Simulating jet and W/Z+jets production with SHERPA

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- The SHERPA approach
- W^{\pm} and Z production @ Tevatron
- jet production @ Tevatron
 - Consistency checks
 - SHERPA vs. NLO/other MC vs. DATA

Combine LO Matrix Elements and Parton Showers according to CKKW

S. Catani, F. Krauss, R. Kuhn, B. Webber, JHEP 0111:063,2001 F. Krauss, JHEP 0208:015,2002

Motivation:

- ME are exact at some given order in the coupling constant and good for description of exclusive quantities (e.g. jet cross sections)
- PS resums $\alpha_S^n log^{2n}(s/Q_0^2)$ to all orders and is good for inclusive quantities

Aim:

- Combine strengths of both approaches
- Good description of soft and hard region
- Avoid double counting of equivalent phase space configurations
- Universality of fragmentation (energy independent)

The SHERPA approach

Solution:

- Divide multi-jet phase space into two regimes (Durham measure Q_{cut})
 - Jet production by ME (if available)
 - Jet evolution down to fragmentation scale by the PS
- Reweight ME's to get exclusive samples at a resolution scale $Q_{\rm cut}$
 - \Rightarrow This allows to add samples of different jet multiplicities
- Veto on PS configurations that have already been taken into account by a higher order ME

Method:

Select a jet multiplicity with probability:

$$P_n = \frac{\sigma_n}{\sum_{i=0}^N \sigma_i}$$

where σ_n is the *n*-jet matrix element taken at resolution scale Q_{cut} . Use Q_{cut} as scale for PDF's and α_S .

• Generate final state momenta p_i according to the ME

The SHERPA approach

Method:

- k_T cluster backwards initial and final state particles until a core $2 \rightarrow 2$ process remains, this results in a chain of resolutions for 1,2,...n jets
- Recalculate α_S at each vertex in the tree at the corresponding k_T scale
- Apply Sudakov weights
 - $\Delta_{q,g}(Q_{\mathrm{cut}},Q_{\mathrm{prod}})$ for outgoing partons
 - $\Delta_{q,g}(Q_{\mathrm{cut}},Q_{\mathrm{prod}})/\Delta_{q,g}(Q_{\mathrm{cut}},Q_{\mathrm{dec}})$ for lines between $Q_{\mathrm{prod}} > Q_{\mathrm{dec}}$
 - Reject events with a combined coupling and Sudakov weight smaller than random number $R \in [0, 1]$
 - Start the initial or final state parton shower for each parton of the event, starting at the scale where it was produced
 - Veto on emissions above the scale $Q_{\rm cut}$



The SHERPA approach

SHERPA specifics:

• Jet measure: $Q_{ij}^2 = \min(p_{\perp i}^2, p_{\perp j}^2) \cdot R_{ij}^2$ or $Q_{iB}^2 = p_{\perp i}^2$

$$R_{ij}^2 = 2 \left[\cosh(\eta_i - \eta_j) - \cos(\phi_i - \phi_j) \right]$$

• For the highest multiplicity ME the scale $Q_{\rm cut}$ in the PDF's and Sudakovs is replaced by the smallest nodal scale of the clustering

Now some applications...

W/Z +jets production @ Tevatron: consitency checks

F. Krauss, A. Schälicke, S. S. and G. Soff, Phys. Rev. D 70 (2004) 114009 The p_{\perp} distribution of the W^- in $pp \rightarrow e^- \bar{\nu_e} + X$ @ Tevatron Run II



$p_{\perp Z}$ distribution

 $Q_{\rm cut}$ =15 GeV

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SHERPA vs. NLO: Incl. W^- +jet prod. @ Tevatron Run II



SHERPA vs. PYTHIA and MC@NLO

The p_{\perp} of the three hardest jets in incl. W production @ Tevatron, Run II



- Inclusive quantities (e.g. p_{\perp}, η distributions of the W/Z) agree quite well
- More exclusive quantities (e.g. p_{\perp} distributions of the jets) show differences, increasing with the order of jet

Inclusive W and Z transverse momentum distributions from Tevatron Run I

- MEs with up to four (W) or three (Z) extra jets
- momenta smaller than merging scale are sensitive to the details of the PS





Specific for pure QCD: Final state parton shower start scale is

$$Q^2 = \frac{2stu}{s^2 + t^2 + u^2}$$

- For a given Q_{cut} the phase space for events with less than two jets above the cut is not filled.
 - \rightarrow can be compensated by a special two scale treatment for the hard $2 \rightarrow 2$ process.
- Sherpa can take into accout up to 4-jet production matrix elements

$p\bar{p} \rightarrow jets @ \sqrt{s} = 1.8 \text{ TeV}$: consistency checks



Effect of varying jet separation: differential jetrates

 $Q_{cut}=20~{
m GeV}$

 $Q_{cut}=30~{
m GeV}$

 $Q_{cut} = 40 \text{ GeV}$

Inclusive jet production: $p\bar{p} \rightarrow jets @ \sqrt{s} = 1.8$ TeV: consistent



 Ratio of inclusive 3-jet to inclusive 2-jet cross section

$$R_{32} = \frac{\sigma_3}{\sigma_2} = \frac{\sigma(p\bar{p} \to 2jets + X)}{\sigma(p\bar{p} \to 3jets + X)}$$

- Tevatron, Run I measurement:
 D0: Phys. Rev. Lett. 86, 1955 (2001)
- $E_{T,jet} > 40 \text{ GeV}, |\eta_{jet}| < 3$
- Jets analysed using a midpointcone algorithm (R=0.7)



Impact of Q_{cut} / N_{max} variations on Sherpa predictions:



Effects of hadronization and Jet-algorithm







Variation of the Factorization and Renormalization scales:



Comparison to data: Azimutal Dijet Decorrelations

 $\hfill \bullet$ Measurement of the $\Phi\mbox{-angle}$ between the two hardest jets of an inclusive sample



• Different regions of $\Delta \phi$ are sensitive to different multiplicities of hard QCD jet production

• Measured at Tevatron ($\sqrt{s} = 1.96$ TeV): DØ hep-ex/0409040

• $p_T^{jet} > 40 GeV$, $|y^{jet}| < 0.5$

Comparison to data: Azimutal Dijet Decorrelations





 $Q_{cut} = 40 GeV$ $N_{max} = 4$

Conclusion/Outlook

Conclusion

- The implemented CKKW-merging works quite well for W/Z +jets and pure jet production Also validated:
 - W/Z @ LHC: F. Krauss, A. Schälicke, S. Schumann and G. Soff, hep-ph/0503280
 - W⁺W⁻ +jets T. G., F. Krauss, A. Schälicke, S. Schumann and J. Winter, hep-ph/0504032
- further validation is in progress
- SHERPA is able to reproduce the shapes of corresponding NLO calculations

SHERPA sources

- SHERPA: T. G., S. Höche, F. Krauss, A. Schälicke, S. Schumann and J. Winter, JHEP 0402:056,2004
- current version SHERPA α -1.0.6 is available under

http://www.physik.tu-dresden.de/~krauss/hep