

High-Energy pA Collisions and Cosmic Ray Airshowers

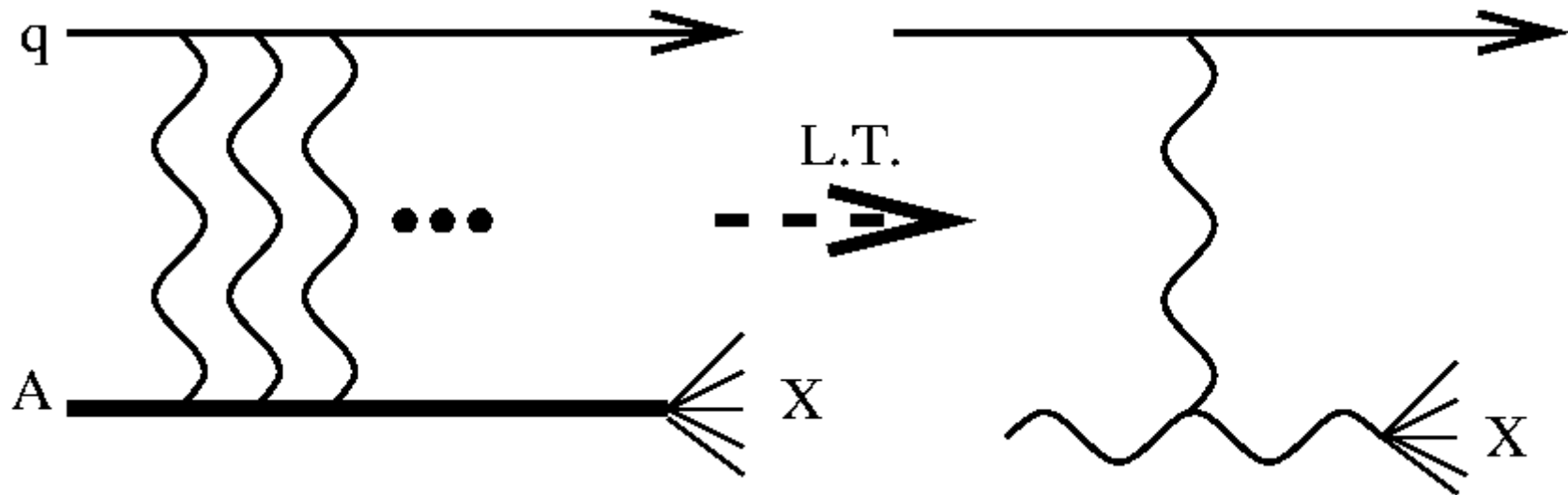
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hep-ph/0408073

- qA Scattering
- Forward Hadron Production
- $Q_s(x)$
- Application to Cosmic Ray Airshowers
- Summary

Is there sensitivity to the evolution of $Q_s(x)$?

Quark-Nucleus Elastic Scattering



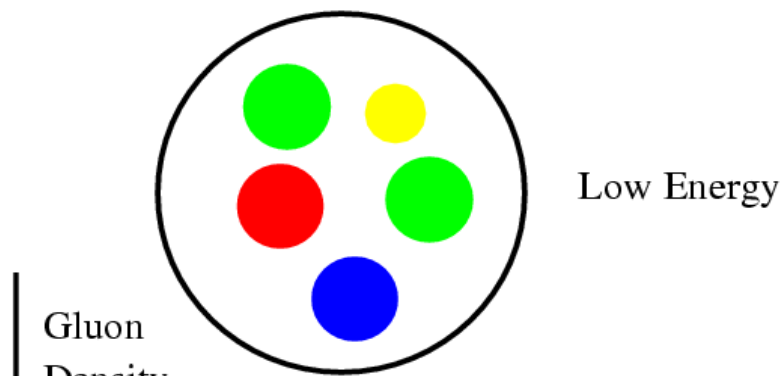
Quark Scattering Amplitude :

$$\langle q \text{ out} | p \text{ in} \rangle = \bar{u}(q) \tau(q, p) u(p)$$

with

$$\tau(q, p) = 2\pi \delta(p^- - q^-) \gamma^- \int d^2 x_t [V(x_t) - 1] e^{i x_t (q_t - p_t)}$$

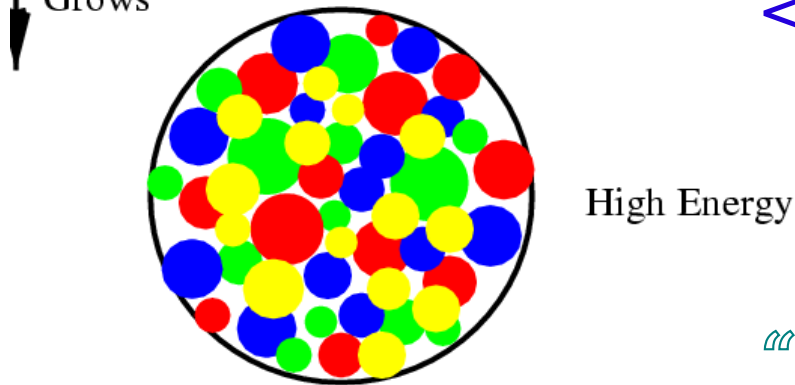
$$V(x_t) = \mathcal{P} \exp \left(-ig^2 \int_{-\infty}^{\infty} dx^- \frac{1}{\partial_t^2} \rho^a(x^-, x_t) t^a \right)$$



<-- Dilute Parton Gas

Low Energy

<-- 'Saturated' Classical Color Field



High Energy

“McLerran-Venugopalan Model”

- Small-x gluons evolve slowly
- Color averaging with static random sources :

$$\langle O \rangle = \int \mathcal{D}\rho \ O[\rho] \exp \left(- \int d^2x_t dx^- \text{tr} \rho^2 / \mu^2 \right)$$

μ^2 is average color charge squared per area

- Field of nucleus : $[D_\mu, F^{\mu\nu}] = \delta(x^-) \delta^{\nu+} g \rho$

Color averaging with a Gaussian leads to

$$\langle V(x_t) \rangle_\rho = \exp -g^4 C \chi \int d^2 z_t G_0^2(x_t - z_t)$$

$$\langle V(x_t) V^\dagger(\bar{x}_t) \rangle_\rho = \exp -g^4 C \chi \int d^2 z_t [G_0(x_t - z_t) - G_0(\bar{x}_t - z_t)]^2$$

with $C = (N_c^2 - 1)/4N_c$, $G_0(x_t) = - \int \frac{d^2 k_t}{(2\pi)^2} \frac{\exp i k_t x_t}{k_t^2}$

and

$$\chi(x^-, x_t) = \int_{x^-}^{x_A^-} dz^- \mu^2(z^-, x_t) \rightarrow \frac{1}{\pi R_A^2} \frac{N_c}{N_c^2 - 1} \int dx g_A(x)$$

Defining the saturation momentum

$$Q_s^2 = 4\pi^2 \alpha_s^2 \frac{N_c^2 - 1}{N_c} \chi$$

--->

$$\frac{d\sigma^{qA}}{dq^- d^2 q_t d^2 b} = \delta(q^- - p^-) C(q_t)$$

$$C(q_t) = \int \frac{d^2 r_t}{(2\pi)^2} e^{iq_t r_t} \left\{ \exp \left[-2Q_s^2 \int_{\Lambda} \frac{d^2 p_t}{(2\pi)^2} \frac{1}{p_t^4} (1 - e^{ip_t r_t}) \right] - 2 \exp \left[-Q_s^2 \int_{\Lambda} \frac{d^2 p_t}{(2\pi)^2} \frac{1}{p_t^4} \right] + 1 \right\}$$

--->

$$d\sigma^{\text{el}}/d^2b = \left[1 - e^{-Q_s^2/4\pi\Lambda^2} \right]^2, \quad d\sigma^{\text{tot}}/d^2b = 2 \left[1 - e^{-Q_s^2/4\pi\Lambda^2} \right]$$

AD + JJM: PRL 89 (2002)

Limits for q_t cross section :

$$C(q_t) =$$

$$q_t \gg Q_s : \quad \frac{1}{2\pi^2} \frac{Q_s^2}{q_t^4} \left[1 + \frac{4}{\pi} \frac{Q_s^2}{q_t^2} \log \frac{q_t}{\Lambda} + \mathcal{O} \left(\frac{Q_s^2}{q_t^2} \right) \right]$$

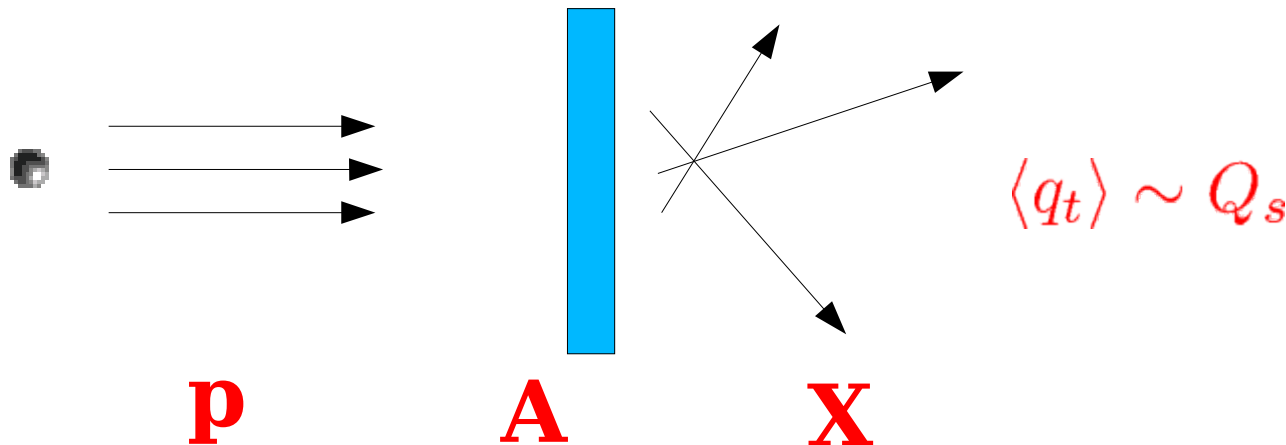
$$q_t \lesssim Q_s : \quad \frac{1}{Q_s^2 \log Q_s/\Lambda} \exp \left(-\frac{\pi q_t^2}{Q_s^2 \log Q_s/\Lambda} \right)$$

Shattering the proton

Probability for quark to be scattered to $q_t \sim 0$ (with color exchange !):

$$\int_0^\Lambda d^2 q_t \frac{d\sigma^{\text{in}}}{d^2 b d^2 q_t} \simeq 1 - \exp\left(-\frac{\pi\Lambda^2}{Q_s^2 \log Q_s/\Lambda}\right) \simeq \frac{\pi\Lambda^2}{Q_s^2 \log Q_s/\Lambda}$$

--> **suppression of “beam-jet remnants”
(soft physics) in the BBL**



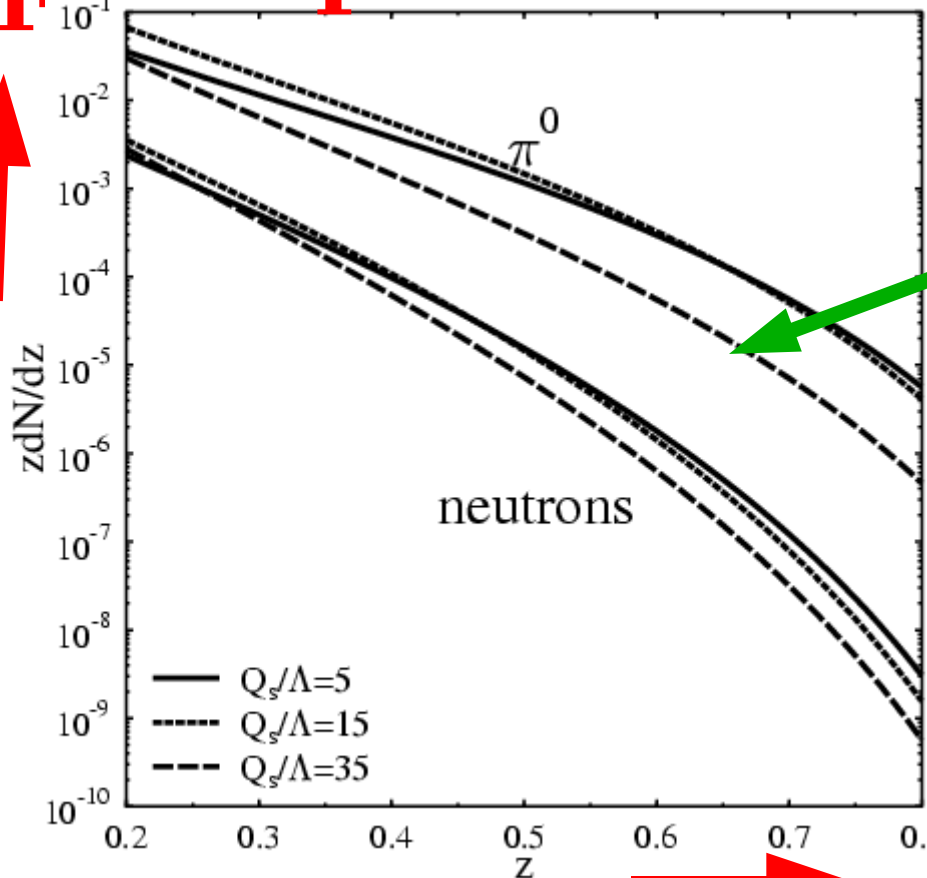
All partons resolved at scale Q_s , coherence of proton destroyed completely.

If one assumes indep. fragm. of scattered partons :

$$x_F \frac{d\sigma^{pA \rightarrow hX}}{dx_F d^2k_t d^2b} = \int_{x_F}^1 dx \frac{x}{x_F} f_{q/p}(x, Q_s^2) D_{h/q}\left(\frac{x_F}{x}, Q_s^2\right) \frac{d\sigma^{qA}}{d^2q_t d^2b}$$

“Limiting Fragmentation” curve in BBL

$x_F \frac{dN}{dx_F}$



Long. distribution steepens

Gerland, A.D., Strikman
PRL 90 (2003)

see also

Frankfurt, Guzey, McDermott,
Strikman: PRL 87 (2001)

x_F

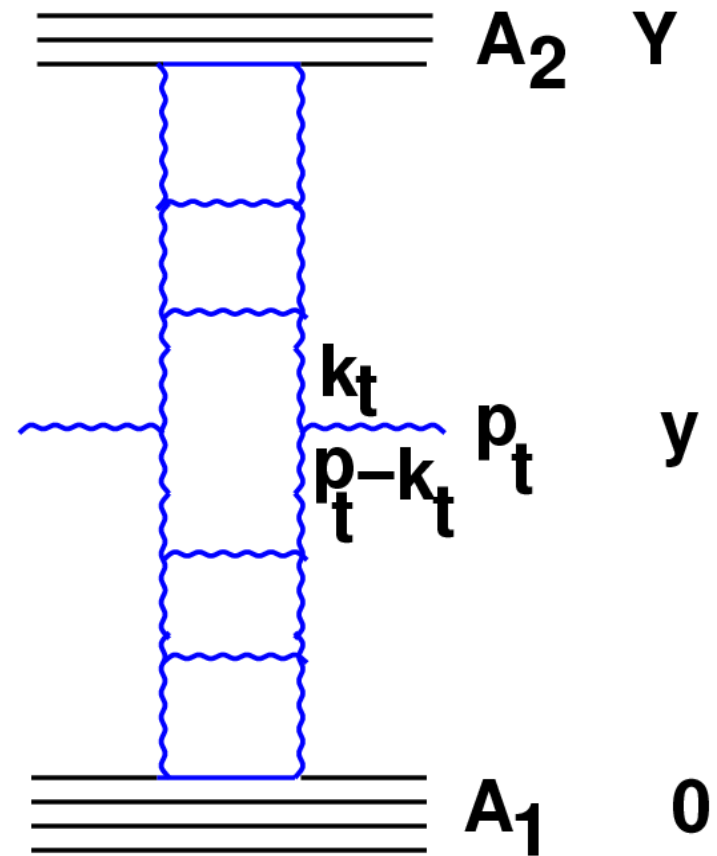
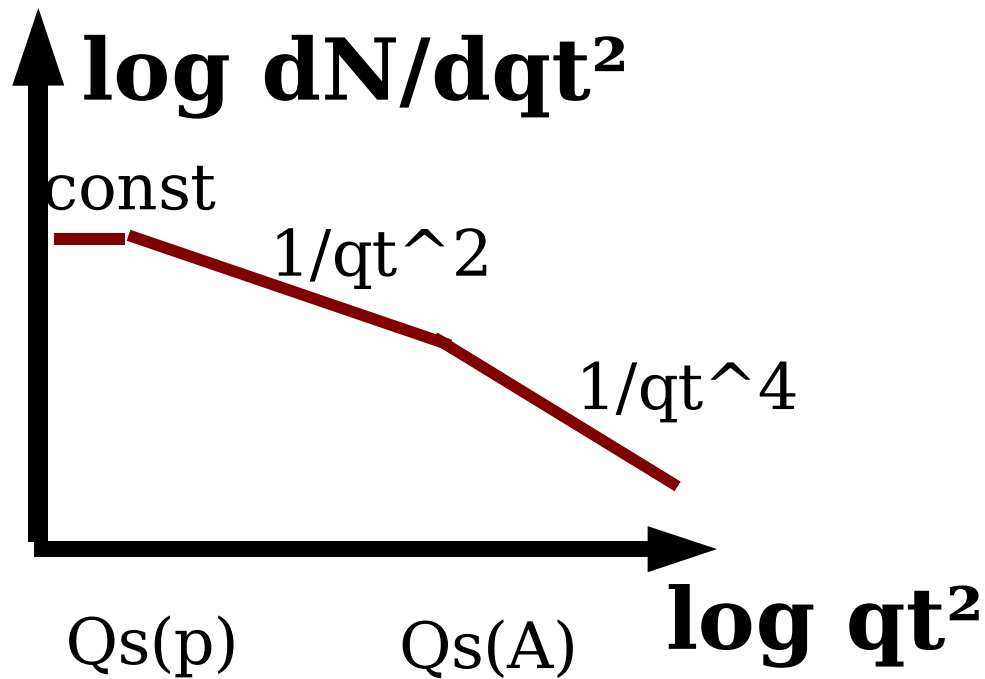
Gluon radiation

$$\frac{dN^g}{dydq_t^2} = 4\pi\alpha_s \frac{N_c}{N_c^2 - 1} \frac{1}{q_t^4}$$

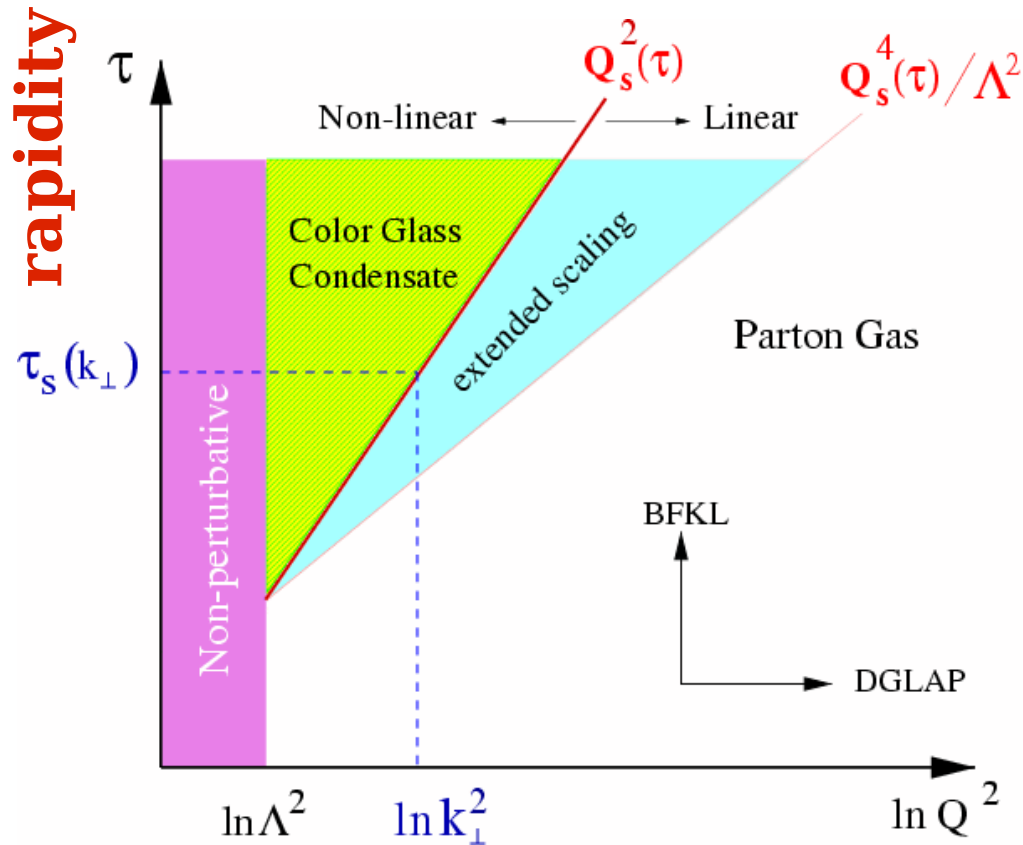
$$x_2 g_A(x_2, q_t^2) x_1 g(x_1, q_t^2)$$

$$xg(x, q_t^2) \propto \frac{1}{\alpha_s} \min(q_t^2, Q_s^2(x)) (1-x)^4$$

(Kharzeev, Levin, Nardi, NPA 2004)



Evolution of Q_s



Iancu, Venugopalan
hep-ph/0303204

determine $Q_s(x)$ from BFKL

fixed coupl. BFKL:

$$Q_s^2(y, A) = Q_0^2 \exp c\bar{\alpha}_s y \rightarrow Q_s^2(x) = Q_0^2 (x_0/x)^\lambda$$

GB-W: $\lambda \sim 0.28$ Initial condition : $Q_0^2 \sim A^{1/3} \log A$

Factorized A- and x-dependence !

running coupl BFKL : $\bar{\alpha}_s(Q^2) = b_0 / \log Q^2 / \Lambda^2$

$$Q_s^2 = \Lambda^2 \exp \left(\log(Q_0^2 / \Lambda^2) \sqrt{1 + 2c \bar{\alpha}_s y} \right)$$

$$y \rightarrow 0 : \rightarrow Q_0^2 \exp \left(\bar{\alpha}_s c y \log Q_0^2 / \Lambda^2 \right) \stackrel{!}{=} Q_0^2 \exp \lambda y$$

$$\rightarrow Q_s^2 = \Lambda^2 \exp \sqrt{\log(Q_0^2 / \Lambda^2) (2\lambda y + \log Q_0^2 / \Lambda^2)}$$

$$y \rightarrow \infty : \rightarrow \Lambda^2 \exp \sqrt{2\lambda y \log Q_0^2 / \Lambda^2}$$

	RHIC	LHC	GZK
yP	10.7	17.3	26.1
Qs r.c.	1.1 GeV	2.4 GeV	5.9 GeV
Qs f.c.	1.4 GeV	4.5 GeV	19.2 GeV

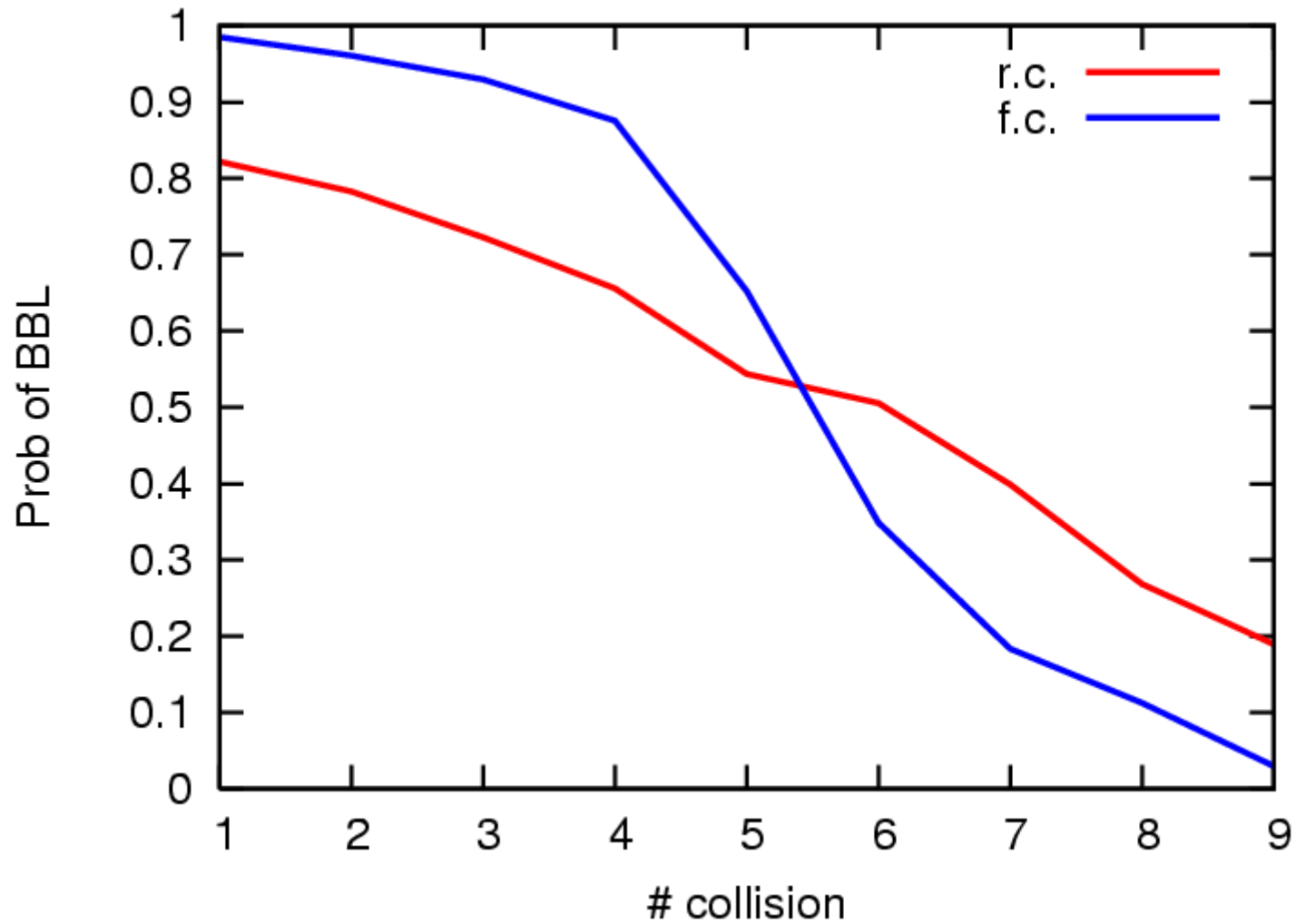
$\lambda = \mathbf{0.28}$; **central "p+N"**; $(Q_0/\Lambda)^2 \sim N_{val}/3$

BBL 1.0 : hA MC for high energies

- Distributions of produced partons in full phase space
(complete event, central+forward regions)
- Conserves energy and momentum
- Connects partons (val q and bremsstr. g) by strings
--> “absorption” of collinear g
- Fragmentation via standard PYTHIA / JETSET
- Realistic transv. density profile for target nucleus
- Low-energy / peripheral collisions handled by any
other model for L.T. + soft hadronic interactions,
for example Lund string models (here: SIBYLL)

here: $Q_s(b, x_F=0.001) > \sim 1 \text{ GeV}$

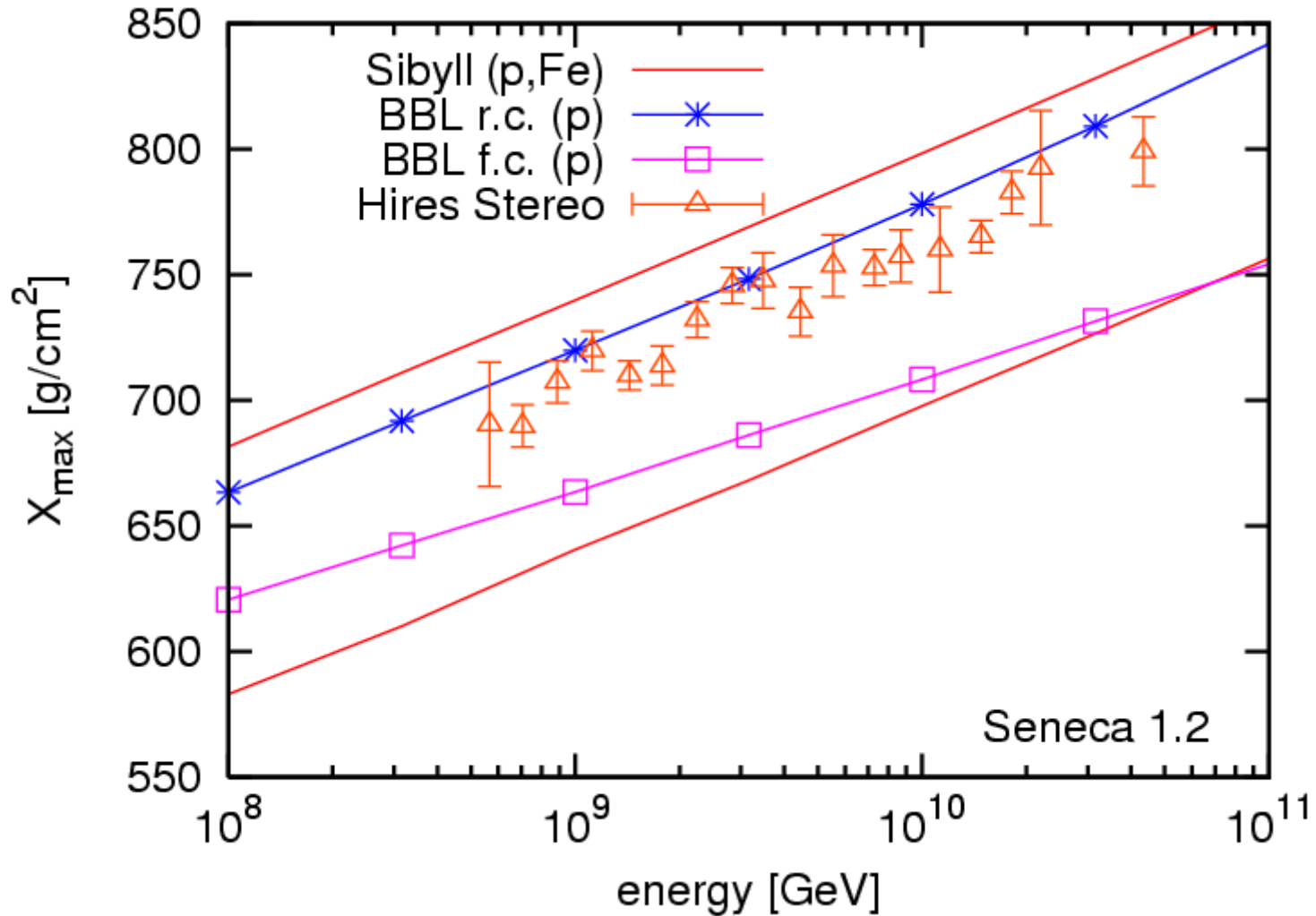
Area of “black disc” seen by leading hadron in n-th collision within one airshower



Cosmic Ray Airshowers

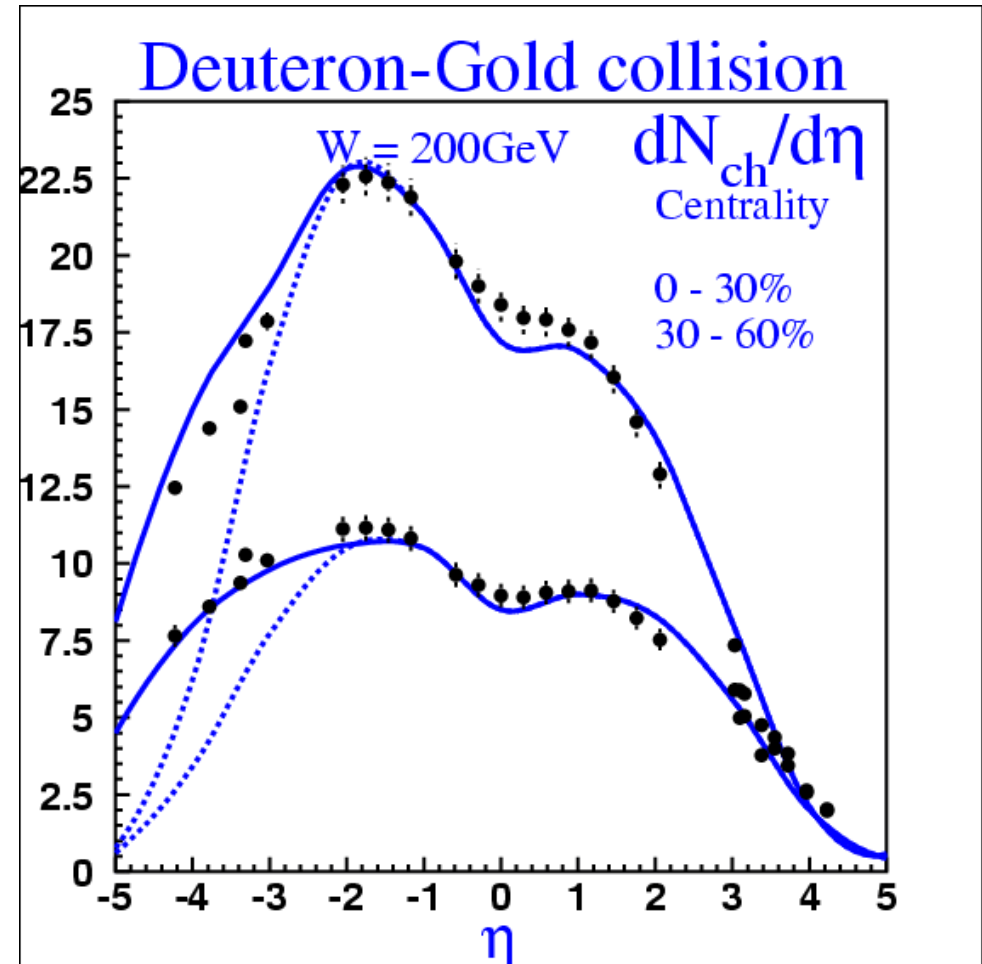
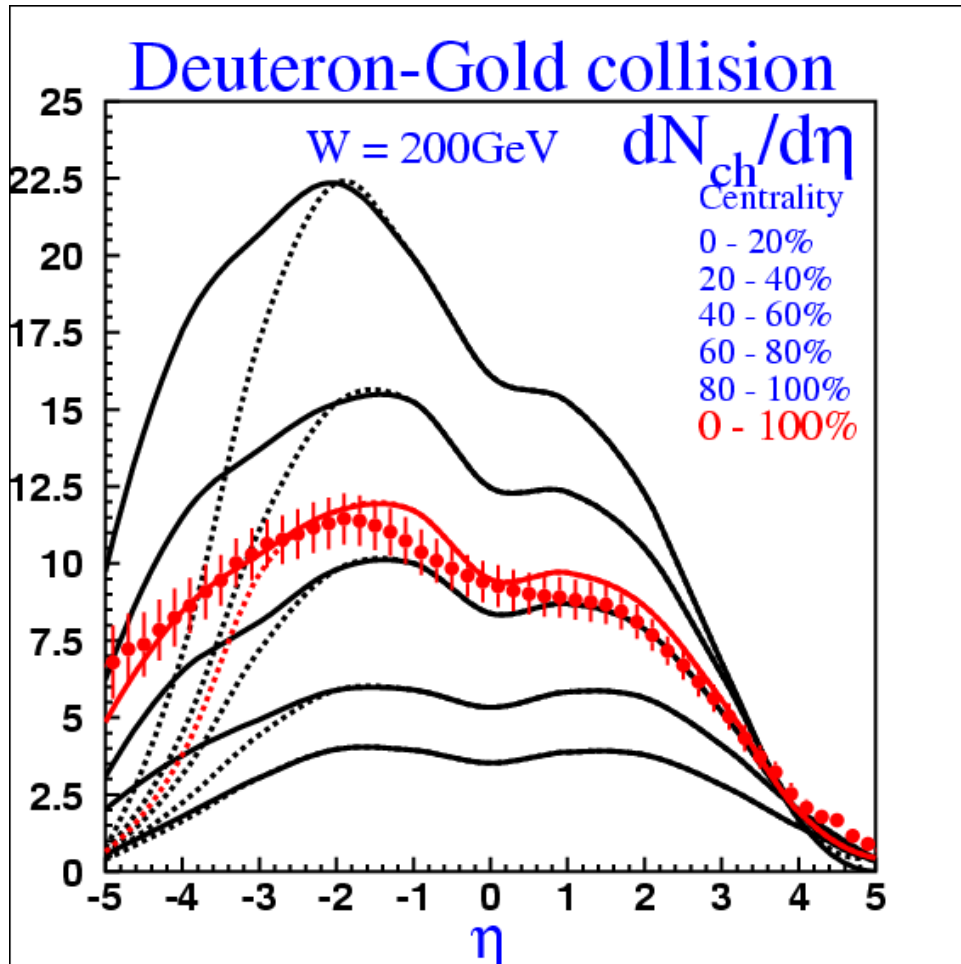
hep-ph/0408073

“penetration depth”



=> Sensitive to evolution of Q_s !!

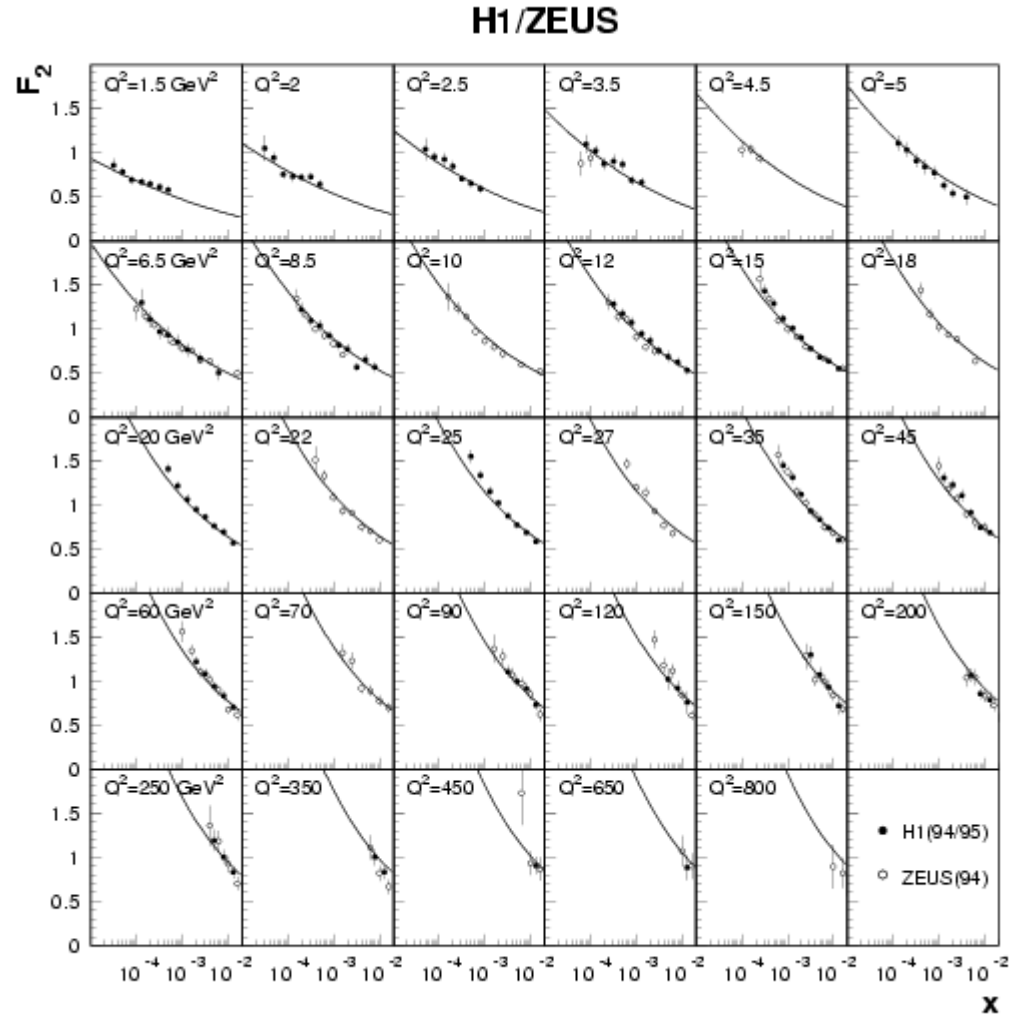
Fixed-Coupl / GB-W evolution at HERA & RHIC



KLN: [hep-ph/0212316](https://arxiv.org/abs/hep-ph/0212316)

GB-W: hep-ph/9903358

DIS on protons: F_2



=> for RHIC and HERA energies, f.c. evolution of Q_s works fine

Summary

- **High-energy pA is of great interest for understanding the high-density non-linear regime of QCD at small x: gluon saturation, unitarity ?**
- **Would be very interesting physics @ LHC!
(especially in the forward region)**
- **Cosmic ray airshowers are sensitive to QCD evolution scenario. Indications for a less rapid growth of $Q_s(x)$ as compared to RHIC or HERA.**