Tommaso Lari Università and INFN Milano On behalf of the BSM convenors

# Beyond the Standard Model WG Status Report (experimental side)

Outline

Introduction

• SUSY BSM

• Non-SUSY BSM



# Working group Topics



http://allanach.home.cern.ch/allanach/lesHouches/susy.html

- Alternative models for Higgs and EWSB (Grojean and Ferrag) Choudhury, Ferrag (8 people in 2<sup>nd</sup> period, 1 not at Les Houches)
- Signature of SUSY breaking scenarii (Lari and Muanza)
- Boudjema, Choudhury, Dittmaier, Galanti, Godbole, Heldmann, Hugonie, Lafaye, Laplace, Lari, Lykken, Mangeol, Penaranda Rivas, Polesello, Prieur, Raklev, Richardson, Rizzi, Schumacher, Spira, Sridhar, Tompkins, Zhukov (30 people in 2<sup>nd</sup> period, 7 not at Les Houches)
- SUSY Les Houches Accord and SPS-like studies (Skands)

Skands (3 people in 2<sup>nd</sup> period)

• Extra-dimensions (Ferrag and Lykken)

Choudhury, Ferrag, Lykken, Przysieniak (6 people in 2<sup>nd</sup> period, 1 not at Les Houches)

- Collider physics and cosmology (Allanach) Lari (6 people in 2<sup>nd</sup> period, 1 not at Les Houches)
- MC and new tools for the new physics (Skands) Skands (2 people in 2<sup>nd</sup> period)

VERY preliminary. List of sub-topics and people to be finalized in these first days.

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# Some general considerations



- By definition, discovery of BSM physics means observing a deviation from the predictions of the SM.
- A good understanding of the signals produced by the SM physics at the LHC is thus necessary to claim discovery of BSM physics (and after discovery to study it).
  - Understanding the detector performance
  - Validate MC tools for LHC energy
  - Use as much as possible the data to estimate the background.
- During early data taking ATLAS and CMS BSM people would actually work on understanding SM physics
  - As day 0 approachs, emphasis on commissioning, background estimation, detailed detector simulation, grid distributed analysis, etc. increases
  - But still ongoing studies on model signatures and new analysis strategies
- I do expect that also here in Les Houches there will be quite some interactions between SM and BSM groups.

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# Supersimmetry



Still the most studed BSM class of models. Among SUSY models, R-parity conserving mSUGRA is probably the most popular.

### **Typical scenario:**

- Production of coloured s-particles, decay into lighter gaugini.
- Stable and weakly interacting Lightest Supersymmetric Particle, to provide a Dark Matter candidate
- Coloured particles mass below 1 TeV (no fine-tuning)

Signatures: Squark and gluino decay into (undetected) LSP produce jets, missing energy, leptons...

### Possible LHC SUSY timeline:

- <u>Phase 1:</u> Discovery (excess of jets and missing energy)
- <u>Phase 2:</u> Masses and decays no mass peak since two undetected LSPs, but if a long enough decay chain can be identified, kinematic endpoints can provide all the masses of the (s)particles involved.
- <u>Phase 3:</u> 2<sup>nd</sup> generation studies: more mass combinations, more decay chains, mass peaks once LSP mass is known, spins, model parameters.

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# SUSY search strategies (1)



What can be seen and at which scale (Heldmann, Hugonie, Savina, ...) SM background to SUSY searches

• Best strategy for mSUGRA is usually jets  $+ E_T^{miss} + n$ -leptons.

• The "Effective Mass":

 $M_{eff} = \Sigma |p_T^i| + E_T^{miss}$ .

discriminates SM and SUSY and has a maximum strongly correlated with the mass of the s-particles produced in the pp collision.

Other MSSM models may have different signatures. Long-lived NLSP decaying in gravitino may give excess of taus or photons, secondary vertices, quasi-stable charged sleptons, ...

Correlation  $(M_{eff}-M_{SUSY})$  also less good in general MSSM.





# SUSY search



• A parameter scan is performed to evaluate the discovery potential and the trigger efficiency of different signatures.

• Natural mSUGRA models ( $m_{SUSY} < 1 \text{ TeV}$ ) may be discovered with a few weeks of data (once calorimeter calibration is understood)

• Caveats:

- Statistical errors only.

-SM background with shower MC (multi-jet xSection too low by orders of magnitude)

Matrix-element MC providing more accurate multi-jet background can be used to re-evaluate discovery potential and benchmark points backgrounds. The background would be eventually be measured from data. How? To which precision? Heldmann, Hugonie, Savina, ... SM background to SUSY searches



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### SUSY mass spectroscopy

g

 $\widetilde{\chi}^{0}$ 

#### Reconstruction of cascade decays p Galanti, Heldmann, Lari, Mangeol, Polesello, Zhukov, ... Precision measurements and new techniques for parameter extraction

### • After discovery: reconstruction of SUSY masses.

• Two undetected LSP: no mass from one specific decay. Measure mass combinations from kinematic endpoints/thresholds. With <u>long enough decay chain</u>, enough relations to get all masses.

### • A point in parameter space is chosen, and decay chains are reconstructed.

•Analysis should be applicable whenever the specific decay do exist.

•Leptonic (e/ $\mu$ ) decay of  $\chi_2^0$  "golden channel" to start reconstruction. But Higgs and  $\tau$  decays can also be used.

• Both ATLAS and CMS have studied in great detail some points favoured by cosmology – at low SUSY scale.

ATLAS Phys. TDR, ATLAS-PHYS-2004-007, CMS-NOTE-2004-029

• Masses can be extracted also by combination of informations from different events (mass relation method, ...)

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# Model-independent masses

• Combine measurements from edges from different jet/lepton combinations to obtain 'model-independent' mass measurements.

• LSP mass poorly determined, and all other masses strongly correlated with it. A Linear Collider input would help a lot!



Reconstruction of cascade decays Galanti, Heldmann, Lari, Mangeol, Polesello, Zhukov, ... Precision measurements and new techniques for parameter extraction LHC/ILC connection Boudjema

masses (GeV)	LHCC5	SPS1a
$m(\tilde{\chi}^{0}_{1})$	122	96
$m(\tilde{l}_R)$	157	143
$m(\tilde{\chi}^{0}_{2})$	233	177
$m(\mathbf{\tilde{q}}_{L})$	687-690	537-543

Sparticle	Expected precision (100 fb <sup>-1</sup> )
$\widetilde{q}_{L}$	± 3%
<b>χ</b> <sup>0</sup> 2	± 6%
Ĩ <sub>R</sub>	± 9%
$\widetilde{\chi}^{0}_{1}$	± 12%

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# Mass peaks



Reconstruction of cascade decays Galanti, Heldmann, Lari, Mangeol, Polesello, Zhukov, ... Precision measurements and new techniques

for parameter extraction

Once  $m(\chi_1^0)$  has been measured, the momentum of the  $\chi_2^0$  can be reconstructed from the approximate relation

 $p(\chi_2^0) = (1-m(\chi_1^0)/m(ll)) p_{ll}$ valid m(ll) near the edge.

The  $\chi_2^0$  can be combined with jets (b-jets) to reconstruct the squark (gluino, sbottom) mass peaks from  $\tilde{g} \rightarrow b\tilde{b} \rightarrow bb\chi_2^0$  and  $\tilde{q} \rightarrow q\chi_2^0$ 

Many other measurements possible:  $\tau\tau$  invariant mass edge  $q_R$  – LSP mass difference Heavy gaugino mass edges







### From masses to model parameters



Precision measurements and new techniques for parameter extraction

SUSY LHA and SPS-like studies Skands

NMSSM DM and colliders

Requirements on LHC and LC data to match precision data on dark matter Allanach

From a given set of measurements one scans the parameter space and finds the points campatible with data. These points are fed to relic density calculators to get constraints on relic density.

ATLAS measu	TLAS measurements Errors			SUSY LHA to interface codes essential here!		
Variable	Value (GeV)	Stat. (GeV)	Scale (GeV)	Total		
$m_{\ell\ell}^{max}$	77.07	0.03	0.08	0.08	Repeat for other benchmark points/models?	
$m_{\ell \ell q}^{max}$	428.5	1.4	4.3	4.5	Micromotics 4.4 (Polonger et	
m <sup>low</sup>	300.3	0.9	3.0	3.1	wicromegas 1.1 (belanger et	
$m_{\ell q}^{high}$	378.0	1.0	3.8	3.9	al.)+ ISASUGRA 7.69 ≝	
$m_{\ell \ell q}^{min}$	201.9	1.6	2.0	2.6		
meth	183.1	3.6	1.8	4.1	$Q_{1}h^{2} = 0.1921 \pm 0.0053$	
$m(\ell_L) - m( ilde{\chi}_1^0)$	106.1	1.6	0.1	1.6	$\chi^{11} = 0.00000$	
$m_{\mathcal{U}}^{max}( ilde{\chi}_4^0)$	280.9	2.3	0.3	2.3	$\log_{10}(\sigma /\text{nb}) = -8.17\pm0.04$	
$m_{\tau\tau}^{max}$	80.6	5.0	0.8	5.1	$S_{10}(O_{\chi p}, p_0) = O_{10}(O_{\chi p}, p_0)$	
$m(g) - 0.99 \times m(\chi_1^{\circ})$	500.0	2.3	6.0	6.4	800 [ <b>Δ<sup>2</sup>χ</b> ]	
$m(q_R) - m(\chi_1)$	424.2	10.0	4.2	10.9	300 fb-1	
$m(g) - m(b_1)$	103.3	1.5	1.0	1.8		
$m(g) - m(b_2)$	70.6	2.5	0.7	2.6		
Parameter	Expect	ed precis	sion (300	fb <sup>-1</sup> )	400 - / ]	
m <sub>o</sub>		± 2%				
m <sub>1/2</sub>		± 0.6%				
tan(B)		+ 9%				
A <sub>0</sub>		± 16	70		$0.17 \ 0.175 \ 0.185 \ 0.19 \ 0.195 \ 0.2 \ 0.205 \ 0.215 \ 0.225 \ 0.215 \ 0.22 \ 0.215 \ 0.22 \ 0.215 \ 0.225 \ 0.215 \ 0.225 \ 0.215 \ 0.225 \ 0.215 \ 0.225 \ 0.215 \ 0.225 \ 0.215 \ 0.225 \ 0.215 \ 0.215 \ 0.225 \ 0.215 \ 0.215 \ 0.225 \ 0.215 \ 0.$	
. Lari BSM report		Les noucnes 3 May 2005			May 2005 Physics at TeV colliders	



# SUSY and Cosmology

Only tiny mSUGRA space allowed by LEP and cosmology ( $\chi$  relic density <= DM abundace).

- **<u>Bulk</u>**: low susy masses, most studied in the past.
- Focus Point: large scalar mass (> 3 TeV), large mixing in neutralino sector. Higgsino component of  $\chi^{0}_{1}$  gives rapid s-channel annihilation in early universe. In this region, large differences between mass spectra and relic density predicted by RGE codes (ISAJET, SOFTSUSY, ...) Also sensitive to top mass value.
- <u>Coannihilation</u>:  $\tau$  and  $\chi$  close in mass, relic density reduced by  $\tau \chi \rightarrow SM$ .
- <u>Higgs funnel</u>: At large tanβ, neutralino annihilation through Higgs resonance.
  Looks like mSUGRA is too constrained.
  Search for cosmologically motivated points with relaxed universality or in NMSSM?

What can be seen and at which scale Heldmann, Hugonie, Savina, ... NMSSM DM and colliders

#### Baer et al. hep-ph/0305191









# **RPV SUSY**



What can be seen and at which scale (Heldmann, Hugonie, Savina, ...)

Discriminating between models

#### SUSY LHA - New ingredients: RPV

CMS Study: Trigger rate vs SUSY selection efficiency, varying trasverse energy cut  $E_T^{min}$ 



Neutralino decay: <u>less missing energy</u>, <u>more jets</u>. Overall somewhat more difficult to see.

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Gravitino LSP. NLSP can be stau or neutralino. Lifetime can be substantial.





# Split SUSY



N. Arkani-Hamed and S.Dimopoulos, hep-th/0405159. A. Romanino and G.F.Giudice, Nucl. Phys. B699 - 65.

Split SUSY Lari,Savina

- Ignore hierachy problem (also there for cosmological constants, one may invoke huge number of vacua and antropic principle)
- Keep SUSY (unification of coupling constants, dark matter...)
- Scalar particles are (VERY) heavy
- Gluino is long-lived (decays to gaugini through virtual squarks) from a narrow resonance to cosmological lifetimes
- If gluino prompt decay: like mSUGRA with heavy scalars (focus-point)
- If gluino lifetime in ps us range: secondary vertices
- If quasi-stable gluino: neutral and charged R-hadrons produced
  - Charge-exchange reaction every nuclear int. length: charge state changes in calorimeter
  - EM+nuclear interaction: no shower, but more energy loss than heavy muon
  - Energy profile in calorimeter, time-of-flight in muon chambers, ... : very typical signature (almost no background expected)
  - LHC sensible up to 1.7 TeV mass

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# SUSY Higgs sector

- 2 doublets, 5 physical states: h<sup>0</sup>,H<sup>0</sup>,A<sup>0</sup>,H<sup>±</sup> (mix if CPV)
- h light, SM-like. m < 133 GeV
- Lots of free parameters in MSSM
- Often assume heavy SUSY states (no Higgs decay into SUSY nor Higgs production in SUSY decays)



- Dittmaier, Penaranda Rivas, Schumacher SUSY Models with an Heavy Higgs Invisible Higgs and CP violation in the Higgs sector
- Define banchmark scenarios. Example (Carena et al., Eur.Phys.J.C26,601):
  - MASSH maximum h mass allowed by theory
  - Nomixing small h mass (difficult for LHC)
  - gluophobic reduces hg coupling (and LHC production xSection)
  - Small α reduces hbb and hττ couplings (harms some discovery channels)
- Parameter scans performed on two free parameters  $(m_{A_{c}} \tan \beta)$ 
  - SM xSection + MSSM correction factors
  - Higgs decays (FeynHiggs)
  - Efficiency and background from MC studies of different channels
  - Corrections from Higgs width and overlap of states

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![](_page_17_Picture_0.jpeg)

**<u>CPV Higgs</u>**. Neutral Higgs states mix. Smaller mass for the lightest state allowed by LEP (much below Z mass). For low mass observation by LHC to be studied yet. Higgs in cascade decays. Peak in bb invariant mass distribution – with SUSY cuts may be much easier to

see than SM Higgs.

INFN

di Fisica Nuclear

![](_page_17_Figure_3.jpeg)

![](_page_17_Picture_4.jpeg)

CPV Higgs states observable.

![](_page_18_Picture_0.jpeg)

# Non-SUSY BSM

![](_page_18_Picture_2.jpeg)

• Of course, lots of ideas....

Leptoquarks, black holes, Left-Right Symmetric Model, excited quarks and leptons, compositness, ...

- Many models are built to solve the hierarchy problem as a guideline.
- Focus here on
  - <u>Little Higgs</u>: the SM is part of a symmetry group broken at a few TeV scale. Delays the fine-tuning problem to that scale by introducing new particles that cancel the quadratic divergences to the Higgs mass (a new heavy quark, new gauge bosons, heavy Higgs)
  - <u>Extra dimensions</u>: gravity is strong at the TeV scale (gravitons, excitations of SM particles if they can propagate in extra dimensions)
  - <u>Higgsless models</u>

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![](_page_19_Picture_0.jpeg)

# Little Higgs Models (LH)

### Higgs as a Goldstone boson

Known and new Higgs, gauge bosons coming from breaking a SU(5) symmetry at scale v (few TeV). Divergent contribution to the Higgs mass from top, W, Z and Higgs loops are canceled by the new particles:

•Heavy gauge bosons  $Z_H$ ,  $W_H$ ,  $A_H$ m < 6 TeV  $(m_h/200 \text{ GeV})^2$ • Heavy quark T (electroweak singlet)  $v\sqrt{2} < m < 2 \text{ TeV} (m_h/200 \text{ GeV})^2$ • New Higgs bosons  $\Phi^0 \Phi^+ \Phi^{++}$ m < 8 TeV  $(m_h/200 \text{ GeV})^2$ 

"Littlest Higgs model" (T. Han et al., Phys. Rev. D67, 095004) used for a detailed ATLAS study (G. Azuelos et al., hep-ph/0402037). Also under study by CMS.

CMS study for generic heavy gauge bosons is also relevant (M. Dittmar et al., hep-ph/0307020).

![](_page_20_Picture_0.jpeg)

# LH: T Quark Search

![](_page_20_Picture_2.jpeg)

![](_page_20_Figure_3.jpeg)

![](_page_21_Picture_0.jpeg)

# LH: New gauge bosons

![](_page_21_Picture_2.jpeg)

Lots of models with heavy W/Z bosons.
Following discovery of ee/μμ resonance discriminating among them would required detailed measurements of width, asymmetries, cross section lineshape, etc.

• Discovery:

 $A_{\rm H}/Z_{\rm H} \rightarrow ee, \mu\mu \quad W_{\rm H} \rightarrow e\nu, \mu\nu$ 

- Up to ~5 TeV, except for small  $\cot\theta$  (Z<sub>H</sub>, W<sub>H</sub>) and  $\tan\theta$ `~1.3 (
- CMS reach similar
- Cross section, width measure  $\theta$
- <u>Specific of LH models</u> (assuming m<sub>h</sub> =120 GeV):

$$Z_H \rightarrow Zh \rightarrow llbb$$

$$W_{H} \rightarrow Wh \rightarrow l\nu bb$$

$$W_{\rm H}/Z_{\rm H} \rightarrow W/Z h \rightarrow qq\gamma\gamma$$

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![](_page_21_Figure_14.jpeg)

![](_page_22_Picture_0.jpeg)

### **Extra Dimensions**

Model independent constraints on new gauge bosons Universal extra Dimensions

> Factorized metric  $ds^2 = dr^2 + dt^2 + du^2$

#### Large xTra Dim

Radius R >> TeV<sup>-1</sup> Modify Newton's Law below R Lower Planck scale to TeV

Only gravitons in xtraDim (SM fields does not show Characteristic excited states at scale R<sup>-1</sup> << TeV<sup>-1</sup>) Signatures:

• (near-)continuum of graviton states

• Direct production, virtual effects observable

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#### TeV<sup>-1</sup> scale xTra Dim

Radius R  $\sim$  TeV<sup>-1</sup> May come with others large xTra Dim.

SM fields allowed in xTra Dim Tower of KK excitations at TeV scale for each particle in bulk.

Signatures: Excited states of gauge bosons Excited states of fermions if live in bulk.

![](_page_22_Picture_14.jpeg)

Non-factorizable metric  $ds^2 = f(u)(dr^2 + dt^2) + du^2$ 

### **Randall-Sundrum**

Radius M<sup>-1</sup><sub>Planck</sub> But phenomenology at TeV scale.

Graviton discrete excitations Also new scalar field (radion)

![](_page_23_Figure_0.jpeg)

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![](_page_24_Picture_0.jpeg)

# TeV<sup>-1</sup> Extra dimension(s)

![](_page_24_Picture_2.jpeg)

#### Model independent constraints on new gauge bosons

- One of the extra dimensions may have smaller size (TeV<sup>-1</sup>): all SM fields (Universal Extra Dimension) or just gauge boson may propagate in it.
- Tower of excited KK states with mass

 $m_k^2 = m_0^2 + k^2 M_C^2$ 

- Gauge bosons KK: probably only first resonance observable (EW data constraints), discovery with ee, μμ, ev, μν
- Precision measurements with electrons

 $Z^{(1)}/\gamma^{(1)}$ : G.Azuelos and G.Polesello, in hep-ph/0204031 $W^{(1)}$  and  $g^{(1)}$  can also be seen by ATLASSensitivity to peak (100 fb<sup>-1</sup>):5.8 TeVReach with interference, el. (100 fb<sup>-1</sup>):9.5 TeVUltimate with interference, e+ $\mu$ , 300 fb<sup>-1</sup>: 13.5 TeV

![](_page_24_Figure_10.jpeg)

![](_page_25_Picture_0.jpeg)

# **Discrimination of Models**

![](_page_25_Picture_2.jpeg)

	process	$\sigma \times BR(Z^* \to e^+e^-)$ (fb)
	$Z^{(1)}/\gamma^{(1)}$	4.05
	$Z^{(1)}/\gamma^{(1)}$ -M2	11.75
· Crass santian width resonance shane	Z'	4.65
• Cross section, width, resonance shape	$qq \rightarrow G^*$	0.20
Not shown: asymmetries	$gg \to G^*$	0.13
• Discrimination $Z^{(1)}/Z'/G^*$ possible	$qq \rightarrow e^+e^-$	4.83
$W^{(1)}/W$ difficult		

![](_page_25_Figure_4.jpeg)

![](_page_26_Picture_0.jpeg)

# **Universal Extra Dimensions**

T. Appelquist, HC Cheng and BA Dobrescu, PR D64 (2001) 035002

### >All SM particles in bulk

- $\Rightarrow$  conservation of momentum in extra dimensions
  - $\Rightarrow$  conservation of KK number
    - $\Rightarrow$  pair production of KK states
      - $\Rightarrow$  lower collider bounds: ~ 350-400 GeV
- LKP quasi-stable (decay only via graviton emission)

### Universal Extra Dimensions

![](_page_26_Figure_10.jpeg)

![](_page_26_Picture_11.jpeg)

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![](_page_26_Picture_14.jpeg)

![](_page_27_Figure_0.jpeg)

![](_page_28_Picture_0.jpeg)

# **Randall-Sundrum model**

Events/0.2

12

10

8

6

Spin-1

gg

![](_page_28_Picture_2.jpeg)

6000

- Model parameters from resonance mass, width and x-section
- May be possible to observe second resonance (spaced as Bessel function zeros)
- Spin measurement possible over most of parameter space

![](_page_28_Figure_6.jpeg)

![](_page_29_Picture_0.jpeg)

# Higgless models

![](_page_29_Picture_2.jpeg)

Warped space, with boundary conditions that break the symmetry on the TeV brane and on the Planck brane:

> C. Csáki et al., hep-ph/0310355, C. Csáki, hep-ph/0412339

The model explains:

- γ : massless photon (flat wavefunction in bulk)

- W, Z : lowest KK states of massive gauge bosons

- correct ration of W/Z mass

![](_page_29_Figure_9.jpeg)

Figure 3: The symmetry breaking structure of the warped higgsless model.

#### Important constraints:

- S parameter from LEP:
  - $\rightarrow$  weak coupling of Z' to fermions (and possibly light Z')
- unitarity in VB scattering
   → resonances in WZ scattering
   distinguishable from QCD-like
   chiral Lagrangian model resonances.

![](_page_29_Figure_15.jpeg)

FIG. 1. Diagrams contributing to the  $W^{\pm}Z \rightarrow W^{\pm}Z$  scattering process: (a), (b) and (c) appear both in the SM and in Higgsless models, (d) and (e) only appear in Higgsless models, while (f) only appears in the SM.

Birkedal et al., hep-

Physics at TeV colliders

T. Laring shiens, with the top

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![](_page_30_Picture_0.jpeg)

![](_page_30_Picture_1.jpeg)

### resonance in WZ scattering

![](_page_30_Figure_3.jpeg)

FIG. 4. The number of events per 100 GeV bin in the  $2j + 3\ell + \nu$  channel at the LHC with an integrated luminosity of 300 fb<sup>-1</sup> and cuts as indicated in the figure. The model assumptions and parameter choices are the same as in Fig. 2.

A. Birkedal et al., hep-ph/0412278

![](_page_30_Picture_6.jpeg)

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![](_page_31_Picture_0.jpeg)

![](_page_31_Picture_1.jpeg)

![](_page_31_Picture_2.jpeg)

- Lots of work has been made in preparation of LHC start-up on extensions of the Standard Model...
- Even more remains to be done!
- So... have a good workshop!

![](_page_31_Picture_6.jpeg)

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![](_page_32_Picture_0.jpeg)

![](_page_32_Picture_1.jpeg)

# Backup slides

![](_page_32_Picture_3.jpeg)

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![](_page_33_Picture_0.jpeg)

# Mass Relation Method

![](_page_33_Picture_2.jpeg)

- Hot off the press: new idea for reconstructing SUSY masses!
- 'Impossible to measure mass of each sparticle using one channel alone' (Page 8).
  - Should have added caveat: Only if done event-by-event!
- Remove ambiguities by combining different events analytically → 'mass relation method' (Nojiri et al.).
  - Also allows all events to be used, not just those passing hard cuts

![](_page_33_Figure_8.jpeg)

![](_page_34_Figure_0.jpeg)

![](_page_35_Picture_0.jpeg)

### **CPV Higgs**

![](_page_35_Picture_2.jpeg)

CP conserving at Born level, but CP violation
 via complex A<sub>t</sub>, A<sub>b</sub> M<sub>gl</sub>
 CP eigenstates h, A, H mix to mass eigenstates

H<sub>1</sub>, H<sub>2</sub>, H<sub>3</sub>

➤ maximise effect → CPX scenario (Carena et al., Phys.Lett B495 155(2000)) arg(A<sub>t</sub>)=arg(A<sub>b</sub>)=arg(M<sub>gluino</sub>)=90 degree

> scan of Born level parameters: tan $\beta$  and M<sub>H+-</sub>

no absolute limit on mass of H<sub>1</sub> from LEP
 strong dependence of excluded region
 on value for m<sub>top</sub>

on calculation used FeynHiggs vs CPH

![](_page_35_Figure_9.jpeg)

![](_page_35_Figure_10.jpeg)

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![](_page_36_Figure_0.jpeg)

MC studies of ATLAS searches for  $m_{\rm H} < 70$  GeV

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M<sub>H3</sub>: 140 to 180 GeV

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### Supersymmetry – Spin Measurement INFN

![](_page_37_Picture_1.jpeg)

#### A.J. Barr, hep-ph/0405052

Evidence for supersymmetry (vs extra dimensions, for example)

![](_page_37_Figure_4.jpeg)

![](_page_37_Picture_5.jpeg)

500

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![](_page_38_Picture_0.jpeg)

![](_page_38_Picture_1.jpeg)

![](_page_38_Picture_2.jpeg)

ADD model: Arkani-Hamed, Dimopoulos and Dvali.

- N. Arkhani-Hamed et al., Phys. Lett. B429, 263
- N. Arkhani-Hamed et al., Phys. Rev. D59, 086004
- I. Antoniadis et al., Phys. Lett. B436, 257
- $\delta$  new dimensions of size TeV<sup>-1</sup> << R<sub>0</sub> < 0.2 mm
- Gravity propagates in the whole space (bulk)  $\rightarrow$  increases as R<sup>-(2+ $\delta$ )</sup> for R < R<sub>0</sub> and is strong at scale M<sub>D</sub> (~ TeV).
- $M_D^{\delta+2} R_0^{\delta} = M_{Planck} \rightarrow R_0 \sim 1 \text{ mm} (\delta=2) \text{ or } 10 \text{ fm} (\delta=6)$
- Direct tests of Newton's law exclude  $\delta=1$ ,  $\delta=2$  marginal ( $R_0 < 190 \mu m$ )
- Stringent (but model-dependent) astrophysical limits
- Low-energy Kaluza-Klein graviton excitations. Universal and weak coupling to SM particles. Large number of states (- continuum).

![](_page_38_Picture_13.jpeg)

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![](_page_39_Figure_0.jpeg)

Scalar particles out of reach.  $\chi\chi$  production (4.5 pb) difficult to separate from SM background gg production (0.6 pb) and decay into gauginos can be observed. Two mass differences from neutralino leptonic decays. Reconstruction of gaugino MSSM Parameters (M<sub>1</sub>,M<sub>2</sub>, $\mu$ ,tan $\beta$ ) to be demonstrated yet. T. Lari BSM report

Sleptons close in mass to neutralinos: slow sleptons  $\chi$  from decay. Still several mass combination can be reconstructed.

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![](_page_40_Picture_0.jpeg)

![](_page_40_Figure_1.jpeg)

**Single production** 

Main background from  $W_TW_T$  scattering

![](_page_40_Figure_4.jpeg)

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