Progress on VirCol A Parton Shower MC based on Antenna Formalism

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QCD at High Energies



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Hard & Soft

Matrix Elements (Fixed Order):

- •Fixed order in α -> Exact interference, helicity, loops, ...
- •At present can do 2->5/6 (less with loops)
- •Perturbative expansion better at higher energy (asymptotic freedom)
- Multiple soft emissions important for full event structure = exclusive observables
- •Widely separated scales -> big logs / big truncation errors.
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- Derived in universal limit of QCD → depend on universal parameters
- •Exponentiate \rightarrow infinite O(α) \rightarrow ideal for widely separated scales (logs resummed)
- Arbitrary number of partons in final state → match to hadronisation descriptions
- •Derived in <u>limit</u> (collinear) of QCD \rightarrow approx. for wide-angle / hard emissions.

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Marriage desireable!!

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Possible Ceremonies



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Vircol – Basic SKETCH

- Perturbative expansion for some observable J, $d\sigma = \Sigma_{m=0} d\sigma_m$; $d\sigma_m = d\Pi_m |M|^2 \delta(J-J(k_1,k_2,...,k_m))$
- Assume we know some Matrix Elements $d\sigma_0, d\sigma_1, \dots d\sigma_n$ (w or w/o loops)
- And we have some approximation $T_{n \rightarrow n+1}$, so that $d\sigma_{n+1} \sim T_{n \rightarrow n+1} d\sigma_n$ (~ parton shower)

• A 'best guess' cross section is then: $d\sigma \sim d\sigma_0 + d\sigma_1 + \dots + d\sigma_n (1 + T_{n \rightarrow n+1} + T_{n \rightarrow n+1} T_{n+1 \rightarrow n+2} + \dots)$ $\Rightarrow d\sigma \sim d\sigma_0 + d\sigma_1 + \dots + d\sigma_n S_n \qquad ; S_n = 1 + T_{n \rightarrow n+1} S_{n+1}$

- For this to make sense, the $T_{n \rightarrow n+1}$ have to at least contain the correct singularities (in order to correctly sum up all logarithmically enhanced terms), but they are otherwise arbitrary.
- We will now reorder this series in a useful way ...

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Reordering example: $h \rightarrow gg$

■ Assume we know ME for $H \rightarrow gg$ and $H \rightarrow ggg$. Then reorder:

Use
$$1=S_n-T_{n \to n+1}S_{n+1}$$

 $d\sigma \sim d\sigma_{gg} + d\sigma_{ggg}S_{ggg}$
 $= S_{gg}d\sigma_{gg} + S_{ggg}(d\sigma_{ggg} - T_{gg \to ggg}d\sigma_{gg})$
 $= S_{gg}d\sigma_{gg} + S_{ggg}d\chi_{ggg}$ (generalises to n gluons)

■ I.e shower off gg and modified ggg matrix element.

 Double counting avoided since singularities/shower subtracted in dχ_{ggg}.

What IS THE Difference?

CKKW (& friends) in a nutshell:

- 1. Generate a n-jet Final State from n-jet (singular) ME.
- 2. Construct a "fake" PS history.
- 3. Apply Sudakov weights on each "line" in history → from inclusive n-jet ME to exclusive n-jet (i.e. probability that n-jet FS remains n-jet above cutoff) → gets rid of double counting when mixed with other ME's (Sudakov wt dampens singularity).
- 4. Apply PS with no emissions above cutoff.

VirCol in a nutshell:

- 1. Subtract PS singularities from n-jet ME (antenna subtraction)
- 2. Generate a n-jet Final State from the subtracted (finite) ME.
- 3. Apply PS \rightarrow Leading Logs resummed.
- + full NLO: divergent part already there = unitarity of shower assumption \rightarrow just include extra finite contribution in $d\sigma_0$: $d\sigma = d\sigma_0^{(0)} + d\sigma_1^{(0)} + sing[d\sigma_0^{(1)}] + F^{(1)} + ...$
- + now NNLO/NLL possible -> talks by Gehrmann, Gehrmann-De Ridder

The ANTENNA Shower

So far, we have written a C++ code that (for the moment) generates a pure gluon cascade ordered in:

 $y_R = 4s_{a1}s_{1b}/s^2_{a1b} = 4p^2_{T;ARIADNE}/s_{a1b}$

- ...with the antenna / subtraction function: $|A(a,b\rightarrow a,1,b)|^2 =$
 - $2(s_{a1b}(s_{a1}+s_{1b})+s_{ab}^2)^2/(s_{a1}s_{1b}(s_{a1b}s_{ab}+s_{a1}s_{1b})s_{a1b})$
 - \rightarrow "usual" collinear limit, but different outside.
- This gives an analytical Sudakov integral
 = [Mathematica output].
- (No Matrix Elements yet ... but work in progress).





- Sudakov → y_R for next branch → select phase space point along iso-y_R contour:
 - 1. Rewrite Antenna function in terms of $y^2 = y_R$; $\xi = (s_{a1} + s_{1b})/R_1^{-1} \quad dy \quad \frac{1}{y} R_{1=y} \quad dy \quad \frac{(y + (1_i - y_i)^2)^2}{1 + y^2 = 4_i - y_i} P_{\frac{1}{y^2} - \frac{1}{i}}$
 - 2. Partial-fraction singular structure + overestimate numerators \rightarrow generate uniform R and solve for ξ_{R} $R = \frac{R_{*R}}{1} d* \frac{A(y;*)}{1; * y + y^{2} = 4} + \frac{B(y;*)}{\frac{B(y;*)}{x^{2}; 1}}$









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A1: Sudakov suppressed, (sum of ordered branching probs → unordered probability)



Q1: Where's the soft singularity?

Q2: Is that a 'dead region' ?



- Q1: Where's the soft singularity?
- Q2: Is that a 'dead region' ?
- A2: Yes, it is cut out by 'unresolution criterion', i.e. that neighbour dipoles remain resolved after branching.
- Due to H→gg in colour singlet state! (pathological)
- Eventually, could be filled by Matrix Elements and/or by changing evolution var.



Preliminary Results

Thrust and 3-jet rate, compared to Q²-ordered and p_T²-ordered PYTHIA showers.



Preliminary: matrix elements should be added and parton-level matched to hadronisation models eventually + all showers only include gluons here...

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Conclusion & Outlook

Construction of VIRCOL shower monte carlo:

- gluons shower MC (based on LO, done!)
- gluons shower MC (based on NLO)
- parton shower MC (LO/NLO(/NNLO))
- parton shower MC (NLL + NLO/NNLO)
- Hadron collider shower MC's
- Higher order Sudakov factor calculations(this will reduce a lot of implicit and explicit uncertainties: e.g. renormalization scale, choice of subtraction function,...)