

# *Lepton Flavour & Number Violation @ LHC*

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- R-parity conservation
  - Sources flavour violation in supersymmetric models
  - Lepton flavour violating decays of supersymmetric particles; Implications for LHC observables
- R-parity violation
  - $\nu$  physics
  - predictions for LSP decays

## Experimental Information

Large mixing angles in neutrino sector

$$\begin{aligned} |\tan \theta_{atm}|^2 &\simeq 1 \\ |\tan \theta_{sol}|^2 &\simeq 0.4 \\ |U_{e3}|^2 &\lesssim 0.05 \end{aligned}$$

Small flavour and CP violation violation in charged lepton sector

$$\begin{aligned} BR(\mu \rightarrow e\gamma) &\lesssim 1.2 \cdot 10^{-11} & BR(\mu^- \rightarrow e^- e^+ e^-) &\lesssim 10^{-12} \\ BR(\tau \rightarrow e\gamma) &\lesssim 1.1 \cdot 10^{-7} & BR(\tau \rightarrow \mu\gamma) &\lesssim 6.8 \cdot 10^{-8} \\ BR(\tau \rightarrow ll') &\lesssim O(10^{-7}) \quad (l, l' = e, \mu) \\ |d_e| &\lesssim 10^{-27} \text{ e cm}, \quad |d_\mu| \lesssim 1.5 \cdot 10^{-18} \text{ e cm}, \quad |d_\tau| \lesssim 1.5 \cdot 10^{-16} \text{ e cm} \end{aligned}$$

possible SUSY contributions to magnetic moments of leptons

$$|\Delta a_e| \leq 10^{-12}, \quad 0 \leq \Delta a_\mu \leq 43 \cdot 10^{-10}, \quad |\Delta a_\tau| \leq 0.058$$

# Sources of Flavour Violation

Sleptons:

$$M_{\tilde{l}}^2 = \begin{pmatrix} M_{L,ij}^2 + \frac{v_d^2 Y_{ki}^{E*} Y_{kj}^E}{2} + D_L \delta_{ij} & \frac{v_d A_{ij} - \mu v_u (Y_{ij}^E)^*}{\sqrt{2}} \\ \frac{v_d A_{ij}^* - \mu v_u Y_{ij}^E}{\sqrt{2}} & M_{E,ij}^2 + \frac{v_d^2 Y_{ik}^E Y_{jk}^{E*}}{2} + D_R \delta_{ij} \end{pmatrix}$$

Sneutrinos:

$$M_{\tilde{\nu}}^2 = M_{L,ij}^2 + D_\nu \delta_{ij}$$

where

$$D_L = \frac{(g'^2 - g^2)(v_d^2 - v_u^2)}{8}, \quad D_R = \frac{g'^2(v_d^2 - v_u^2)}{4}$$

$$D_\nu = \frac{(g^2 + g'^2)(v_d^2 - v_u^2)}{8}$$

Without loss of generality:  $Y_{ij}^E = Y_i^E \delta_{ij}$ ,  $Y_i^E$  real

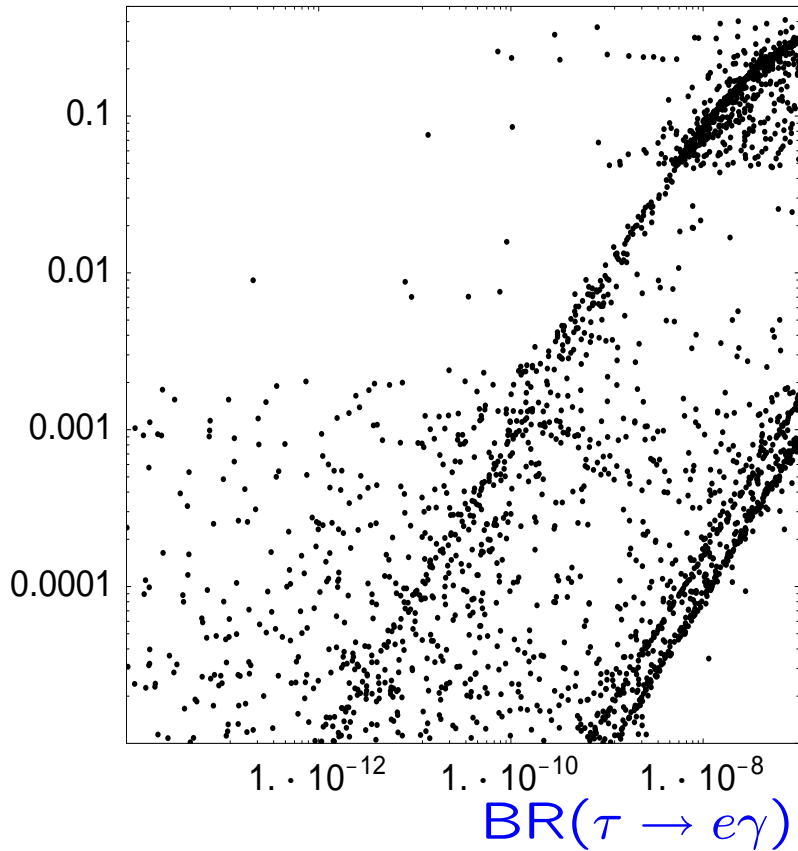
# Lepton flavour violating SUSY decays

Lepton flavour violating couplings:

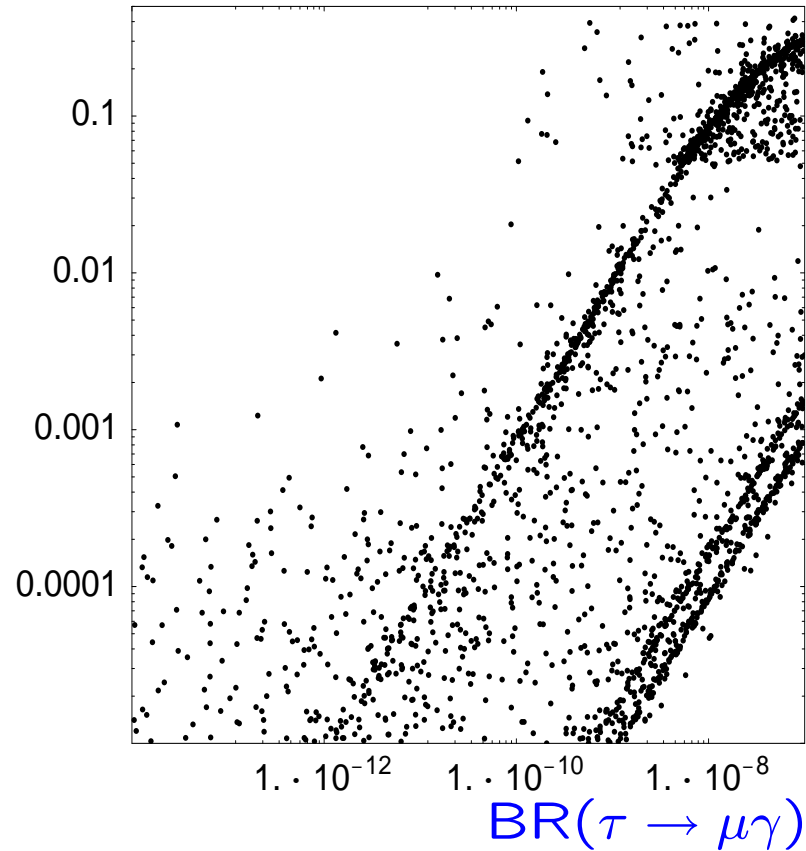
- $\tilde{l}_i - l_j - \tilde{\chi}_k^0 \Rightarrow \tilde{l}_i \rightarrow l_j \tilde{\chi}_k^0, \tilde{\chi}_k^0 \rightarrow l_j \tilde{l}_i$
- $\tilde{\nu}_i - l_j - \tilde{\chi}_k^+$
- $\tilde{l}_i - \tilde{\nu}_j^\dagger - W, \tilde{l}_i - \tilde{\nu}_j^\dagger - H^+$
- $\tilde{l}_i - \tilde{l}_j^\dagger - Z, \tilde{l}_i - \tilde{l}_j^\dagger - (h^0, H^0, A^0)$
- $\tilde{\nu}_i - \nu_j - \tilde{\chi}_k^0, \tilde{l}_i - \nu_j - \tilde{\chi}_k^+$

$$\tilde{\chi}_2^0 \rightarrow \tilde{l}_i l_j \rightarrow l_k l_j \tilde{\chi}_1^0$$

$$\text{BR}(\tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 e^\pm \tau^\mp)$$



$$\text{BR}(\tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 \mu^\pm \tau^\mp)$$



Variations around SPS1a

## Implications for LHC

studies by:

- F. Paige, I. Hinchliffe, PRD 63 (2001) 115006
- J. Hisanao, R. Kitano, M. Nojiri, PRD 65 (2002) 116002
- D.F. Carvalho et al., PLB 618 (2005) 162
- A. Bartl et al., hep-ph/0510074
- talk by R. Rückl

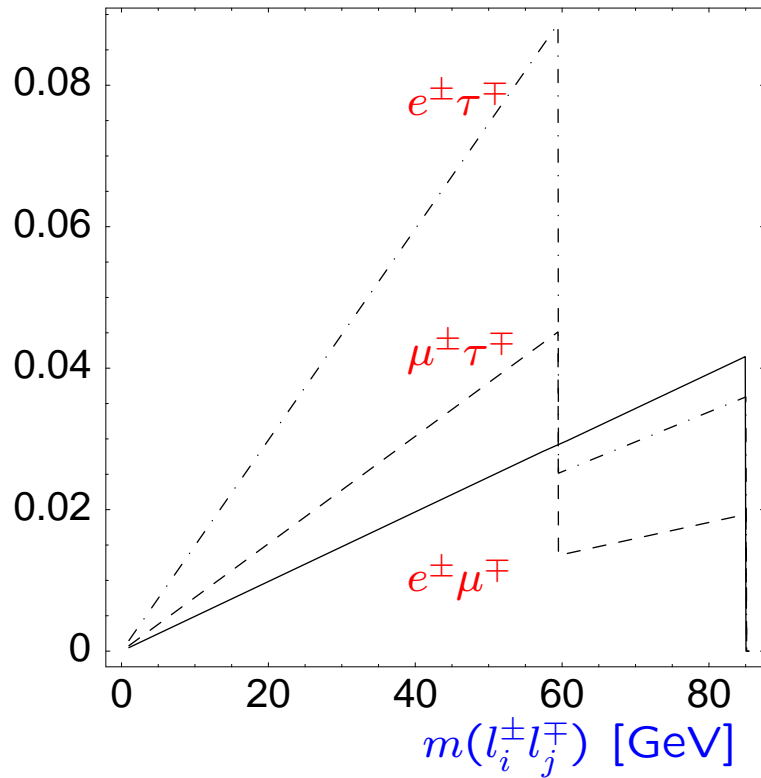
Edge variables positions changes only slightly:  $\pm(1-2) \%$

However, new combinations:  $m_{ll}^{max} \rightarrow m_{e\mu}^{max} m_{e\tau}^{max} m_{\mu\tau}^{max}$   
similarly for  $m_{llq}^{max}$  and  $m_{llq}^{min}$

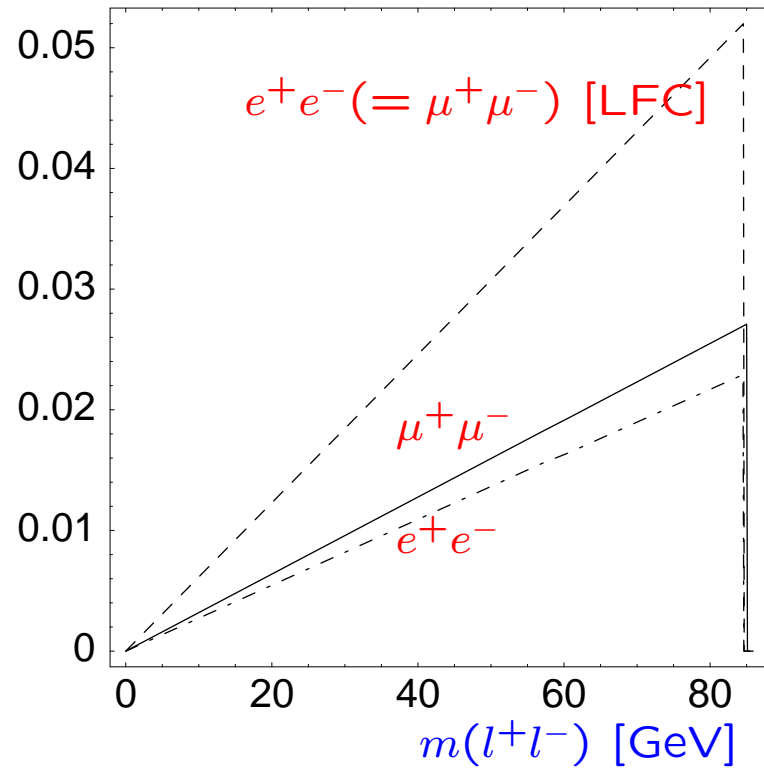
Note:  $\text{BR}(\tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 e^\pm \tau^\mp)$ ,  $\text{BR}(\tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 \mu^\pm \tau^\mp) \simeq \text{BR}(\tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 e^\pm e^\mp)$   
or even larger

$\Rightarrow$  pairing of different lepton flavours necessary

$$100/\Gamma_{tot}d\Gamma(\tilde{\chi}_2^0 \rightarrow l_i^\pm l_j^\mp \tilde{\chi}_1^0)/dm(l_i^\pm l_j^\mp)$$



$$100/\Gamma_{tot}d\Gamma(\tilde{\chi}_2^0 \rightarrow l^+ l^- \tilde{\chi}_1^0)/dm(l^+ l^-)$$



## Bilinearly broken R-parity

Is defined as  $\text{MSSM} + \epsilon_i \hat{L}_i \hat{H}_u + B_i \epsilon_i \tilde{L}_i H_u$

$B_i \epsilon_i$  induces sneutrinos vevs  $v_i$

$(\lambda'_{ijk} \simeq (\epsilon_i/\mu) h_{jk}^D, \lambda_{ijk} \simeq (\epsilon_{[i}/\mu) h_{j]k}^E + v'_i)$

Induced **mixings**: (leptons, charginos), (neutrinos, neutralinos),  
(Higgs bosons, sleptons)

**Solves neutrino problems**:

Atmospheric at tree level, solar at loop level

Negligible flavour violating decays of leptons:

$\text{BR}(\mu \rightarrow e\gamma) < 10^{-17}$ ,  $\text{BR}(\tau \rightarrow e\gamma, \mu\gamma) < 10^{-16}$ .

**Leads to predictions for collider physics**



# Parameters controlling $\nu$ -Physics

If  $m_{\nu,Loop} \ll m_{\nu,Tree}$

$\Delta m_{atm}^2$	$M_2 / \det(\mathcal{M}_{\chi^0})  \vec{\Lambda} ^2$
$\tan^2 \theta_{atm}$	$ \Lambda_3 / \Lambda_3 ^2$
CHOOZ	$ \Lambda_1  / \sqrt{\Lambda_2^2 + \Lambda_3^2}$
$\tan^2 \theta_{sol}$	$ \epsilon_1 / \epsilon_2 ^2$
$m_{sol}^2 / m_{atm}^2$	$ \vec{\epsilon} ^2 /  \vec{\Lambda} $

where

$$\Lambda_i = \mu v_i + v_d \epsilon_i$$

# Approximate Couplings

smallness of R-parity coupling  $\Rightarrow$  expansion of all R-parity violating couplings:

$$c = f(\mu, M_k, A_l, M_{\tilde{j}})\Lambda_i + g(\mu, M_k, A_l, M_{\tilde{j}})\epsilon_i$$

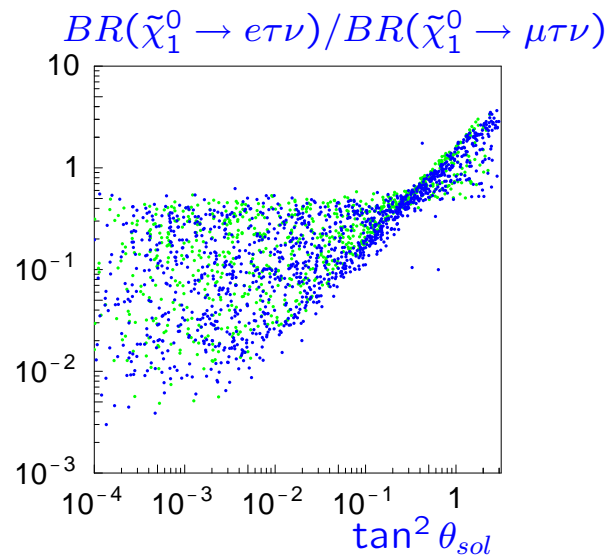
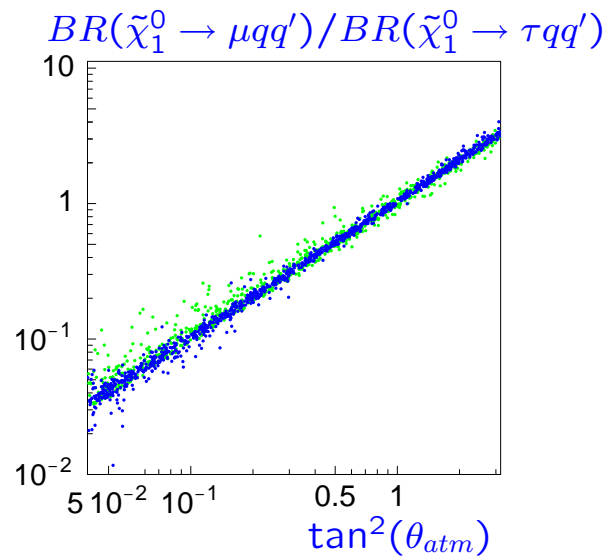
e.g.  $\tilde{\chi}_1^0$ - $W^\pm$ - $l_i$  couplings:

$$O_{Ri} = \frac{gh_{ii}^E v_d}{2\text{Det}_+} \left[ \frac{gv_d N_{12} + M_2 N_{14}}{\mu} \epsilon_i \right. \\ \left. + g \frac{(2\mu^2 + g^2 v_d v_u) N_{12} + (\mu + M_2) g v_u N_{14}}{2\mu \text{Det}_+} \Lambda_i \right]$$

$$O_{Li} = \frac{g\Lambda_i}{\sqrt{2}} \left[ -\frac{g' M_2 \mu}{2\text{Det}_0} N_{11} + g \left( \frac{1}{\text{Det}_+} + \frac{M_1 \mu}{2\text{Det}_0} \right) N_{12} \right. \\ \left. - \frac{v_u}{2} \left( \frac{g^2 M_1 + g'^2 M_2}{2\text{Det}_0} + \frac{g^2}{\mu \text{Det}_+} \right) N_{13} \right. \\ \left. + \frac{v_d (g^2 M_1 + g'^2 M_2)}{4\text{Det}_0} N_{14} \right]$$

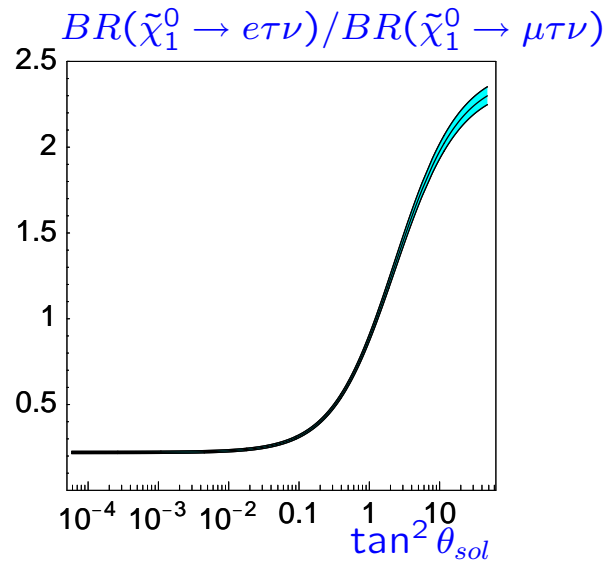
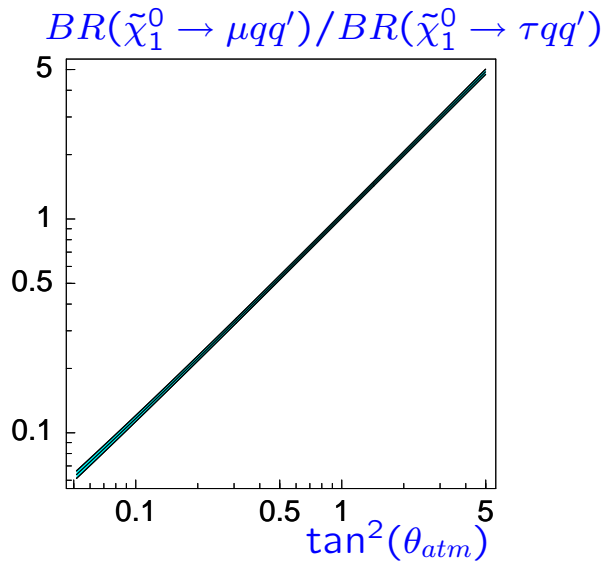
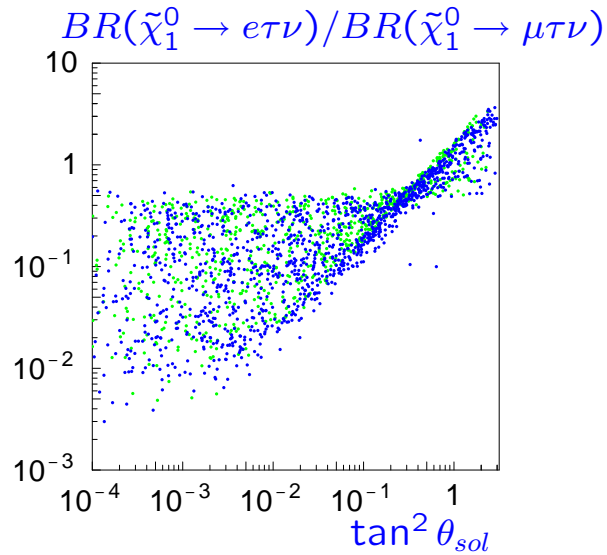
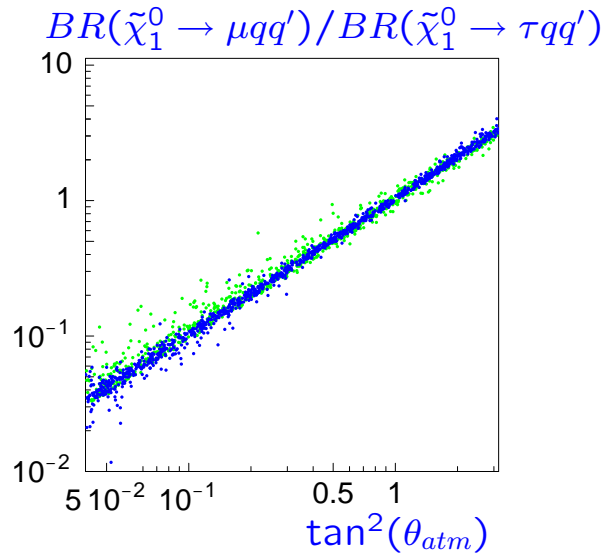
$$|O_{Ri}| \ll |O_{Li}|, \quad \left| \frac{O_{L2}}{O_{L3}} \right|^2 = \left| \frac{\Lambda_2}{\Lambda_3} \right|^2 \simeq \tan^2 \theta_{atm}$$

# Correlations



Summing over all neutrinos.

# Correlations



Assumptions:

- spectrum, mixing angles within 10 percent
- statistical error:  $10^5 \chi_1^0$

Parameters:

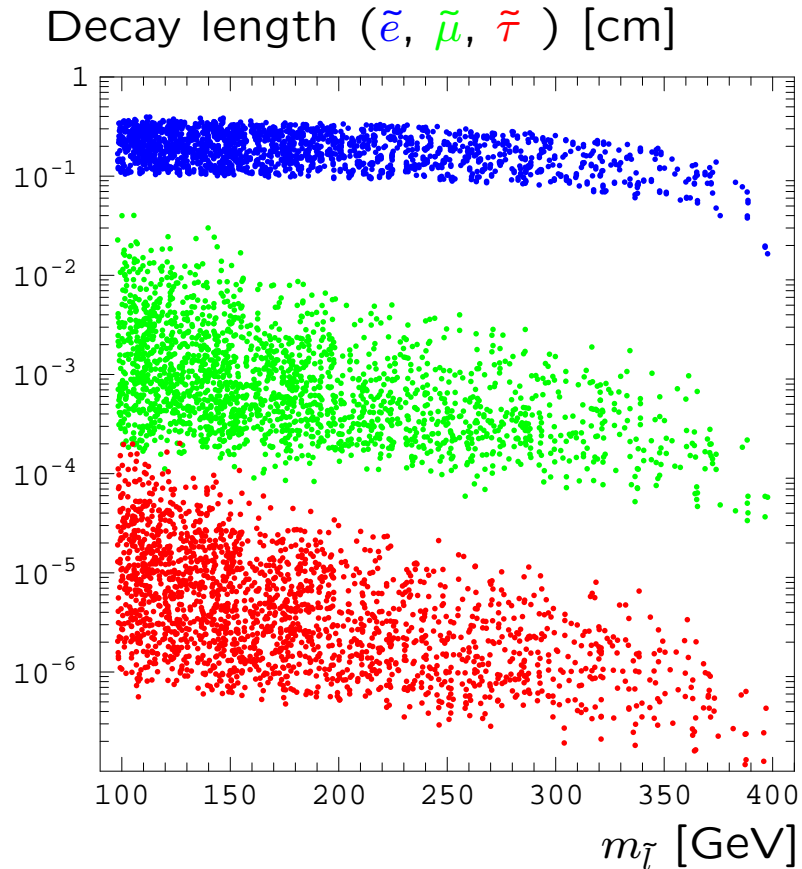
$$M_2 = 120 \text{ GeV}, \mu = 500 \text{ GeV}$$

$$\tan \beta = 5, m_0 = 500 \text{ GeV}$$

$$A = -500 \text{ GeV}$$

Summing over all neutrinos.

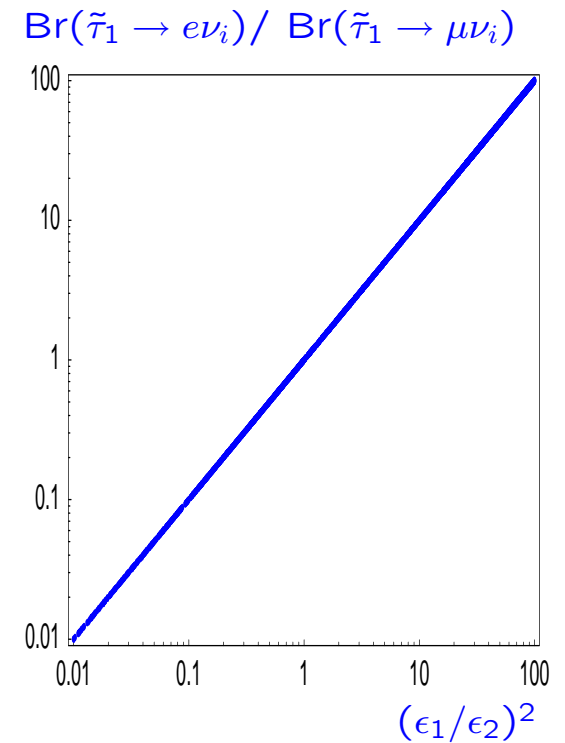
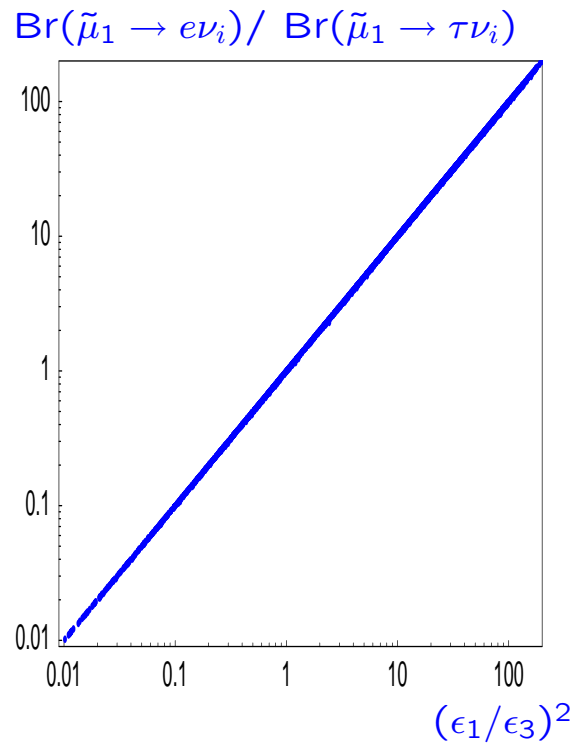
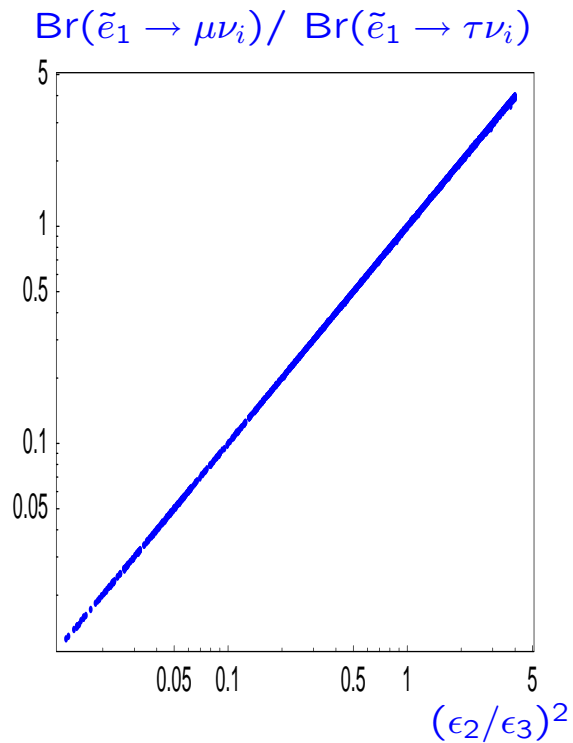
## Charged Scalar LSP



$\Rightarrow \tilde{e}, \tilde{\mu}, \tilde{\tau}$  can be separated  
in this model.

Moreover

$$\frac{\Gamma(\tilde{\tau})}{\Gamma(\tilde{\mu})} \simeq \left( \frac{Y_{\tau}}{Y_{\mu}} \right)^2 \frac{m_{\tilde{\tau}}}{m_{\tilde{\mu}}}$$



Cross check possible:  $(\epsilon_1/\epsilon_3)^2 / (\epsilon_1/\epsilon_2)^2 \equiv (\epsilon_2/\epsilon_3)^2$   
 $\Rightarrow$  Measure 2 ratios, 3rd is fixed.

# Gravitino Dark Matter

GMSB: light gravitino LSP,  $\tilde{\chi}_1^0$  of  $\tilde{l}_R$  NLSP

Standard thermal history of the universe:

$$\Omega_{3/2} h^2 \simeq 0.11 \left( \frac{m_{3/2}}{100 \text{ eV}} \right) \left( \frac{100}{g_*} \right) \quad (g_* \simeq 90 - 140)$$

Current data:  $\Omega_M h^2 \simeq 0.134 \pm 0.006$  ,  $\Omega_B h^2 \simeq 0.023 \pm 0.001$

$\Rightarrow m_{3/2} \simeq 100 \text{ eV}$  if DM candidate, warm dark matter

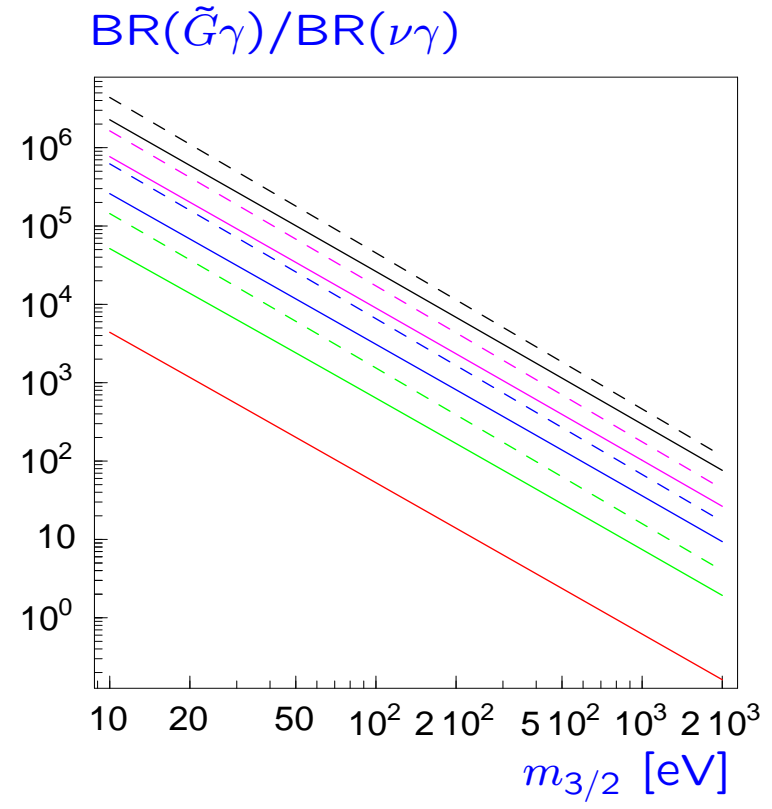
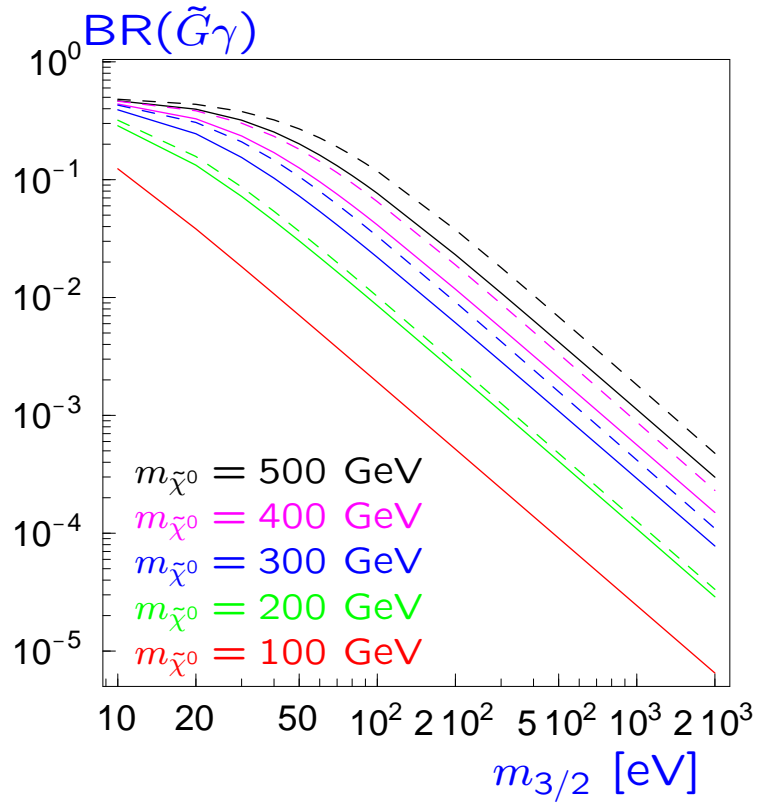
constraints from Lyman- $\alpha$  forest:  $m_{WDM} \gtrsim 550 \text{ eV}$

(M. Viel et al., arXiv:astro-ph/0501562)

$\Rightarrow$  assume additional entropy production, e.g. non-standard decays of messenger particles

(E. Baltz, H. Murayama, astro-ph/0108172; M. Fujii and T. Yanagida hep-ph/0208191)

# GMSB signals



$n_5 = 1, \tan \beta = 10$



## Comments

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$$\frac{m_{\tilde{\tau}_1}}{m_{\tilde{\chi}_1^0}} \propto \frac{1}{\sqrt{n_5}}$$

⇒ for  $n_5 \geq 3$  hardly points with  $\tilde{\chi}_1^0$  NLSP

- $\tilde{l}_R$  NLSPs:  $\text{BR}(l\nu) > \text{BR}(l\tilde{G})$

- $n_5 = 2$ :  $\text{BR}(\tilde{G}\gamma)$  reduced by a factor 2-3

- $\tilde{G}$  decays via R-parity violating couplings, however:

$$\Gamma(\tilde{G}) \simeq 3.5 \cdot 10^{-16} \frac{m_\nu [\text{eV}]^3}{0.05 \text{eV}} \frac{m_{3/2}^3}{M_{Pl}^2} \Rightarrow \tau(\tilde{G}) \sim O(10^{31}) \text{Hubble times}$$