## Low Mass Dimuon Production in Proton-Nucleus Collisions with the NA60 Apparatus

Outline:

- Introduction

Motivation
NA60 apparatus; setup in 2002; data reconstruction and selection
Detector performance: Phase space coverage
Dimuon mass resolution and
Signal-to-background ratio

- Monte Carlo generation and comparison to data
- Extraction of physics results and discussion

Elementary (pp) production cross-sections of $\eta, \rho, \omega$ and $\phi$
Nuclear dependence of the $\eta, \rho / \omega$ and $\phi$ production cross-sections

## Motivation

The study of low mass dilepton production in nuclear collisions provides a window of opportunity to learn about several interesting physics topics:

- Medium effects on the mass and width of the $\rho$ vector meson (less for $\omega$ and $\phi$ )
$\rightarrow$ which might be due to the approach to chiral symmetry restoration
- Thermal virtual photon production from the earliest stages of the collision $\rightarrow$ which would constitute direct evidence of a quark-gluon-plasma phase
- Strangeness enhancement, through the $\phi$ meson
$\rightarrow$ link to general strangeness enhancement in deconfined phase
Such "new physics" studies must be built on top of a solid understanding of low mass dilepton production in proton-nucleus collisions, which provide a reference baseline with respect to which the heavy-ion specific phenomena can be extracted.

NA60 is presently taking a high statistics proton-nucleus data sample with seven different nuclear targets, at 400 and 158 GeV . The present talk reports on results obtained from a much smaller data sample collected in 2002.

## The NA60 Apparatus during the 2002 Proton Run

The muon spectrometer:

- Hadron absorber ( $13+7 \lambda_{1}$ )
- 4+4 MWPCs with 3 individual planes each (rotated by $60^{\circ}$ )
- 4 trigger hodoscopes
- Toroidal magnet

The vertex region:


## provides:

- Tracking of charged particles
- Target identification
- Improved dimuon kinematics via track "matching"
provides:
- Muon identification
- Highly selective dimuon trigger
- Track reconstruction of two muons however: affected by multiple scattering and energy loss induced by the hadron absorber


Strip segmentation adapted to the highly inhomogeneous particle production across the sensor surface.

## Data Taking, Reconstruction and Event Selection

## 1) Data Taking

- 400 GeV proton beam incident on Be , In and Pb targets ( $2 \%, 0.9 \%$ and $1.2 \% \lambda_{\text {Int }}$ respectively)
- All targets were simultaneously placed in the beam to reduce systematic effects in the extraction of the nuclear dependence of the production cross-sections
- During 4 days in 2002: 600000 dimuons collected (at "low" beam intensity: 1-3 $10^{8}$ protons/burst)

2) Data Reconstruction and Event Selection

- Reconstruct tracks in the muon spectrometer and build dimuons of all charge combinations
- Reconstruct the charged particle tracks in the vertex telescope
- Reconstruct the primary interaction vertex
- Select events with only one reconstructed vertex in the target region to reject pile-up
- Match the two muons to vertex telescope tracks, in coordinate and momentum space

Like Sign (LS) dimuons:
used to evaluate the "combinatorial background"

Opposite Sign (OS) dimuons: used for the physics data analysis from $\pi$, K decays through a mixed event technique

- Select matched dimuons in a well defined phase space window
- After the full reconstruction, vertex selection and phase space cuts we are left with $\sim 15000$ OS dimuons


## Target Identification

- Z-vertex resolution ~600-900 $\mu \mathrm{m}$ depending on the target position
$\Rightarrow$ allows us to clearly separate the individual targets ( 2 mm thick, 8 mm interspacing)
- Vertexing algorithm tuned through MC simulation
$\Rightarrow$ in only $\sim 2 \%$ of all generated events the collision vertex is reconstructed in a wrong target


The use of 3 target materials with very different mass numbers ( $\mathrm{Be}, \mathrm{In}, \mathrm{Pb}$ ) allows us to extract the nuclear dependence of the particle production cross-sections.

## Phase Space Window \& Acceptances

> Dimuon phase space $3.3<\mathrm{y}_{\mathrm{lab}}<4.2$ $\left|\cos \theta_{\mathrm{cs}}\right|<0.5$
> $\mathrm{~m}_{\mathrm{T}}>0.4+0.7(\mathrm{y}-4.2)^{2} \mathrm{GeV}$


- The phase space window was tuned to keep most of the dimuons collected in the $\omega$ and $\phi$ mass windows.
- Apart from the dimuon selection cuts ( $\mathrm{y}_{\mathrm{lab}}, \cos \theta$ and $\mathrm{m}_{\mathrm{T}}$ ) we also apply an angular single muon cut to stay away from the "beam-hole" of the strip sensors: $\eta(\mu)<4.2$
- Acceptances: $\rho \sim 3.3 \%, \omega \sim 3.6 \%, \phi \sim 6.5 \%$ (the exact value depending on the target position)
- The dipole magnetic field in the vertex region improves significantly the acceptance for low mass and low $\mathrm{p}_{\mathrm{T}}$ opposite sign dimuons



## Mass Resolution and Signal / Background

- Measuring the muons before they suffer multiple scattering and energy loss in the hadron absorber, thanks to our silicon vertex telescope, allows us to achieve a mass resolution on the $\omega$ and $\phi$ resonances of around 30 MeV . That is exactly the value expected from our MC simulations.
- Through the matching procedure the signal to background ratio improves by a factor of 4 .



## MC Generation of Light Meson Decays

- The $\eta, \eta^{\prime}, \rho, \omega$ and $\phi$ mesons were generated with
- $\mathrm{p}_{\mathrm{T}}$ distributions: $\left.\frac{1}{p_{T}} \cdot \frac{d N}{d p_{T}}=\frac{1}{m_{T}} \cdot \frac{d N}{d m_{T}}=m_{T} K_{1}\left(\frac{m_{T}}{T}\right)\right]$
- y distributions:

$$
\frac{d N}{d y} \propto \frac{1}{\cosh ^{2}\left(a\left(y-y^{*}\right)\right)}
$$

the $y$ width scales with $y_{\max }=\ln (\sqrt{ } / \mathrm{m})$

- mass distributions: Gounaris-Sakurai parameterisation

| $\begin{aligned} & \stackrel{\rightharpoonup}{\circ} \\ & \stackrel{y}{\circ} \\ & \stackrel{1}{\mathrm{~N}} \end{aligned}$ | $\eta$ | $\mu^{+} \mu^{-}$ |
| :---: | :---: | :---: |
|  | P | $\mu^{+} \mu^{-}$ |
|  | $\omega$ | $\mu^{+} \mu^{-}$ |
|  | ¢ | $\mu^{+} \mu^{-}$ |



- $\rho$ mass line shape modified to include phase space effects

$$
\frac{\mathrm{d} R(M)}{\mathrm{d} M}=\frac{\alpha^{2} m_{\rho}^{4}}{3(2 \pi)^{4}} \frac{\left(1-\frac{4 m_{\pi}^{2}}{M^{2}}\right)^{3 / 2} \sqrt{1-\frac{4 m_{\mu}^{2}}{M^{2}}}\left(1+\frac{2 m_{\mu}^{2}}{M^{2}}\right)}{\left(M^{2}-m_{\rho}^{2}\right)^{2}+M^{2} \Gamma_{\mathrm{tot}}^{2}}
$$



- Dimuon mass distributions of Dalitz decays:

Kroll-Wada form multiplied with transition form factors (Lepton-G data)

- Branching ratios from PDG04;
we use $\omega \rightarrow$ ee instead of $\omega \rightarrow \mu \mu$ because it is known more accurately ("lepton universality")


## MC Simulation of Hard Processes

## Events generated with Pythia 6.2

Open Charm ( $D \bar{D}$ ): semi-muonic decays of two $D$ mesons

- Generation done with CTEQ6L PDFs
- Branching Ratios taken from PDG04
- Normalisation:
$\sigma(\mathrm{c} \overline{\mathrm{c}}) \sim 20 \mu \mathrm{~b}$ (from a compilation of charm measurements) Linear A-dependence, including nuclear effects on the PDFs


## Drell-Yan (DY)

- Generated with MRS-A Low $Q^{2}$ to obtain events with low masses
- K factor $=1.4$ (to reproduce NA3 data: $\mathrm{p}-\mathrm{Pt}$ at 400 GeV ) $\sigma(\mathrm{DY})_{\mathrm{pp}}=17 \mathrm{nb}, \sigma(\mathrm{DY})_{\mathrm{pn}}=15 \mathrm{nb}$
- Linear scaling with the number of nucleons:
 $\sigma(\mathrm{DY})_{\mathrm{p}-\mathrm{A}}=\mathrm{Z} \cdot \sigma_{\mathrm{pp}}+(\mathrm{A}-\mathrm{Z}) \cdot \sigma_{\mathrm{pn}}$


## Reconstruction:

- All generated muon pairs are immersed in an underlying hadronic event, using VENUS, to correctly reproduce the reconstruction efficiencies.
- Particle tracking through the apparatus done via GEANT.
- The reconstruction was done with the same settings as the real data
$\rightarrow$ gives particle acceptances and detector smearing effects.


## Reconstructed MC vs. Data

- Before extracting physics results from the data using our MC simulations, we must ensure that the data's kinematical distributions are reproduced.
- Among other variables, we compare the rapidity, the decay angle and the transverse momentum distributions of various mass windows, where the comparison is performed on the raw data level.
- Within the statistics available in the $\omega$ and $\phi$ mass windows we see good agreement between reconstructed MC and data.

rapidity:



## Fitting the OS Dimuon Mass Spectrum

- Background fixed by a mixed event technique using single muons from the measured like-sign dimuons.
- Open charm and Drell-Yan production cross-sections fixed from previous measurements (describes nicely the region between the $\phi$ and the $\mathrm{J} / \psi$ peaks).
- The $\omega$ and $\phi$ cross-sections can be extracted from the resonance peaks.
- The good mass resolution allows us to extract the $\rho$ normalisation independently of the $\omega$.
- The $\eta$ cross-section is essentially determined from the mass region below 0.45 GeV , where its Dalitz decay is the dominating process (the $\eta$ 2-body peak does not have enough statistics to influence the fit).
- From a simultaneous fit of the 3 data samples ( $\mathrm{Be}, \mathrm{In}$ and Pb ), we can extract the dependence of the $\eta, \rho / \omega$ and $\phi$ cross-sections with $A$.

$$
\sigma_{p A}=\sigma_{0} \cdot A^{\alpha}
$$

$\sum_{i=B e, I n, P b} \frac{\mathrm{~d} N_{i}^{O S}}{\mathrm{~d} M}=\sum_{i=B e, I n, P b} \frac{\mathrm{~d} N_{i}^{B G}}{\mathrm{~d} M}+\mathcal{L}_{i}\left(\frac{\mathrm{~d} \sigma_{p A_{i}}^{D \bar{D}}}{\mathrm{~d} M}+\right.$
$B^{\eta_{D}} \sigma_{0}^{\eta} \mathrm{A}_{i}^{\alpha^{\eta}} \frac{\mathrm{d} N_{i}^{\eta_{D}}}{\mathrm{~d} M}+B^{\eta_{D}^{\prime}} \sigma_{0}^{\eta^{\prime}} \mathrm{A}_{i}^{\alpha^{\eta^{\prime}}} \frac{\mathrm{d} N_{i}^{\eta_{D}^{\prime}}}{\mathrm{d} M}+B^{\omega_{D}} \sigma_{0}^{\omega} \mathrm{A}_{i}^{\alpha^{\omega}} \frac{\mathrm{d} N_{i}^{\omega_{D}}}{\mathrm{~d} M}+$
$\left.B^{\eta} \sigma_{0}^{\eta} \mathrm{A}_{i}^{\alpha^{\eta}} \frac{\mathrm{d} N_{i}^{\eta}}{\mathrm{d} M}+B^{\rho} \sigma_{0}^{\rho} \mathrm{A}_{i}^{\alpha^{\rho}} \frac{\mathrm{d} N_{i}^{\rho}}{\mathrm{d} M}+B^{\omega} \sigma_{0}^{\omega} \mathrm{A}_{i}^{\alpha^{\omega}} \frac{\mathrm{d} N_{i}^{\omega}}{\mathrm{d} M}+B^{\phi} \sigma_{0}^{\phi} \mathrm{A}_{i}^{\alpha^{\phi}} \frac{\mathrm{d} N_{i}^{\phi}}{\mathrm{d} M}\right)$

- Fit parameters: $\sigma_{0}{ }^{\eta}, \sigma_{0}{ }^{\rho}, \sigma_{0}{ }^{\omega}, \sigma_{0}{ }^{\phi}$,

$$
\alpha^{\eta}, \alpha^{\omega}, \alpha^{\phi}
$$

- Assuming: $\alpha^{\rho}=\alpha^{\omega} ; \alpha^{\eta}=\alpha^{\eta}$ and $\sigma^{\eta^{\prime}}=0.15 \cdot \sigma^{\eta}$ [Eur. Phys. J. C4 (1998) 231]
- The fit is performed in the mass window $0.2-1.1 \mathrm{GeV}$.

Fitting $\mathrm{p}-\mathrm{Be}, \mathrm{p}-\mathrm{In}$ and $\mathrm{p}-\mathrm{Pb}$ simultaneously




- The fitting procedure (7 free parameters) describes the low mass dimuon spectra of the three data samples without additional sources (like in HELIOS-1 and CERES).
- From these fits we can derive the number of $\omega$ 's and $\phi$ 's present in our data samples:
- The $\phi$ peak increases relatively to the $\omega$, from $\mathrm{p}-\mathrm{Be}$ to $\mathrm{p}-\mathrm{Pb}$.

|  | $\mathrm{N}^{\omega}$ | $\mathrm{N}^{\phi}$ |
| :--- | :---: | :---: |
| Be | 966 | 575 |
| $\ln$ | 676 | 464 |
| Pb | 660 | 511 |

## Results I: Nuclear dependence of production cross-sections

The fit gives the nuclear dependence of the $\eta, \omega$ and $\phi$ cross-sections:

$$
\sigma_{p A}=\sigma_{0} \cdot A^{\alpha}
$$

- The $\eta$ and $\phi$ production cross-sections scale faster with A than the $\omega$.

| $\alpha^{\eta}$ | $0.93 \pm 0.02$ |
| :---: | :---: |
| $\alpha^{\omega}$ | $0.82 \pm 0.01$ |
| $\alpha^{\phi}$ | $0.91 \pm 0.02$ |

(statistical error only)
This should be kept in mind when interpreting data collected in heavy ion collisions.

- No previous measurements are worth comparing to, except with HERA-B, which measured


$$
\alpha(\phi)=1.01 \pm 0.01 \pm 0.06
$$

in $p-C, T i, W$ at 920 GeV , in the $\phi \rightarrow \mathrm{K}^{+} \mathrm{K}^{-}$decay channel
Phase space domain:

$$
\begin{aligned}
& 0.5<\mathrm{p}_{\mathrm{T}}^{2}<12.1(\mathrm{GeV} / \mathrm{c})^{2} \\
& 2.95<\mathrm{y}_{\text {lab }}<4.2 \\
& \text { (i.e. }-0.85<\mathrm{y}^{*}<0.4 \text { ) }
\end{aligned}
$$

note that $\alpha$ decreases with $\mathrm{X}_{\mathrm{F}}$ and increases with $\mathrm{p}_{\mathrm{T}}$

## Results II: Elementary $4 \pi$ production cross-sections

We have extracted the absolute production cross-sections of the $\eta, \rho, \omega$ and $\phi$ mesons in elementary p-nucleon collisions at $\mathrm{E}_{\mathrm{lab}}=400 \mathrm{GeV}$.
The extrapolation to $4 \pi$ requires assuming certain kinematical distributions outside of our phase space window. For the decay angle distributions of the 2-body decays we have two reasonable options: $1+\cos ^{2} \theta$ or uniform

Looking at the results:

| $\sigma_{0}[\mathrm{mb}]$ | $1+\cos ^{2} \theta$ | uniform |
| :---: | :---: | :---: |
| $\rho$ | $11.6 \pm 1.0$ | $8.9 \pm 0.7$ |
| $\omega$ | $10.5 \pm 0.6$ | $8.0 \pm 0.5$ |
| $\phi$ | $0.53 \pm 0.05$ | $0.40 \pm 0.03$ |
| $\eta$ | $9.5 \pm 0.6$ | $10.2 \pm 0.6$ |
| (statistical errors only) |  |  |

(a) $\left(\sigma_{0}^{\rho} / \sigma_{0}^{\omega}\right)_{\mu \mu}=1.1 \pm 0.1$
(b) $\left(\sigma_{0}^{\rho}+\sigma_{0}^{\omega}\right)_{\mu \mu}=\left\{\begin{array}{l}22.1 \pm 1.2 \mathrm{mb} \text { for } 1+\cos ^{2} \theta \\ 16.9 \pm 0.9 \mathrm{mb} \text { for uniform }\end{array}\right.$
(c) $\left(\frac{\sigma_{0}^{\eta}}{\sigma_{0}^{\rho}+\sigma_{0}^{\omega}}\right)_{\mu \mu}= \begin{cases}0.43 \pm 0.04 & \text { for } 1+\cos ^{2} \theta \\ 0.60 \pm 0.05 & \text { for uniform }\end{cases}$
(statistical errors only)

How do these results compare to previous measurements?

## NA27





NA27 measured the elementary $\eta, \rho, \omega$ and $\phi$ full phase space production cross-sections in pp @ 400 GeV
[Z. Phys. C50 (1991) 405]
Phase space coverage: $x_{F}>0$

## HELIOS-1



- HELIOS-1 measured di-electron and dimuon spectra in p-Be @ 450 GeV
[Z. Phys. C68 (1995) 47.]
- The $\eta$ yield in the dilepton spectra was fixed from the independent $\mathrm{I}^{+} \mathrm{I} \gamma$ measurement

Comparison to NA60:

- similar phase space coverage:

$$
\begin{array}{rlr}
+0.25<\mathrm{y}^{*}<+1.50 & \mathrm{e}^{+} \mathrm{e}^{-} \\
-0.25<\mathrm{y}^{+}<+1.25 & \mu^{+} \mu^{-} \\
-0.75<\cos \theta<0.75 & \\
\mathrm{~m}_{\mathrm{T}}>0.25 \mathrm{GeV} & \mathrm{e}^{+} \mathrm{e}^{-} \\
\mathrm{m}_{\mathrm{T}}>0.4 \mathrm{GeV} & \mu^{+} \mu^{-}
\end{array}
$$

- same mass resolution
- higher background


## CERES-TAPS



- CERES-TAPS measured $\eta$ and $\omega$ production in $p$-Be and $p-A u$ collisions @ 450 GeV
[Eur. Phys. J. C4(1998) 249].
- Phase space coverage: $3.1<y<3.7$
- Published particle ratios (no absolute cross-sections) in their phase space window.


## Absolute cross-sections: $\left(\sigma_{0}{ }^{\rho}+\sigma_{0}{ }^{\omega}\right), \sigma_{0}{ }^{\eta}$ and $\sigma_{0}{ }^{\phi}$

To compare our $\rho$ and $\omega$ cross-sections with measurements done in independent decay channels, we must take into account the interference effect in our data ( $\rightarrow$ overlapping mass; measurement in the same decay channel). HELIOS-1 found in their analysis a negative interference effect, giving a total $\sigma^{\rho / \omega} 15 \%$ smaller than their sum, measured in independent channels, $\sigma^{\rho}+\sigma^{\omega}$.


| NA60 |  |  |
| :--- | :---: | :---: |
|  |  |  |
|  | $1+\cos ^{2} \theta$ | uniform |
| $\sigma_{0}{ }^{\eta} \quad[\mathrm{mb}]$ | $9.5 \pm 0.6$ | $10.2 \pm 0.6$ |
| $\sigma_{0}{ }^{\rho+\sigma_{0}{ }^{\omega}[\mathrm{mb}]}$ | $25.4 \pm 1.3$ | $19.4 \pm 1.0$ |
| $\sigma_{0}{ }^{\phi}$ | $[\mathrm{mb}]$ | $0.53 \pm 0.05$ |

- The cross-sections $\sigma_{0}{ }^{\rho}+\sigma_{0}{ }^{\omega}$ measured by NA27 and NA60 agree perfectly if the $1+\cos ^{2} \theta$ decay angle distribution is used for NA60's extrapolation to $4 \pi$.
- The $\eta$ cross-sections of NA27 and NA60 also agree very well with $1+\cos ^{2} \theta$.
- The $\phi$ cross-section measured by NA60 is lower than NA27's, but maybe the NA27 value is slightly overestimated.
- The comparisons of absolute cross-sections indicate that the $1+\cos ^{2} \theta$ decay angle distribution is the more appropriate one.



## The $\eta /(\rho+\omega)$ cross-section ratio in $p$-nucleon collisions

In order to compare apples with apples:

1. Correct the measurements in leptonic decay channels (HELIOS-1 and NA60) for the $\rho / \omega$ interference, and use the results obtained with $1+\cos ^{2} \theta$.
2. Extrapolate the HELIOS-1 and CERES-TAPS measurements to elementary p-nucleon collisions (using our $\alpha^{\eta}$ and $\alpha^{\rho}=\alpha^{\omega}$ ).
3. Extrapolate the CERES-TAPS measurement to full phase space.

Note: CERES-TAPS assumed $\sigma^{\rho}=\sigma^{\omega}$ for the calculation of $\eta /(\rho+\omega)$.
$\rightarrow$ The NA60 value agrees with these previous measurements.

The open (closed) symbols show the measurements before (after) extrapolating to $4 \pi$ and $p$-nucleon collisions

## Summary and Outlook

- Although the proton 2002 run had limited statistics, we achieved
- a good mass resolution ( $\sim 30 \mathrm{MeV}$ at 1 GeV )
- a good signal-to-background ratio
- This performance allowed us to
- clearly separate the $\omega$ and $\phi$ peaks and
- estimate the $\rho$ normalisation independently of the $\omega$
- Having 3 target materials with very different mass numbers, we extracted the nuclear dependence of the production cross-sections for the $\eta, \omega$ and $\phi$ mesons.
- The extracted elementary proton-nucleon $4 \pi$ cross-sections for the $\eta, \rho, \omega$ and $\phi$ mesons are in good agreement with existing measurements (NA27, HELIOS-1 and CERES)
- The observed faster scaling of the $\eta$ and $\phi$ mesons with respect to the $\omega$ should be taken into account when interpreting the heavy-ion data.


## Outlook:

NA60 is currently collecting a large data sample with a proton beam at 400 GeV incident on 7 nuclear targets ( $\mathrm{Be}, \mathrm{Al}, \mathrm{Cu}, \mathrm{In}, \mathrm{W}, \mathrm{Pb}, \mathrm{U}$ ) to collect further reference data.

