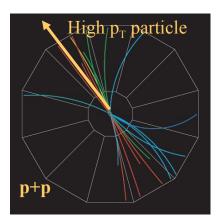
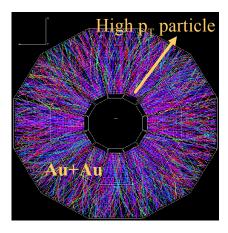
# Dihadron fragmentation in vacuum and in matter

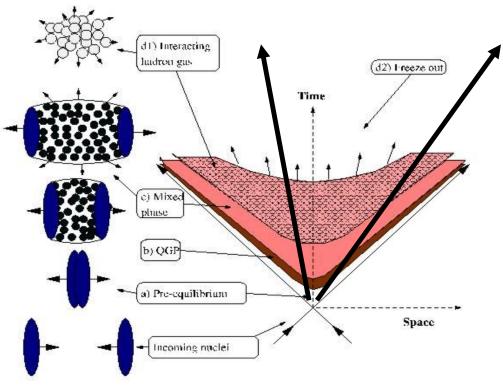
Hard Probes 2004 Ericeira, Portugal Nov. 4 - 10 Abhijit Majumder Nuclear Theory Group, LBNL

in collaboration with Xin-Nian Wang

## **HEAVY-ION COLLISIONS AND JETS**

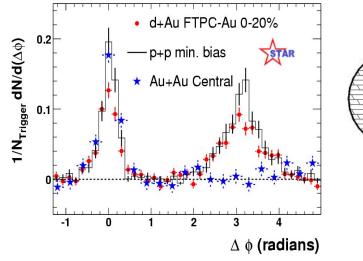


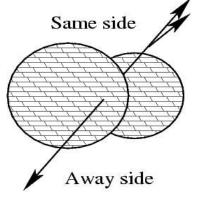




Select a leading particle  $4 < p_t < 6 \text{ GeV/c}$ ,  $|\eta| | < 0.75$ Associate all other particles ( $0.15 < p_t < 4 \text{ GeV/c}$ ,  $|\eta| < (1.1)$  with the leading particle.

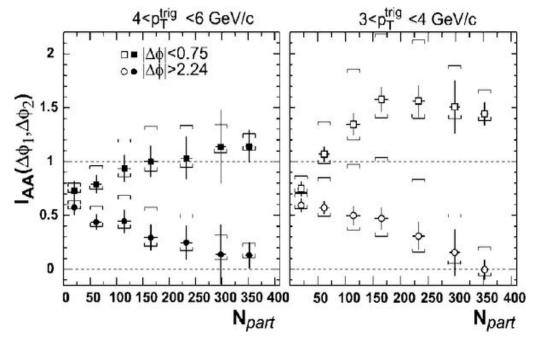
## **RESULTS FROM STAR, heavy-ion collisions**



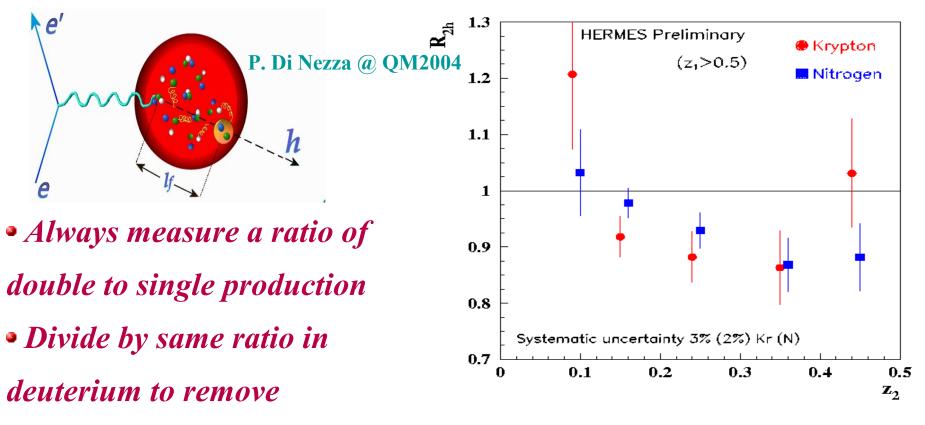


• d+Au central very similar to p+p

- Away side, central Au+Au suppressed
- Near side, central Au+Au unchanged
- Can explain the integral of the near side vs Npart.



### **RESULTS FROM HERMES, DIS on cold nuclei**



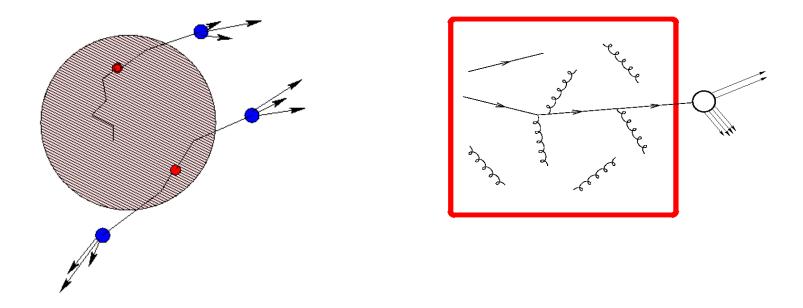
detector systematics

$$R_{2h} = \frac{No. of events with at least 2 hadrons with z}{No. of events with at least one hadron with same ratio on deuterium 4}$$

## **TWO POSSIBILITIES !**

#### **PARTONIC ENERGY LOSS :**

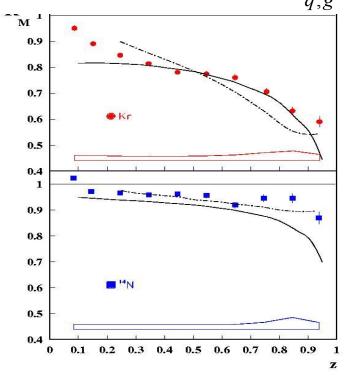
• High energy partons are created over the entire collision zone

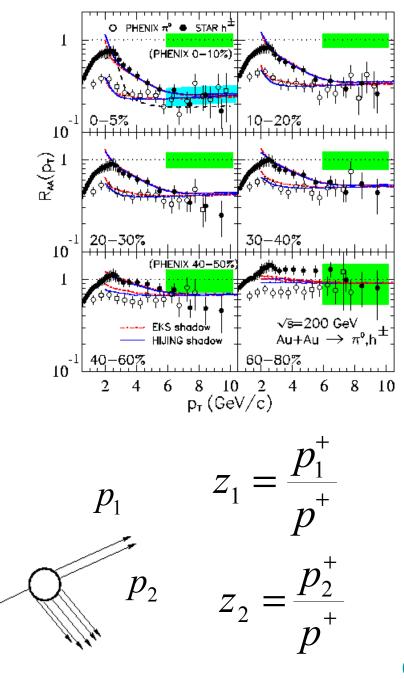


- Lose energy by partonic interaction, medium may be hadronic or partonic
- Emerge as partons and then fragment

 Partonic energy loss models explain single inclusive suppression pretty well !
 (GLV, BDMPS, WGZ, SW) To explain double inclusive spectra requires a new phenomenological object: Dihadron fragmentation function! D<sup>h1h2</sup><sub>a,g</sub>(z1, z2)

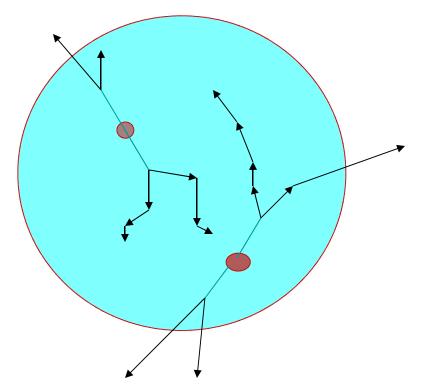
p





#### HADRONIC ENERGY LOSS:

- Fragmentation occurs inside the hot medium
- Hadrons become independent due to scattering
- Each hadron suffers the same Energy Loss on average



Hadronic scattering models can explain mean single supp. ! 11/22/Greiner et. al. @QM2004, V. Koch (unpublished!)

- Probability of observing two hadrons factorizes P(1,2) = P(1)P(2)
- Each probability is suppressed compared to p+p: P(h) = s p(h)
- Thus the conditional probability is also suppressed compared to p+p collisions

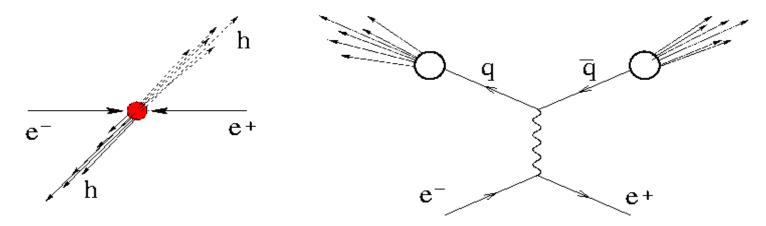
$$\frac{P(1,2)}{P(1)} = \frac{P(1)P(2)}{P(1)} = P(2) = sp(2) = s\frac{p(1)p(2)}{p(1)}$$

P(h) is for A+A; p(h) is for p+p

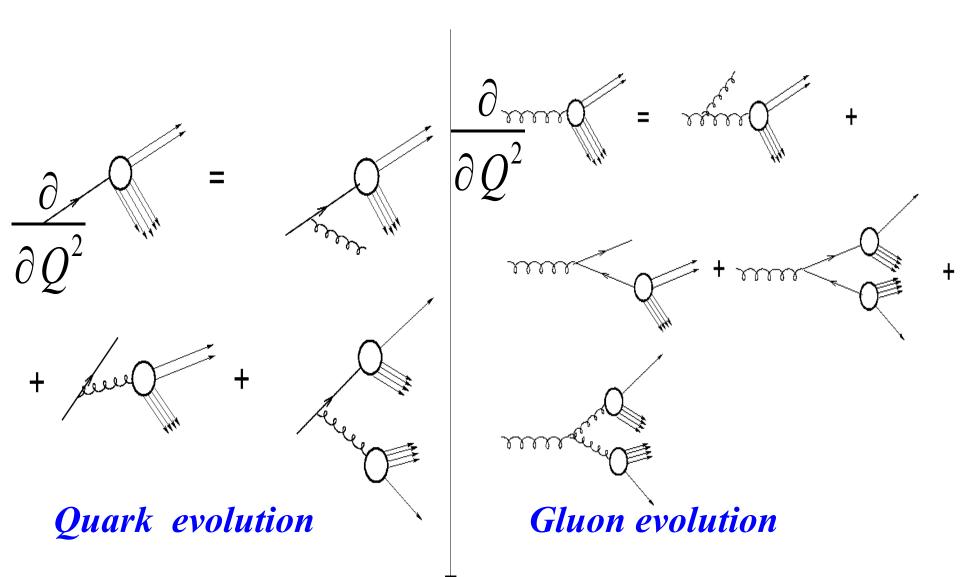
• Hadronic absorption models cannot explain the double inclusive spectrum

Modification of Dihadron fragmentation function!

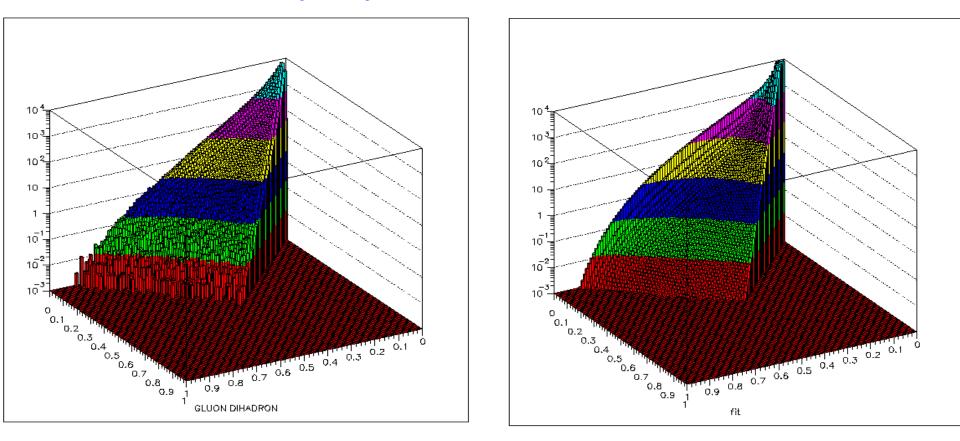
- Fragmentation functions are universal
- We have a definition in terms of operators
- Start with simple system :  $e^+e^-$ , can factorize  $\lambda^2 \ll \mu^2 \ll Q^2$
- Derive evolution (vacuum splitting functions same)
- Check with data and repeat in medium (similar to single)



Evolving to a higher scale Q solving DGLAP equations
Set of coupled differential equations containing the following processes:

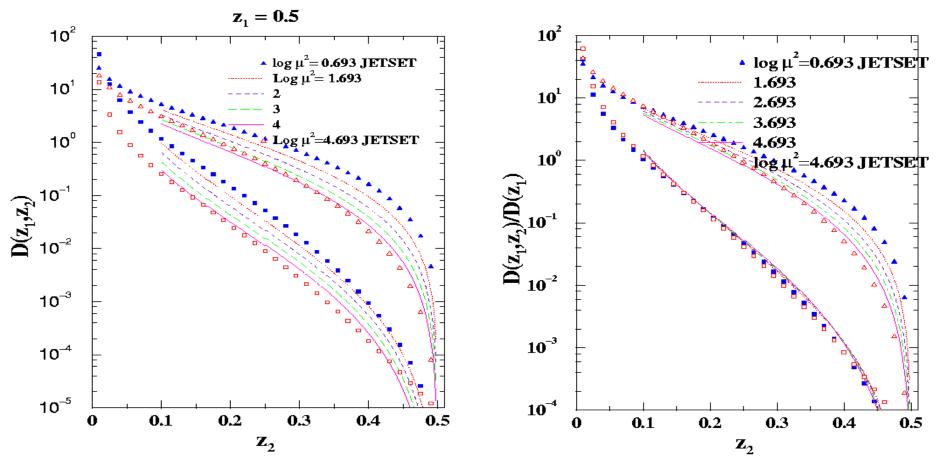


## **Results from Event generators: JETSET** *insist z*<sub>1</sub> > *z*<sub>2</sub>, *fit a function to it !*



$$D(z_{1}, z_{2}) = N z_{1}^{\alpha_{1}} z_{2}^{\alpha_{2}} (z_{1} + z_{2})^{\alpha_{3}} (1 - z_{1})^{\beta_{1}} (1 - z_{2})^{\beta_{2}} (1 - z_{1} - z_{2})^{\beta_{3}}$$

11/22/04

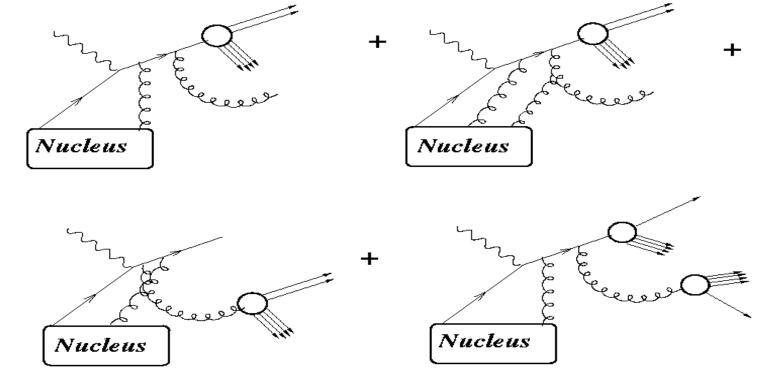


Quark and Gluon evolution fits event generator data very well! Thus we can understand evolution of FF from QCD.

the double to single ratio shows little change

## **Medium modification**

- Apply to DIS of Nuclei (HERMES expt. at DESY)
- A parton in a nucleon is struck by EM probe
- Parton scatters in medium and then exits & fragments
- Fragmentation function is medium modified. The medium modification also has new set of diagrams!



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**DIS** followed by di-hadron fragmentation from a large nuclei may be generally expressed as

$$\frac{d^{-2}W^{\mu\nu}}{dz_{1}dz_{2}} = \int dx f_{q}^{A}(x)H^{-\mu\nu}\tilde{D}^{h_{1}}$$

nucleons

$$\tilde{D} = medium \ modified \ fragmentation \ fun$$

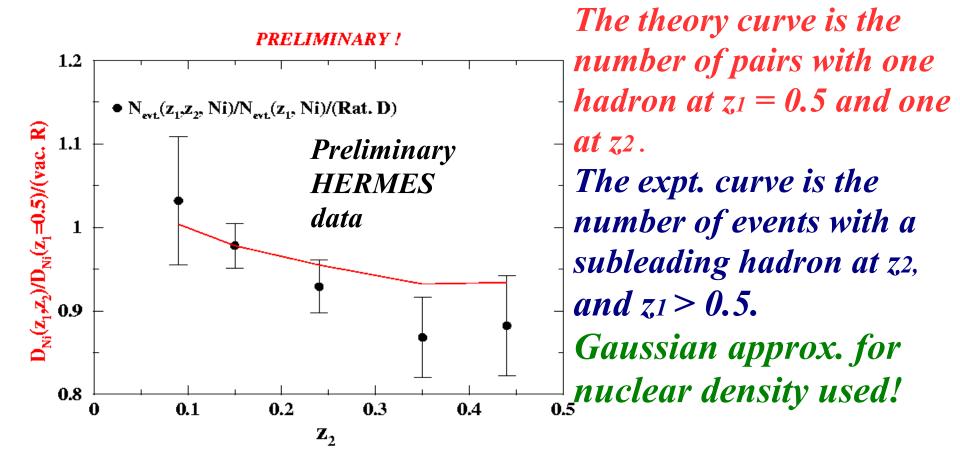
$$\tilde{D}(z_1, z_2, \mu^2) = D(z_1, z_2, \mu^2) + \frac{\alpha_s}{2\pi} \int_0^{\mu^2} \frac{dl^2}{l^2} \int \frac{dy}{y^2} \left( \frac{1+y^2}{1-y} T_{gg}(x, y, Q^2, l^2) + V.C. \right) D(z_1/y, z_2/.$$

$$\int_0^{\varphi} \int T_{gg} \to \int dy \ dy \ dy \ dy \ dy \ 2 \langle A | \overline{\psi}(y) F(y_1) F(y_2) + this \ \sim A^{\frac{1}{3}}$$

$$Luo, Qiu \ and \ Sterman \ PRD \ 50, \ 1951 \ (1994).$$

$$Wang \ and \ Guo \ NP \ A696, \ 788 \ (2001).$$
14

14



Theory curve: (FF(2h)/FF(1h) in A) / (FF(2h)/FF(1h) in vac.)

**Expt ratio** =  $\frac{No. of events with at least 2 hadrons with <math>z_1}{No. of events with at least one hadron with same ratio on deuterium}$ 

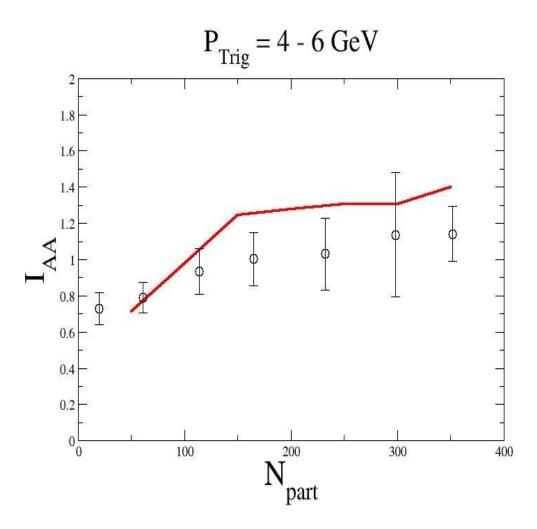
11/22/04

## **Dihadron results for hot medium**

Very preliminary estimate for the same side two body correlation

*Results include the effect of trigger bias.* 

Initiating parton in a heavy-ion collision has higher energy than that in p-p collision..



## Summary & Conclusions!

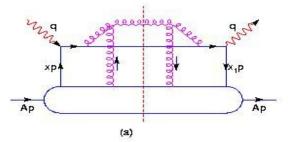
- We have defined a new phenomenological object in QCD:
   `` The Dihadron Fragmentation function ''
- Demonstrated its factorization at LO in  $e^+e^-$
- Derived it evolution equation (has extra components)
- Matched results with JETSET!!
- Allowed a physical understanding of change with scale
- Extended formalism to medium modification in **DIS**
- Extended formalism to medium modification in heavy-ion collision

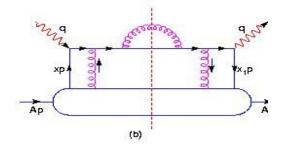
Back up slides ...

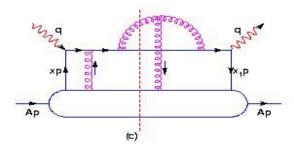
## Multiple higher twist diagrams need to be evaluated

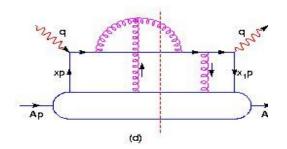
Multiple scattering from soft gluons lead to LPM interference

Assume a Gaussian density distribution for nucleons in a medium sized nucleus

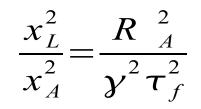








$$T_{qg}^{A} = C A^{1/3} (x G^{N}(x)) (1 - e^{-x_{L}^{2}/x_{A}^{2}})$$



 $\tau_f = Formation time$ 

 $\gamma = boost$ 

R = Nuclear size

Fragmentation functions come from the soft matrix elements  

$$D_{q}(z) = \frac{z^{3}}{2} \int \frac{d^{-4}p}{(2\pi)^{4}} Tr \left[ \frac{y^{+}}{2p_{h}^{+}} \delta\left(z - \frac{p_{-h}^{+}}{p^{+}}\right) \hat{T}_{q}(p_{-}, \frac{p_{h}^{-}}{p_{-}}) \hat{T}(p_{-}, \frac{p_{h}^{+}}{p_{-}}) \right]$$
where  $\hat{T}_{q}(p_{-}, p_{-h}) = \int d^{-4}x e^{-ip \cdot x} \sum_{S-1} \left| 0 \right| \psi(0) |p_{-h}, S-1 \right| \left| p_{-h}, S-1 \right|$ 

$$D_{q}(z_{1}, z_{2}) = \int \frac{dq^{-2}}{8(2\pi)^{2}} \frac{z^{4}}{4z_{1}z_{2}} \int \frac{d^{-4}p}{(2\pi)^{4}} Tr \left[ \frac{y^{+}}{2p_{-h}^{+}} \delta\left(z_{1} + z_{2} - \frac{p_{-h}^{+}}{p^{+}}\right) \hat{T}_{q}(p_{-}, \frac{q_{-}}{2}) \right]$$
A. Mueller, PRD 18, 3705 (1978).  

$$= \frac{z^{4}}{z_{1}z_{2}} P_{h} = p_{h}^{-4} + p_{2}^{+} \int \frac{dq^{-2}}{4(2\pi)^{2}} \hat{T}(p_{-}, p_{-}, p_{-}) = \frac{z^{4}}{4(2\pi)^{2}} P_{h} = p_{h}^{-4} + p_{2}^{-4} \int \frac{dq^{-2}}{4(2\pi)^{2}} \hat{T}(p_{-}, p_{-}, p_{-}) = \frac{z^{4}}{2} P_{h} = p_{h}^{-4} + p_{2}^{-4} \int \frac{dq^{-2}}{4(2\pi)^{2}} \hat{T}(p_{-}, p_{-}, p_{-}) = \frac{z^{4}}{2} P_{h} = p_{h}^{-4} + p_{2}^{-4} \int \frac{dq^{-2}}{4(2\pi)^{2}} \hat{T}(p_{-}, p_{-}, p_{-}) = \frac{z^{4}}{2} P_{h} = p_{h}^{-4} + p_{2}^{-4} \int \frac{dq^{-2}}{4(2\pi)^{2}} \hat{T}(p_{-}, p_{-}, p_{-}) = \frac{z^{4}}{2} P_{h} = p_{h}^{-4} + p_{2}^{-4} \int \frac{dq^{-2}}{4(2\pi)^{2}} \hat{T}(p_{-}, p_{-}, p_{-}) = \frac{z^{4}}{2} P_{h} = p_{h}^{-4} + p_{2}^{-4} + p_{2}^{-4} \int \frac{dq^{-2}}{4(2\pi)^{2}} \hat{T}(p_{-}, p_{-}, p_{-}) = \frac{z^{4}}{2} P_{h} = p_{h}^{-4} + p_{2}^{-4} + p_{2}^{-4} \int \frac{dq^{-2}}{4(2\pi)^{2}} \hat{T}(p_{-}, p_{-}, p_{-}) = \frac{z^{4}}{2} P_{h} = \frac{z^{4}}{2} P_$$