

Naked Cronin effect in Au+Au collisions

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Hard Probes 2004 - Ericeira, Nov 3-10, 2004

- ★ **Intro:** overview of theory models
- ★ **Glauber-Eikonal model** (pQCD multiple parton scatterings)
 - ➔ calibrated in p+p collisions
 - ➔ d+Au @ $\eta = 0 \Rightarrow$ OK!
- ★ **Naked Cronin effect** = Au+Au without medium effects
 - ➔ onset and magnitude of hadron quenching
 - ➔ from RHIC to SPS
- ★ **Remarks** on baryon anomaly and HERMES data

Based on: A.A., M.Gyulassy, PLB 586 (04) 244 & J.Phys.G 30 (04) s969
A.A., nucl-th/0405046

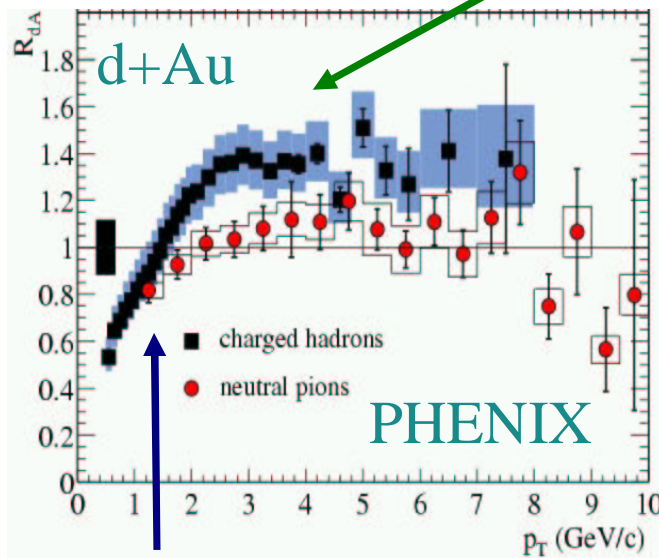
Cronin ratio of inclusive spectra

$$R_{AB}(p_T) = \frac{d^2 N_{AB}^h / dp_T d\eta}{T_{AB} d^2 \sigma_{NN}^h / dp_T d\eta}$$

Binary collision scaling

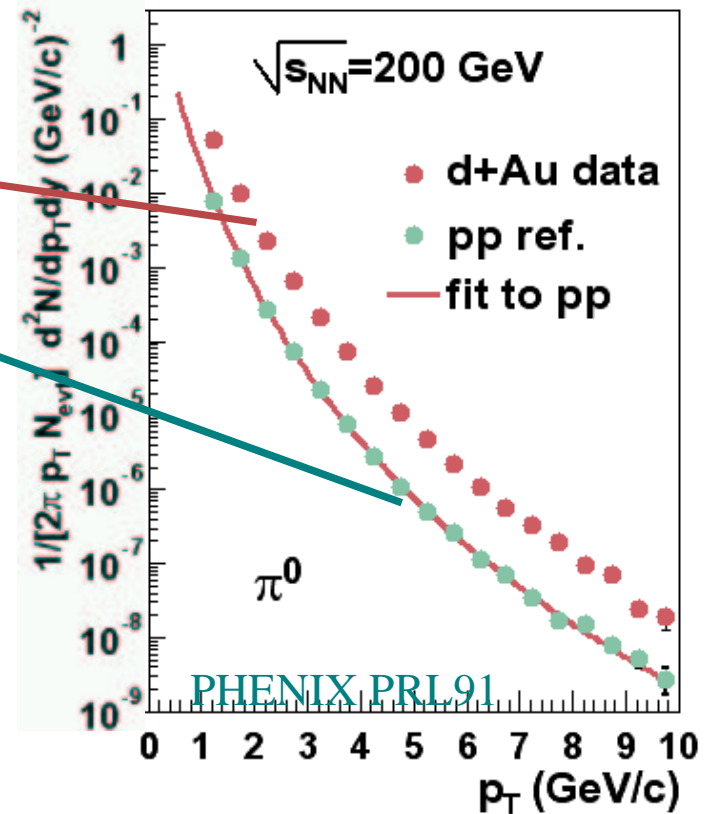
p+p reference

Enhancement at moderate-pT



Suppression at low pT

(same pattern in Fermilab p+A data)



Note: baryon/meson anomaly

Theoretical models for hA

★ Parton rescatterings

→ Glauber-Eikonal models (pQCD)

"Exact" Kuhn '76; Kryzwicki '79; Lev-Petersson '83;
 A.A., Gyulassy; Cattaruzza, Treleani;

Approx. X.N. Wang et al.; Barnafoldi et al.;
 Vitev, Gyulassy; Fai, Zhang;

→ Colour dipole models

Vintage Kopeliovich et al.

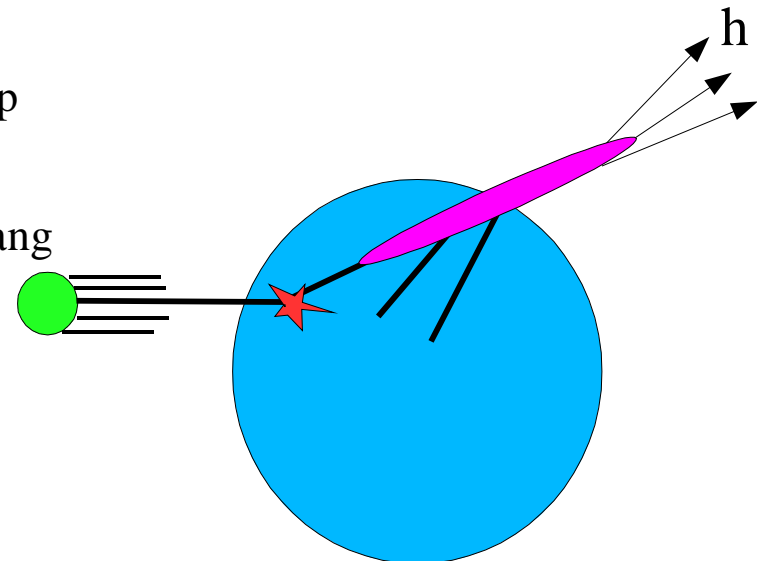
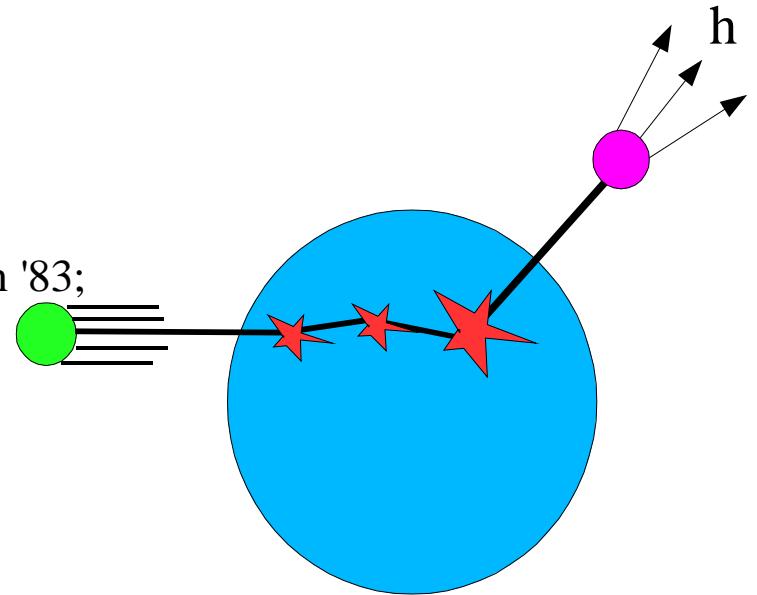
→ Golec Biernat-Wuesthoff σ_{dip}

CGC inspired Jalilian-Marian
 Kharzeev, Kovchegov, Tuchin

→ CGC inspired parametrization of σ_{dip}

★ Final state parton recombination (parton shower and pickup from medium)

Hwa, Yang



First act: p+p collisions

Single scattering
= LO in pQCD

renormalization scale

fragmentation scale

$$\frac{d\sigma^{pp}}{d^2p_t} = \int dx \sum_i f_{i/h}(x, Q_p^2) \underbrace{\int dx' \sum_j f_{j/A}(x', Q_p^2) \frac{d\sigma_{pQCD}^{ij}}{d^2p_T}}_{= d\sigma^{iN} / d^2p_T} \int dz D_{i \rightarrow h}(z, Q_h^2)$$

K-factor (simulates NLO)

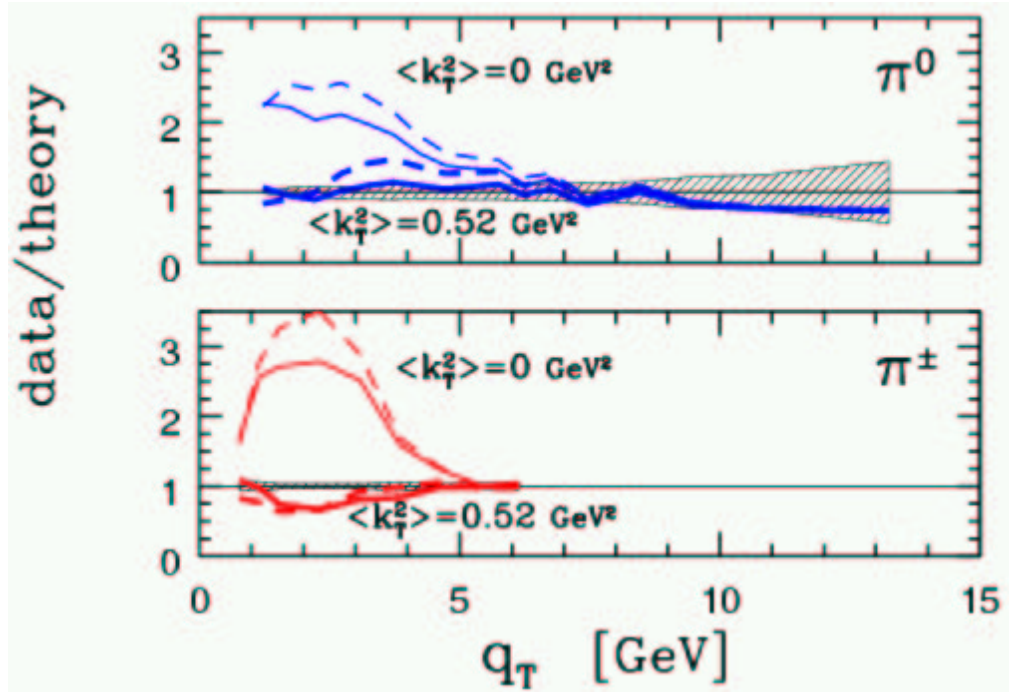
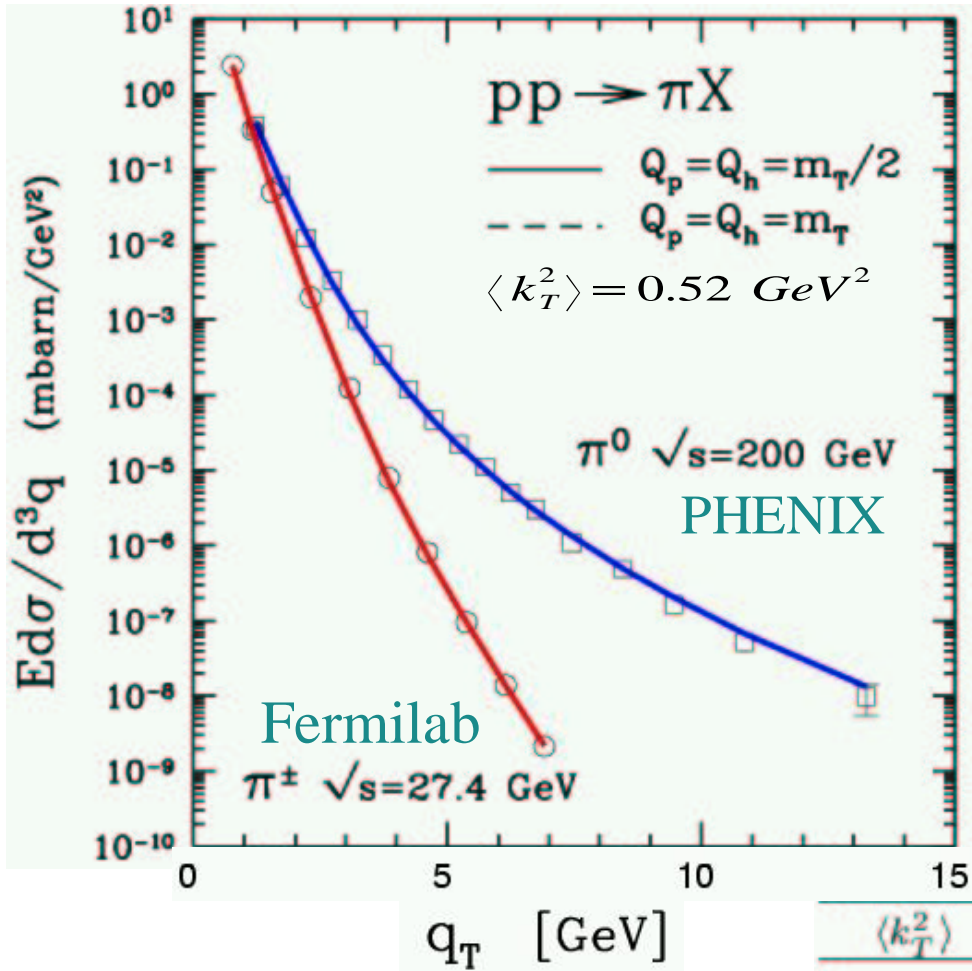
where

$$\frac{d\sigma^{iN}}{d^2p_T} \approx \int_{\frac{4p_T^2}{s}}^1 dx' f_{j/A}(x', Q^2) \frac{C^{ij}}{(p_T^2 + p_0^2)^2} \propto \frac{1}{(p_T^2 + p_0^2)^{2+n}} = \text{parton-nucleon cross section}$$

IR regulator

- CHOOSE the scales: $Q_p = Q_h = \sqrt{p_T^2 + p_0^2} / 2$
(or $Q_p = Q_h = \sqrt{p_T^2 + p_0^2}$)
- FIT $K=K(s)$ to the high- p_T tail of the hadron spectrum
(χ^2 fit, sensitive to the choice of scales)
- FIT intrinsic $\langle k_T^2 \rangle = 0.52 \text{ GeV}^2$ to the moderate- p_T
- FIT $p_0 = p_0(s)$ to the low- p_T hadron spectrum

Results of the fit



Fit procedure: Eskola, Honkanen '02

$\langle k_T^2 \rangle$	$Q_p = Q_h$	$\sqrt{s} = 27.4 \text{ GeV}$	$\sqrt{s} = 200 \text{ GeV}$
0.52 GeV ²	$m_T/2$	$p_0 = 0.70 \pm 0.1 \text{ GeV}$ $K = 1.07 \pm 0.02$	$p_0 = 1.0 \pm 0.1 \text{ GeV}$ $K = 0.99 \pm 0.03$
	m_T	$p_0 = 0.85 \pm 0.1 \text{ GeV}$ $K = 4.01 \pm 0.08$	$p_0 = 1.2 \pm 0.1 \text{ GeV}$ $K = 2.04 \pm 0.12$
0 GeV ²	$m_T/2$	$p_0 = (0.70 \pm 0.1 \text{ GeV})$ $K = 3.96 \pm 0.11$	$p_0 = (1.0 \pm 0.1 \text{ GeV})$ $K = 1.04 \pm 0.06$
	m_T	$p_0 = (0.85 \pm 0.1 \text{ GeV})$ $K = 13.4 \pm 0.4$	$p_0 = (1.2 \pm 0.1 \text{ GeV})$ $K = 2.04 \pm 0.12$

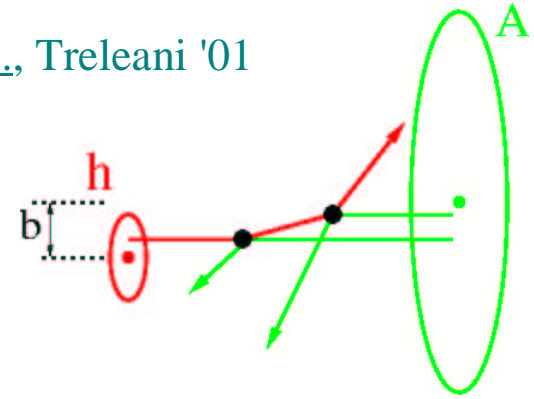
other energies \Rightarrow log fit:

$$p_0 = 0.151 + 0.200 \log \sqrt{s}$$

Second act: h+A collision and Cronin effect

Multiple parton scattering Calucci, Treleani '90-'91 & A.A., Treleani '01

Assuming: generalized collinear factorization
 factorization of the n -body cross-section
 only elastic parton scatterings



$$\frac{d\sigma^{iA}}{d^2p_t} = \sum_{n=1}^{\infty} \frac{1}{n!} \int d^2b \int d^2k_1 \cdots d^2k_n$$

$$\times \underbrace{\frac{d\sigma^{iN}}{d^2k_1} T_A(b) \times \cdots \times \frac{d\sigma^{iN}}{d^2k_n} T_A(b)}_{\text{n-fold parton rescattering}} e^{-\sigma^{iN}(p_0)T_A(b)} \times \delta^{(2)}\left(\sum \mathbf{k}_i - \mathbf{p}_t\right)$$

unitarity factor (probability conserv.)

and:
$$\frac{d\sigma_{pA}^h}{d^2p_t} = \sum_i f_{i/p} \otimes \frac{d\sigma^{iA}}{d^2p_t} \otimes D_{i \rightarrow h} + A \sum_j f_{j/A} \otimes \frac{d\sigma^{jP}}{d^2p_t} \otimes D_{j \rightarrow h}$$

● pA = unitarized multiple parton scatterings on free nucleons

● **Spectra in pp coll. as limiting case (high- p_T or $A \rightarrow 1$)**

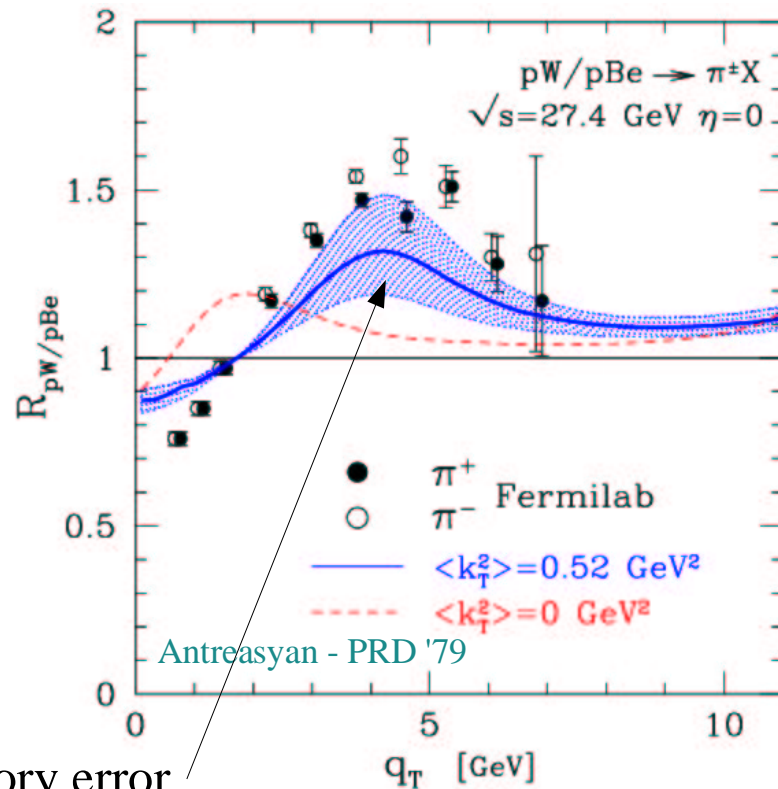
● **No extra free parameters**

$$\frac{d\sigma_{pA}^h}{d^2p_T^h} \xrightarrow{p_T \rightarrow \infty} A \frac{d\sigma_{pp}^h}{d^2p_T^h}$$

h+A : midrapidity pions

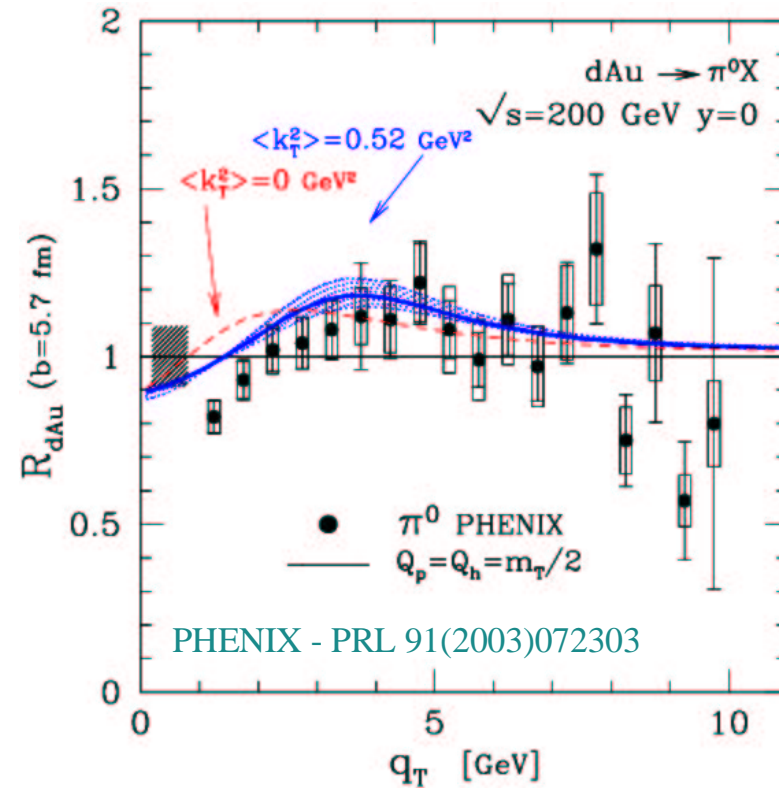
Fermilab $\sqrt{s} = 27.4 \text{ GeV}$

$p_0 = 0.7 \text{ GeV} \pm 10\%$



Phenix $\sqrt{s} = 200 \text{ GeV}$

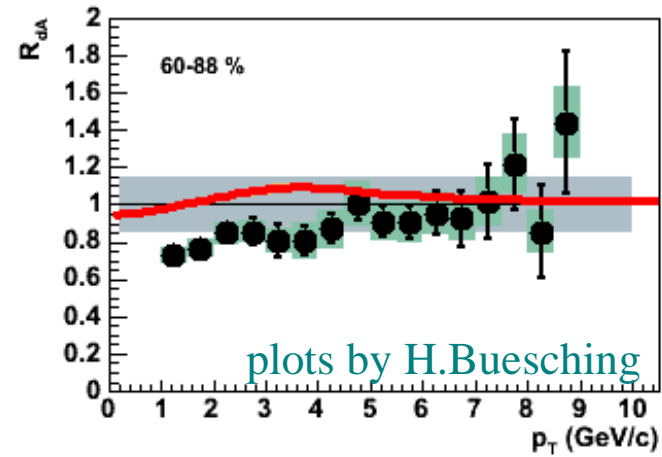
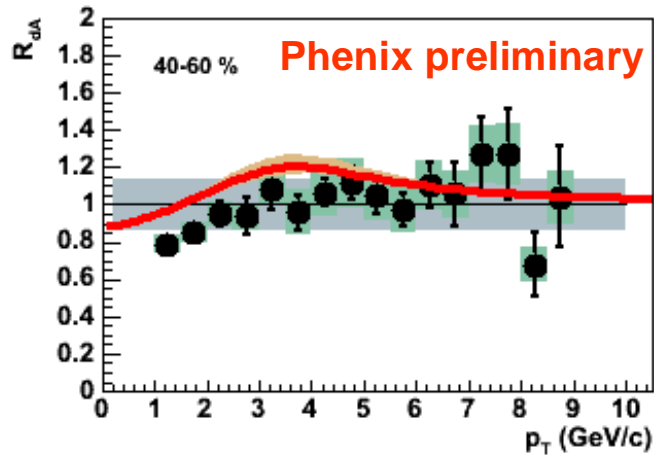
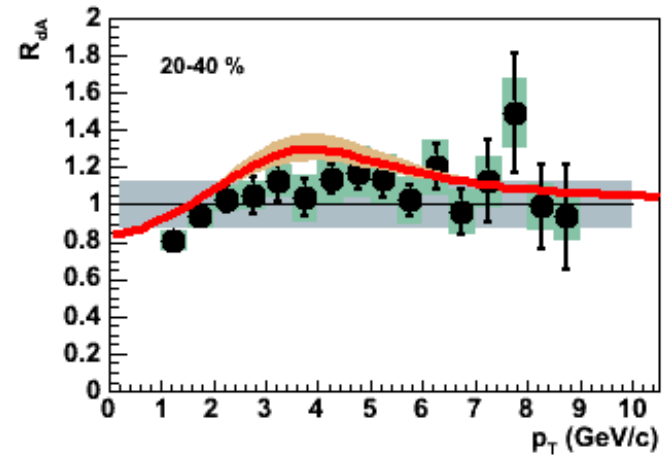
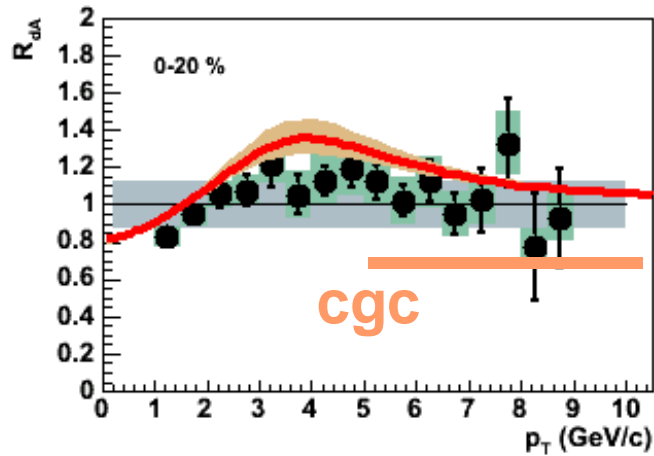
$p_0 = 1.0 \text{ GeV} \pm 10\%$



theory error
due to fit of p_0

Glauber-Eikonal model is OK at $\eta=0$

Check: centrality dependence at 200 GeV



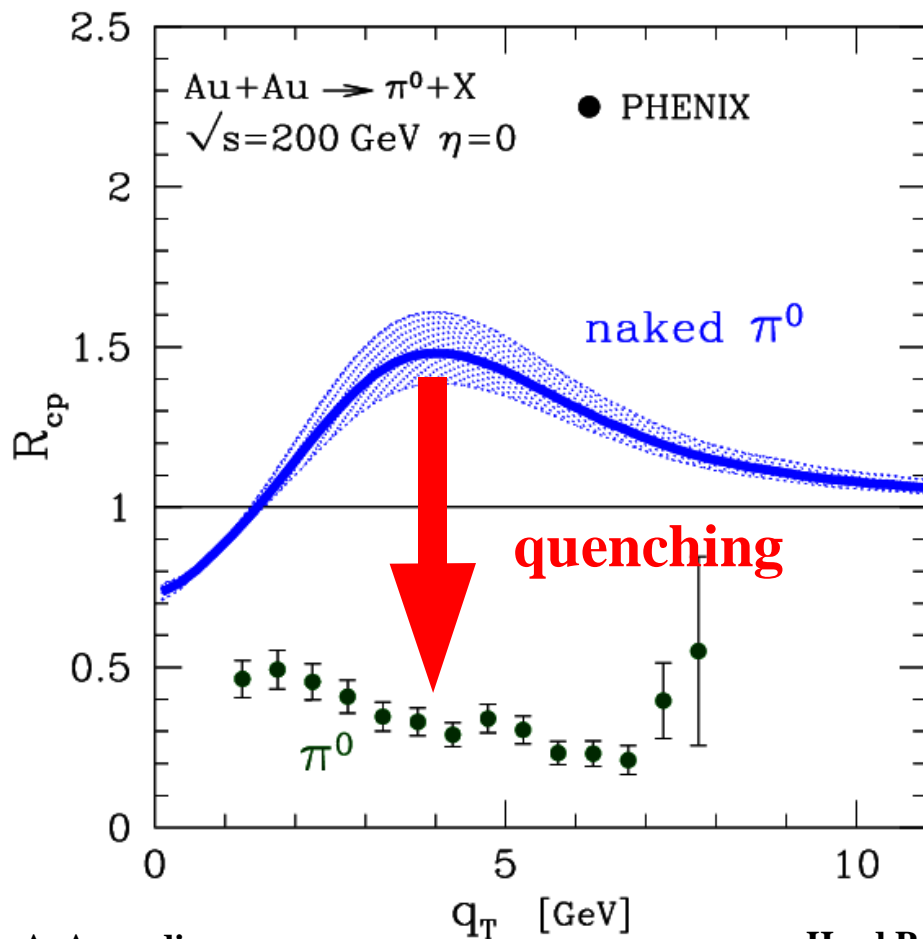
Glauber-Eikonal model is OK at $\eta=0$

Third act: A+A collision

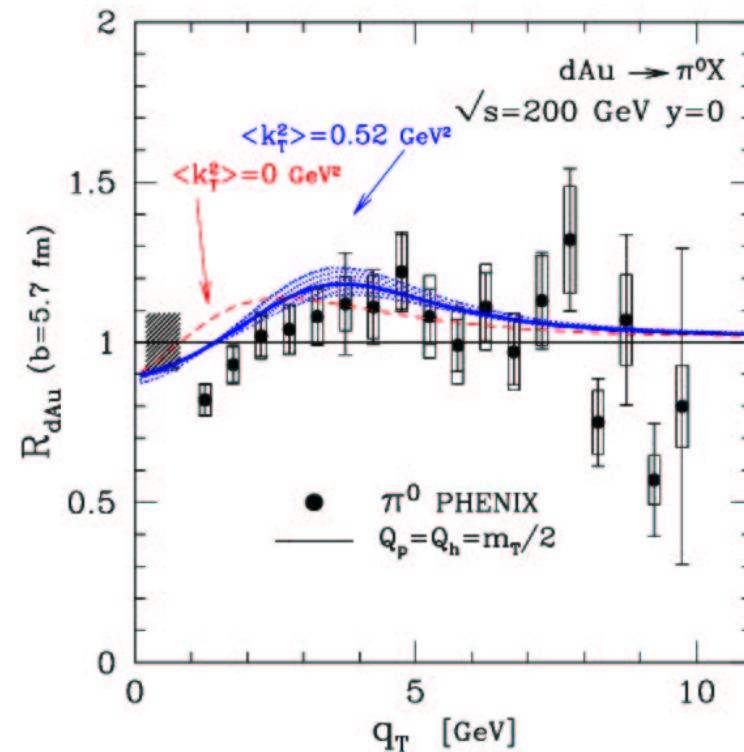
minimal extension
of GE model:

$$\frac{d\sigma_{AB}^h}{d^2bd^2p_t} = \sum_i f_{i/A} \otimes \frac{d\sigma^{iB}}{d^2bd^2p_t} \otimes D_{i \rightarrow h} + \sum_j f_{j/B} \otimes \frac{d\sigma^{jA}}{d^2bd^2p_t} \otimes D_{j \rightarrow h}$$

1) Phenix π^0 - $\sqrt{s}=200$ GeV ($p_0 = 1.0 \text{ GeV} \pm 10\%$)

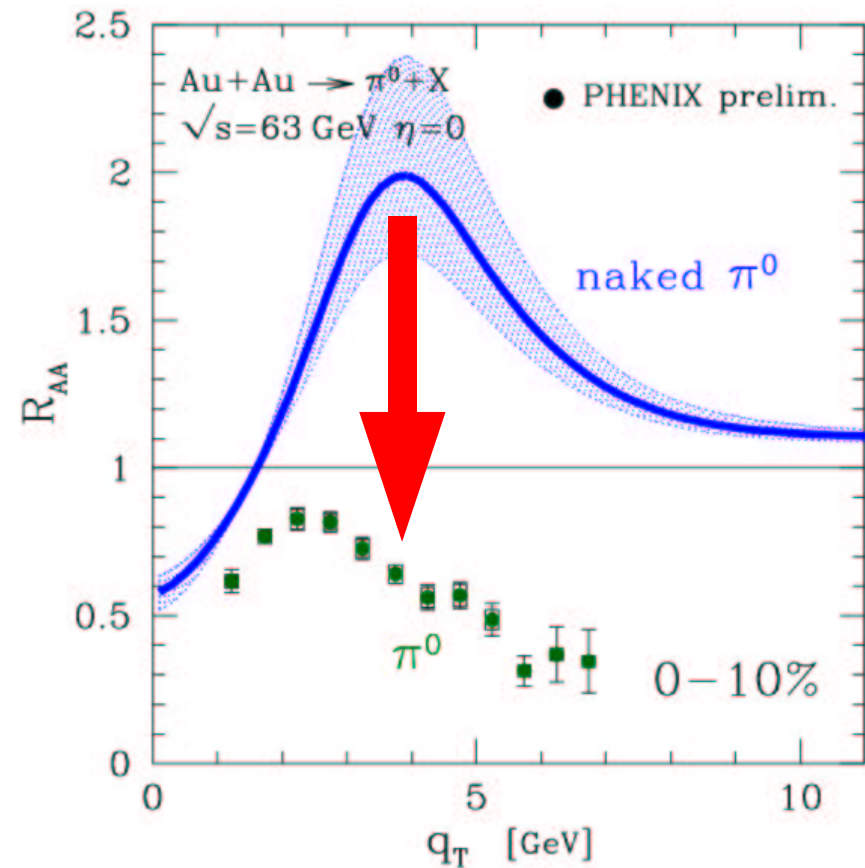
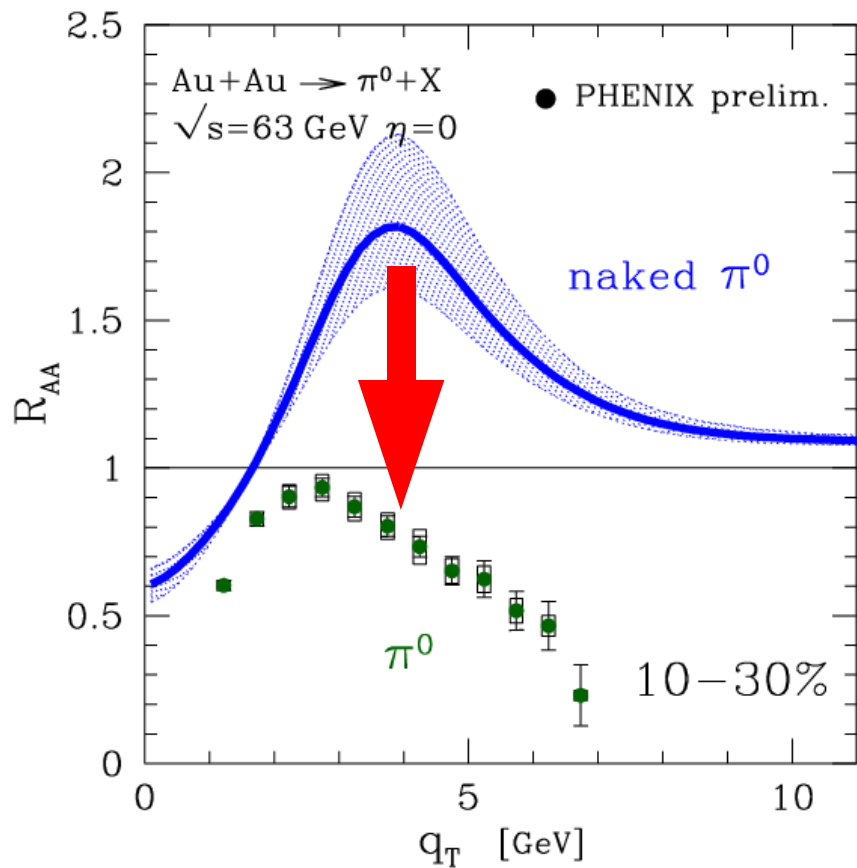


Note: existence of quenching from
d+Au data alone - for magnitude
we need theory computation



2) Phenix π^0 - $\sqrt{s}=62.4$ GeV ($p_0 = 0.82$ GeV $\pm 10\%$)

At 62 GeV we do not have data on d+Au \Rightarrow **theory is needed**

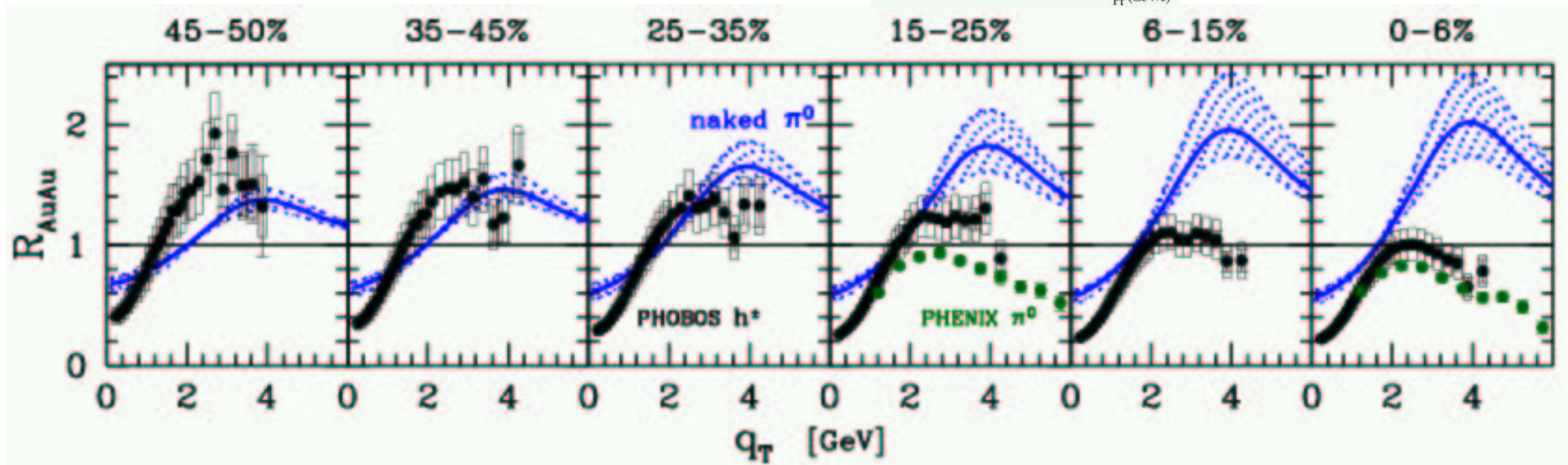
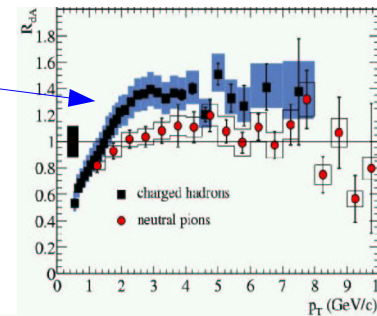


- ★ As expected, less suppression in peripheral collisions.
- ★ Where does quenching begin? **theory vs. data is needed**
 \Rightarrow let's look at **PHOBOS h^\pm data**

3) Phobos h^\pm - $\sqrt{s}=62.4$ GeV

Attention: baryon-meson anomaly
computations are for π^0 's

($p_0 = 0.82$ GeV \pm 10%)



no quenching

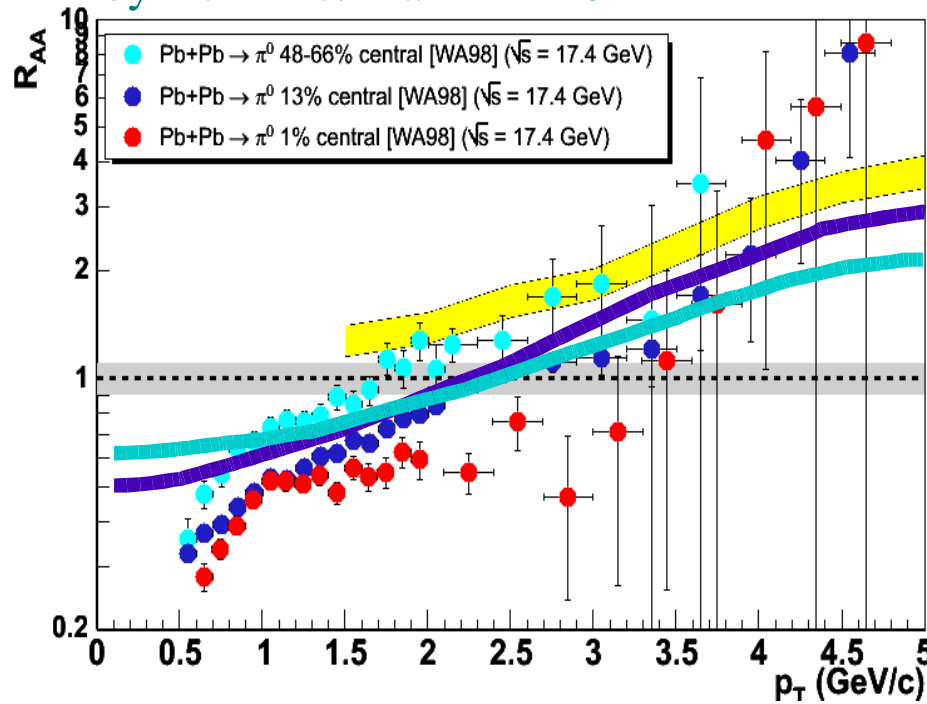
quenching
begins

more and more quenched

Quenching begins at around 40-45% centrality

4) WA98 π^0 - $\sqrt{s}=17.4$ GeV ($p_0 = 0.59$ GeV \pm 10%) work in progress

reanalysis of p+p reference for WA98 data
by D.D'Enterria - PLB'04



indications of:

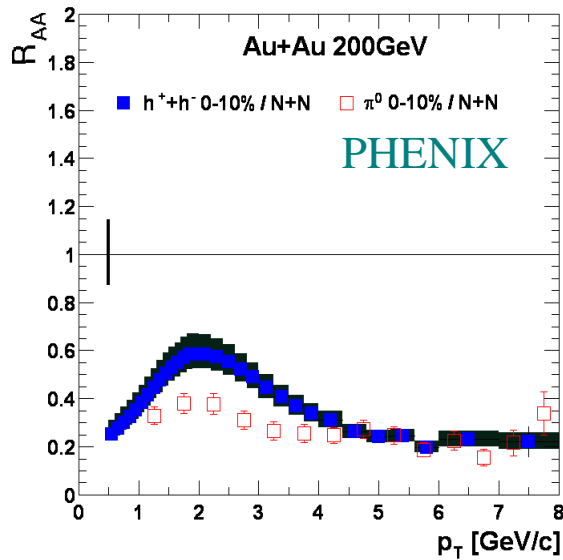
**very small or no quenching
in peripheral collisions**

**Moderate quenching
in central collisions**

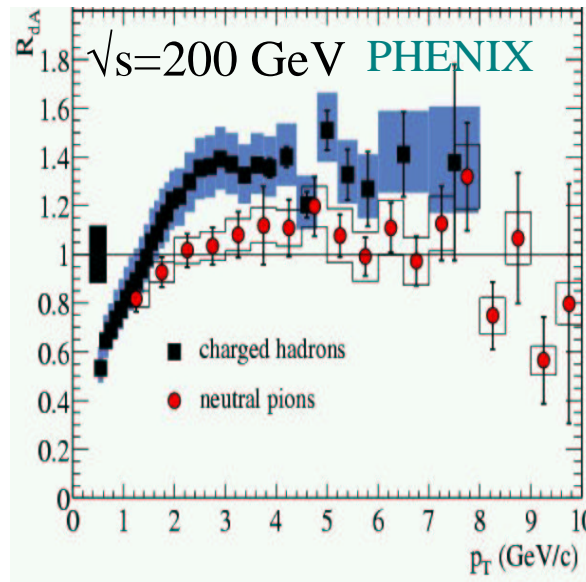
- ★ **large experimental uncertainty** due to p+p normalization
 - ➔ 25% systematic error on p+p baseline not shown
 - ➔ chosen p+p parametrization underestimates data at $p_T \gtrsim 3$ GeV
- ★ **large theory uncertainty** due to $K=1.1$ log-extrapolation
 - ➔ theory curves to be taken as approximate lower limits
 - ➔ also, scales are quite low for perturbative computations...

Baryon anomaly, a pervasive theme

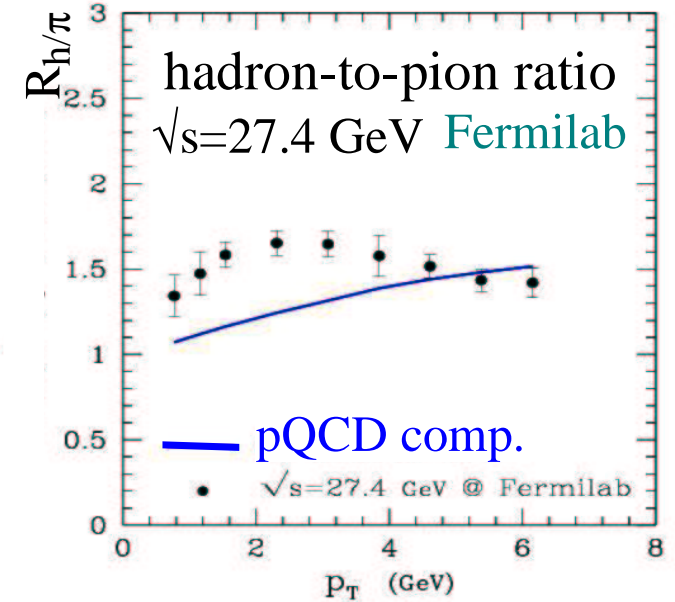
in A+A



in h+A



in p+p



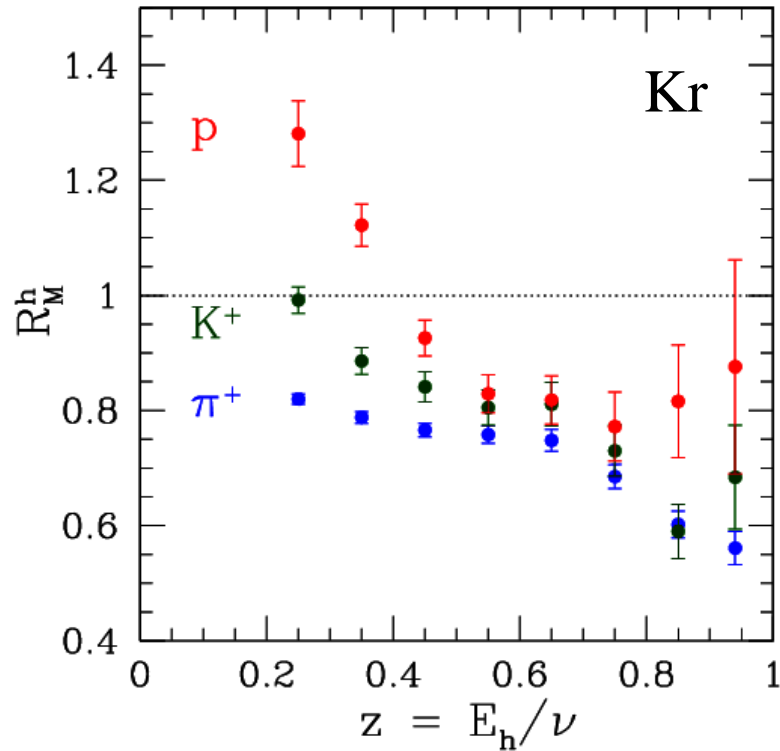
Baryon anomaly is hard to understand in pQCD
 \Rightarrow Cronin computations only for π^0

Baryon anomaly also in DIS on nuclei!

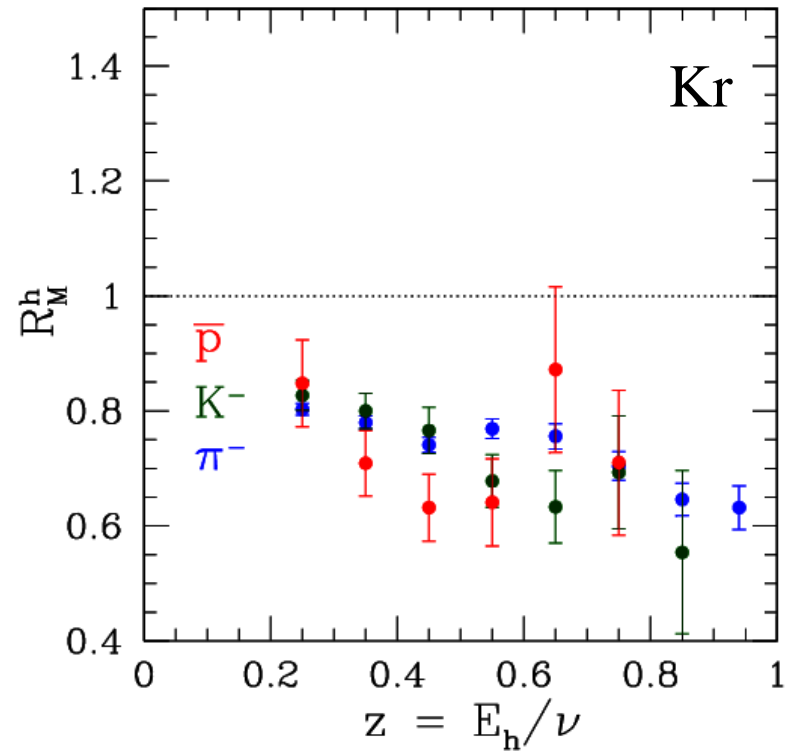
HERMES $E_{\text{lab}}=27$ GeV - Phys.Lett.B577(03)37

hadron "quenching" in e+A:

$$R_M^h(z, \nu, p_t^2, Q^2) = \frac{\left. \frac{N_h(z, \nu, p_t^2, Q^2)}{N_e(\nu, Q^2)} \right|_A}{\left. \frac{N_h(z, \nu, p_t^2, Q^2)}{N_e(\nu, Q^2)} \right|_D}$$



baryons are anomalous...



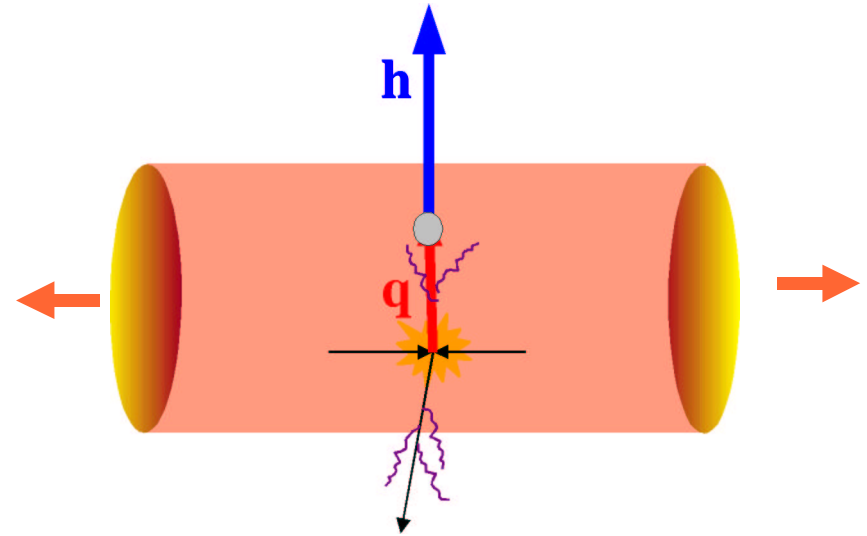
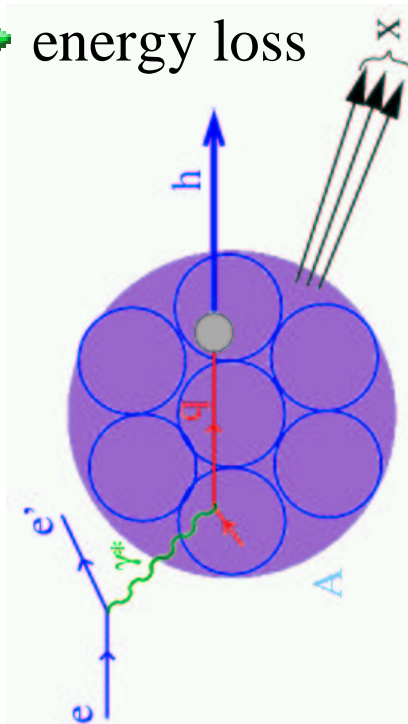
... but antibaryons are not!

Lessons from HERMES data

★ cold nuclear matter effect in clean environment

➔ hadron formation time

➔ energy loss



$$E_q = \nu = E_e - E_{e'} \approx 13 \text{ GeV}$$

on average

$$E_h = z \nu \approx \mathbf{2 - 12 \text{ GeV}}$$

$$E_q = p_T / z$$

$$E_h = p_T \approx \mathbf{2 - 12 \text{ GeV}}$$

★ HERMES kinematics is relevant to RHIC mid-rapidity

Conclusions

★ **Glauber-Eikonal model** describes fairly well Cronin effect in mid-rapidity d+Au on a broad energy range $\sqrt{s}=20-200$ GeV

⇒ **baseline computation for Au+Au collisions**
("Naked Cronin effect")

★ **d+Au data vs. theory**

➔ no initial-state effect at mid-rapidity

★ **Au+Au data vs. theory**

➔ quenching at $\sqrt{s}=62.4$ GeV starts at around 40% centrality

➔ indications moderate quenching in SPS central collisions?

★ **read HERMES papers:** nDIS is relevant to RHIC!

★ Future:

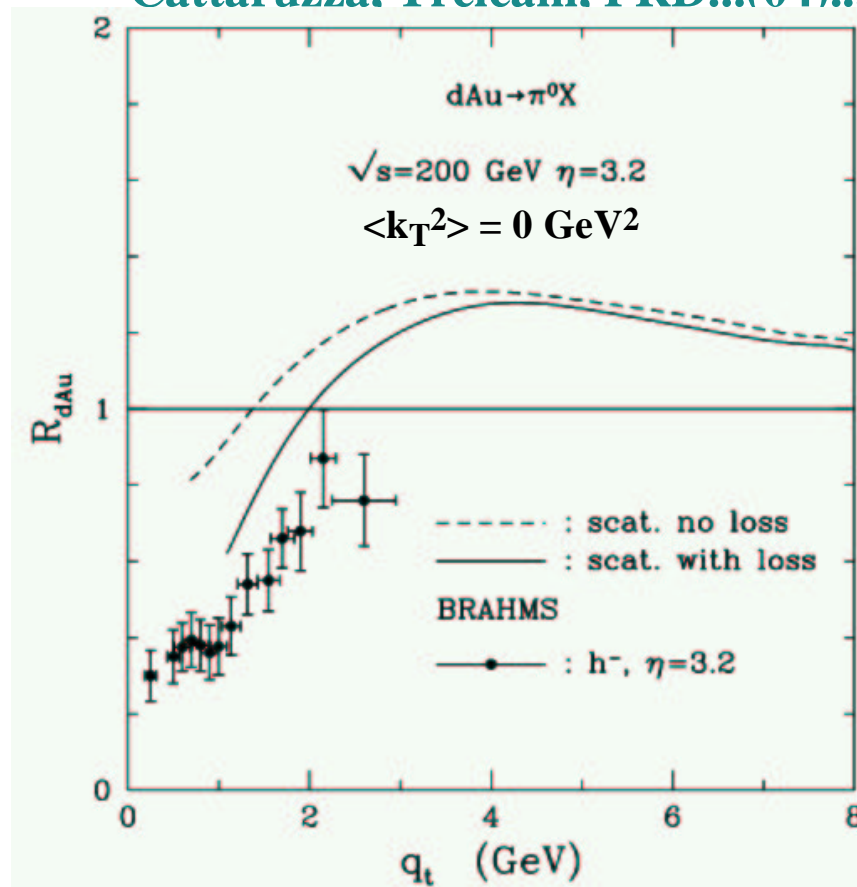
➔ careful check of assumptions

➔ inclusion of elastic energy loss & coherent multi-scatterings

The End

Glauber-Eikonal with elastic energy loss

Cattaruzza. Treleani. PRD...(04)...



Elastic energy loss:

parton energy is conserved,
but x is not!

- ◆ $x \rightarrow x + \Delta x$
- ◆ negligible at high- p_T
- ◆ **potentially large correction at low- p_T**

Geometric shadowing & Cronin effect

Integrated parton yield
(dominated by low- p_T)

$$\frac{d\sigma^{iA}}{d^2bd\eta} \approx 1 - e^{-\underbrace{T_A(b)\sigma^{iN}(\eta;p_0,K)}_{\text{opacity } \chi=\chi_A^i(b,\eta)}} \leq 1$$

unitarity

opacity $\chi=\chi_A^i(b,\eta)$

(average no. of scatterings)

Two limits:

1) low-opacity

$$\chi \ll 1 \implies \frac{d\sigma^{iA}}{d^2bd\eta} \approx T_A \sigma^{iN}$$

collision scaling
(single semihard scattering)

2) high-opacity

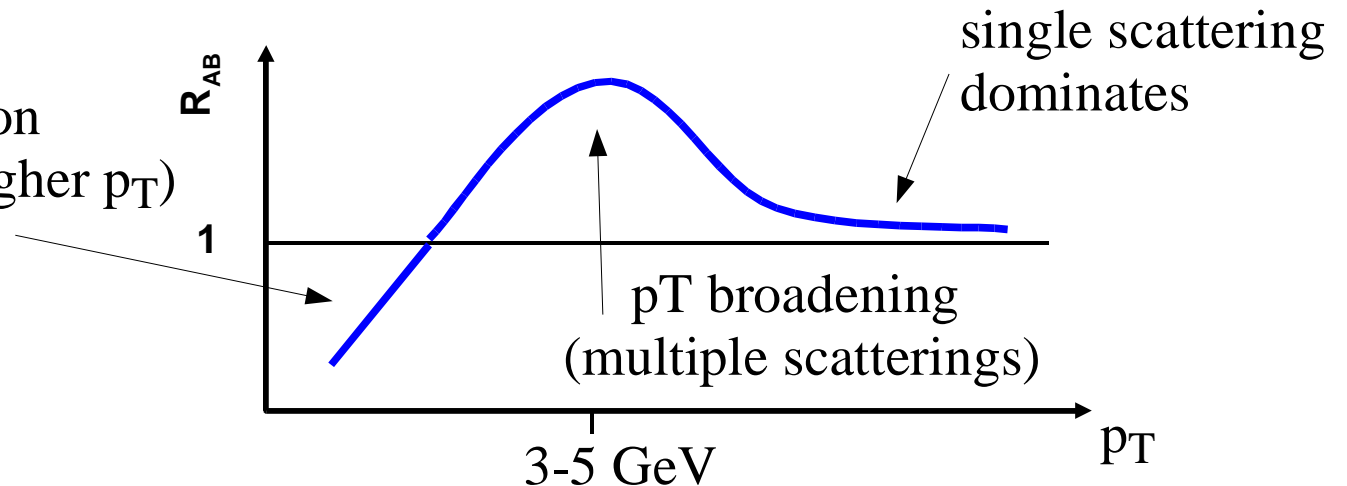
$$\chi \gtrsim 1 \implies \frac{d\sigma^{iA}}{d^2bd\eta} \ll 1 \lesssim T_A \sigma^{iN}$$

geometric shadowing

Sum of 2 effects:

a) momentum conservation
(spectrum shifted to higher p_T)

b) **geometric shadowing**



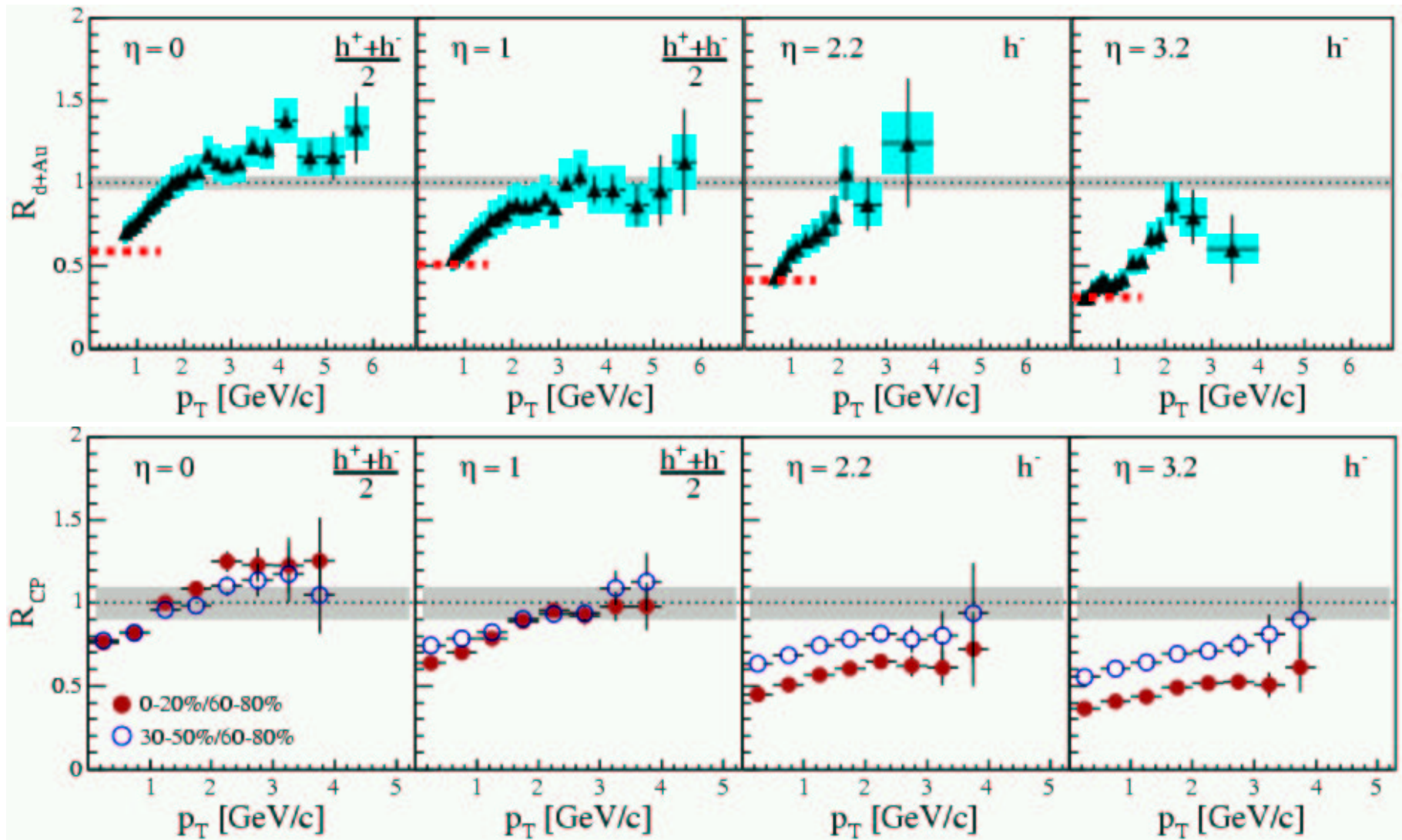
NOTE: "Dynamical" shadowing NOT included (no CGC, no geom. scaling, no EKS98, ...)

Hermes kinematics

		EMC		HERMES	
		h		h	π, K
E_{beam}	GeV	100	200	27.5	27.5
Q_{min}^2	GeV ²	2	2	1	1
W_{min}^2	GeV ²	4	4	4	4
y_{max}		0.85	0.85	0.85	0.85
x_{min}		0.02	0.02	0.06	0.02
x_{max}		1	1	1	1
z_{min}		0.2	0.2	0.2	0.2
z_{max}		1	1	1	1
ν_{min}	GeV	10	30	7	7
ν_{max}	GeV	85	170	23.4	23.4
$E_{h\text{min}}$	GeV	3	3	1.4	2.5
$E_{h\text{max}}$	GeV	85	170	23.4	15.0

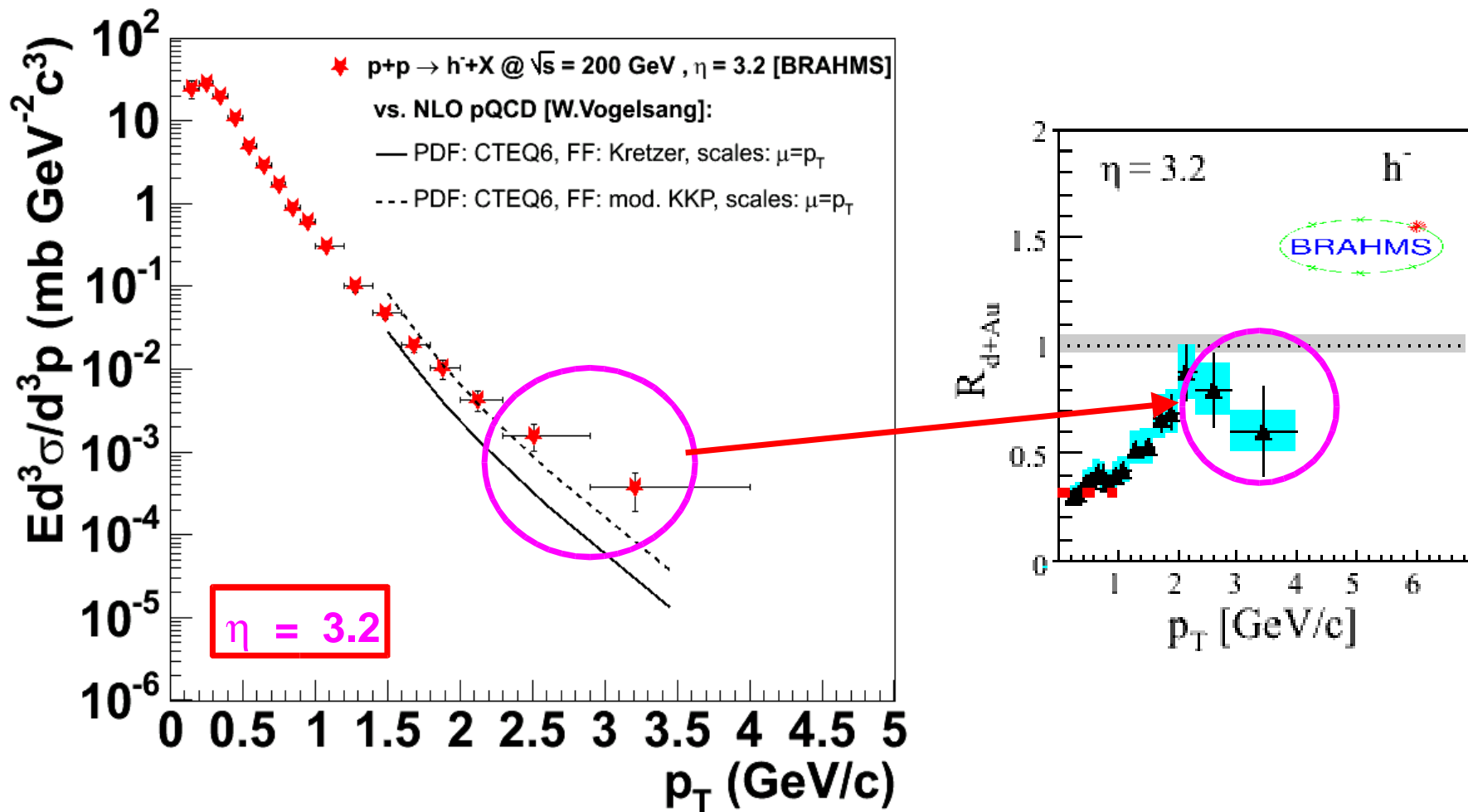
Table 3: Kinematic cuts of the EMC and HERMES experiments.

Mystery no. 1



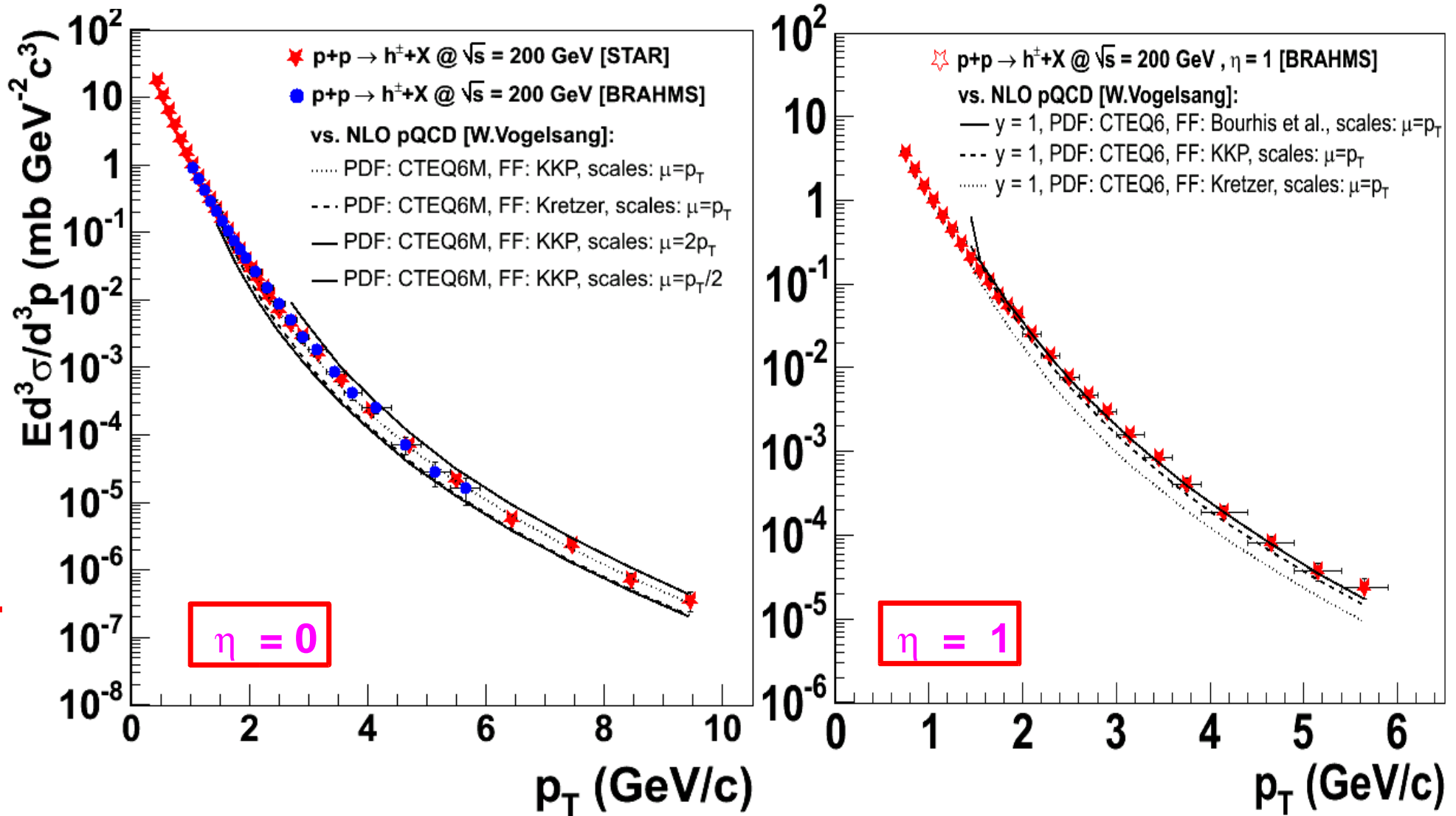
- ◆ **Discrepancy between R_{dAu} and R_{CP}**
- ◆ **Hard to understand theoretically**

High p_T p+p reference spectrum vs pQCD ($\eta = 3.2$)



- Two highest p_T points (there where the R_{dAu} "suppression" appears) are enhanced w.r.t. NLO (which describes all other rapidities, even $y = 3.8$!).
- Let's be provocative ... Is there d+Au suppression or (genuine" ?) p+p enhancement

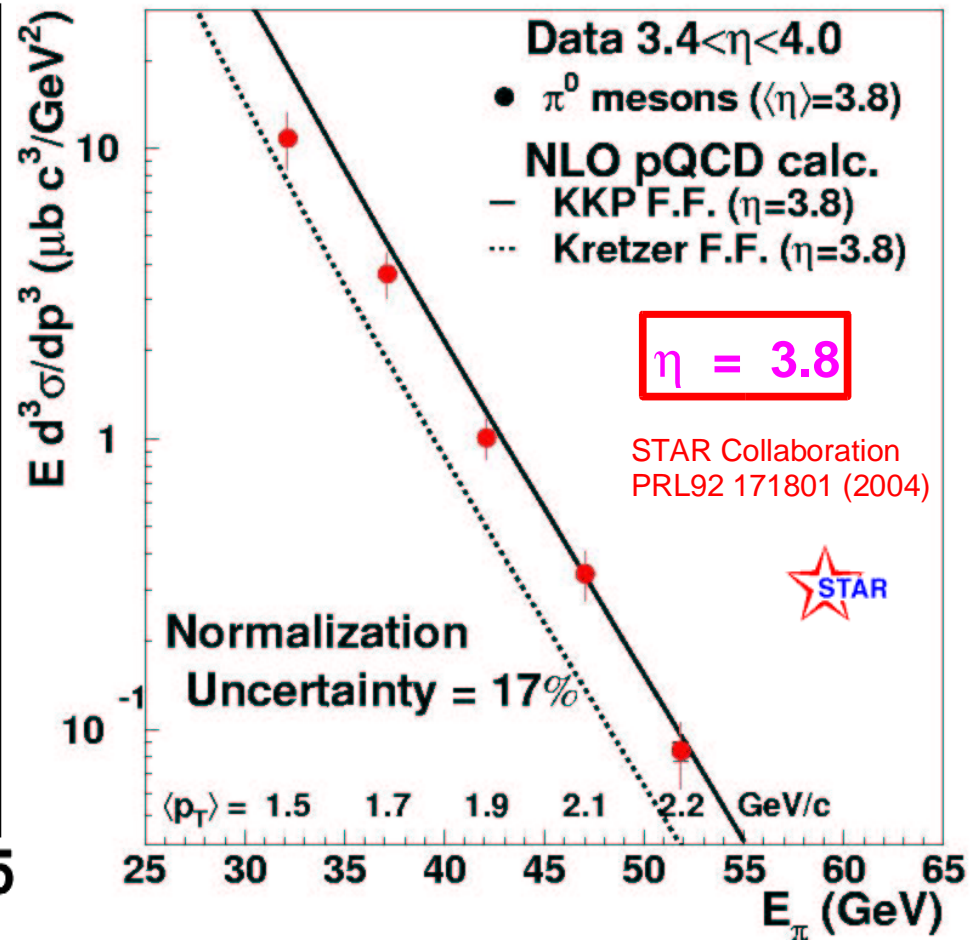
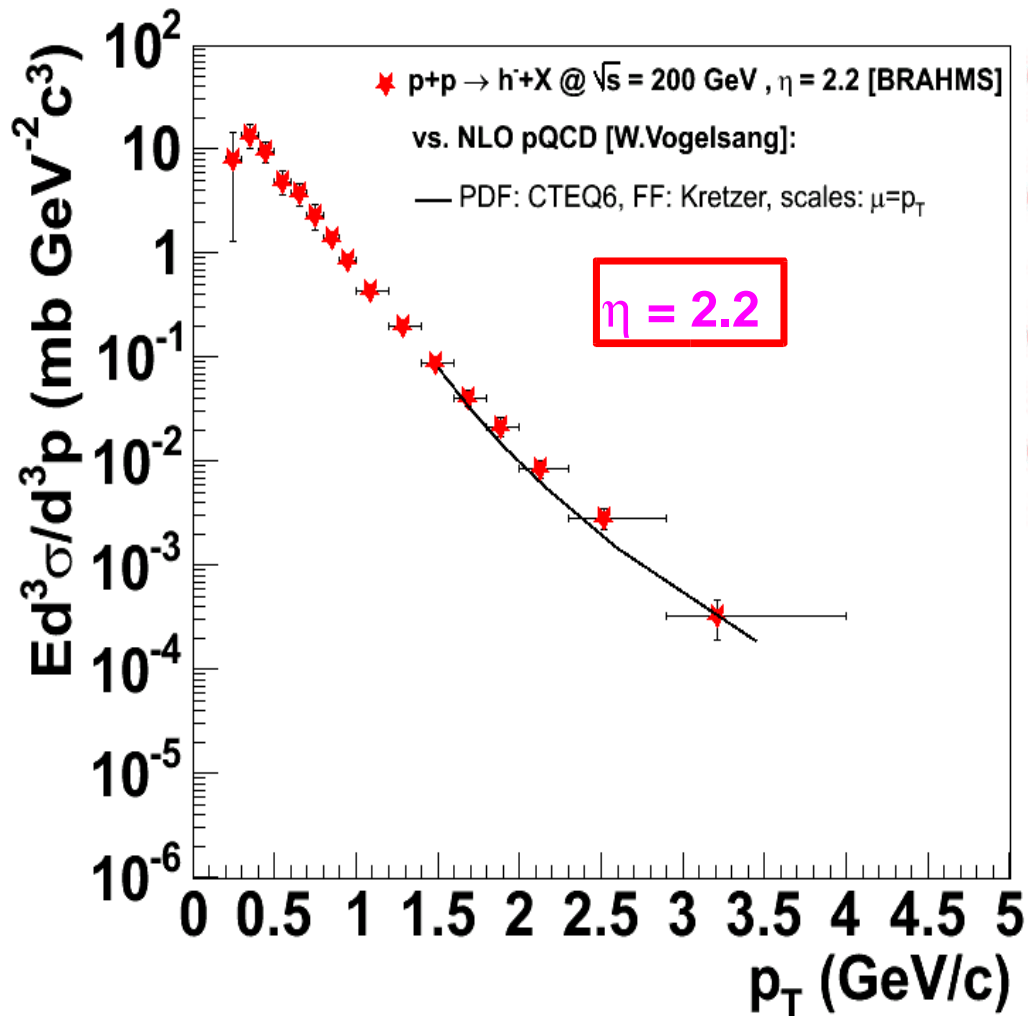
High p_T p+p reference spectra vs pQCD (mid-rapidity)



Good agreement with NLO pQCD

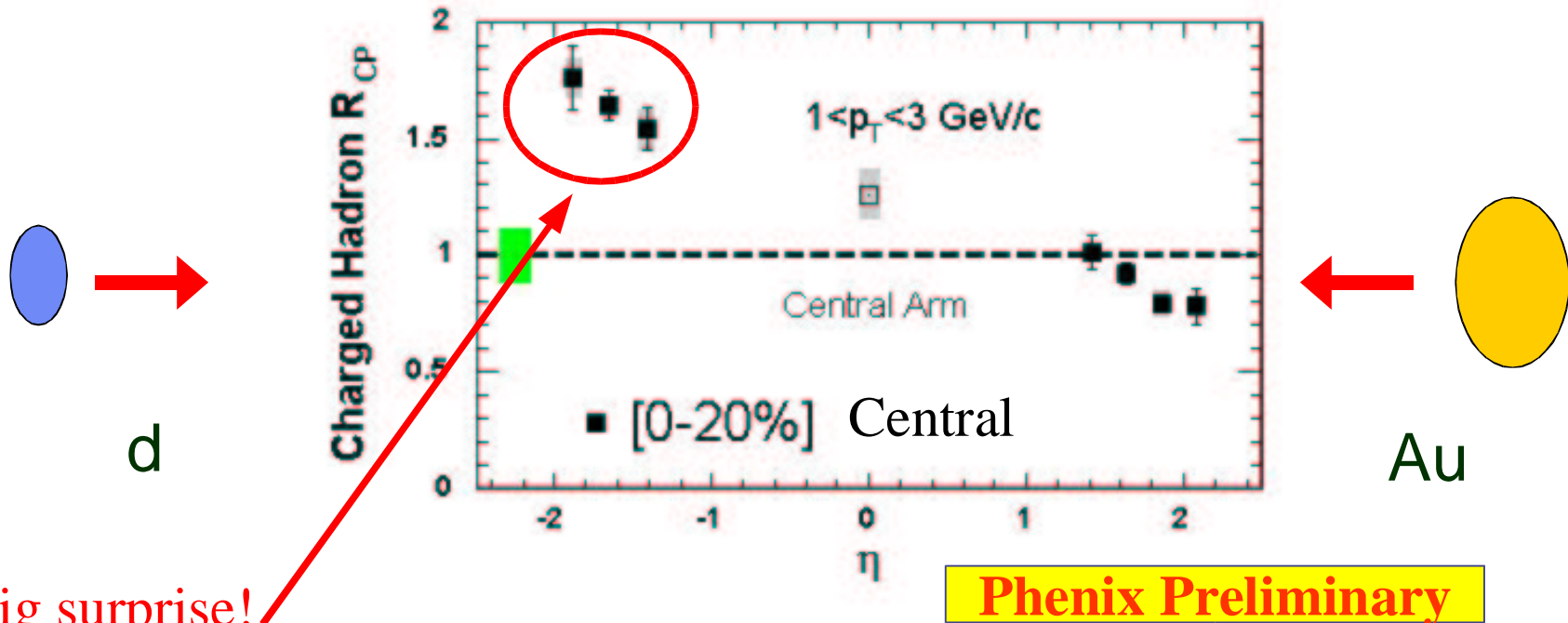
[calculations by W. Vogelsang]

High p_T p+p reference spectra vs pQCD (forward)



 Good agreement with NLO pQCD [calculations by W. Vogelsang]

Mystery no. 2: increasing Cronin at $\eta < 0$



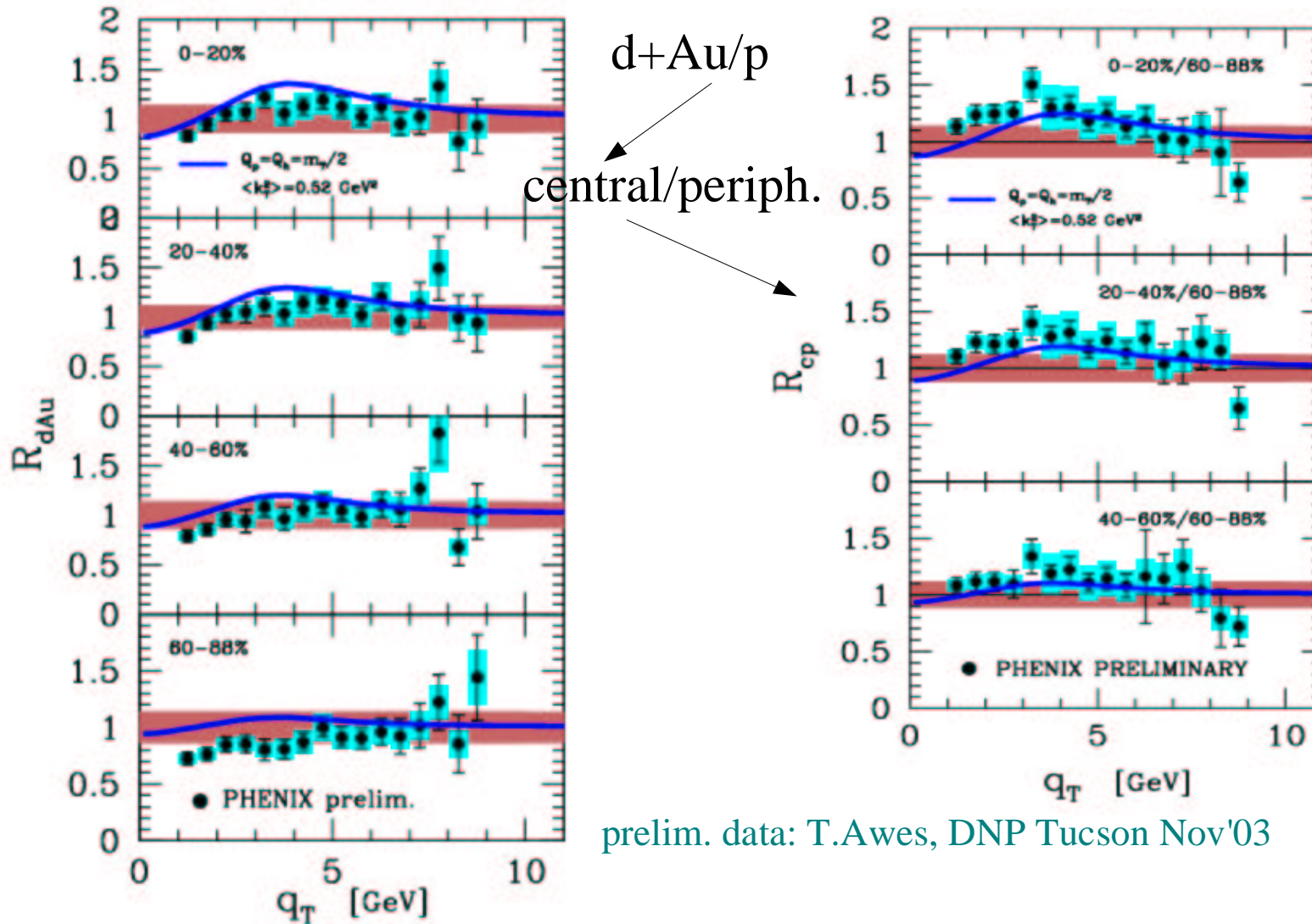
Big surprise!

The deuterium is a smallish object,
it should cause only a small Cronin!
Soft physics? Particle mix?
Anti-shadowing of Au partons?

no model to my knowledge can explain this

Centrality dependence at $\eta=0$

If dynamical shadowing at work, \Rightarrow stronger suppression in central



prelim. data: T.Awes, DNP Tucson Nov'03

no dynamical shadowing at $\eta=0$

GE model computation

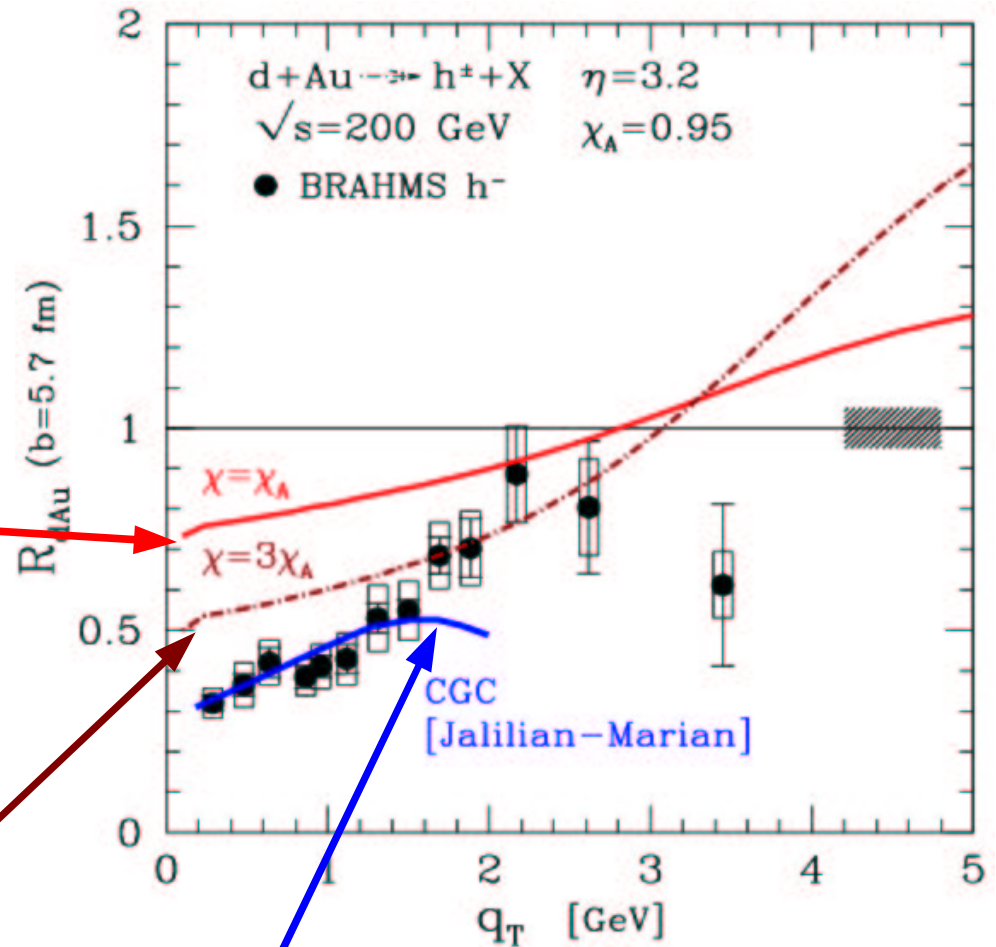
- No available data for pp coll's at forward rapidity \Rightarrow we take **same parameters as for $\eta=0$** :
 - $p_0(\eta>0) = p_0(\eta=0) = 1.0$ GeV
 - $K(\eta>0) = K(\eta=0) = 1$

↓

η	0	1.8	3.2
χ	0.74	0.85	0.95

Are we underestimating opacity?
 $\chi_A=0.95$ obtained with p_0 and K fitted in pp at $\eta=0$: **may change with η** :

With $\chi=3\chi_A$, we are not so bad



CGC à la Jalilian-Marian
 underestimates data at $p_T > 1.5$ GeV

CGC and p+A collisions

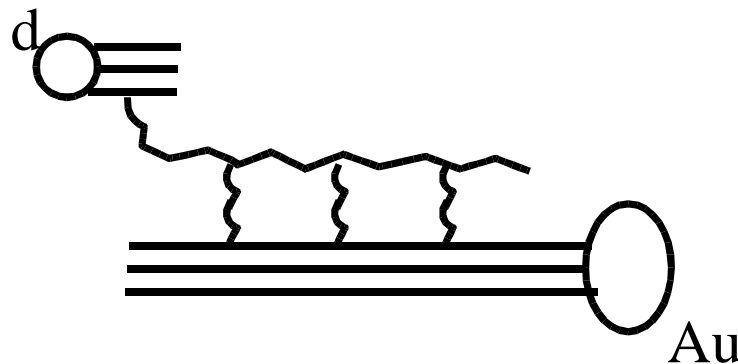
For gluon and valence quark production (Dumitru, Gelis, JalilianMarian; Blaizot et al.)

$$\frac{d\sigma^{d+A \rightarrow g(q_V)+X}}{d^2p_T dy} \propto \frac{1}{p_T^2} \int \frac{d^2k_T}{(2\pi)^2} f_{g(q_V)}^p(x_1, \vec{p}_T - \vec{k}_T) k_T^2 C(x_2; \vec{k}_T)$$

where $C(x_2; k_T) = \int d^2r_T e^{i\vec{r}_T \cdot \vec{k}_T} \langle U^\dagger(0) U(\vec{r}_T) \rangle_{x_2}$ (fundam. or adjoint Wilson line).

If W is Gaussian, then Glauber multiscattering interpretation:

$$C(x_2; k_T) = \int d^2r_T e^{i\vec{r}_T \cdot \vec{k}_T} e^{\tilde{\sigma}(r_T; b)}$$



MV model: local gaussian correlations

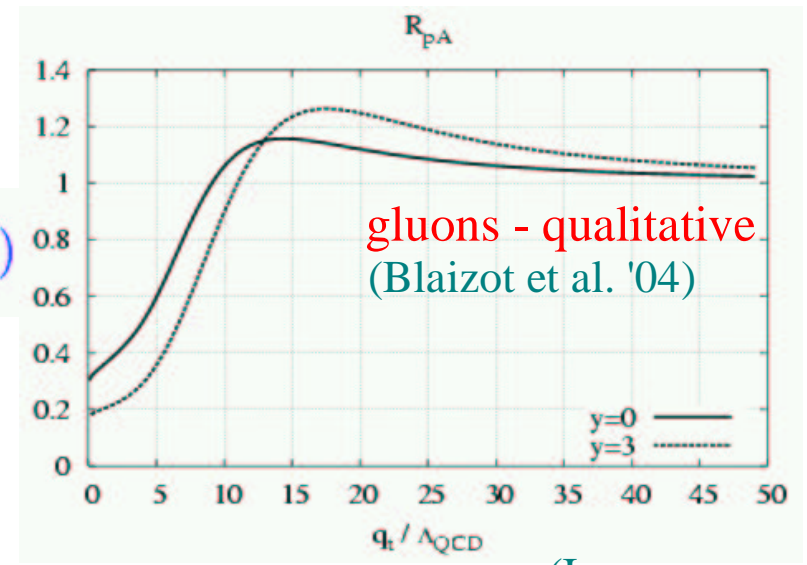
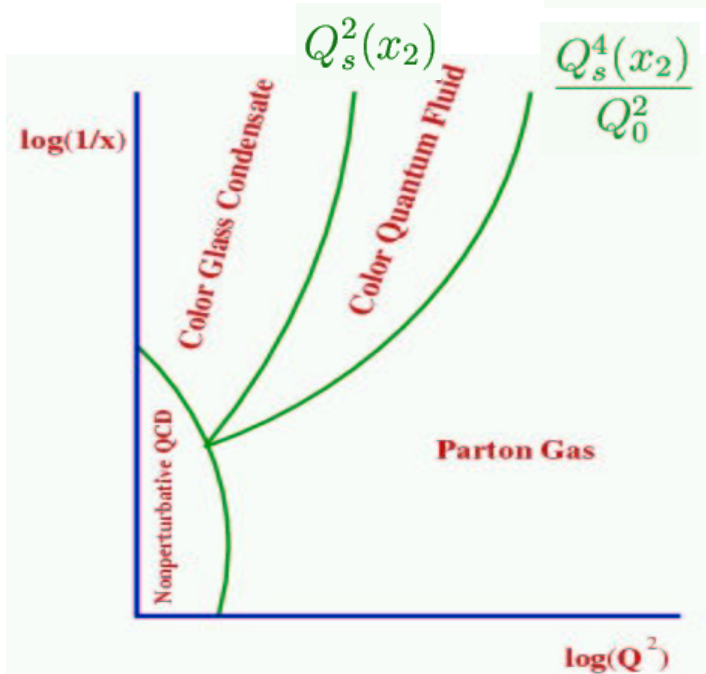
$$\lambda_y(y_T - x_T) = \lambda_y = Q_s^2(x_2) = \alpha_s^2 \frac{8\pi N_c}{N_c - 1} x_2 G(x_2; Q^2)$$

$$\tilde{\sigma}(r_T) = r_T^2 Q_s^2(x_2) + O(r_T^4) \approx \tilde{\sigma}_{pQCD} !! \implies$$

CGC = MV + quantum evolution:

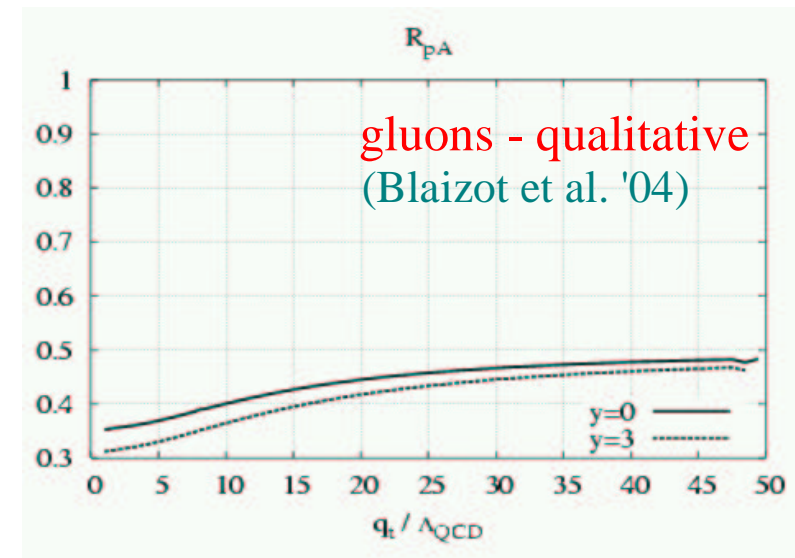
non-local gaussian correlations \implies approx. solution of RG equation: [Itakura, McL](#) (Iancu,

In saturation region: $\lambda_y(y_T - x_T) \propto \int d^2k_T e^{i(\vec{y}_T - \vec{x}_T) \cdot \vec{k}_T} k_T^2 \ln\left(1 + \left(\frac{Q_s^2(x_2)}{k_T^2}\right)^\gamma\right)$



gluons - qualitative
(Blaizot et al. '04)

(Iancu,
Itakura, McL)



gluons - qualitative
(Blaizot et al. '04)