

# Probing Flavor Structure in Supersymmetric Theories

Shaaban Khalil

Ain Shams University, Cairo, Egypt

# I. Introduction

- The recent results on the mixing-induced asymmetries of  $B \rightarrow \varphi K$  and  $B \rightarrow \eta' K$  are:

**Belle**

$$\begin{cases} S_{\varphi K} = 0.44 \pm 0.27 \pm 0.05 \\ S_{\eta' K} = 0.62 \pm 0.12 \pm 0.04 \end{cases}$$

**BaBar**

$$\begin{cases} S_{\varphi K} = 0.50 \pm 0.25^{+0.07}_{-0.04} \\ S_{\eta' K} = 0.30 \pm 0.14 \pm 0.02 \end{cases}$$

- The direct CP violation in  $B^0 \rightarrow K^- \pi^+$  and  $B^- \rightarrow K^- \pi^0$ :

$$A_{K^- \pi^+}^{\text{CP}} = -0.113 \pm 0.019 \quad 4.2 \sigma \text{ deviation from zero}$$

$$A_{K^- \pi^0}^{\text{CP}} = 0.04 \pm 0.04. \quad \text{is quite small}$$

- These observations are considered as signals to new physics.

- To accommodate the CP asymmetries of B decays, SUSY models with flavor non-universal soft breaking terms are favored.
- The squark mixings are classified as
  - i) LL and RR mixings given by  $(\delta^{u,d}_{LL})_{ij}$  and  $(\delta^{u,d}_{RR})_{ij}$ .
  - ii) LR and RL mixings given by  $(\delta^{u,d}_{LR})_{ij}$  and  $(\delta^{u,d}_{RL})_{ij}$ .
- Constraints are more stringent on the LR (RL) mass insertions than the LL (RR) mass insertions.
- MIs between 1st & 2nd generations are severely constrained more than MIs between 1st or 2nd & 3rd generations.
- This gives the hope that SUSY contributions to the B-system.

## II. Squark Mixing: LL versus LR

- The EDM constraints severely restrict the LL and RR mass insertions:

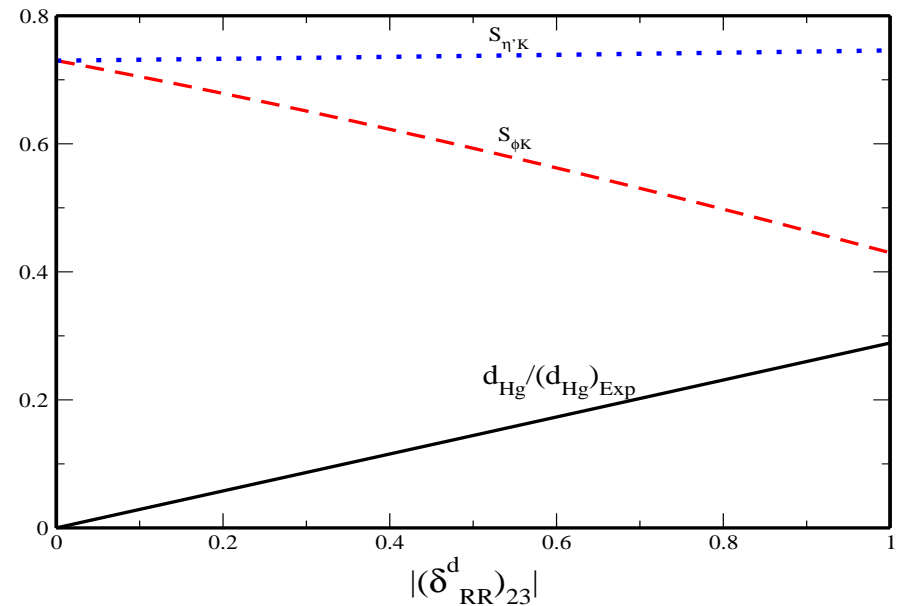
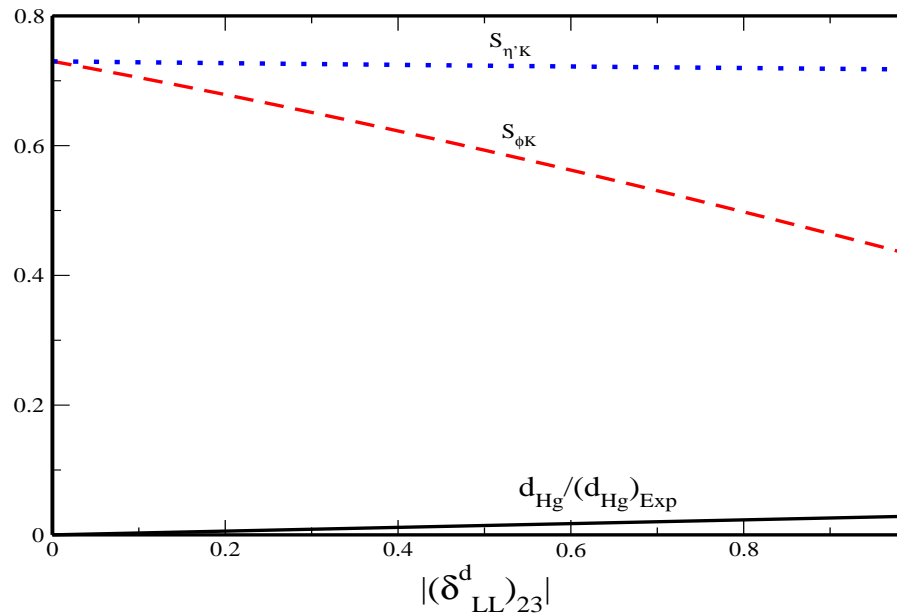
$$\begin{aligned}(\delta_{LR}^d)_{22\text{eff}} &\approx (\delta_{LR}^d)_{22} + (\delta_{LL}^d)_{23} (\delta_{LR}^d)_{33} (\delta_{RR}^d)_{32} \\ &\approx 10^{-2} (\delta_{LL}^d)_{23} (\delta_{RR}^d)_{32} \quad (\text{for neg. } (\delta_{LR}^d)_{22}).\end{aligned}$$

- From Hg EDM:  $\text{Im}(\delta_{LR}^d)_{22} < 5.6 \times 10^{-6}$ .
- The SUSY contributions to the decay amplitudes  $B \rightarrow \varphi K$  and  $\eta' K$ :

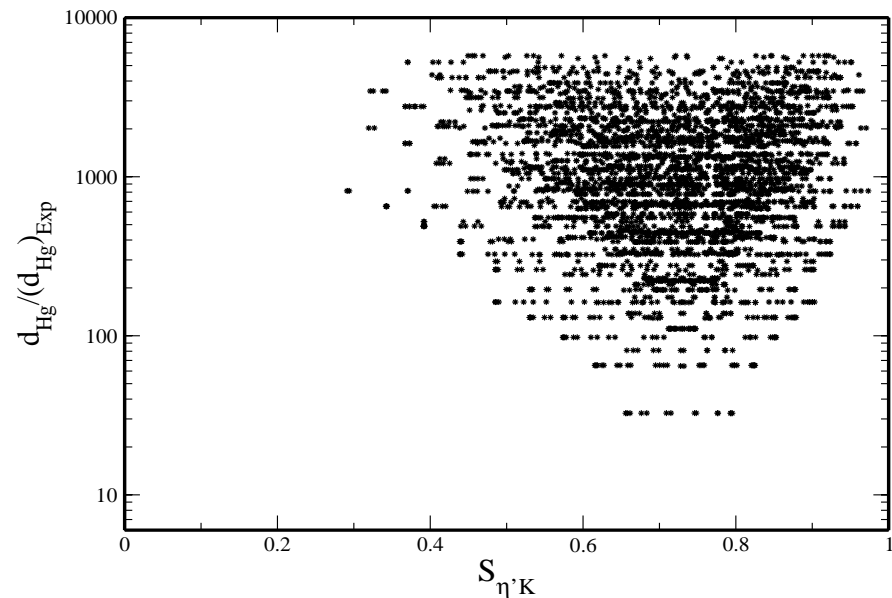
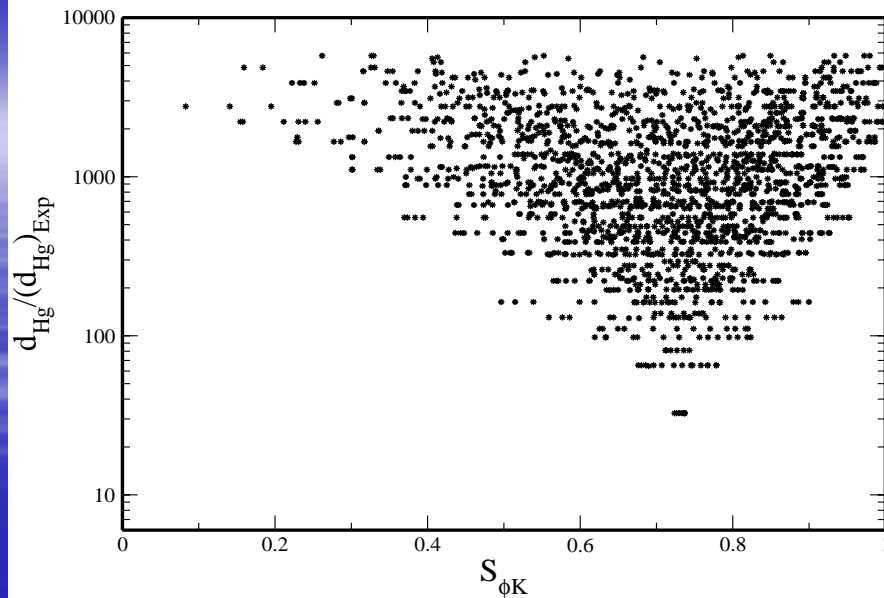
$$R_\varphi = -0.14 e^{-i0.1} (\delta_{LL}^d)_{23} - 127 e^{-i0.08} (\delta_{LR}^d)_{23} + (L \leftrightarrow R),$$

$$R_{\eta'} = -0.07 e^{-i0.24} (\delta_{LL}^d)_{23} - 64 (\delta_{LR}^d)_{23} - (L \leftrightarrow R).$$

- The CP asymmetries  $S_{\phi K}$  and  $S_{\eta K}$  can not be accommodated



- Combining the effects of LL and RR MIs, we can fit the exp. data of  $S_{\phi K}$  and  $S_{\eta K}$ .
- However, the Hg EDM exceeds with many order of magnitudes its exp. Bound.



- Thus, SUSY models with dominant LR and RL may be the most favorite scenario.
- However, it is difficult to arrange for  $(\delta_{LR}^d)_{23} \approx O(10^{-2})$  whilst  $(\delta_{LR}^d)_{12}$  remains small:

From  $\Delta M_K$  and  $\epsilon/\epsilon$ :  $\text{Re}(\delta_{LR}^d)_{12} < O(10^{-4})$  &  
 $\text{Im}(\delta_{LR}^d)_{12} < O(10^{-5})$ .



- The MI  $(\delta_{LR}^d)_{ij}$  are given by

$$(\delta_{LR}^d)_{ij} \approx [V_L^{d+} \cdot (Y^d A^d) \cdot V_R^d]_{ij} .$$

- The factorizable A-term is an example of a specific texture to satisfy this hierarchy between  $(\delta_{LR}^d)_{12}$  and  $(\delta_{LR}^d)_{23}$  .

- With intermediate/large  $\tan\beta$ , an effective  $(\delta_{LR}^d)_{23}$  can be obtained:  $(\delta_{LR}^d)_{23 \text{ eff}} = (\delta_{LR}^d)_{23} + (\delta_{LL}^d)_{23} (\delta_{LR}^d)_{33}$

- For negligible  $(\delta_{LR}^d)_{23}$  , we find

$$(\delta_{LR}^d)_{23 \text{ eff}} \approx (\delta_{LL}^d)_{23} \frac{m_b}{\tilde{m}} \tan \beta$$

- Thus if  $(\delta_{LL}^d)_{23} \approx 10^{-2}$ , we get  $(\delta_{LR}^d)_{23\text{eff}} \approx O(10^{-2} - 10^{-3})$ .
- These contributions are considered as LL (or RR).
- The main effect is still due to the Wilson coefficient  $C_{8g}$  of the chromomagnetic operator.
- Although  $(\delta_{LL}^d)_{23} \approx 10^{-2}$  is not enough to explain the CP asymmetries of B-decays.
- Still it can induce an effective LR mixing that accounts for these results.



# Suggested supersymmetric flavor model

- As an example, we consider the following SUSY model:

$$M_1 = M_2 = M_3 = M_{1/2}$$

$$A^u = A^d = A_0$$

$$M_U^2 = M_D^2 = m_0^2$$

$$m_{H1}^2 = m_{H2}^2 = m_0^2$$

- The masses of the squark doublets are given by

$$M_Q^2 = \begin{pmatrix} m_0^2 & & \\ & m_0^2 & \\ & & a^2 m_0^2 \end{pmatrix}$$

$$\tan \beta = 15, m_0 = M_{1/2} = A_0 = 250 \rightarrow a \leq 5$$

- The Yukawa textures play an important role in the CP and flavour supersymmetric results.

$$Y^u = 1/v \sin\beta \text{diag}(m_u, m_c, m_t)$$

$$Y^d = 1/v \sin\beta V_{CKM}^+ \cdot \text{diag}(m_d, m_s, m_b) \cdot V_{CKM}$$

- Although, it is hierarchical texture, it leads to a good mixing between the second and third generations.
- The LL down MIs are given by:

$$(\delta_{LL}^d)_{ij} = 1/m_q^2 [V_L^{d+} (M^d)^2_{LL} V_L^d]_{ij}$$

With  $a=5 \rightarrow (\delta_{LL}^d)_{23} \sim 0.08 e^{0.4i}$

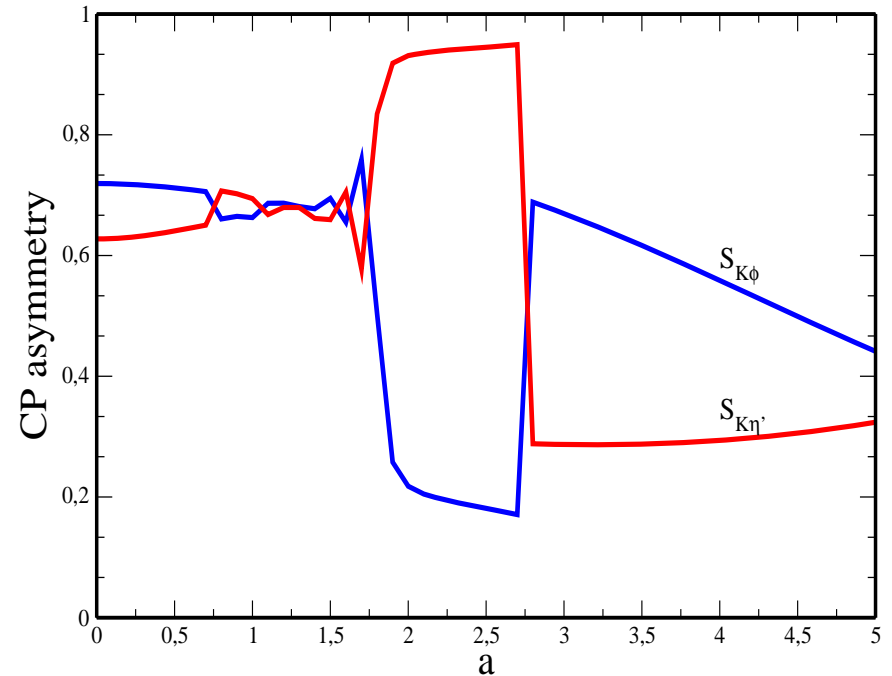
- Thus, one gets:  $(\delta_{LR}^d)_{23\text{eff}} \sim O(10^{-2}-10^{-3})$ .
- The corresponding single LR MI is negligible due to the degeneracy of the A-terms.

# Contribution to $S_{\phi K}$ and $S_{\eta K}$

- Gluino exchanges give the dominant contribution to the CP asymmetries:  $S_{\phi K}$  and  $S_{\eta K}$ .

$$A(B \rightarrow \phi K) \sim -i \frac{G_F}{\sqrt{2}} m_B^2 F_+^{B \rightarrow K} f_\phi H_{8g} (C_{8g} + \tilde{C}_{8g})$$

$$A(B \rightarrow \eta K) \sim -i \frac{G_F}{\sqrt{2}} m_B^2 F_+^{B \rightarrow K} f_\eta H_{8g} (C_{8g} - \tilde{C}_{8g})$$



- 1  $\sigma$  constraints on  $S_{\eta K}$  leads to a lower bound on  $a$ :  $a \geq 3$ .
- With a large  $a$ , it is quite possible to account simultaneously for the experimental results  $S_{\phi K}$  and  $S_{\eta K}$ .

# Contribution to $B \rightarrow K \pi$

- The direct CP asymmetries of  $B \rightarrow K \pi$  are given by:

$$A_{K\pi^+}^{\text{CP}} \sim 2 r_T \sin \delta_T \sin(\theta_P + \gamma) + 2 r_{EW}^C \sin \delta_{EW}^C \sin(\theta_P - \theta_{EW}^C),$$
$$A_{K\pi^0}^{\text{CP}} \sim 2 r_T \sin \delta_T \sin(\theta_P + \gamma) - 2 r_{EW} \sin \delta_{EW} \sin(\theta_P - \theta_{EW}).$$

- The parameters  $\theta_P$ ,  $\theta_{EW}^C$ ,  $\theta_{EW}$  and  $\delta_T$ ,  $\delta_{EW}$ ,  $\delta_{EW}^C$  are the CP violating and CP conserving phases respectively.
- The parameter  $r_T$  measures the relative size of the tree and QCD penguin contributions.
- The parameter,  $r_{EW}$ ,  $r_{EW}^C$  measure the relative size of the electroweak and QCD contributions.

$$P e^{\theta_P} = P^{sm} (1 + k e^{\theta'_P}),$$

$$r_{EW} e^{\delta_{EW}} e^{\theta_{EW}} = r_{EW}^{sm} e^{\delta_{EW}} (1 + l e^{\theta'_{EW}}), r_{EW}^C e^{\delta_{EW}^C} e^{\theta_{EW}^C} = (r_{EW}^C)^{sm} e^{\delta_{EW}^C} (1 + m e^{\theta'_{EW}^C}),$$

$$r_T e^{\delta_T} = (r_T e^{\delta_T})^{sm} / |1 + k e^{\theta'_P}|$$

- $k, l, m$  are given by ( $m_g = m_q = 500, M_2 = 200, \mu = 400$ ):

$$k e^{\theta'_P} = -0.0019 \tan \beta (\delta_{LL}^u)_{32} - 35.0 (\delta_{LR}^d)_{23} + 0.061 (\delta_{LR}^u)_{32}$$

$$l e^{\theta'_{EW}} = 0.0528 \tan \beta (\delta_{LL}^u)_{32} - 2.78 (\delta_{LR}^d)_{23} + 1.11 (\delta_{LR}^u)_{32}$$

$$m e^{\theta'_{EW}^C} = 0.134 \tan \beta (\delta_{LL}^u)_{32} + 26.4 (\delta_{LR}^d)_{23} + 1.62 (\delta_{LR}^u)_{32}$$

$$a=5, m_0 = M_{1/2} = A_0 = 250 \text{ GeV} \rightarrow (\delta_{LR}^d)_{23} \sim 0.006 x e^{-2.7i}$$

Thus:  $k \sim 0.2, l \sim 0.009, m \sim 0.16 \rightarrow r_{EW} \sim 0.13, r_{EW}^C \sim 0.012,$   
 $r_T \sim 0.16 \rightarrow A_{K^-\pi^+}^{CP} \sim -0.113$  and  $A_{K^-\pi^0}^{CP} < A_{K^-\pi^+}^{CP}$



# Conclusions

- We studied the possibility of probing SUSY flavor structure using K, B CP asymmetries constraints & EDM.
- One possibility: large LR and/or RL sector.
- Second possibility: large LL combined with a very small RR and also intermediate or large  $\tan \beta$ .
- Large LR requires a specific pattern for A terms.
- Large LL seems quite natural and can be obtained by a non-universality between the squark masses.
- As an example, a SUSY model with a non-universal left squark masses is considered.
- one gets effective  $(\delta_{LR}^d)_{23}$  that leads to a significant SUSY contributions to the CP asymmetries of B decays.