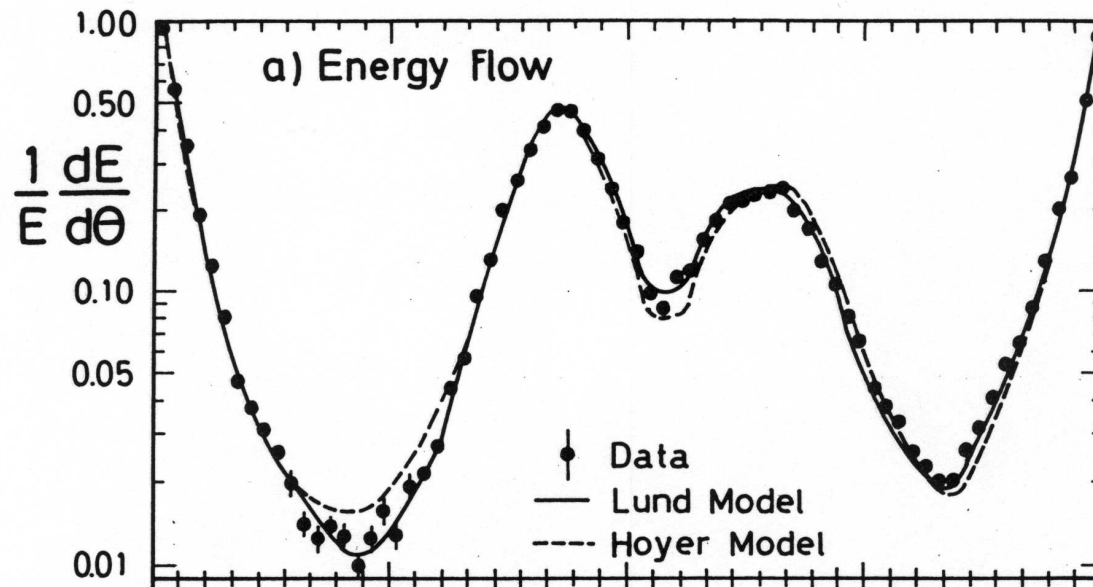
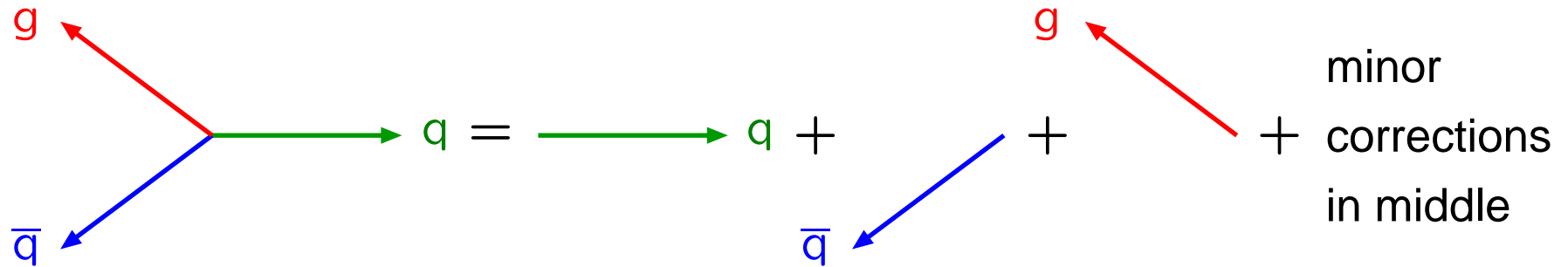
Force ratio gluon/ quark = 2, cf. QCD $N_C/C_F = 9/4, \rightarrow 2$ for $N_C \rightarrow \infty$

No new parameters introduced for gluon jets!, so:

- Few parameters to describe energy-momentum structure!
 - Many parameters to describe flavour composition!

Independent fragmentation

Based on a similar iterative ansatz as string, but

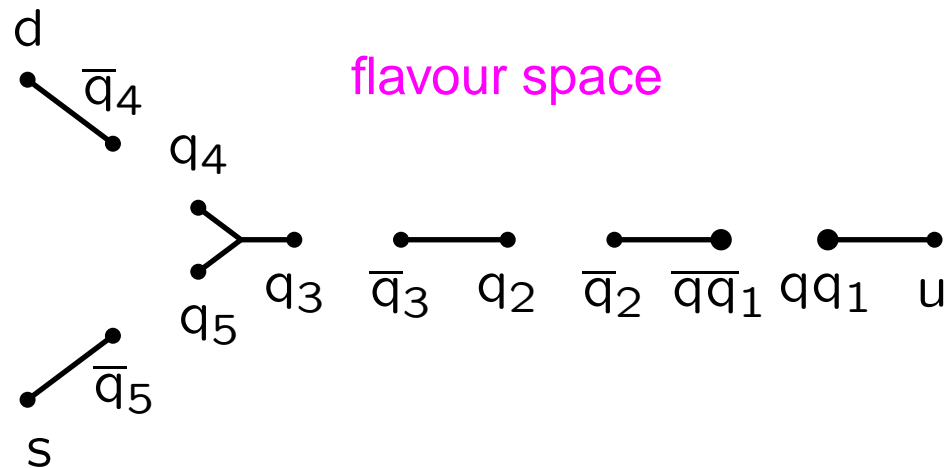
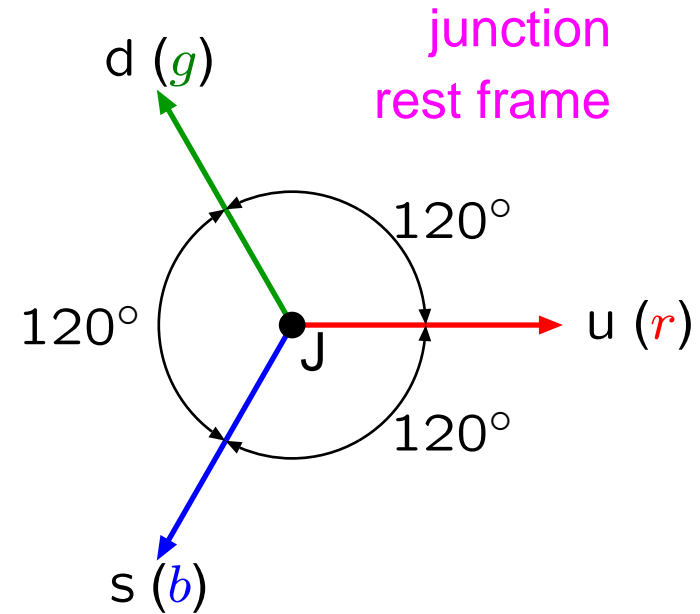
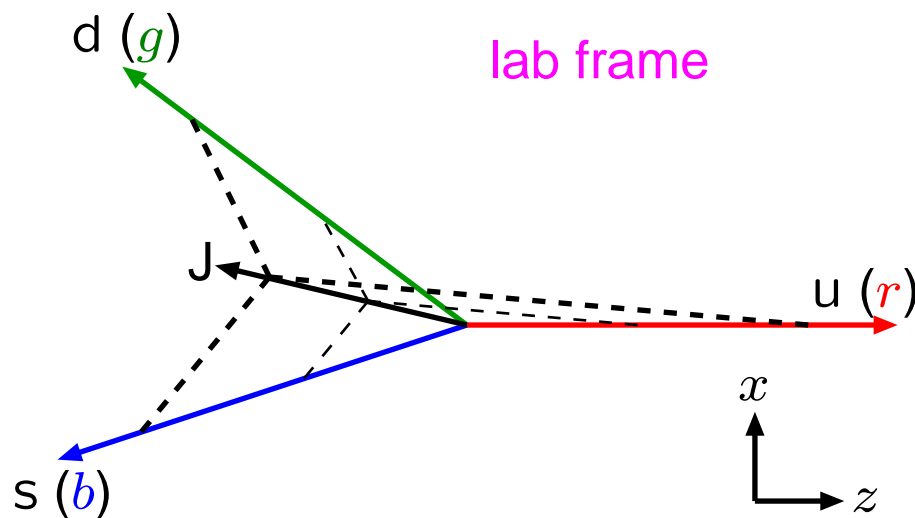


String effect
(JADE, 1980)
 \approx coherence in
nonperturbative
context

Further numerous and detailed tests at LEP favour string picture ...
... but much is still uncertain when moving to hadron colliders.

Lund news: fragmentation of junction topology

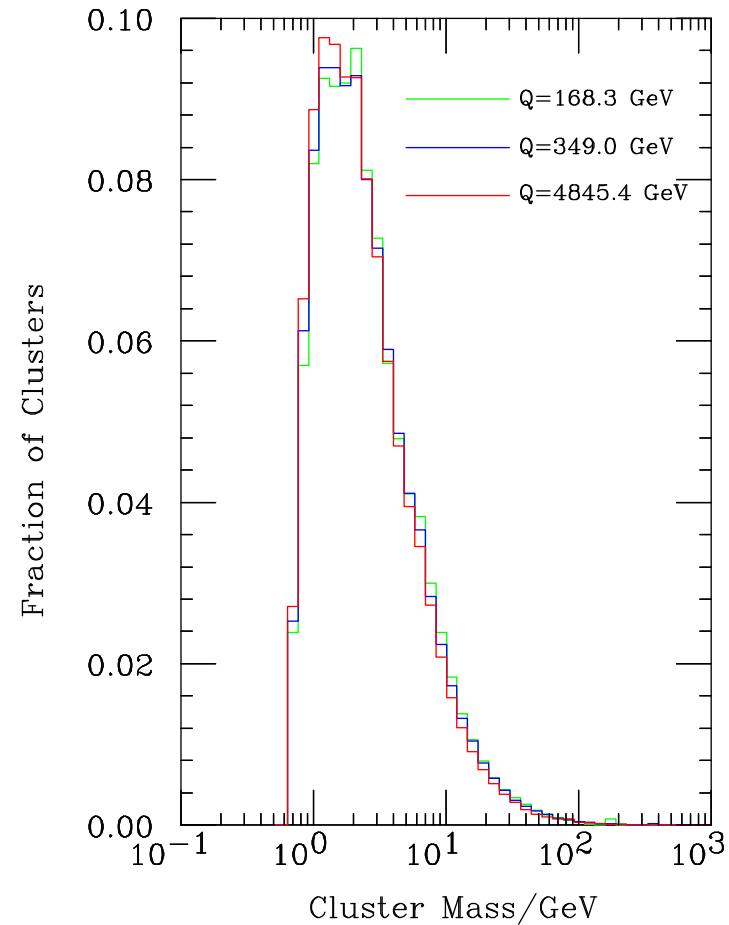
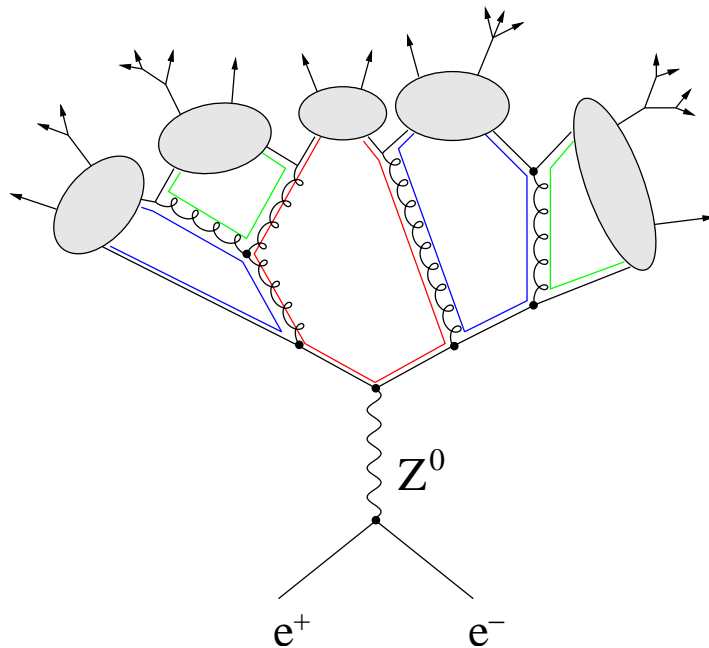
Encountered in R -parity violating SUSY decays $\tilde{\chi}_1^0 \rightarrow uds$,
 or when 2 valence quarks kicked out of proton beam



More complicated
 (but \approx solved) with
 gluon emission and
 massive quarks

The HERWIG Cluster Model

“Preconfinement”:
colour flow is local
in coherent shower evolution



- 1) Introduce forced $g \rightarrow q\bar{q}$ branchings
- 2) Form colour singlet clusters
- 3) Clusters decay isotropically to 2 hadrons according to phase space weight $\sim (2s_1 + 1)(2s_2 + 1)(2p^*/m)$
simple and clean, but ...

1) Tail to very large-mass clusters (e.g. if no emission in shower);
 if large-mass cluster \rightarrow 2 hadrons then
 incorrect hadron momentum spectrum, crazy four-jet events
 \Rightarrow split big cluster into 2 smaller along “string” direction;
 daughter-mass spectrum \Rightarrow iterate if required;
 \sim 15% of primary clusters are split, but give \sim 50% of final hadrons

2) Isotropic baryon decay inside cluster

\Rightarrow splittings $g \rightarrow qq + \bar{q}\bar{q}$

3) Too soft charm/bottom spectra

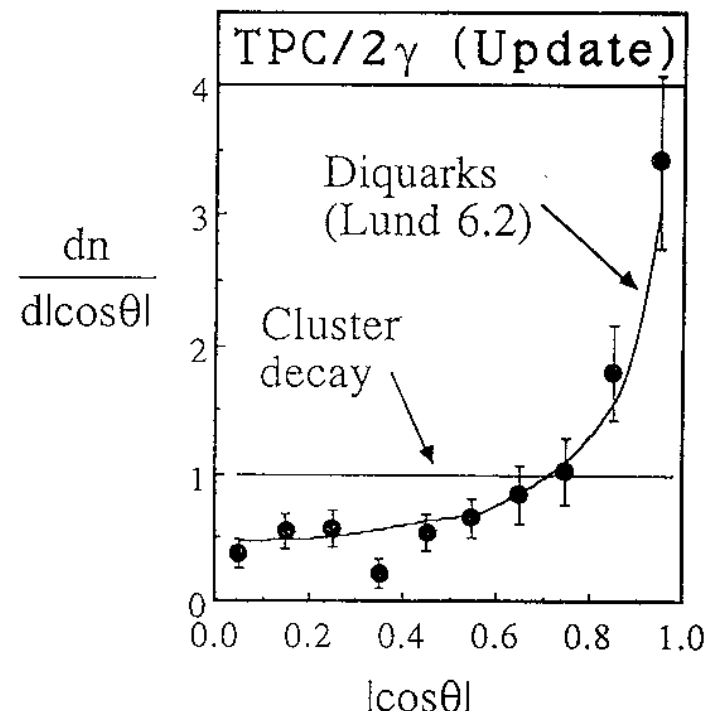
\Rightarrow anisotropic leading-cluster decay

4) Charge correlations still problematic

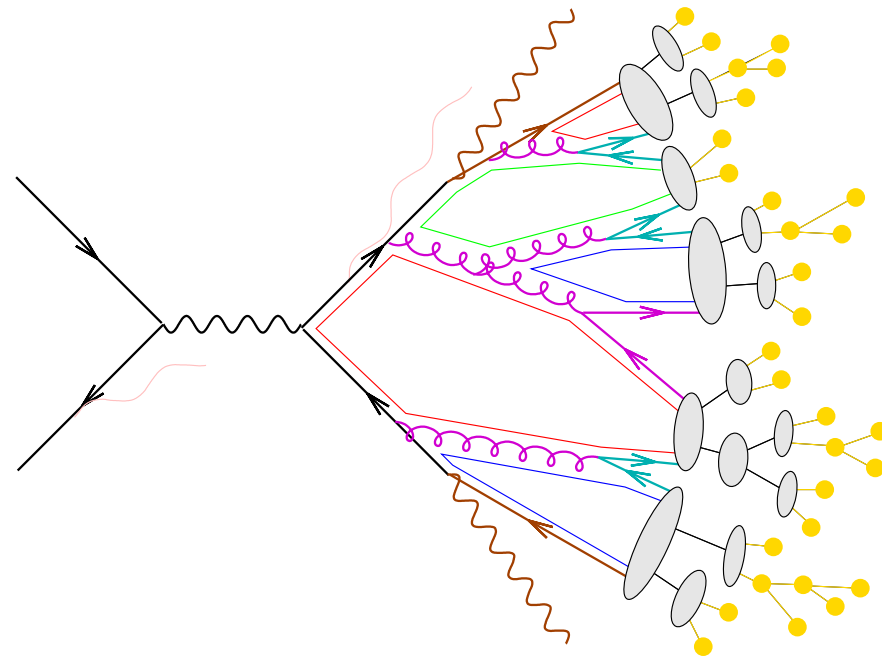
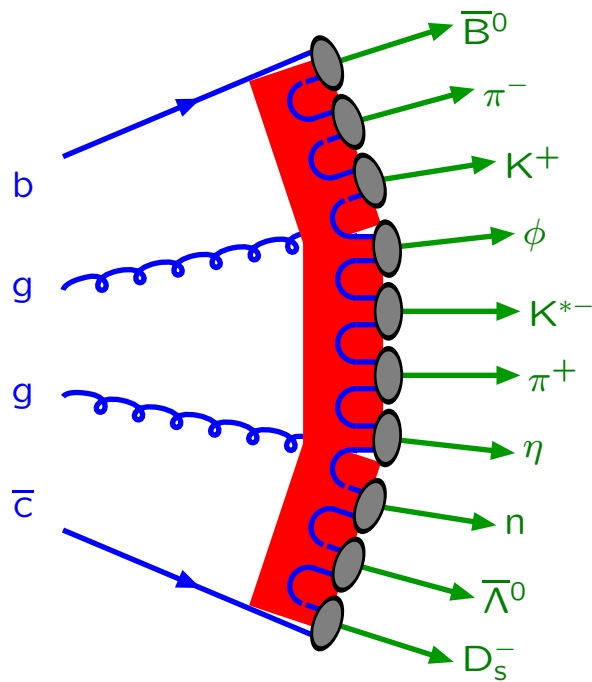
\Rightarrow all clusters anisotropic (?)

5) Sensitivity to particle content

\Rightarrow only include complete multiplets



String vs. Cluster



program	PYTHIA	HERWIG
model	string	cluster
energy-momentum picture	powerful	simple
parameters	predictive	unpredictive
flavour composition	few	many
parameters	many	few

“There ain’t no such thing as a parameter-free *good* description”

Local Parton–Hadron Duality

Analytic approach:

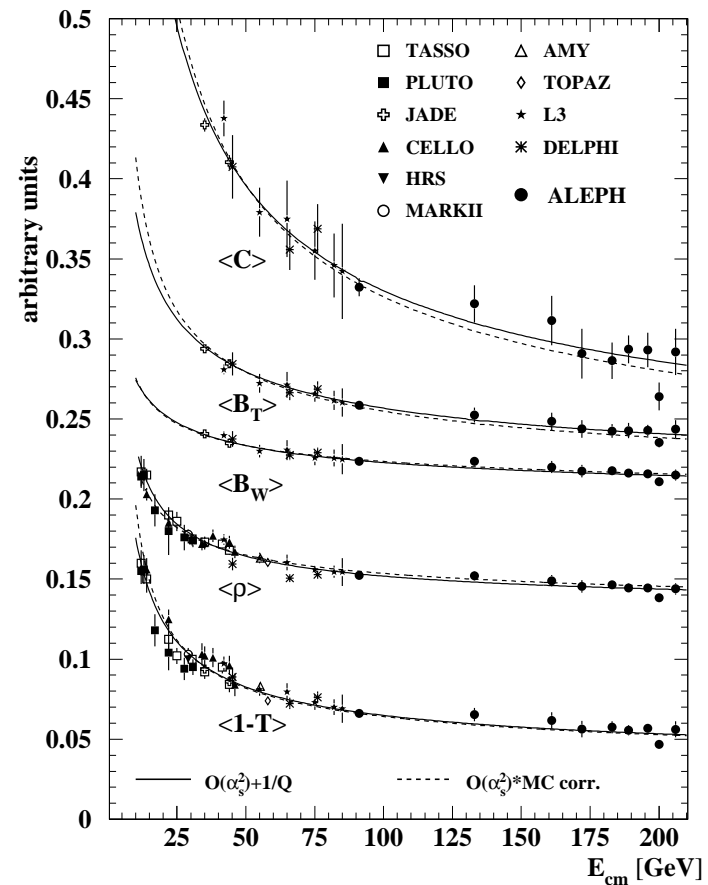
Run shower down to $Q \approx \Lambda_{\text{QCD}}$
(or m_{hadron} , if larger)

“Hard Line”: each parton \equiv one hadron

“Soft Line”: local hadron density
 \propto parton density

describes momentum spectra dn/dx_p
and semi-inclusive particle flow,
but fails for identified particles

+ “renormalons” (power corrections)
 $\langle 1 - T \rangle = a \alpha_s(E_{\text{cm}}) + b \alpha_s^2(E_{\text{cm}})$
 $+ c/E_{\text{cm}}$



Not Monte Carlo, not for arbitrary quantities

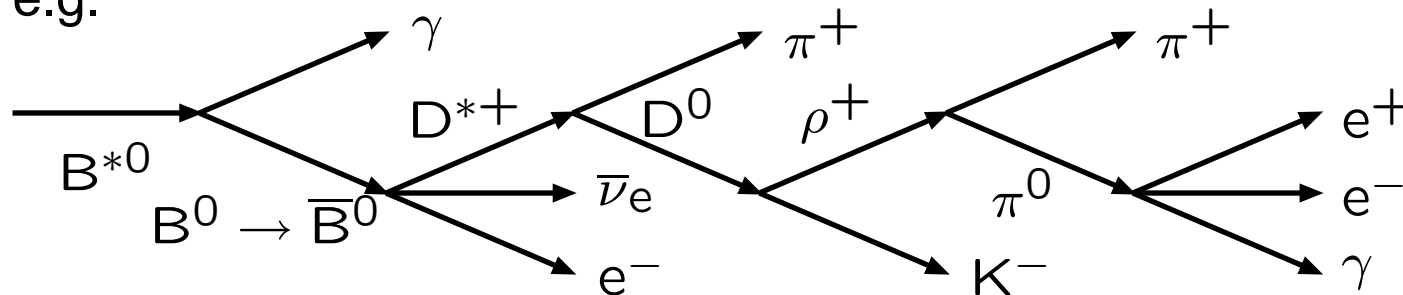
Decays

Unspectacular/ungrateful but necessary:

this is where most of the final-state particles are produced!

Involves hundreds of particle kinds and thousands of decay modes.

e.g.



- $B^{*0} \rightarrow B^0 \gamma$: electromagnetic decay
- $B^0 \rightarrow \bar{B}^0$ mixing (weak)
- $\bar{B}^0 \rightarrow D^{*+} \bar{\nu}_e e^-$: weak decay, displaced vertex, $|\mathcal{M}|^2 \propto (p_{\bar{B}} p_{\bar{\nu}})(p_e p_{D^*})$
- $D^{*+} \rightarrow D^0 \pi^+$: strong decay
- $D^0 \rightarrow \rho^+ K^-$: weak decay, displaced vertex, ρ mass smeared
- $\rho^+ \rightarrow \pi^+ \pi^0$: ρ polarized, $|\mathcal{M}|^2 \propto \cos^2 \theta$ in ρ rest frame
- $\pi^0 \rightarrow e^+ e^- \gamma$: Dalitz decay, $m(e^+ e^-)$ peaked

Dedicated programs, with special attention to polarization effects:

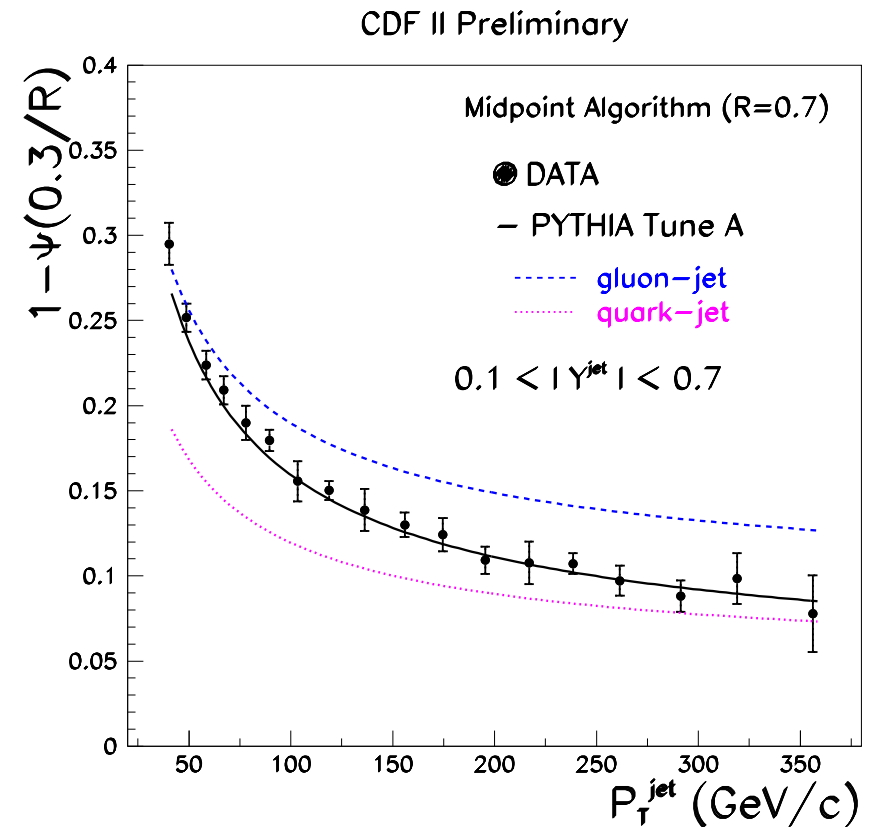
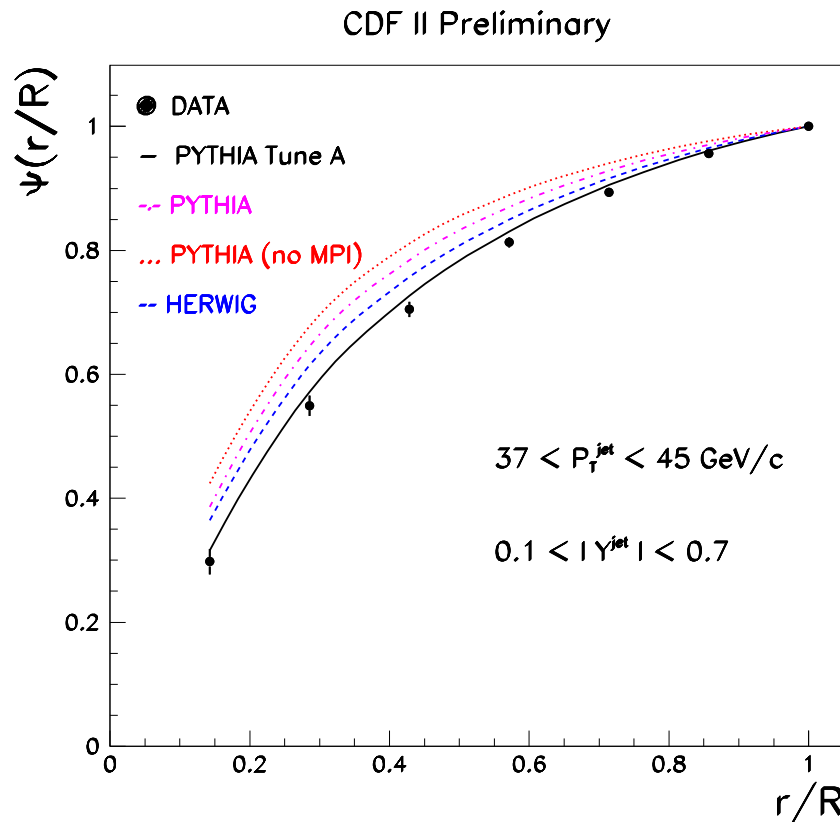
- EVTGEN: B decays
- TAUOLA: τ decays

Jet Universality

Question: are jets the same in all processes?

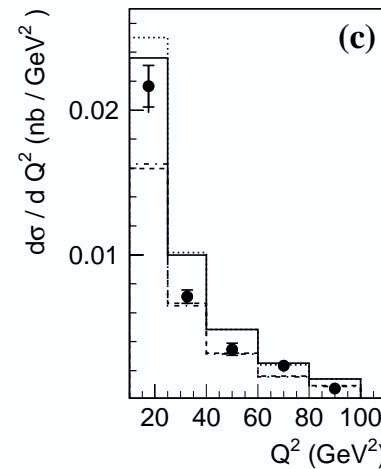
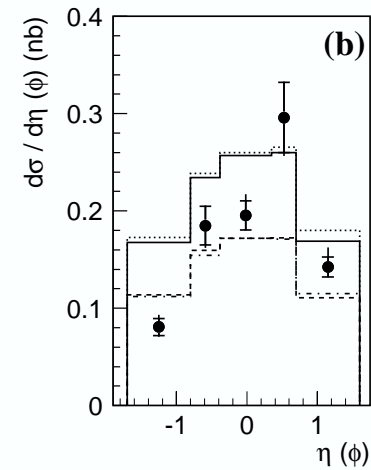
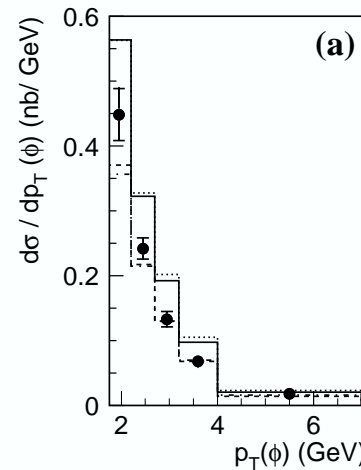
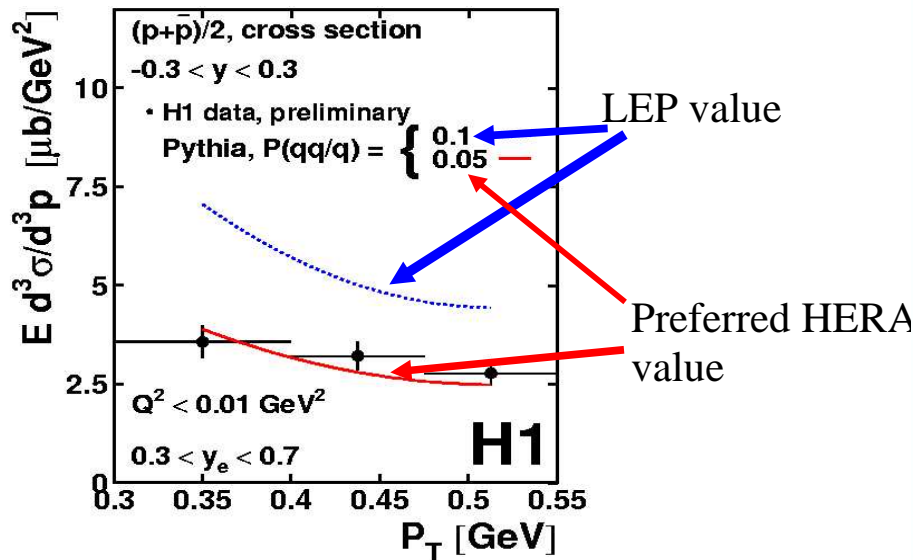
Answer 1: no, at LEP mainly quarks jets, often b/c,
at LHC mainly gluons, if quarks then mainly u/d.

Answer 2: no, perturbative evolution gives calculable differences.



Answer 3: (string) hadronization mechanism assumed universal, but is not quite.

$E d^3\sigma/d^3p$: Dependence on proton P_T



so discrepancies $\begin{matrix} P_{qq}/P_q & = & 0.1 \text{ at LEP,} & = & 0.05 \text{ at HERA} \\ P_s/P_u & = & 0.3 \text{ at LEP,} & = & 0.2 \text{ at HERA} \end{matrix}$

- Reasons? HERA dominated by “beam jets”, so
- Less perturbative evolution \Rightarrow strings less “wrinkled”?
 - Many overlapping strings \Rightarrow collective phenomena?

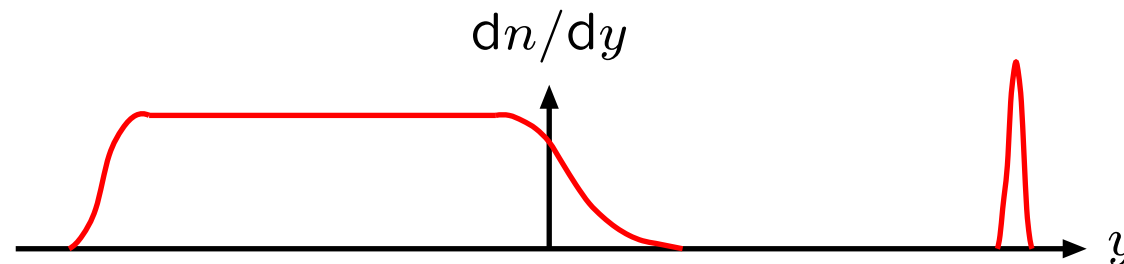
Other program tasks/elements

- Diffractive physics (\approx rapidity-gap physics)

$$\sigma_{\text{el}} \approx 25 \text{ mb} \quad pp \rightarrow pp$$

LHC: $\sigma_{\text{diff}} \approx 25 \text{ mb} \quad pp \rightarrow pX, pp \rightarrow X_1X_2, \text{ etc}$

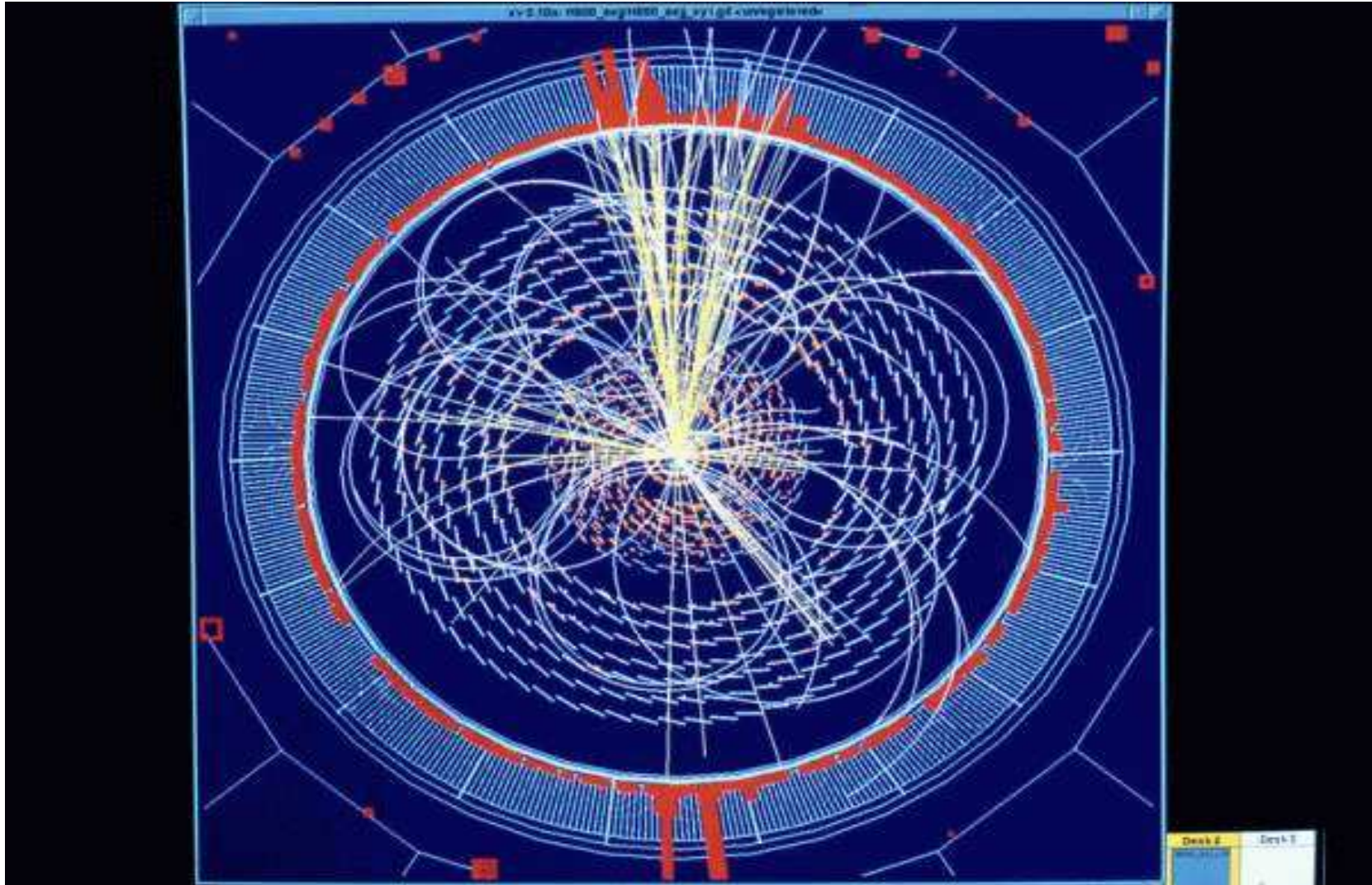
$$\sigma_{\text{inel,nondiff}} \approx 50 \text{ mb} \quad pp \rightarrow X \text{ (without obvious subdivision of } X\text{)}$$



- Colour reconnection: how well can we trust “perturbatively” calculable colour flow in soft region?
- Bose-Einstein: must we use amplitudes to describe production of identical particles? ($\sim 50 \pi^+$, $\sim 50 \pi^-$, $\sim 70 \pi^0$ per event)
- Event measures; jet clustering routines; other utilities

... and more

Event Generator Practicalities



Event generation structure

1) Initialization step

- select process(es) to study
- modify physics parameters: m_t, m_h, \dots
 - set kinematics constraints
- modify generator performance
 - initialize generator
 - book histograms

2) Generation loop

- generate one event at a time
- analyze it (or store for later use)
 - add results to histograms
 - print a few events

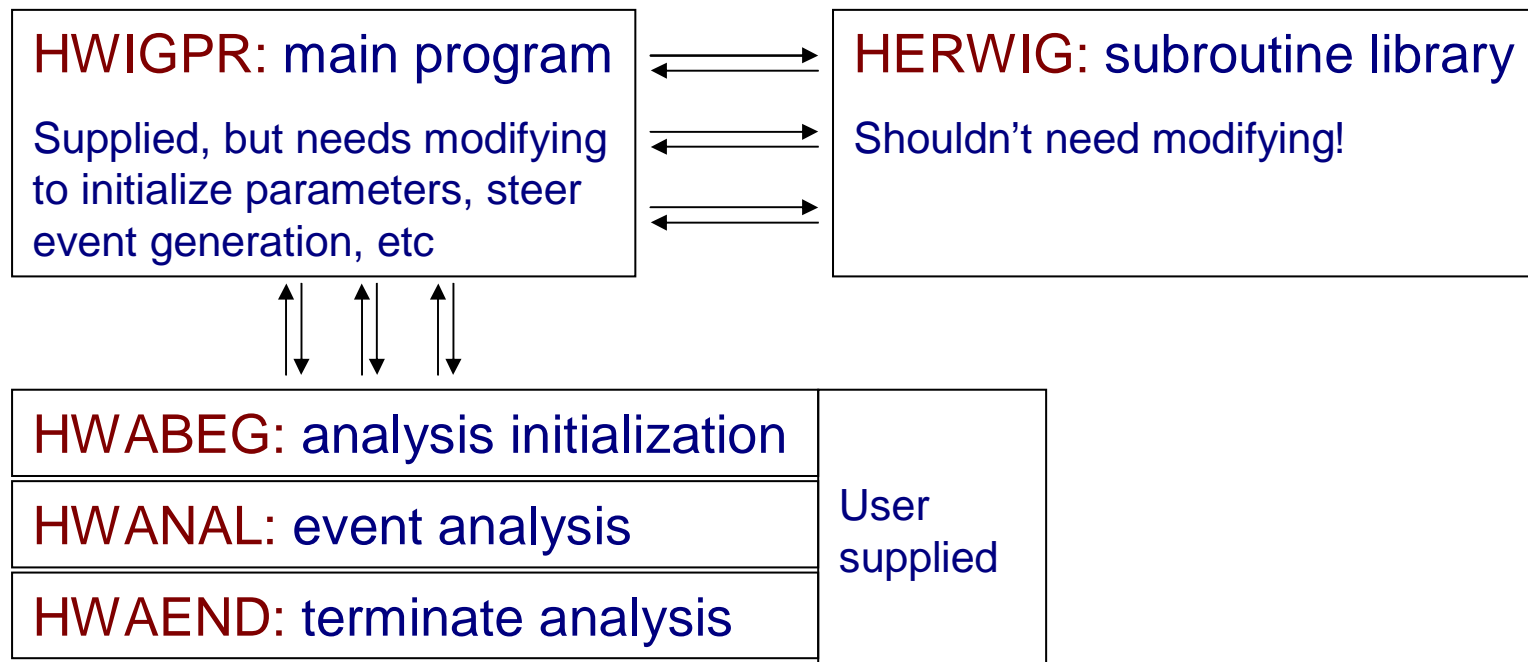
3) Finishing step

- print deduced cross-sections
- print/save histograms etc.

How to run event generators

Often forced to use what is allowed by constricted collaboration framework, but for maximal power and minimal bugs run raw generator:

- HERWIG, ISAJET: supplied but modifiable main program, calling user-written routines



- PYTHIA: generator is subroutine package, user writes main program

```

C...Arithmetic in double precision; integer functions; PYDATA.
  IMPLICIT DOUBLE PRECISION(A-H, O-Z)
  INTEGER PYK,PYCHGE,PYCOMP
  EXTERNAL PYDATA
C...The event record and other common blocks.
  COMMON/PYJETS/N,NPAD,K(4000,5),P(4000,5),V(4000,5)
  COMMON/PYDAT2/KCHG(500,4),PMAS(500,4),PARF(2000),VCKM(4,4)
  COMMON/PYSUBS/MSEL,MSELPD,MSUB(500),KFIN(2,-40:40),CKIN(200)
  COMMON/PYPARS/MSTP(200),PARP(200),MSTI(200),PARI(200)
C...Physics scenario.
  MSEL=0          ! Mix subprocesses freely
  MSUB(102)=1     ! g + g -> h0
  MSUB(123)=1     ! f + f' -> f + f' + h0
  MSUB(124)=1     ! f + f' -> f" + f"' + h0
  PMAS(25,1)=300D0 ! Nominal Higgs mass.
C...Run parameters.
  NEV=1000        ! Number of events
  ECM=14000D0     ! CM energy of run
  CKIN(1)=200D0   ! Minimum Higgs mass.
  CKIN(2)=400D0   ! Maximum Higgs mass.
C...Initialize and book histogram(s).
  CALL PYINIT('CMS','p','p',ECM)
  CALL PYBOOK(1,'Higgs mass distribution',80,200D0,400D0)
C...Generate events and look at first few.
  DO 200 IEV=1,NEV
    CALL PYEVNT
    IF(IEV.LE.1) CALL PYLIST(1)
C...Find Higgs and fill its mass. End event loop.
  DO 150 I=7,9
    IF(K(I,2).EQ.25) CALL PYFILL(1,P(I,5),1D0)
  150  CONTINUE
  200  CONTINUE
C...Final output.
  CALL PYSTAT(1)  ! Print cross section table
  CALL PYHIST     ! Print histogram(s)
  END

```



```

C...Test program to generate ttbar events at Tevatron using PYTHIA
C...internal ttbar production subprocesses.
C...Ref: PYTHIA Tutorial, Fermilab, Dec 2004.
C ----- PREAMBLE: COMMON BLOCK DECLARATIONS ETC -----
C...All real arithmetic done in double precision.
      IMPLICIT DOUBLE PRECISION(A-H, O-Z)
C ----- PYTHIA SETUP -----
C...Number of events to generate
      NEV=100
C...Select type of events to be generated: ttbar (using PYGIVE)
C...And use the new world average mt.
      CALL PYGIVE('MSEL=6')
      CALL PYGIVE('PMAS(6,1)=178.0')
C...Initialize PYTHIA for Tevatron ppbar collisions
      ECM=1960D0
      CALL PYINIT('CMS','p','pbar',ECM)
C...Initialize user stuff, e.g. book histograms etc.
      CALL MYSTUF(0,NEV)
C ----- EVENT LOOP -----
      DO 1000 IEV=1,NEV
C...Generate event
          CALL PYEVNT
C...Print out the event record of the first event
          IF (IEV.EQ.1) CALL PYLIST(2)
C...Do event-by-event user stuff, e.g. fill histograms.
          CALL MYSTUF(1,IEV)
      1000 CONTINUE
C ----- FINALIZATION -----
C...Print some info on cross sections and errors/warnings
      CALL PYSTAT(1)
C...Finalize my user stuff, e.g. close histogram file.
      CALL MYSTUF(2,NEV)
      END

```



On To C++

Currently HERWIG and PYTHIA are successfully being used,
also in new LHC environments, using C++ wrappers

Q: Why rewrite?

A1: Need to clean up!

A2: Fortran 77 is limiting

Q: Why C++?

A1: All the reasons for ROOT, Geant4, ...

(“a better language”, industrial standard, ...)

A2: Young experimentalists will expect C++

(educational and professional continuity)

A3: Only game in town! **Fortran 90**

So far mixed experience:

- Conversion effort: everything takes longer and costs more
(as for LHC machine, detectors and software)
- The physics hurdle is as steep as the C++ learning curve

C++ Players

PYTHIA7 project \implies **ThePEG**

Toolkit for High Energy Physics Event Generation
(L. Lönnblad; S. Gieseke, A. Ribon, P. Richardson)

HERWIG++: complete reimplementaion
(B.R. Webber; S. Gieseke, A. Ribon, P. Richardson,
M. Seymour, P. Stephens, 3 new)

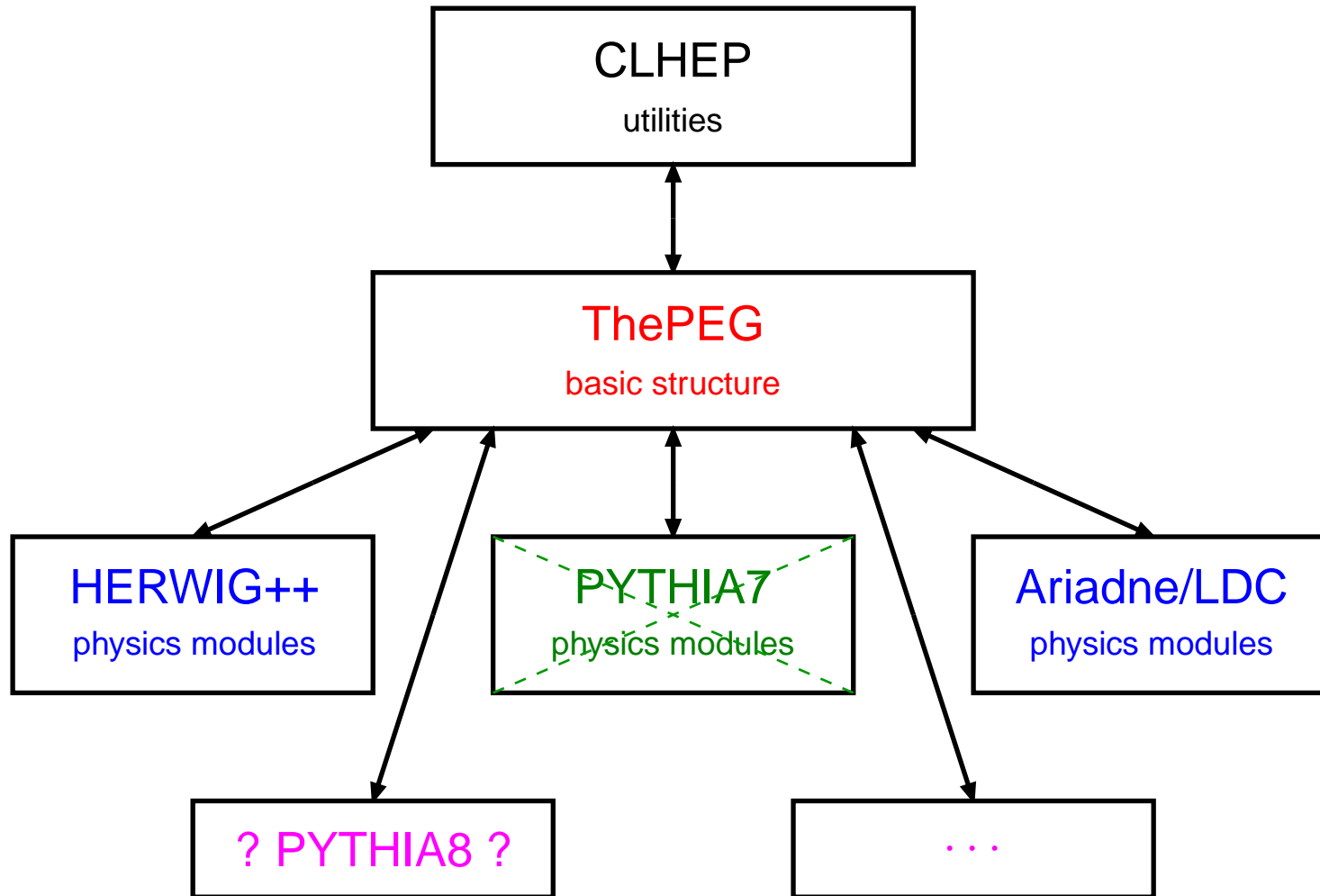
ARIADNE/LDC: to do ISR/FSR showers, multiple interactions
(L. Lönnblad; N. Lavesson)

SHERPA: partly wrappers to PYTHIA Fortran; has CKKW
(F. Krauss; T. Gleisberg, S. Hoeche, A. Schaelicke,
S. Schumann, J. Winter)

PYTHIA8: restart to write complete event generator
(T. Sjöstrand, (S. Mrenna?, P. Skands?))

What is ThePEG?

Toolkit for High Energy Physics Event Generation



not SHERPA

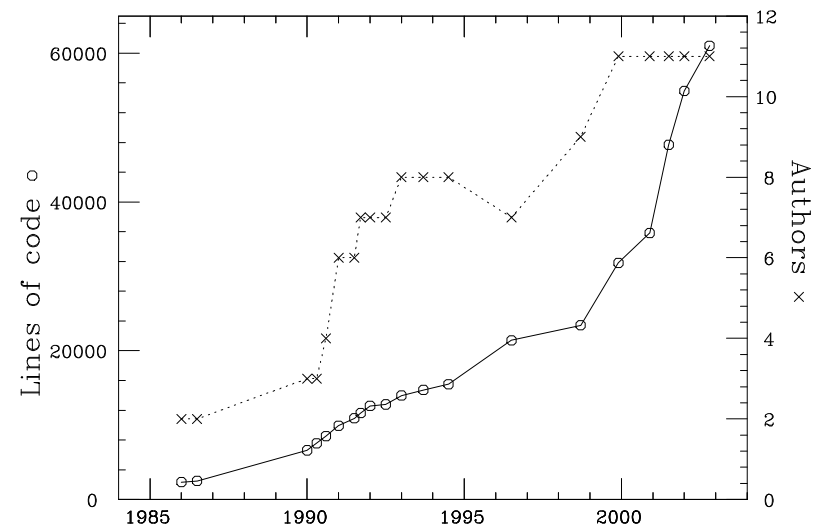
Running ThePEG

- ThePEG defines a set of abstract `Handler` classes for hard partonic sub-processes, parton densities, QCD cascades, hadronization, ...
- These handler classes interact with the underlying structure using a special `Event Record` and a pre-defined set of `virtual` functions.
- The procedure to implement e.g. a new hadronization model, is to write a new (C++) class *inheriting* from the abstract `HadronizationHandler` base class, implementing the relevant virtual functions.
- The end-user will use a setup program to be able to pick objects corresponding to different physics models to build up an `EventGenerator` which then can be run interactively or off-line, or as a special slave program e.g. for Geant4.
- The setup program is used to choose between a multitude of pre-defined generators, to modify parameters and options of the selected models and, optionally, to specify the analysis to be done on the generated events.
- The `Repository` is the central part of the setup phase. It handles a structured list of all available objects and allows the user to manipulate them.

The new generator Herwig++

A completely new event generator in C++

- Aiming at full multi-purpose generator for LHC and future colliders.
- Preserving main features of HERWIG such as
 - angular ordered parton shower
 - cluster hadronization
- New features and improvements
 - covariant shower formulation
 - improved parton shower evolution for heavy quarks
 - consistent radiation from unstable particles (multiscale evolution)



Growth of Fortran HERWIG

Hard interactions

- Basic ME's included in **ThePEG**, such as:

$$e^+e^- \rightarrow q\bar{q}, \text{ partonic } 2 \rightarrow 2,$$

we use them.

- Soft and hard **matrix element corrections** implemented for $e^+e^- \rightarrow q\bar{q}g$.
- **AMEGIC++** will provide arbitrary ME's for multiparton final states via **AMEGICInterface**.
- **LesHouchesFileReader** enables to read in and process *any* hard event generated by parton level event generators (MadGraph/MadEvent, AlpGen, CompHEP, ...).
- CKKW ME+PS foreseen.
- Other authors can easily include their own matrix elements (\rightarrow *safety* of OO code)

New/Future: HELAS like structures are already implemented for decays and spin correlations \longrightarrow allows us to code simple processes efficiently.

. . . and New Decays!

- Better decayers are being developed for almost all decay modes.
- → B decays.
- Spin correlations will be included.
- Major effort ongoing
 - a universal database is being set up.
 - contains 448 particles and 2607 decay modes at present.
 - possibility to generate configuration files for different generators (they need to write their own code however. . .).
- Particle data book as guideline.

→ *look at examples. . .*

Herwig++ Particle Properties DataBase

This is the development version of the Herwig++ particle properties database. This is intended to replace the storage of particle properties as a text file to improve maintainance and accessibility.

This version is for the Herwig authors only and much of the information is preliminary.

The database currently contains 448 particles and 2607 decay modes.

The information is available in a number of forms

- The particles [numerically listed](#) according to the [PDG code](#)
- The particles [listed](#) according to the multiplets taken from the [PDG](#)
- The [decayers](#)
- The [Width Generators](#)
- The [Mass Generators](#)
- The [references](#)
- Generate the [input files](#) for event generation

The contents of the database can be altered by following the links in the particle table or particle descriptions or by selecting an option below

- Add or modify a particle: 0 _____
- Add a decay mode for particle with id: 0 _____
-
-
-
-
-
-
- Set the decay modes for a particle to the charge conjugates of the antiparticle: 0 _____
-
-
-

Simple checks on the contents of the database

-

-
-
-

[Peter Richardson](#)

Last modified: Mon Jan 31 17:56:08 GMT 2005

What's next?

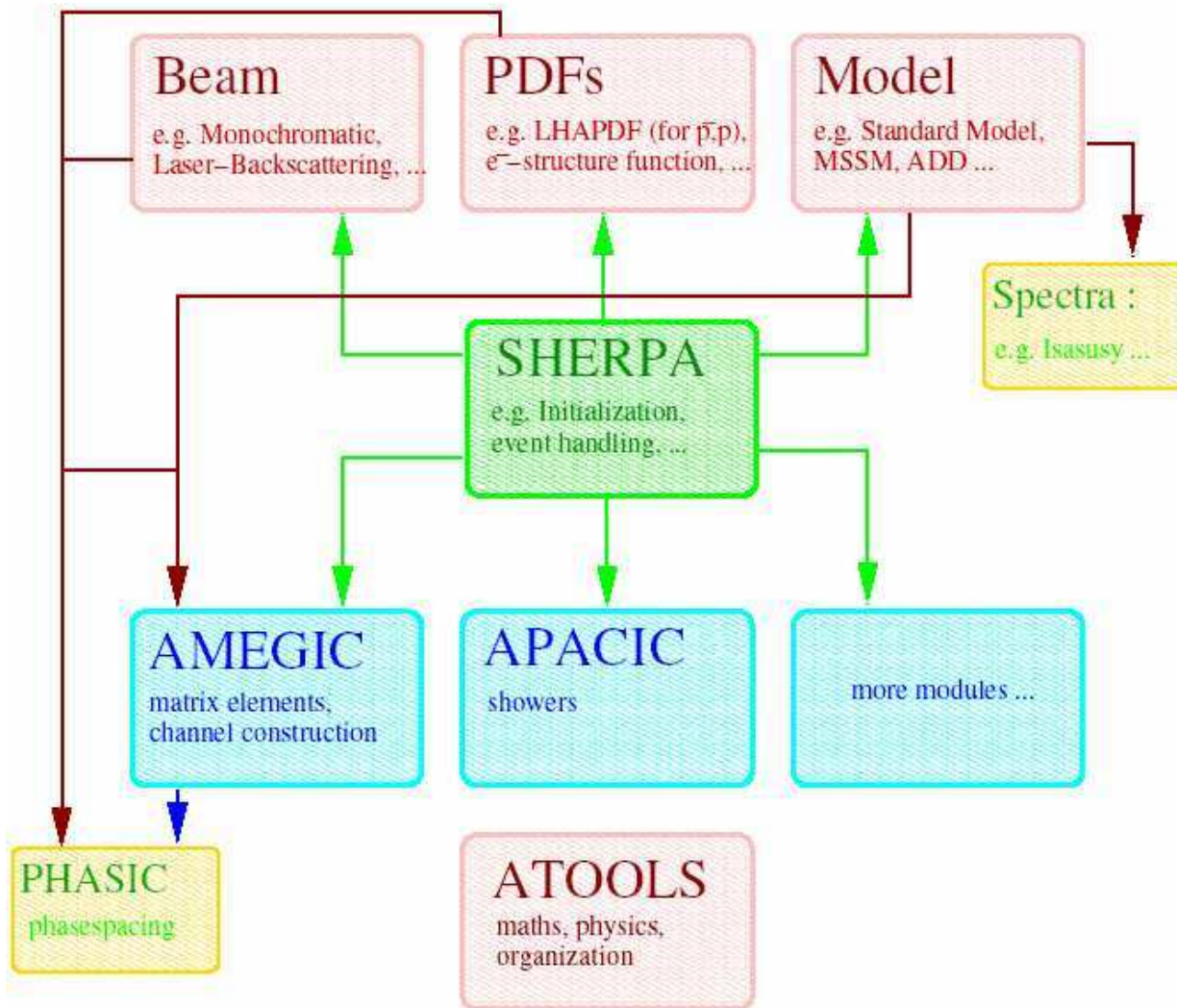
Near Future. . .

- ★ Initial state shower:
 - Complete implementation and tests.
- ★ Refine e^+e^- :
 - Full CKKW ME+PS matching.
 - Precision tune to LEP data should be possible.
- ★ with IS and FS showers running:
 - we can start to test Drell–Yan and jets in pp collisions.
 - cross check with Tevatron data and finally make predictions for the LHC.
- ★ Underlying Event.
- ★ Hadronic Decays: *NEW!* many new decayers, τ -decays, Spin correlations (P Richardson).
- ★ *New Ideas*: soft gluons, improved shower algorithm, NLO, . . .

Schedule?

- Ready for LHC!

SHERPA



Conclusions/Outlook

SHERPA including the ME's of AMEGIC++ and the CKKW prescription to combine them with the PS **is a powerful tool** to attempt the description of present-day Tevatron data and to study the extrapolation to LHC energies.

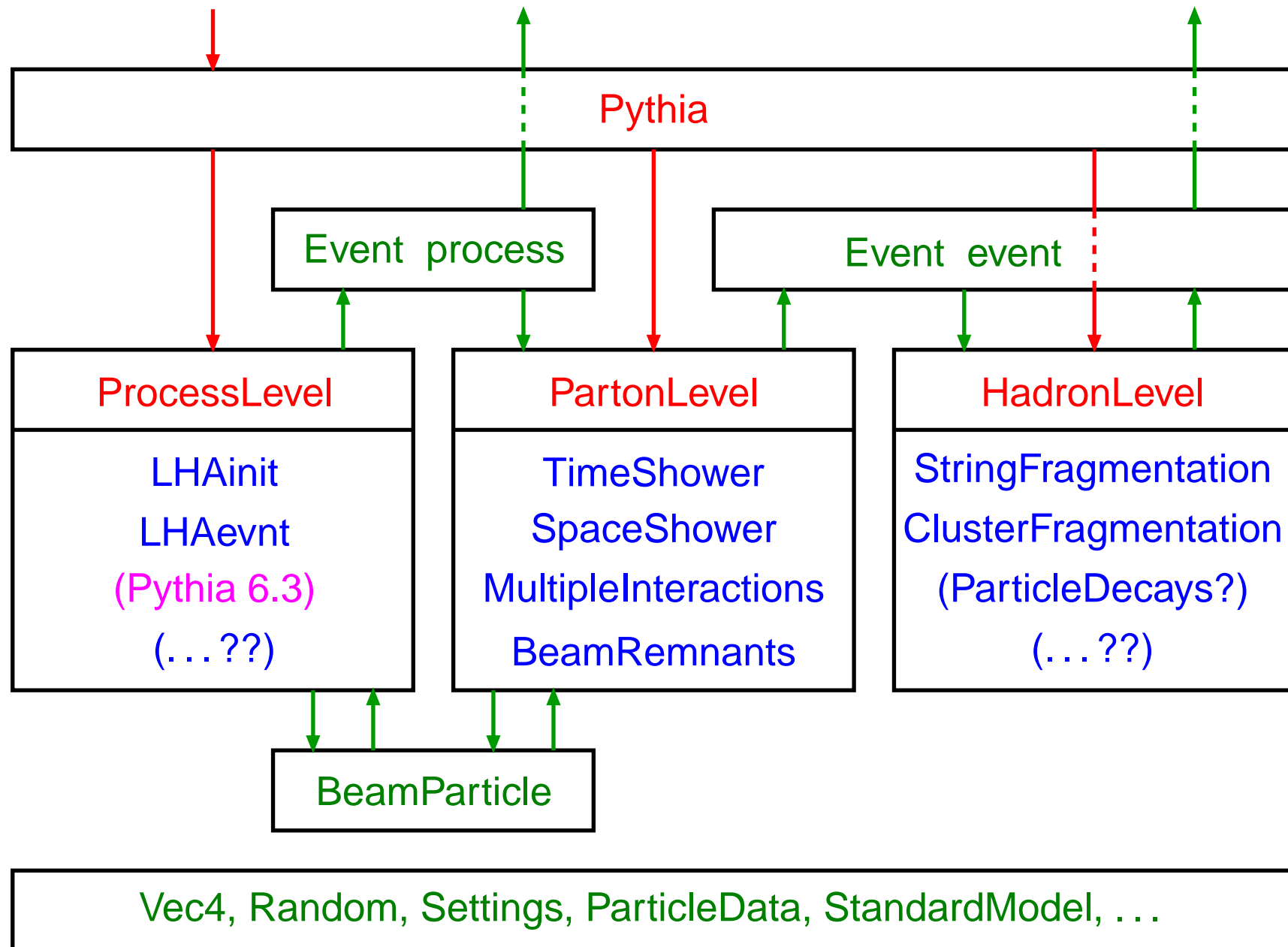
The next release will include:

- The simple hard underlying event model
- Revision of the phase space integration
(enhanced integration performance and unweighting efficiencies)
- Support of the SLHA for MSSM spectrum input

Sources:

- T. Gleisberg, S. Höche, F. Krauss, A. Schälicke, S. S. and J. Winter, JHEP 0402:056,2004
- download (SHERPA α -1.0.4), manual, bug reports etc. under
<http://www.physik.tu-dresden.de/~krauss/hep>

Current PYTHIA8 structure



Current PYTHIA8 status

Existing classes			Missing classes
Process	LHAinit	★	ThePEG input, alternatively Cross section administration
Level	LHAevnt	★★	
Parton	TimeShower	★ ★ ★	Phase space selection Process matrix elements Parton density libraries Resonance decays
Level	SpaceShower	★★	
	MultipleInteractions	★	
	BeamRemnants	★	
Hadron	StringFragmentation	★	...
Level	ClusterFragmentation	★	ME/PS matching
—	Event	★★	Junction fragmentation ParticleDecays
	BeamParticle	★★	
	Vec4, Random	★ ★ ★	Bose-Einstein ...
	Settings	★★	
	ParticleData	★	

⇒⇒ Roughly according to three-year plan so far!

Outlook

Generators in state of continuous development:

- ★ better & more user-friendly general-purpose matrix element calculators+integrators ★
 - ★ new libraries of physics processes, also to NLO ★
 - ★ more precise parton showers ★
 - ★ better matching matrix elements \Leftrightarrow showers ★
 - ★ improved models for underlying events / minimum bias ★
 - ★ upgrades of hadronization and decays ★
 - ★ moving to C++ ★
- \Rightarrow always better, but never enough

But what are the alternatives, when event structures are complicated and analytical methods inadequate?

Final Words of Warning

[. . .] The Monte Carlo simulation has become the major means of visualization of not only detector performance but also of physics phenomena. So far so good. But it often happens that the physics simulations provided by the Monte Carlo generators carry the authority of data itself. They look like data and feel like data, and if one is not careful they are accepted as if they were data.

[. . .] I am prepared to believe that the computer-literate generation (of which I am a little too old to be a member) is in principle no less competent and in fact benefits relative to us in the older generation by having these marvelous tools. They do allow one to look at, indeed visualize, the problems in new ways. But I also fear a kind of “**terminal illness**”, perhaps **traceable to the influence of television** at an early age. There the way one learns is simply **to passively stare into a screen and wait for the truth to be delivered**. A number of physicists nowadays seem to do just this.

J.D. Bjorken

from a talk given at the 75th anniversary celebration of the Max-Planck Institute of Physics, Munich, Germany, December 10th, 1992. As quoted in: Beam Line, Winter 1992, Vol. 22, No. 4