

2004 LHC Days in Split  
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# Importance of $\tau$ 's in the MSSM Higgs Boson Discovery

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# Introduction

In the MSSM, the SM-like lighter scalar  $h$  is expected to be found in several decay channels in the region of large  $m_A$  and  $\tan\beta$

To best disentangle between SM and MSSM, the heavy bosons,  $A$ ,  $H$  and  $H^\pm$  should be looked for !

Couplings to fermions and weak bosons (when compared to SM) lead to specific MSSM Higgs search strategies at large  $\tan\beta$

	$g_{\phi uu}$	$g_{\phi dd}$	$g_{\phi VV}$
$h$	$\cos\alpha/\sin\beta$	$-\sin\alpha/\cos\beta$	$\sin(\beta-\alpha)$
$H$	$\sin\alpha/\sin\beta$	$\cos\alpha/\cos\beta$	$\cos(\beta-\alpha)$
$A$	$\gamma_5/\tan\beta$	$\gamma_5 \tan\beta$	$0$



- suppression of  $HZZ/HWW$  couplings
- absence of  $AZZ/AWW$  couplings
- enhancement of  $\phi dd$  ( $d=\tau, b..$ ) couplings



Importance of decays to  $\tau$ 's,  $H/A/h \rightarrow \tau\tau$   
and production in association with  $b$ 's in  $gg \rightarrow bbH/A$

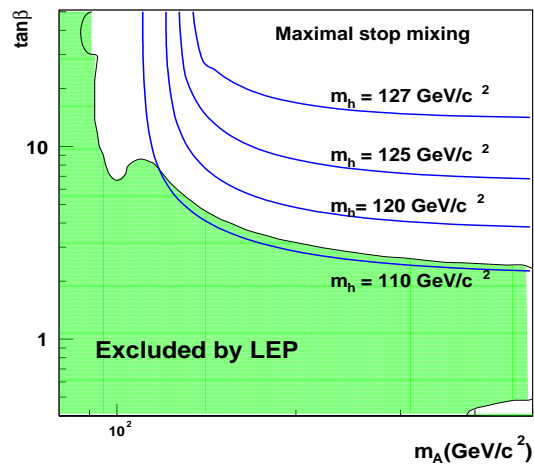


# MSSM parameter space and mass spectrum

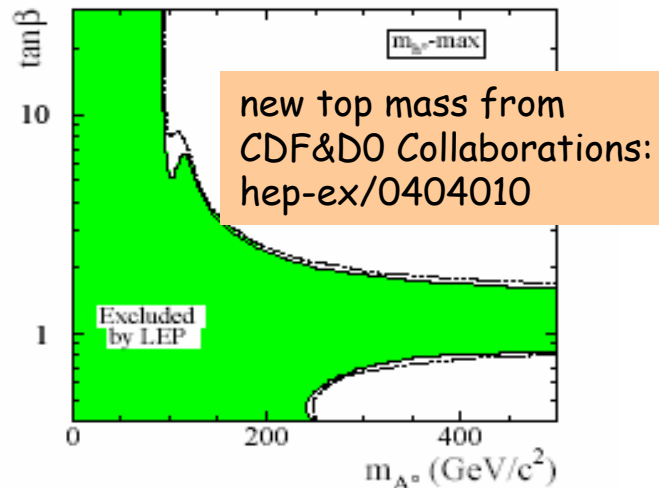
LEP-type SUSY scenario assumed (for most of the LHC studies):

$M_2 = 200 \text{ GeV}/c^2$ ,  $\mu = -200 \text{ GeV}/c^2$ ,  $M_{\text{gluino}} = 800 \text{ GeV}/c^2$ ,  $M_{\text{squark, slepton}} = 1 \text{ TeV}/c^2$   
 No stop mixing ( $X_t = 0$ ) or maximal stop mixing ( $X_t = 2450 \text{ GeV}/c^2$ )

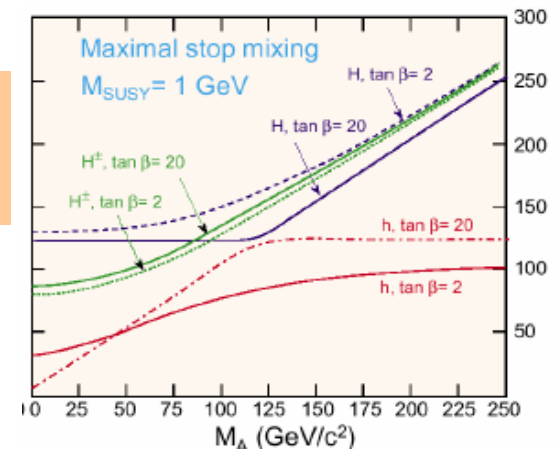
LEP exclusion  
with  $m_{\text{top}} = 175 \text{ GeV}$



LEP exclusion with  
 $m_{\text{top}} = 179.3 \text{ GeV}$



MSSM Higgs boson  
mass spectrum  
with maximal stop mixing



Mass of the lighter scalar  $h$  with two-loop/RGE-improved radiative corrections:

$m_h^{\text{max}} = 113 \text{ (116) GeV}$  for no stop mixing

$m_h^{\text{max}} = 127 \text{ (132) GeV}$  for maximal stop mixing

with  $m_{\text{top}} = 175 \text{ (179.3) GeV}$ ,  $M_S = 1 \text{ TeV}$



# Heavy neutral MSSM Higgs bosons H and A

Production through  $gg \rightarrow H/A$  and  $gg \rightarrow bbH/A$

Associated production  $gg \rightarrow bbH/A$   
dominates at large  $\tan\beta$

$\rightarrow$  b tagging can be used to suppress  
 $Z, \gamma^*, W + \text{jet}$  and QCD multi-jet backgrounds  
leading to real (or fake)  $\tau\tau$  pair production

Branching fractions at large  $\tan\beta$  ( $>10$ ):

$BR(H, A \rightarrow \tau\tau) \sim 10\%$

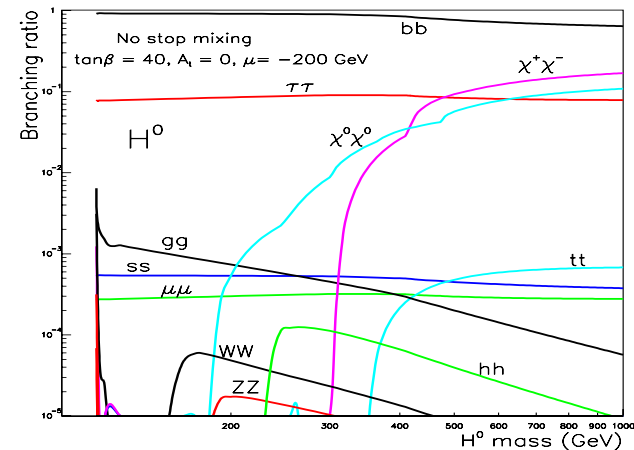
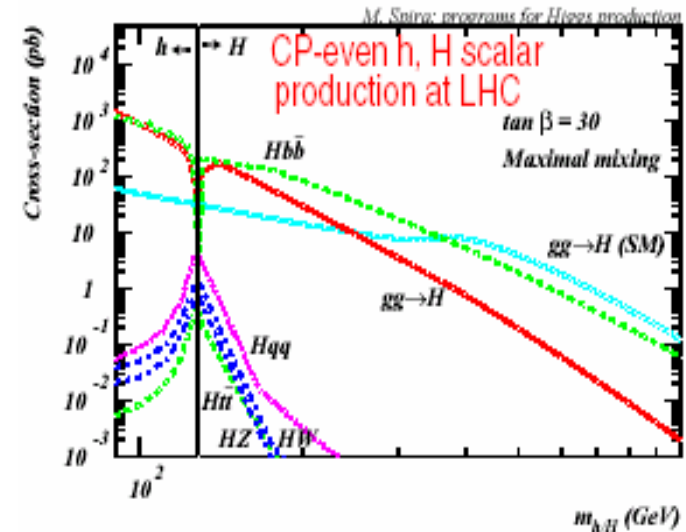
$H, A \rightarrow bb$  dominates

-but background reduction difficult,  
not yet shown to be really useful

$BR(H, A \rightarrow \mu\mu)$  small  $\sim 3 \times 10^{-4}$

- but precise mass measurement possible

At large  $m_A$ , sensitivity to SUSY parameters  
( $\mu$  and  $M_2$ ) due to opening of  $H, A \rightarrow \chi\chi$  decay modes



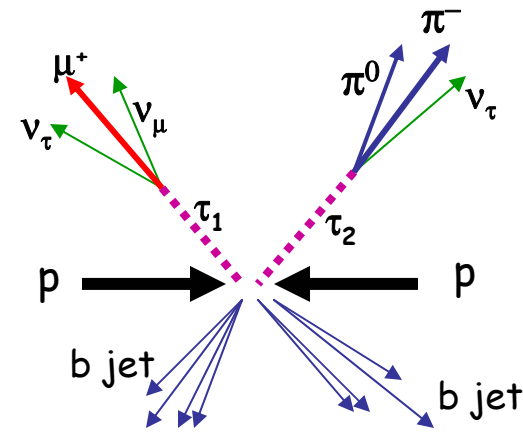
## qq->bbH/A, H/A -> $\tau\tau$ channels

Final states from H,A ->  $\tau\tau$ :

2 leptons,	BR ~ 12.4%
lepton + $\tau$ jet,	BR ~ 45.6%
2 $\tau$ jets,	BR ~ 42%
$\tau$ jet = hadronic $\tau$ decay	

Backgrounds from

$Z, \gamma^* \rightarrow \tau\tau$ ,  $t\bar{t}$ ,  $Wt$ ,  $W$ +jets,  $bb$ , QCD multi-jet events



Challenges with  $\tau\tau$  decay modes:

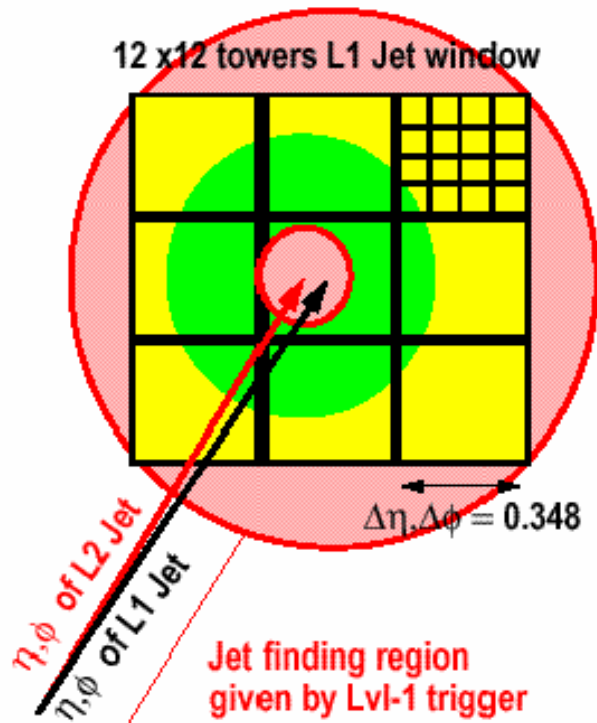
- Suppress the fake  $\tau$ 's from hadronic jets  
(QCD multi-jet,  $W$ +jet backgrounds)
- Trigger on fully hadronic final states, 2  $\tau$  jets (also at low  $m_A$ )
- Reconstruct the Higgs boson mass from  $E_{\tau}^{\text{miss}}$  + leptons, jets

Hadronic jet suppression (at trigger level and offline) is based on narrowness and low multiplicity (1 or 3 prongs) of a  $\tau$  jet

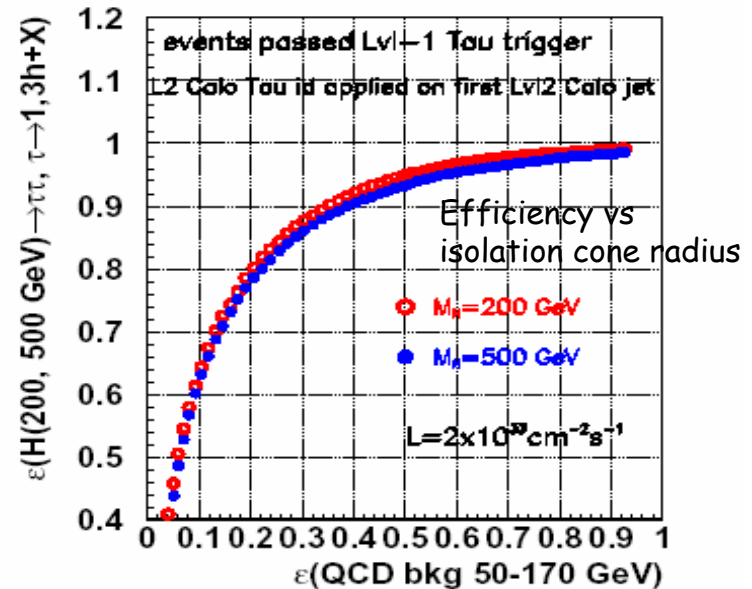
# Hadronic Tau trigger

Requirements: 3(6) kHz output rate at Level-1 at low(high) luminosity  
 Reduction of hadronic QCD events by  $\sim 10^3$  at HLT

Level-1: Narrow hadronic jet in calorimeters



Level-2: Isolation of the jet core ( $\Delta R < 0.13$ ) in the fine-grained EM calorimeter

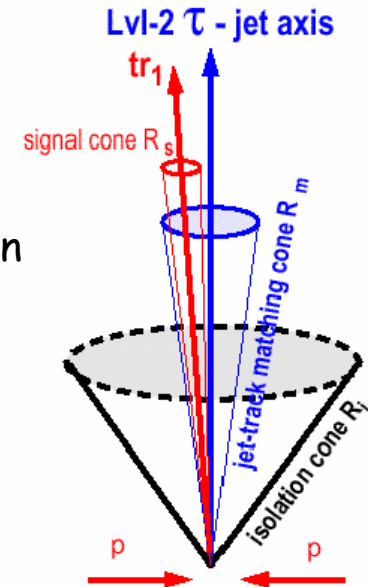


# Level-3 Tau trigger

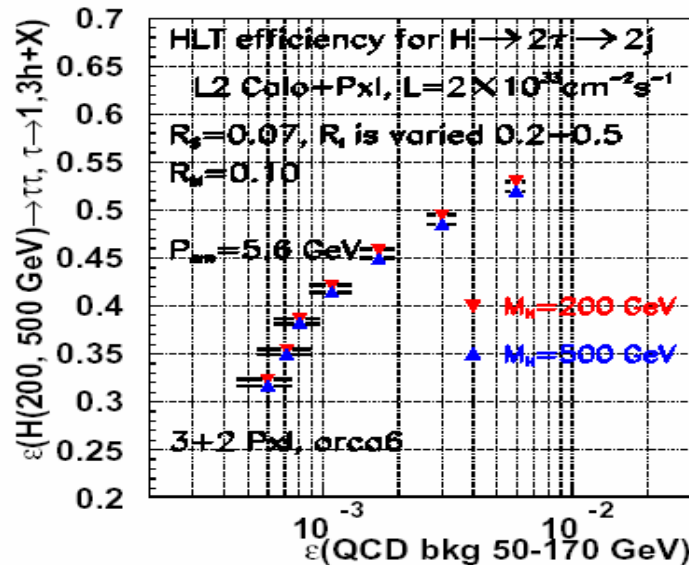
with isolation in the Pixel detector or in the full tracker (Pixel + Silicon)

Method:

- Reconstruction of tracks around the Level-1 jet direction in the Pixel detector or in the Pixel+Silicon tracker
- Small signal cone ( $\Delta R_s = 0.07$ ) around the hardest track
- Larger isolation cone around jet direction



Efficiency (QCD vs  $H \rightarrow \tau\tau \rightarrow 1/3$  prong jets) as a function of the isolation cone size



$H \rightarrow \tau\tau \rightarrow 2\tau$  jets,  $m_H = 200$  GeV:  
Signal efficiency for a QCD background suppression of  $10^3$  at  $L = 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$

HLT path	eff.	cpu[ms]
Calo+Pixel	0.41	59
Calo+Tracker	0.45	130

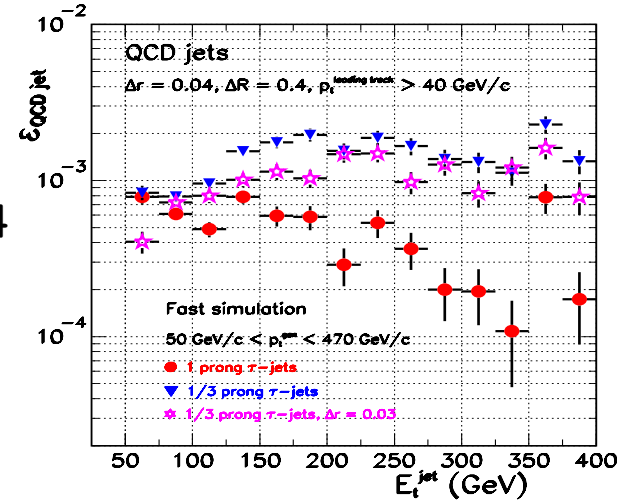
# Hadronic $\tau$ identification

## Algorithm:

- $p_T$  (leading track)  $> 40$  GeV in a jet with  $E_T > 60$  GeV
- two other tracks,  $p_T > 1$  GeV, allowed in a small signal cone of  $\Delta r < 0.04$  around the leading track
- isolation of the signal cone in a larger cone of  $\Delta R = 0.4$

QCD jet suppression  $\sim 1000$   
 $\tau$  jet efficiency  $\sim 30\%$  for  $m_A = 500$  GeV

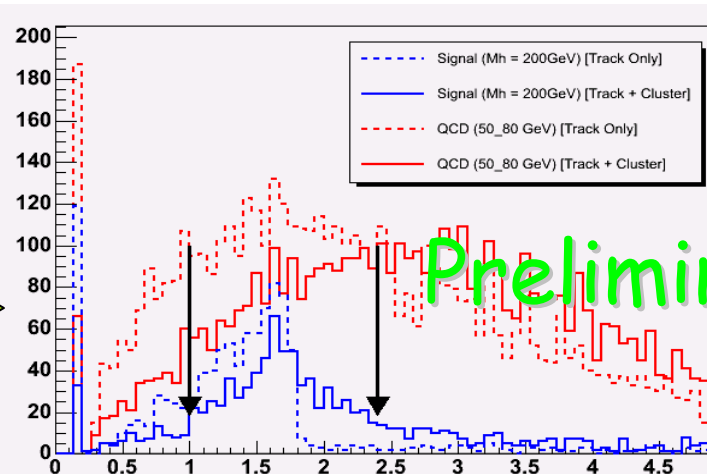
$\tau \rightarrow 3$  prong decays can be used too !



## Further methods exploiting $\tau$ properties

### $\tau$ tagging with mass

Reconstructed  $\tau$  mass for  
 3 prongs and  $h^\pm + n\pi^0$  final states  
 $H \rightarrow \tau\tau \rightarrow 2\tau$  jets,  $m_H = 200$  GeV  
 and for QCD jets





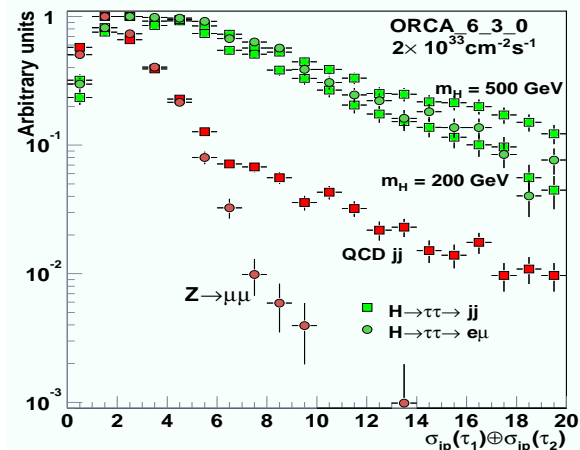
## Methods exploiting $\tau$ lifetime: $c\tau \sim 90 \mu\text{m}$

$\tau$ 's from  $H \rightarrow \tau\tau$  with  $m_H = 200 \text{ GeV}$  travel  $\sim 5 \text{ mm}$  before decaying  $\rightarrow$   
 $Z \rightarrow \ell\ell$  and QCD multijet backgrounds can be suppressed by

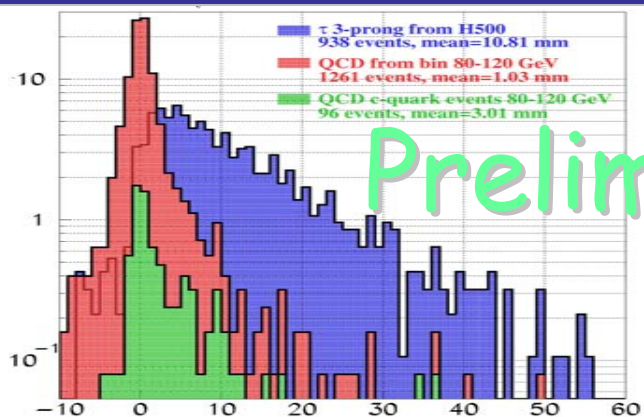
### $\tau$ tagging with impact parameter

combining impact parameter measurements to  
 $((\sigma_{ip}^{\tau_1})^2 + (\sigma_{ip}^{\tau_2})^2)^{1/2}$   
 in two  $\tau$ 's for 1- and 3-prong  $\tau$ 's

$\rightarrow$  QCD di-jet suppression further by  $\sim 9$   
 for  $\sim 60\%$  signal efficiency

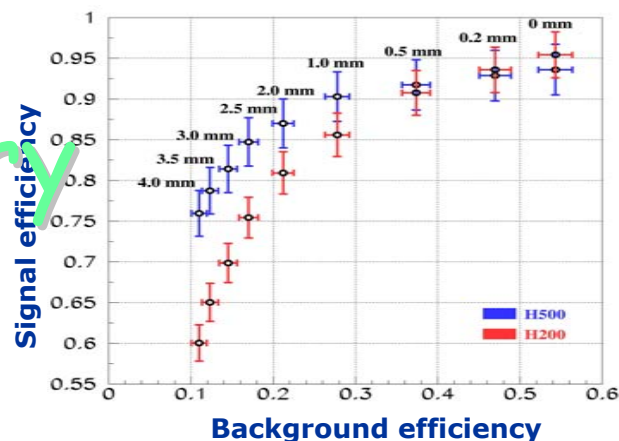


### Secondary vertex reconstruction in 3-prong $\tau$ decays



3D reconstructed flightpath in mm

Preliminary

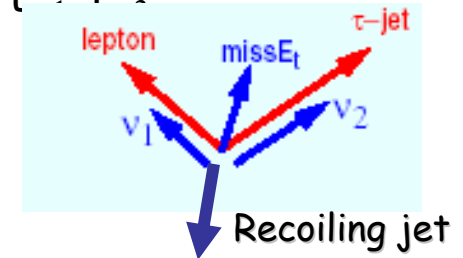


## Higgs boson mass reconstruction in $H, A \rightarrow \tau\tau$

Assume  $\nu$ 's emitted in the two  $\tau$  directions given by the visible decay products (leptons, hadrons), project  $E_{\tau}^{\text{miss}}$  on the two  $\tau$  directions to give  $E_{\tau 1}$ ,  $E_{\tau 2}$

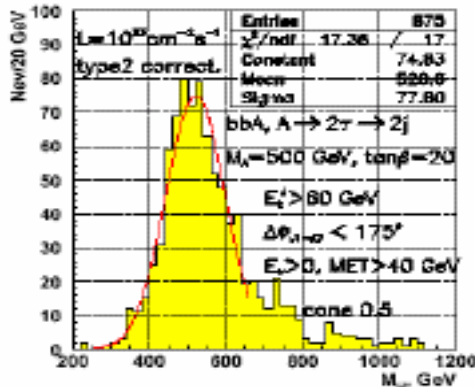
$$m_H = (2E_{\tau 1}E_{\tau 2}(1-\cos\theta_{\tau\tau}))^{1/2}$$

$$E_{\tau 1} = E_{\text{jet}} + E_{\nu} \quad , \quad E_{\tau 2} = E_{\text{lepton}} + E_{\nu}$$

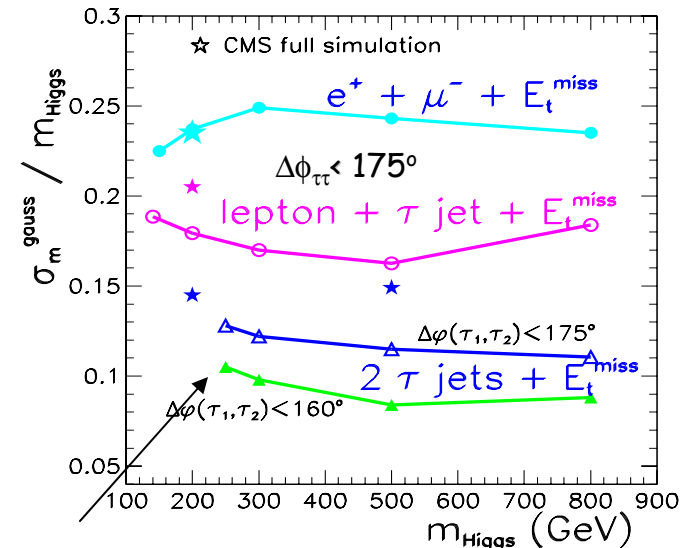


- Resolution and reconstruction efficiency sensitive to  $E_{\tau}^{\text{miss}}$  measurement
- Resolution depends on  $\Delta\theta_{\tau\tau}$  as  $1/\sin\theta_{\tau\tau}$

Mass reconstruction in  $H \rightarrow \tau\tau \rightarrow 2 \tau$  jets:



Resolution for  $m_H = 500 \text{ GeV}$ :  
 $\sigma / \langle m_H \rangle \sim 15\%$   
 with  $\Delta\phi_{\tau\tau} < 175^\circ$



Best resolution obtained with fully hadronic  $\tau$  final states



# Discovery potential for $H/A \rightarrow \tau\tau$

Background suppression with

lepton isolation in  $e\mu$  and  $\ell\ell$  final states:  $bb$  background

hadronic  $\tau$  identification, lepton+jet and 2-jet final states: QCD multi-jets,  $W$ +jets

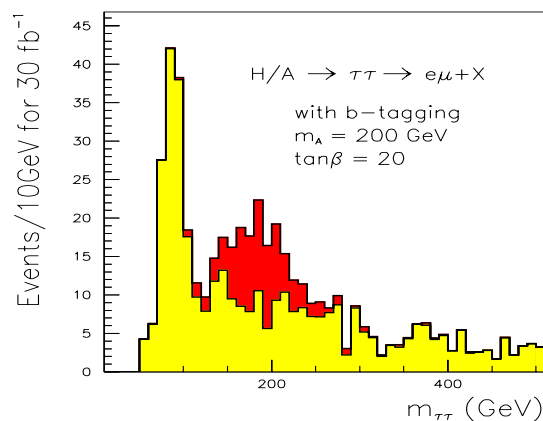
b-jet tagging:  $Z \rightarrow \tau\tau$ , QCD multi-jets,  $W$ +jets

$\tau$  tagging with impact parameter:  $Z \rightarrow \ell\ell$ , QCD multi-jets

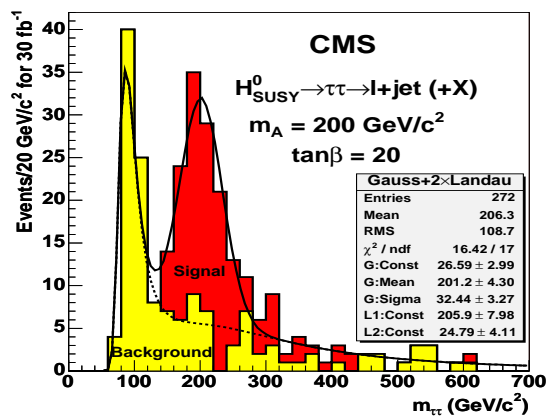
central jet veto:  $t\bar{t}$  and  $Wt$  backgrounds

Signal superimposed on the total background in

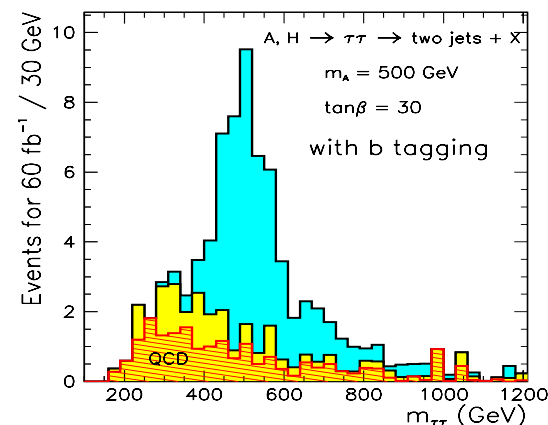
$e\text{-}\mu$  final state  $30\text{fb}^{-1}$



lepton+jet final state  $30 \text{ fb}^{-1}$

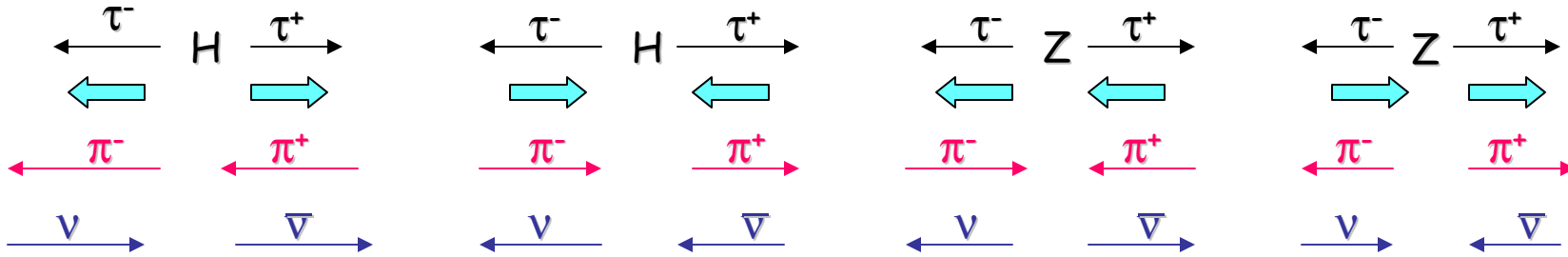


two-jet final state  $60 \text{ fb}^{-1}$



# Helicity correlations in $H, A \rightarrow \tau\tau \rightarrow 2 \tau$ jets

a possibility to suppress further  $Z, \gamma^*$  background?



Two possible spin configurations:

	$H(A) \rightarrow \tau\tau$	$Z \rightarrow \tau\tau$	
$\tau$ helicity	LL	LR	
	RR	RL	
$\tau$ polarization	+-	++	$\tau(+\text{pol}) \rightarrow$ hard $\pi^\pm, \rho_L, a_L$ , soft $\rho_R, a_R$
	-+	--	$\tau(-\text{pol}) \rightarrow$ soft $\pi^\pm, \rho_L, a_L$ , hard $\rho_R, a_R$

Expectation:

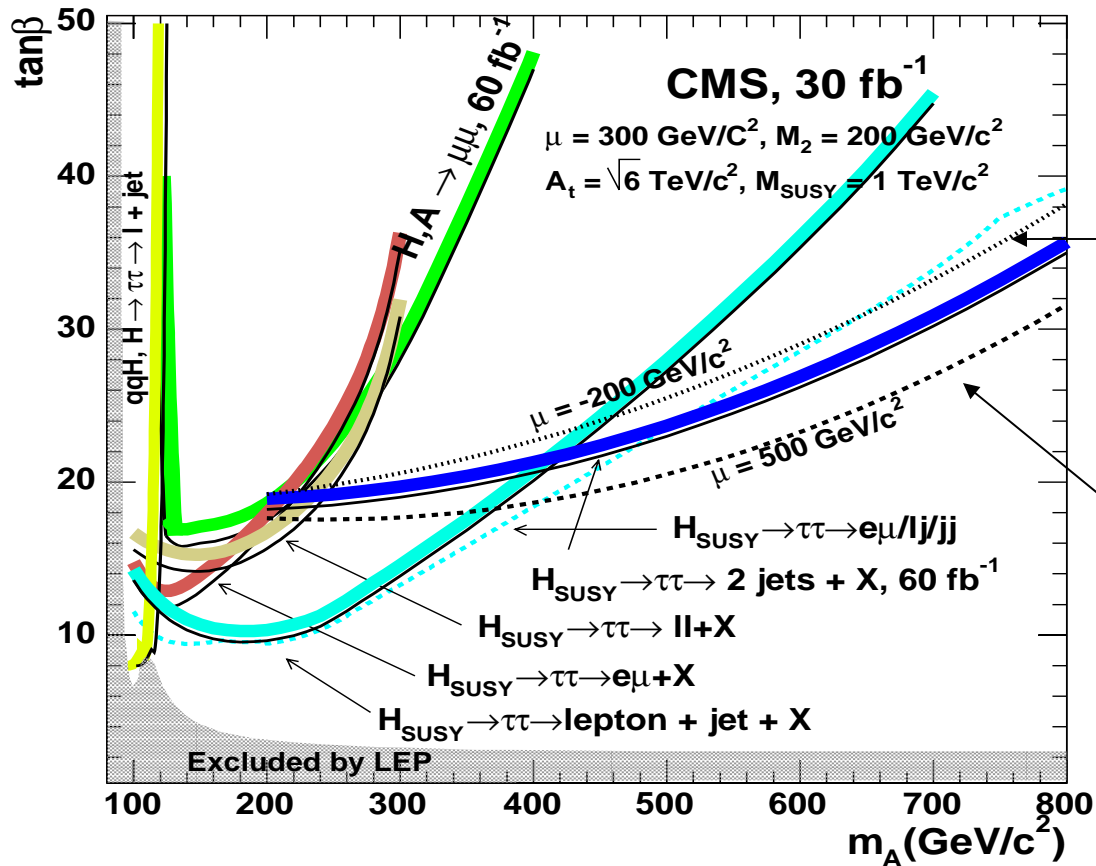
$H(A) \rightarrow \tau\tau$ :  $\Delta E(\tau_1)$  large and  $\Delta E(\tau_2)$  small

$Z \rightarrow \tau\tau$ :  $\Delta E(\tau_1)$  and  $\Delta E(\tau_2)$  small or  $\Delta E(\tau_1)$  and  $\Delta E(\tau_2)$  large

where  $\Delta E(\tau) = E_{\pi^\pm} - \Sigma E_{\pi^0}$



## 5 $\sigma$ discovery potential for the heavy neutral MSSM Higgs bosons



Reach sensitive to SUSY parameters ( $\mu$  and  $M_2$ ) at large  $m_A$  :

Reduction of the  $H \rightarrow \tau\tau$  branching fraction for smaller  $\mu$  and  $M_2$   
 $\rightarrow$  enhancement of  $H, A \rightarrow \chi\chi$  due to lighter gauginos

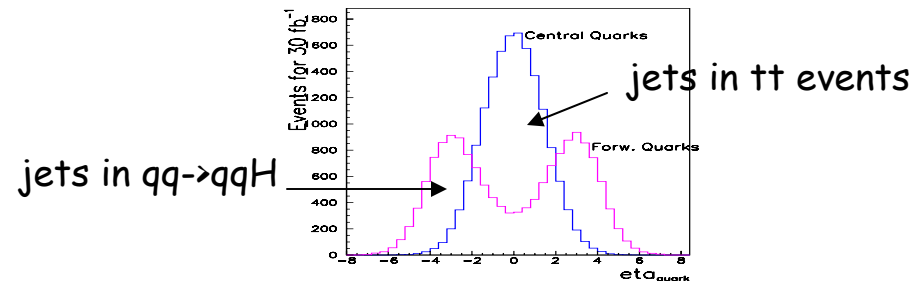
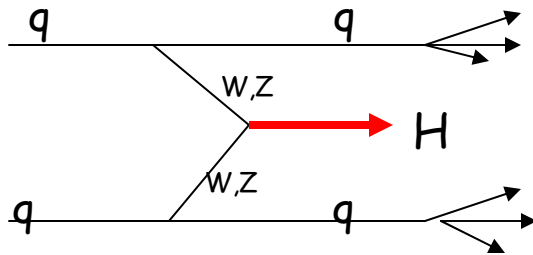
Enhancement of the  $H \rightarrow \tau\tau$  branching fraction for larger  $\mu$  and  $M_2$   
 $\rightarrow$  reduction of  $H, A \rightarrow \chi\chi$  due to heavier gauginos



# H production in weak gauge boson fusion

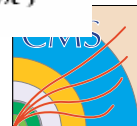
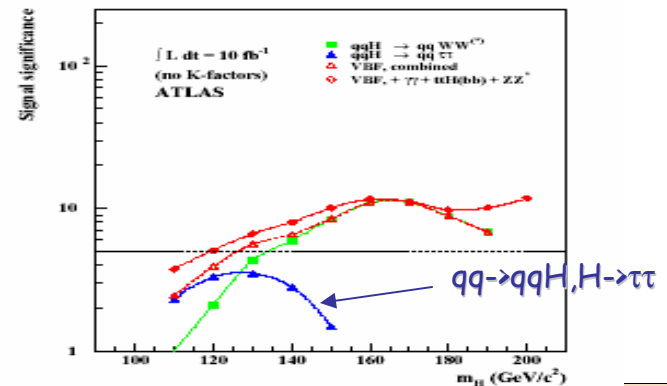
Production characteristics:

- the final state q-jets from  $qq \rightarrow qqH$  are energetic and distributed in the forward regions
- no jets expected in the central region



Forward jet tagging and central jet veto can be used to suppress the QCD multi-jet, W+jet, Z+jet,  $\gamma$ +jet and  $tt$  backgrounds

$qq \rightarrow qqH$  important for the SM Higgs boson searches with  $H \rightarrow WW, WW^*, \gamma\gamma, \tau\tau$  decay modes, in particular in the region of small ( $<120 \text{ GeV}$ ) and large ( $>500 \text{ GeV}$ )  $m_H$

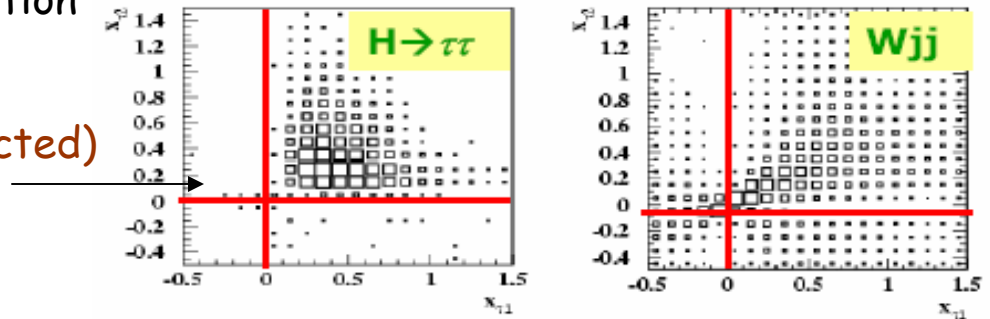


# H → ττ → 2 leptons, lepton + τ jet in weak boson fusion

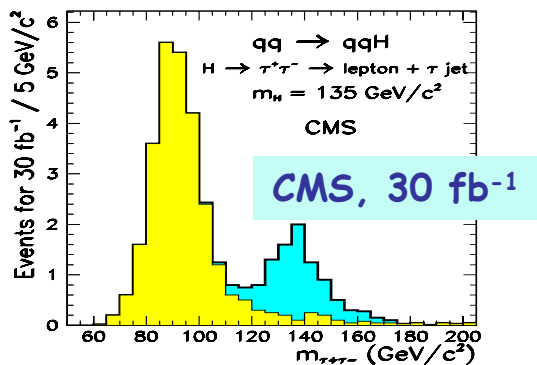
Backgrounds from Z+jets, W+jets, tt, WW suppressed with **forward jet tagging and central jet veto**  
 The electro-weak qq→qqZ/W background is irreducible but initially smaller than the QCD induced ones

Higgs boson mass can be reconstructed from visible τ's (jets and leptons) and E<sub>τ<sup>miss</sup></sub> with collinear neutrino approximation

Further background suppression with  
 $x_{\tau} = p(\tau \text{ decay products})/p(\tau \text{ reconstructed})$   
 $x_{\tau 1} x_{\tau 1} > 0$

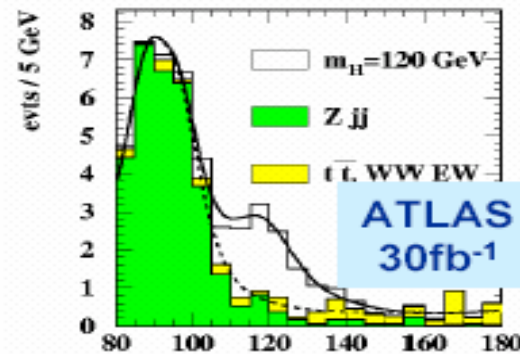


lepton+τ jet final state



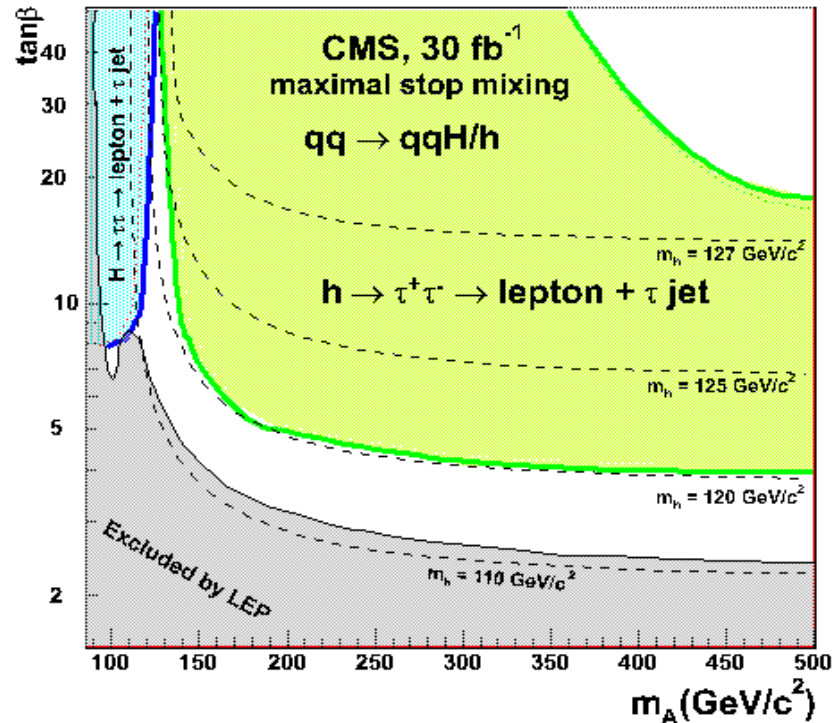
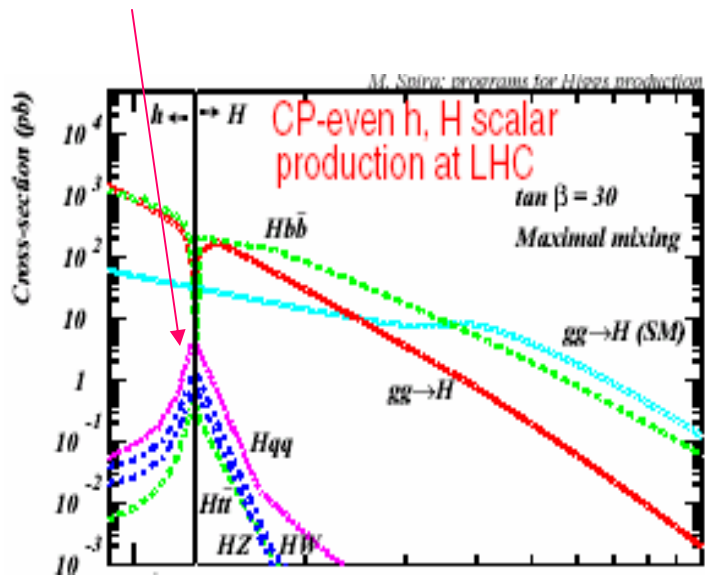
Integrated luminosity  
 ~40 fb<sup>-1</sup> needed for  
 a 5σ-significance  
 around m<sub>H</sub> = 125 GeV  
 in SM

leptonic final state



# $qq \rightarrow qqH/h, H/h \rightarrow \tau\tau$ in MSSM

$qq \rightarrow qqH/h$  significant near the **lower (upper)** mass bound of  $H/h$



In this region,  $H/h$  is SM-like  
 $\rightarrow$  discovery region calculated from the SM sensitivity

Large coverage in MSSM with  $qq \rightarrow qqh, h \rightarrow \tau\tau$  already with 30  $\text{fb}^{-1}$  !

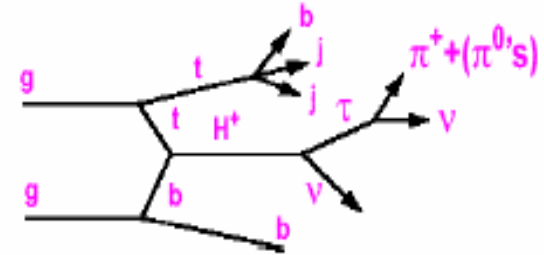




# Charged MSSM Higgs bosons

Production:

- in  $tt$  events with  $t \rightarrow bH^\pm$  if  $m_{H^\pm} < m_{top}$
- through  $gg \rightarrow tbH^\pm$  if  $m_{H^\pm} > m_{top}$



For  $m_{H^\pm} > m_{top}$  no need to detect the associated  $b$  (at large rapidities)  $\rightarrow$   
 $gb \rightarrow tH^\pm$  can be used

Decay channels

$m_{H^\pm} > m_{top}$ :  $BR(H^\pm \rightarrow \tau\nu) \sim 100\%$

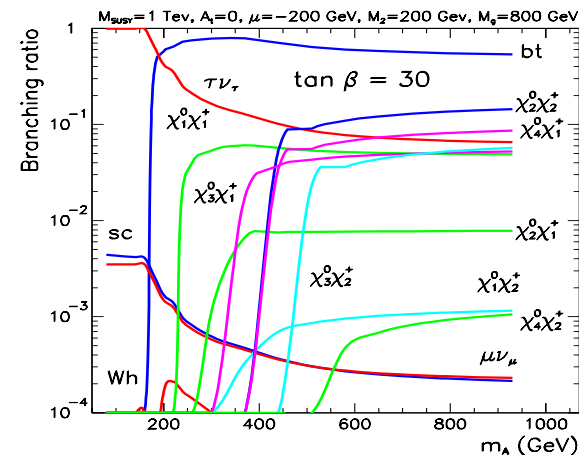
$m_{H^\pm} > m_{top}$  and large  $\tan\beta$  ( $>10$ ):

$H^\pm \rightarrow tb$  dominates

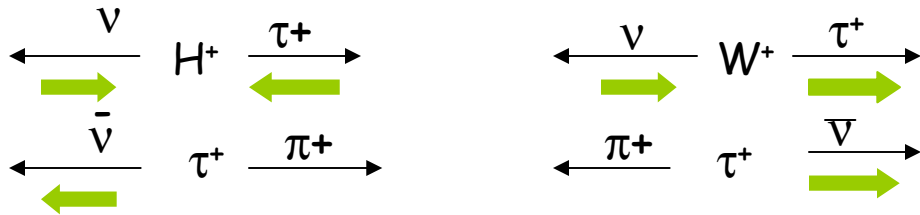
$BR(H^\pm \rightarrow \tau\nu)$  sizeable  $\sim 10\%$

Advantage with  $H^\pm \rightarrow \tau\nu$ ,  $\tau \rightarrow \text{hadrons} + \nu$ :

Helicity correlations can be exploited  
to suppress irreducible backgrounds  
from  $tt$ ,  $Wt$  and  $W + \text{jets}$  with  $W \rightarrow \tau\nu$



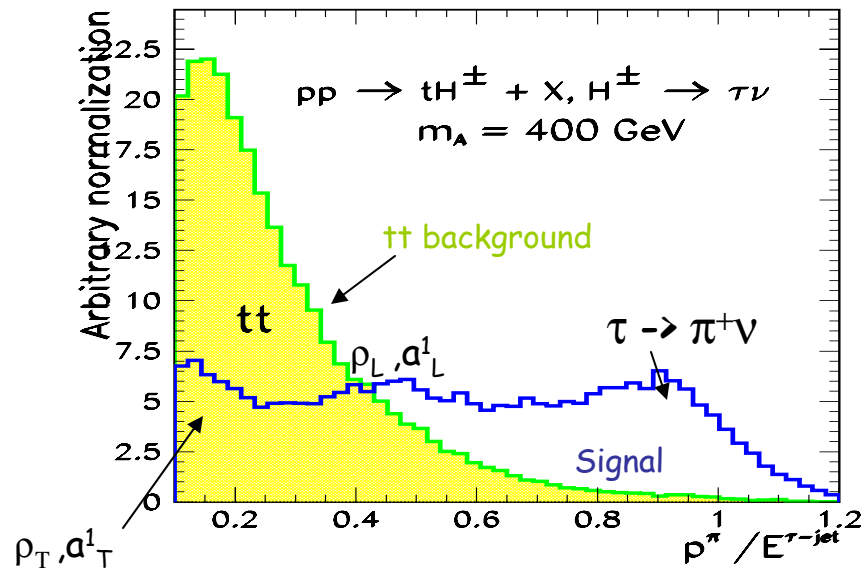
## Helicity correlations in $H^+ \rightarrow \tau\nu$ and $W^+ \rightarrow \tau\nu$



Harder pions from  $H^+ \rightarrow \tau^+\nu$  than from  $W^+ \rightarrow \tau^+\nu$   
 (through  $\tau \rightarrow \pi^+\nu$  and the longitudinal components of  $\rho$  and  $a_1$ )

Suppression of backgrounds  
 with genuine  $\tau$ 's from  $W \rightarrow \tau\nu$   
 with a cut in  $p^\pi/E^\tau_{jet}$

Efficiency with  $p^\pi/E^\tau_{jet} > 0.8$ :  
 Signal,  $m_{H^\pm} = 400$        $\sim 45\%$   
 $t\bar{t}$  background               $\sim 2\%$   
 (fast simulation)



$H^\pm \rightarrow \tau\nu, \tau \rightarrow \text{hadrons}+\nu$  for  $m_{H^\pm} < m_{\text{top}}$   
in tt events with  $t_1 \rightarrow bH^\pm, t_2 \rightarrow \text{lepton} + qq$

Background from tt, Wt, W+jets

Background suppression with  $p_\pi/E^{\tau \text{ jet}}$  cut, lepton isolation, b-tagging,  
top mass reconstruction

$H^\pm$  mass reconstruction not possible, signal as an excess of  $\tau$ 's in tt events

$H^\pm \rightarrow \tau\nu, \tau \rightarrow \text{hadrons}+\nu$  for  $m_{H^\pm} > m_{\text{top}}$   
in fully hadronic events from  $gg \rightarrow tbH^\pm$

Background from tt, Wt, W+jets

Background suppression with  $p_\pi/E^{\tau \text{ jet}}$  cut,  $E_{\text{miss}}$  cut, b-tagging,  
hadronic top mass reconstruction, central jet veto

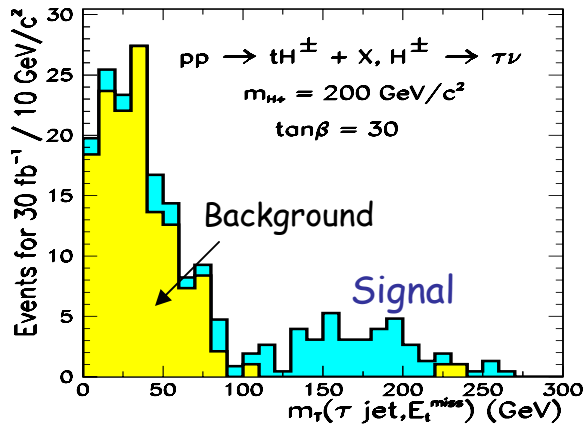
$E_{\text{miss}}$  (mainly) from  $H^\pm \rightarrow \tau\nu$ : transverse mass  $m_T(\tau\text{-jet}, E_{\text{miss}})$   
can be reconstructed with an endpoint at  $m_{H^\pm}$



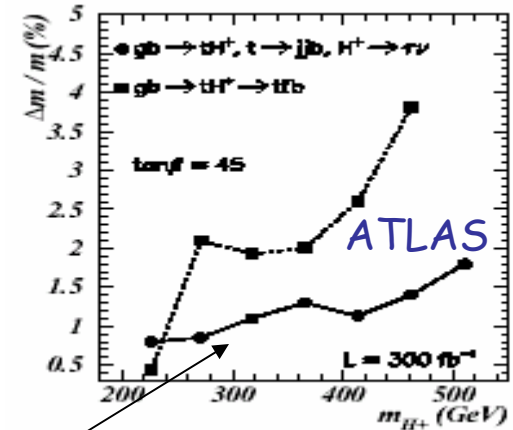
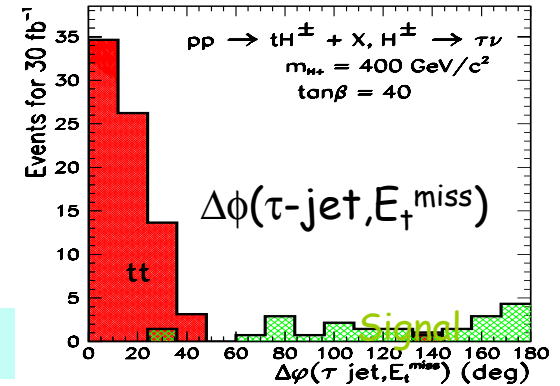
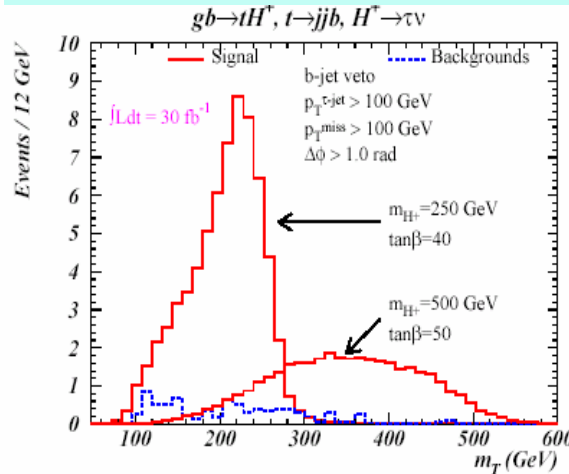
# Transverse mass reconstruction in $tH^\pm, H^\pm \rightarrow \tau\nu$

Quasi two-body decay between the  $\tau$  jet and  $E_+^{\text{miss}}$  in fully hadronic events  $\rightarrow$  almost background-free situation in  $m_T(\tau\text{-jet}, E_+^{\text{miss}})$

CMS,  $\Delta\phi(\tau\text{-jet}, E_+^{\text{miss}}) > 20^\circ$



ATLAS,  $\Delta\phi(\tau\text{-jet}, E_+^{\text{miss}}) > 57^\circ$

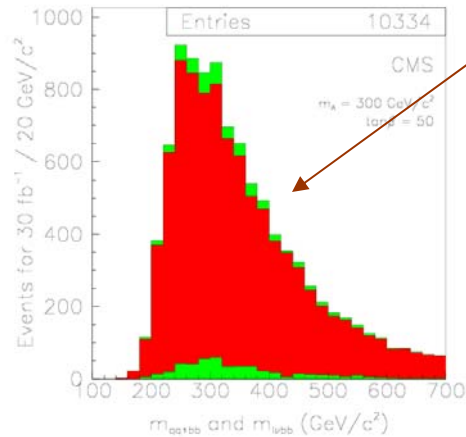
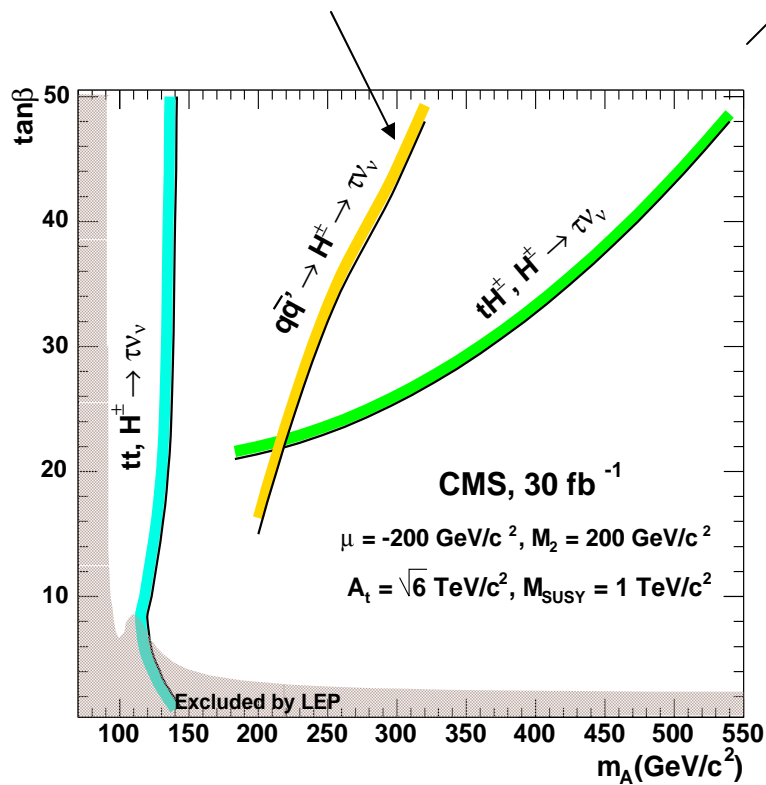
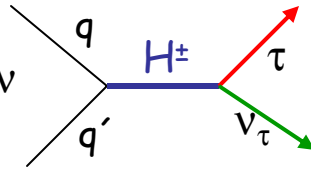


Precision of  $m_{H^\pm}$  measurement from  $m_T(\tau\text{-jet}, E_+^{\text{miss}})$  with likelihood fits:  $\Delta m_{H^\pm} / m_{H^\pm} \sim 1 - 1.5\%$



# Discovery potential for charged Higgs bosons

Production in  $q\bar{q}' \rightarrow H^\pm \rightarrow \tau\nu$   
 large background from  $q\bar{q} \rightarrow W \rightarrow \tau\nu$

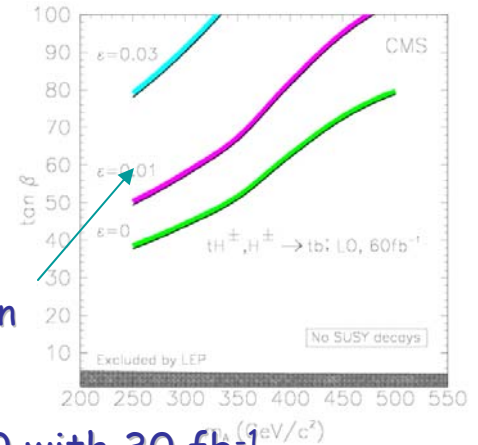


$H^\pm \rightarrow tb$  in  $gg \rightarrow tbH$

- Modest mass resolution
- Difficult background suppression

Background determination  
 uncertainties important

Discovery for  $\tan\beta \sim 50$  with  $30 \text{ fb}^{-1}$



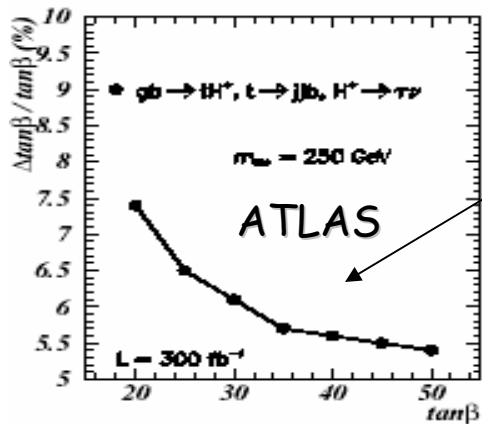
# Measurement of $\tan\beta$ in $H \rightarrow \tau\tau$ from event rates

At large  $\tan\beta$   $\sigma_H \sim \tan^2\beta_{\text{eff}} \times f(m_A)$ ,  $\text{BR}(\tau\tau) \sim \text{constant}$

$$\Delta\tan\beta/\tan\beta = \frac{1}{2} * ((N_S+N_B)^{1/2} / N_S + \Delta\tan\beta_{\text{syst}})$$

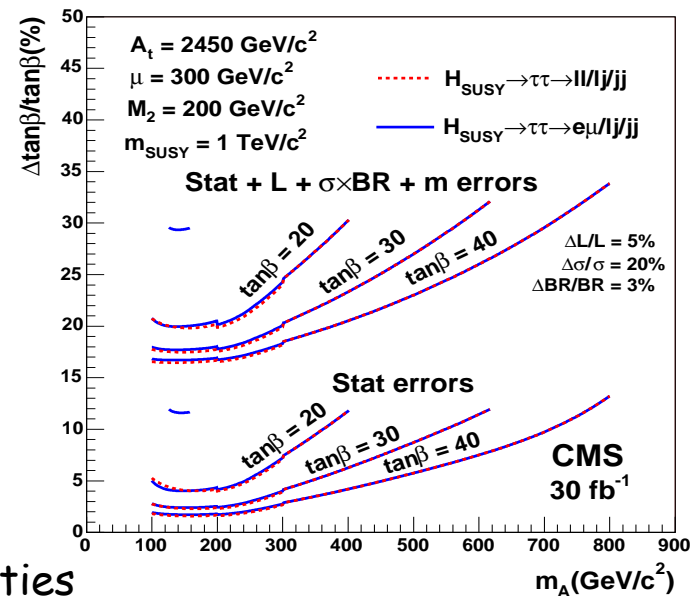
$\Delta\tan\beta_{\text{syst}}$  consists of:

- Theoretical (scale) uncertainty about 20% for  $\sigma(\text{gg} \rightarrow \text{bb}H)$  with NLO calculations
- Luminosity uncertainty  $\Delta L/L \sim 5\%$
- Uncertainty of mass measurement, preliminary estimate 5%
- Uncertainty on event selections  $\sim 3\%$



Theoretical uncertainties not included

Event rates (branching fraction) sensitive to SUSY parameters, uncertainty due to parameter measurement not yet included



## Conclusions

The main discovery potential for the **heavy MSSM Higgs bosons H and A** is through the  **$H/A \rightarrow \tau\tau$**  decay modes, with  $\ell\ell$ ,  $\ell$ +jet and 2-jet final states, in the  $gg \rightarrow bbH/A$  production at large  $\tan\beta$

The most probable discovery mode for the **charged Higgs bosons** is the  **$H^\pm \rightarrow \tau\nu$**  decay in the  $t\bar{t}$  events for  $m_{H^\pm} < m_{\text{top}}$  and in the  $gg \rightarrow t\bar{t}H^\pm$  production for  $m_{H^\pm} > m_{\text{top}}$ .

Advantages with  $\tau$ 's:

- Leptonic and hadronic (including 3-prongs) decay modes can be used
- $\tau$  lifetime can be exploited in  $H/A \rightarrow \tau\tau$  with impact parameter or flightpath measurements
- $\tau$  mass tagging may be used
- Helicity correlations in  $H^\pm \rightarrow \tau\nu$  (and in  $H/A \rightarrow \tau\tau$  ?) can be used to suppress backgrounds from  $W \rightarrow \tau\nu$  (and  $Z \rightarrow \tau\tau$  ?) decays



## Conclusions (cont.)

Requirements for efficient use of  $\tau$ 's in the search of MSSM Higgs bosons:

- Triggering hadronic  $\tau$  decays on lowest possible thresholds is advantageous
- Identification of  $\tau$ 's is (largely) based on isolation:  
high reconstruction efficiencies needed for tracks in jets, i.e.  
**tracker performance is essential**
- Good vertex and impact parameter measurement precisions for  $\tau$  tagging
- $E_{\text{T}}^{\text{miss}}$  measurement precision important for mass ( $m_{\tau}$ ) reconstructions

