# $\Lambda$ -Based Perturbation Theory for Event Shapes in $e^+$ - $e^-$ Annihilation

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#### Talk Plan

• What is  $\Lambda$ -Based Perturbation Theory?

a.k.a. RESIPE (Dhar, Gupta), Method of Effective Charges (Grunberg), RGI perturbation theory...

Event Shapes

taking thrust as an example

• Previous Applications of  $\Lambda$ -Based PT to Event Shapes

means by the DELPHI collaboration, distributions by Burby and Maxwell

Our Analysis

including resummation of logs and power correction fits

Results

fits for  $\Lambda_{\overline{MS}}$  and  $C_1$ 

Work done in collaboration with C. J. Maxwell

# What is $\Lambda$ -based Perturbation Theory?

#### Standard Approach to Perturbation Theory

• Consider a perturbative QCD expansion of an observable normalized as an effective charge (with  $a = \alpha/\pi$ ):

$$R(Q) = a_S(\mu) + r_1(\mu/Q, S) * a_S^2(\mu) + r_2(\mu/Q, S) * a_S^3(\mu) + \cdots$$

- R(Q) is independent of  $\mu$  and S, but only because of cancellation between different orders. So, all finite order predictions depend on  $\mu$  and S.
- Normally have NLO calculation  $\rightarrow \mu$  dependence at  $O(a_S^3)$ .
- Standard fix: choose  $\mu=Q$ , vary to get theoretical error.
- BUT the definition of  $\mu$  depends on S! Specifically our results depend only on the ratio  $\mu/\Lambda_S$ .

# What is $\Lambda$ -based Perturbation Theory?

#### $\Lambda$ -Based Approach to Perturbation Theory

- We take the dimensional transmutation parameter  $\Lambda$ , defined in some scheme (e.g.  $\overline{MS}$ ) as the fundamental parameter to fit to data.
- For each effective charge R we write

$$\frac{dR(Q)}{d\ln Q} = \rho_R(R) = -bR^2(1 + cR + \rho_2^R R^2 + \cdots)$$

where  $\rho$  is the effective charge  $\beta$ -function.

Compare with

$$\frac{da_S(Q)}{d\ln Q} = \beta_S(a) = -ba_S^2(1 + ca_S + c_2^S a_S^2 + \cdots)$$

the  $\beta$ -function for the scheme S.

# What is $\Lambda$ -based Perturbation Theory?

#### $\Lambda$ -Based Approach to Perturbation Theory

•  $\Lambda$  enters as a boundary condition. We can integrate up to find the relation

$$\Lambda_S e^{-r_1(\mu=Q,S)/b} = Q\mathcal{F}(R(Q))\mathcal{G}(R(Q)) \left(\frac{2c}{b}\right)^{c/b}$$

where the LHS and RHS are scheme independent. Explicitly:

$$\mathcal{F}(R) = e^{-1/bR} (1 + 1/cR)^{c/b}$$

$$\mathcal{G}(R) = \exp\left[-\int_0^R dx \left(\frac{1}{\rho_R(x)} + \frac{1}{bx^2(1+cx)}\right)\right]$$

• At NLO  $\rho_R(x) = -bx^2(1+cx)$  so  $\mathcal{G} = 1$ .

### **Event Shapes**

• Eg. thrust

$$\tau \equiv 1 - T \equiv 1 - \max_{\vec{n}} \frac{\sum_{i} |\vec{p_i}.\vec{n}|}{\sum_{i} |\vec{p}|}$$

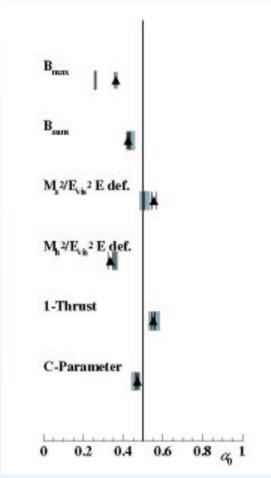
 Standard PT on its own doesn't describe <1-T> very well. Instead we fit for an additional "power correction" giving something like:

$$<\tau> = <1-T> = 0.335\alpha_{\overline{MS}}(Q) + 1.02\alpha_{\overline{MS}}^{2}(Q) + \frac{C_1}{Q}$$

where  $C_1 \sim 1 GeV$ .

- More sophisticated approach: "Dokshitzer-Webber ansatz".
  Relate the power correction to a hypothetical infrared finite coupling → relate power corrections to different observables!
- Gives fairly successful description of data "approximate universality" for new parameter  $\overline{\alpha}_0$ .

# $\Lambda$ -based PT for Event Shape Means @ DELPHI

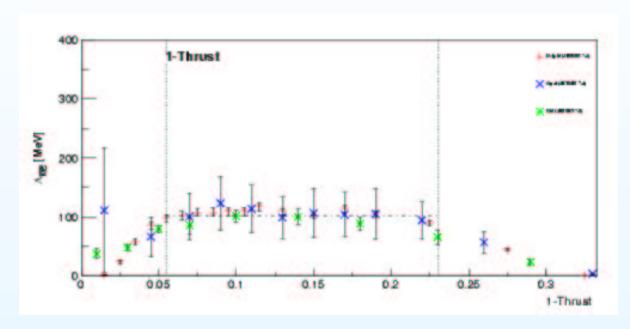


- The DELPHI collaboration have found that Λ-based PT describes the event shape means well without any power corrections.
- This plot shows the effective  $\overline{\alpha}_0$  induced by converting the  $\Lambda$ -based results to the  $\overline{MS}$  scheme.
- Remarkable agreement with experiment compared to universality hypothesis!

A study of the energy evolution of event shape distributions and their means with the DELPHI detector at LEP, DELPHI Collaboration, J. Abdallah, et al, Eur.Phys.J. C29 (2003) 285-312

# $\Lambda$ -based PT for Event Shape Distributions

Also interesting to look at distributions. At NLO:



- Problems:
  - $\circ$  Large logs  $L = \log(1/\tau)$  appear in the 2-jet region  $\tau \to 0$
  - $^{\circ}$  The EC description breaks down as au o 1/3

Direct Extraction of QCD Lambda MS-bar from e+e- Jet Observables, S. J. Burby, C. J. Maxwell, Nucl. Phys. B609 (2001) 193-224

# Our Analysis

We wanted to build on this analysis in 2 ways:

• Large logs in the perturbative series for the distribution give large logs in the perturbative series for  $\rho \rightarrow$  need for resummation:

$$\rho(R, L) = -bR^{2}(\rho_{LL}(RL) + L^{-1}\rho_{NLL}(RL) + \cdots).$$

• To see if the  $\Lambda$ -based approach needs smaller power corrections we wanted to fit for 1/Q shifts:

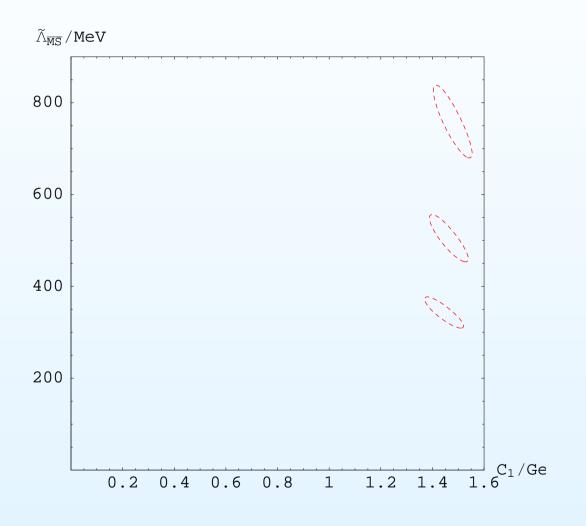
$$\left. \frac{1}{\sigma} \frac{d\sigma}{d\tau} \right|_{\tau} = \left. \frac{1}{\sigma} \frac{d\sigma}{d\tau} \right|_{PT, \ \tau - C_1/Q}$$

Define an effective charge via

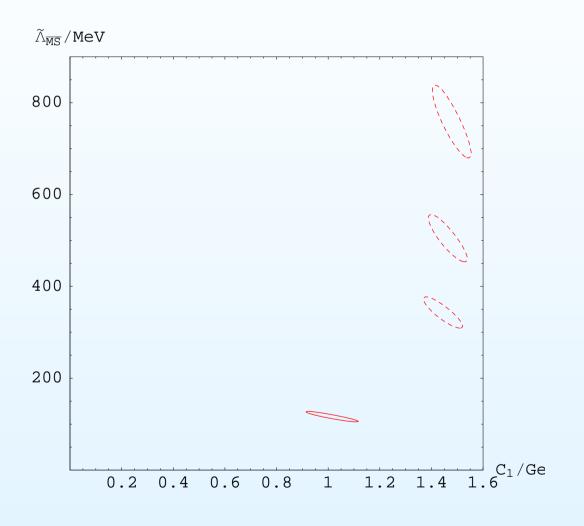
$$R(Q,\tau) = \ln\left(\int_0^{\tau} \frac{1}{\sigma} \frac{d\sigma}{d\tau}\right) / LO$$

# **Our Analysis**

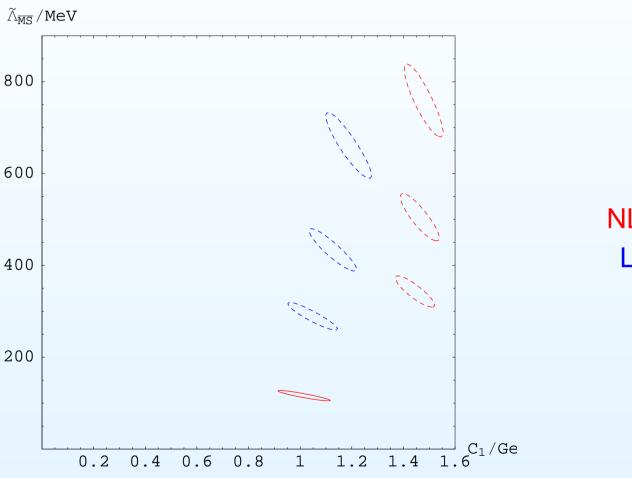
- We need to select a limited fit range to avoid problems in the extreme 2-jet region and as  $\tau \to 1/3$  (made worse by exponentiating).
- Found good fit quality with  $\tau = 0.04 0.16$  at Q = 91.2 GeV with the lower bound scaling  $\propto 1/Q$ .
- Take data from TASSO, JADE, DELPHI, L3, SLD at energies Q=35-189 GeV.
- Note that the "standard approach" in the  $\overline{MS}$  scheme we can fit over a large range of  $\tau!$

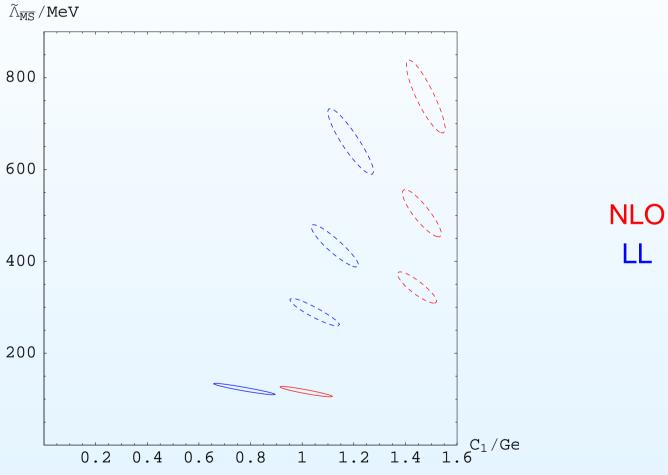


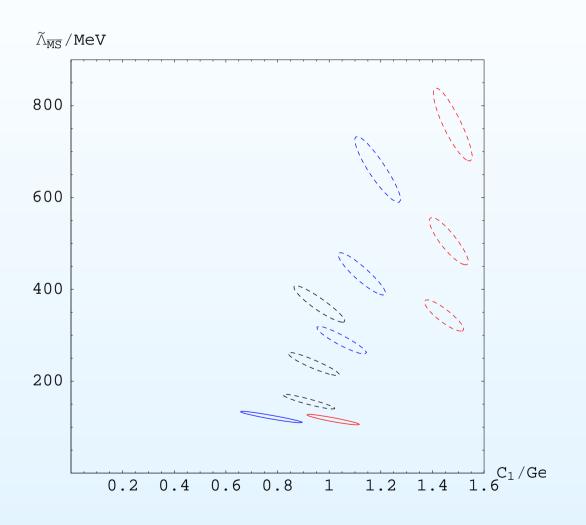
NLO



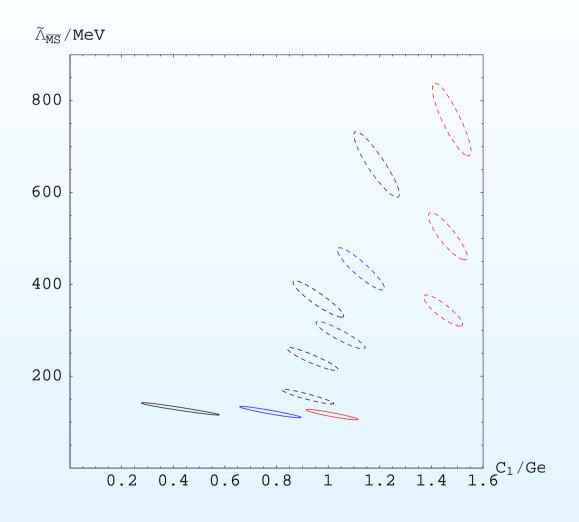
NLO







NLO LL NLL



NLO LL NLL

#### Conclusions

- $\Lambda$ -based PT provides us with a way to make perturbative predictions without needing to fix  $\mu$  by some kinematical argument.
- Analysis of event shape means has suggested this may give more accurate predictions than the standard approach, with significantly reduced power corrections needed to fit the data.
- We have applied this approach to some event shape distributions, and after significantly restricting the fit range we find good agreement with reduced (though still significant) power corrections.
- Adding resummed logs further reduces the fitted power correction.