

Higgs report

Scott Willenbrock

U. of Illinois at Urbana-Champaign

TeV4LHC@BNL

February 5, 2005

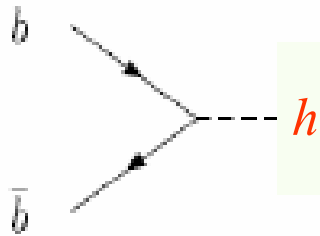
Talks

- Fabio Maltoni $b\bar{b} \rightarrow h, Z$
- Alessandro Vicini NLO EW corrections to $gg \rightarrow h$
- Kirill Melnikov NNLO differential $gg \rightarrow h$
- Marcela Carena \mathcal{CP} violating Higgs in MSSM
- Kamal Benslama H^{++} in the LR symmetric model
- Ariel Schwartzman SVX from D0 to CMS
- Ariel Schwartzman Jet energy resolution for Higgs
- Bruce Mellado Higgs plus 1-jet signatures

Update on $b\bar{b}\rightarrow Z$

Fabio Maltoni
CERN

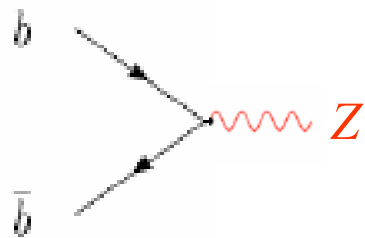
Motivation



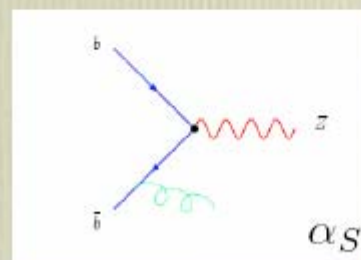
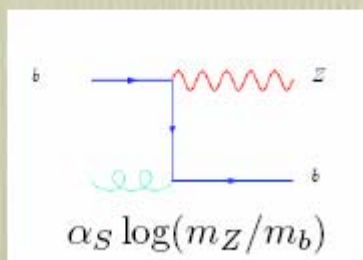
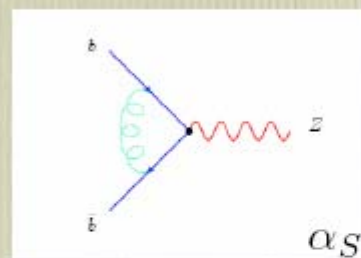
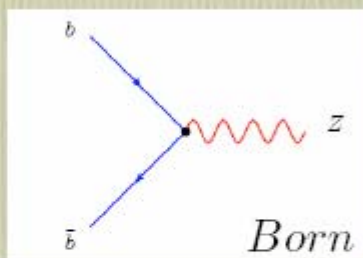
“Inclusive” b tagging

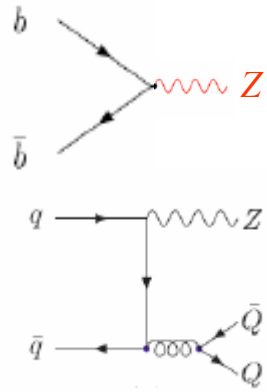
Inclusive approach: look only at the track displacements in the vertex detector and give to the event a probability of containing a b quark.

Test this method on



$b\bar{b} \rightarrow Z$ at NLO





Cross sections (pb) for Z + “inclusive b” Summary

| Process | TeV | LHC |
|----------------------|-------------|-------------|
| $bb \rightarrow Z$ | (30.6) 30.0 | (1860) 1800 |
| $qq \rightarrow Zbb$ | 14.0 | 122 |

$$\mu_F = \mu_R = m_Z$$

Proposal for the TEV₄LHC

Study “inclusive” bottom measurements in W/Z production

theory: we can predict cross sections extremely well

experiment: new approach, maybe better sensitivity

theory : perform the new NLO (and NNLO) calculations for Z and W that are needed

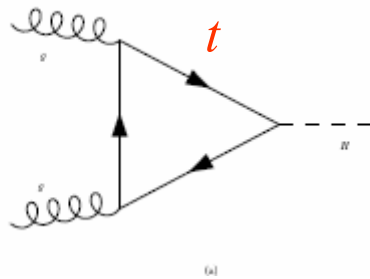
experiment: look at what CMS has done, use CDF and DO data for Wbb and Zbb to test feasibility, find efficiencies, etc...

Two-loop light fermion corrections to Higgs production and decays

Alessandro Vicini

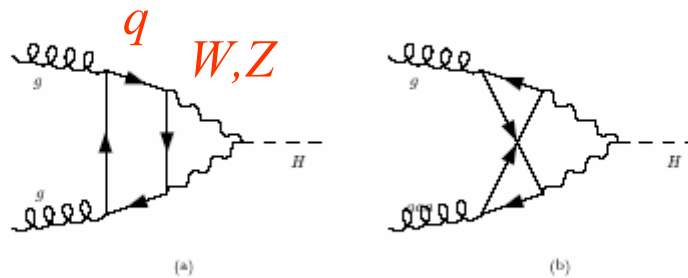
Università degli Studi di Milano and
INFN - Sezione di Milano

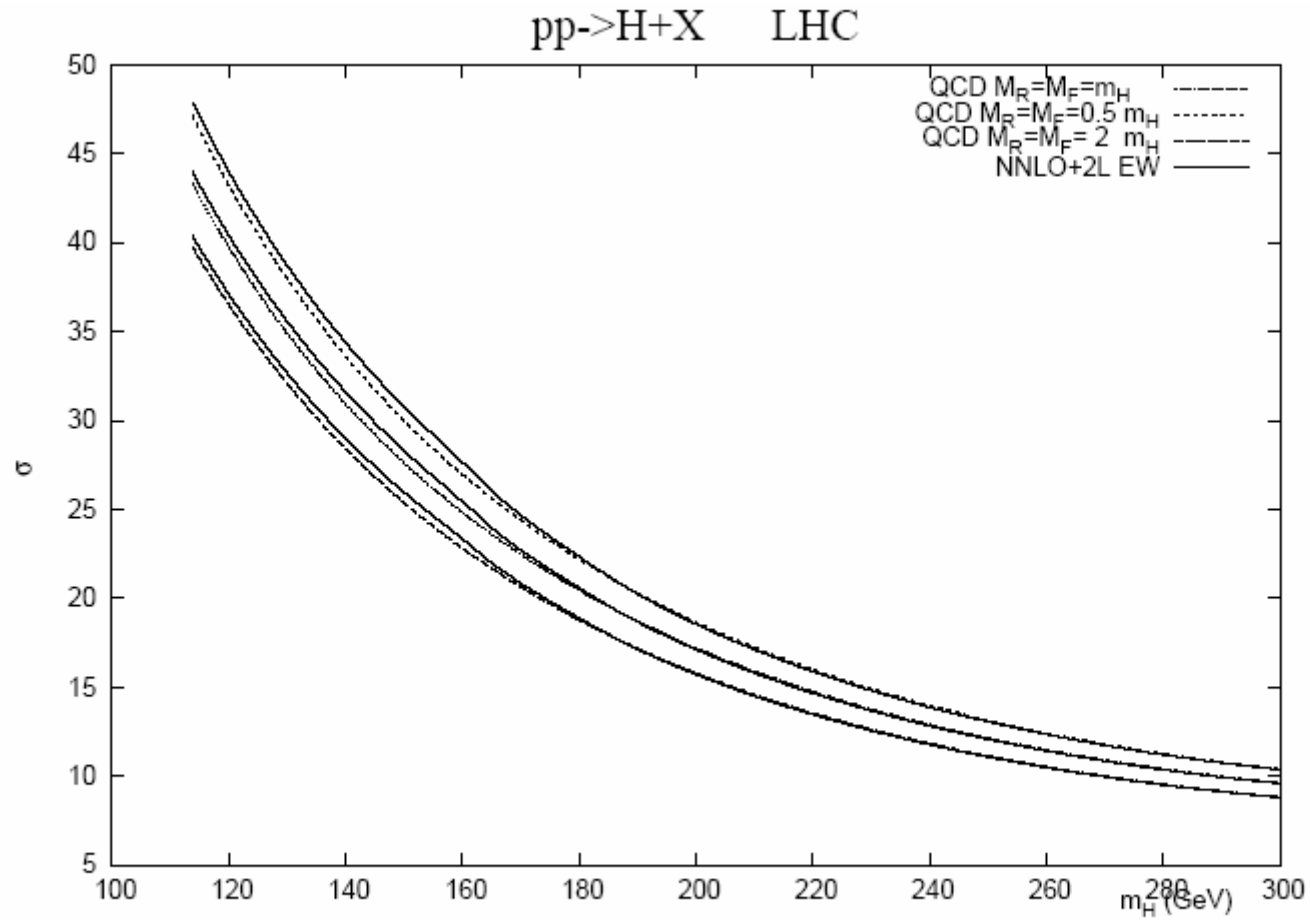
The gluon fusion process:



NNLO QCD (large m_t)

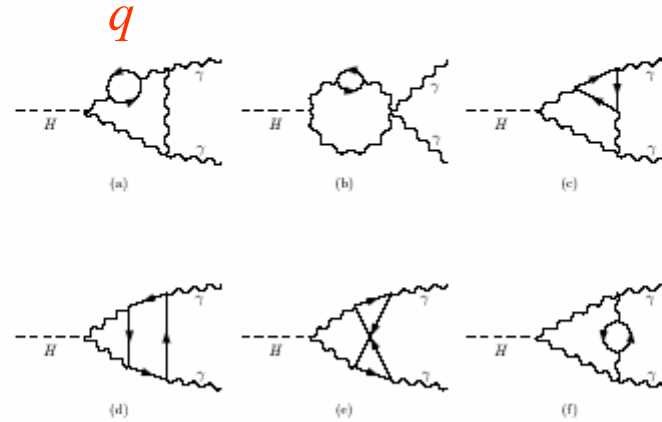
NLO EW correction:



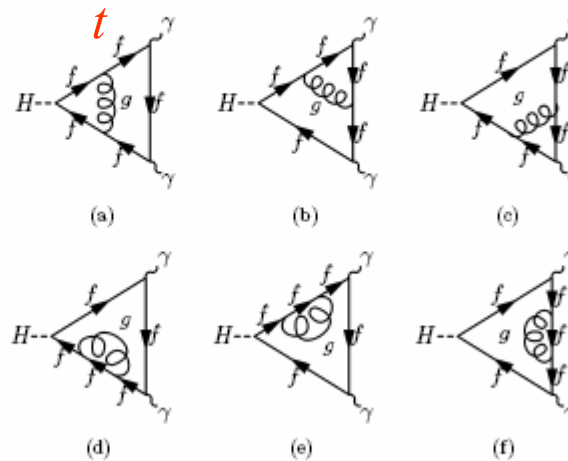


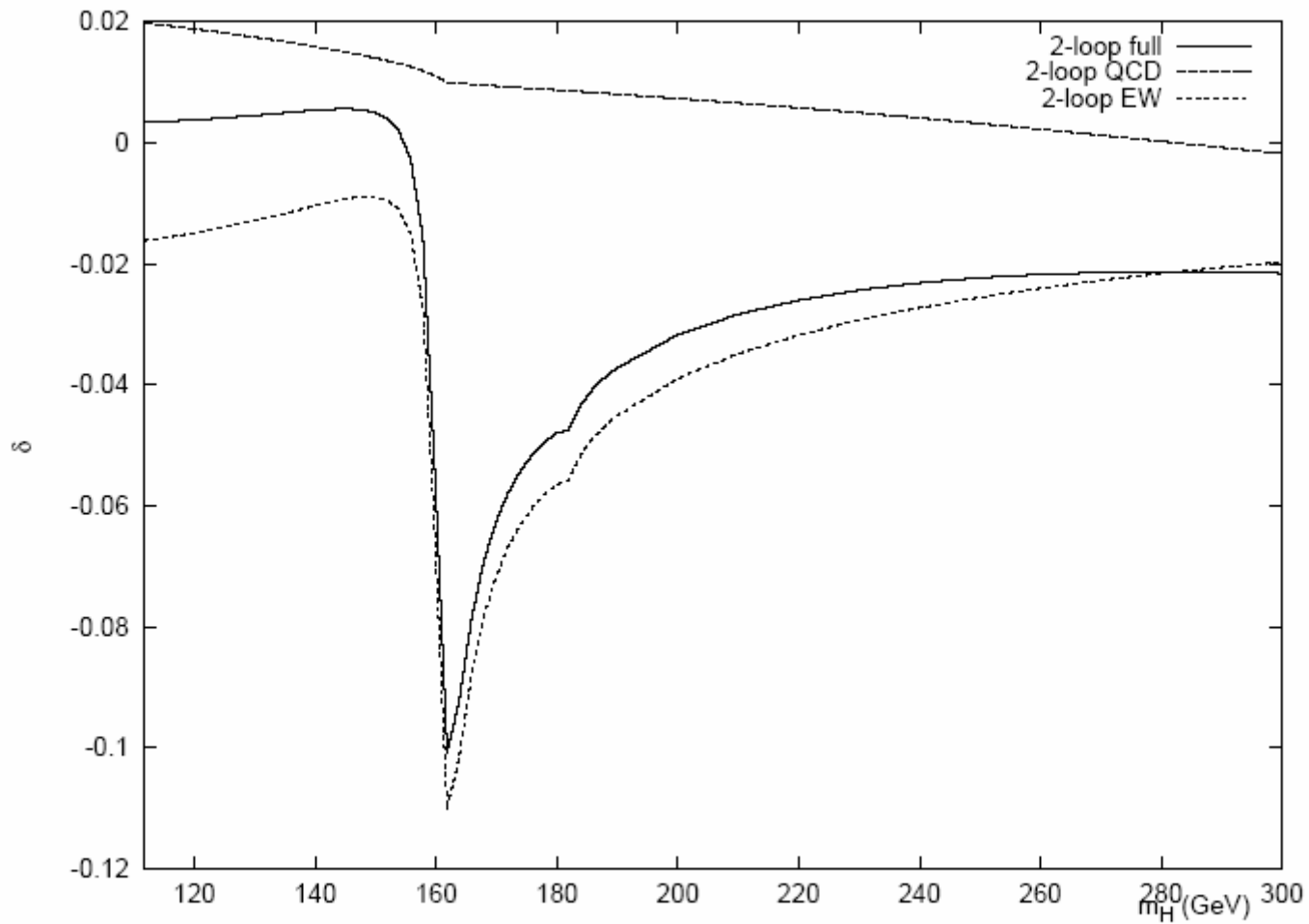
NLO EW correction at most 4%

$H \rightarrow \gamma\gamma$: 2-loop EW light fermion corrections



$H \rightarrow \gamma\gamma$: 2-loop QCD corrections





Relative corrections to the decay width $\Gamma(H \rightarrow \gamma\gamma)$.

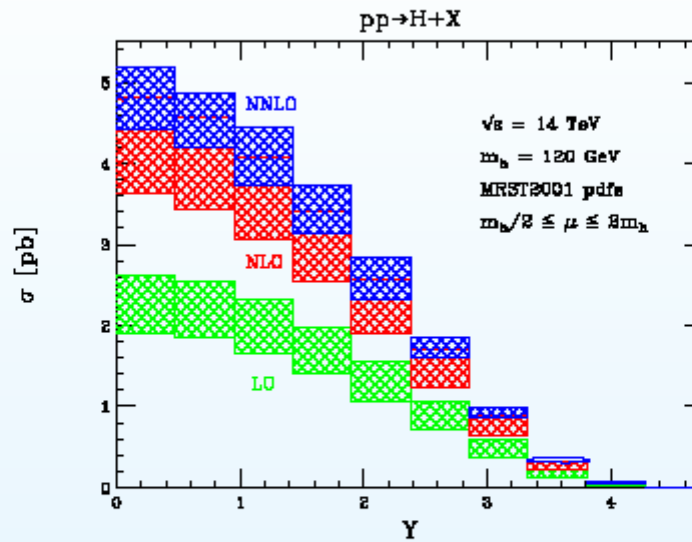
*Higgs boson di-photon signal at the LHC:
realistic K -factors through NNLO in pQCD*

Kirill Melnikov

University of Hawaii at Manoa

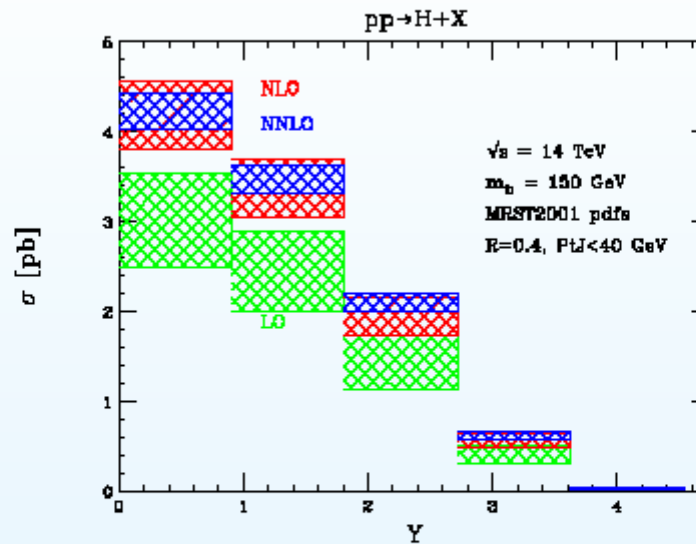
Existence proof: first NNLO calculation of the fully differential cross-section for any process at hadron colliders;

Results: Higgs rapidity distribution



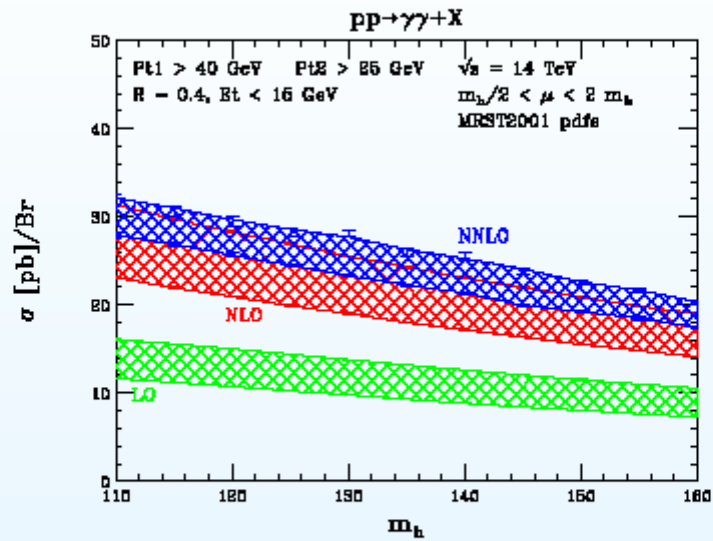
- NNLO corrections important, but things do look convergent;
- Improved stability w.r.t. scale variations;
- Insignificant rapidity dependence of the K -factor.

Results: Higgs rapidity distribution and the jet veto



- Relevant for heavier Higgs discovery ($H \rightarrow WW \rightarrow l\nu l\nu$);
- Only events with $p_{\perp}^J < 40$ GeV are accepted \Leftrightarrow cuts away $pp \rightarrow t\bar{t}$ background;
- Smaller K -factors, compared to inclusive case;
- p_{\perp}^H increases from NLO to NNLO ($p_{\perp}^{H,NLO} = 37.6$ GeV, $p_{\perp}^{H,NNLO} = 44.6$ GeV ; hence, larger fraction of the NNLO cross-section is removed.

Results: realistic di-photon cross-sections



- $p_{\perp}^{\gamma,1} > 40$ GeV and $p_{\perp}^{\gamma,2} > 25$ GeV; $|\eta^{\gamma,1(2)}| < 2.5$.
- Isolation cut: $E_{\perp}^{\text{hadr}} < 15$ GeV for $R < 0.4$.

Results: realistic di-photon signal

- Is it sufficient to know only **inclusive** K -factors?

| m_h , GeV | $\sigma_{\text{NNLO}}^{\text{cut}}/\sigma_{\text{NNLO}}^{\text{inc}}$ | $K_{\text{cut}}^{(2)}/K_{\text{inc}}^{(2)}$ |
|-------------|---|---|
| 110 | 0.559 | 0.93 |
| 115 | 0.589 | 0.95 |
| 120 | 0.601 | 0.95 |
| 125 | 0.632 | 0.98 |
| 130 | 0.669 | 1.00 |
| 135 | 0.668 | 1.00 |

$$K^{(2)} = \frac{\sigma_{\text{NNLO}}}{\sigma_{\text{NLO}}},$$
$$\mu = \frac{m_h}{2}.$$

- The ratio of inclusive to "cut" K -factors is the best guess for the NNLO differential cross-section if only inclusive NNLO K -factor and the differential NLO cross-section were available.
- **The heavier the Higgs, the smaller the impact of the cuts on the K -factor is.**

Conclusions

- New method for NNLO calculations (real radiation) in QCD; applicable to many phenomenologically relevant processes;
- Existence proof: **first NNLO calculation of the fully differential cross-section for any process at hadron colliders**;
- Complete control over the kinematics of the final states allows **arbitrary cuts** to be imposed;
- **State-of-the-art calculation of the Higgs signal in the di-photon channel for the LHC**;
- We plan to extend FEHiP to include $H \rightarrow W^+W^-$ and $H \rightarrow ZZ$ decays in the future. <http://www.phys.hawaii.edu/~kirill/FEHiP.htm>
- We hope FEHiP will be useful in devising strategies to enhance the signal to background ratio and facilitate the Higgs boson discovery.

MSSM Higgs Sector with CP Violation

Marcela Carena
Fermilab

Motivation for CP-Violation in SUSY Models

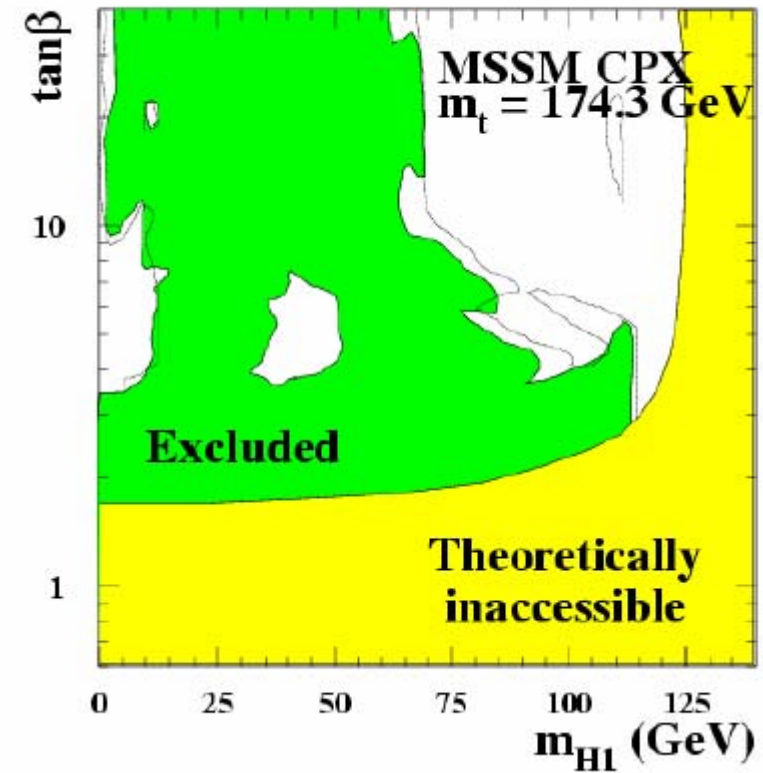
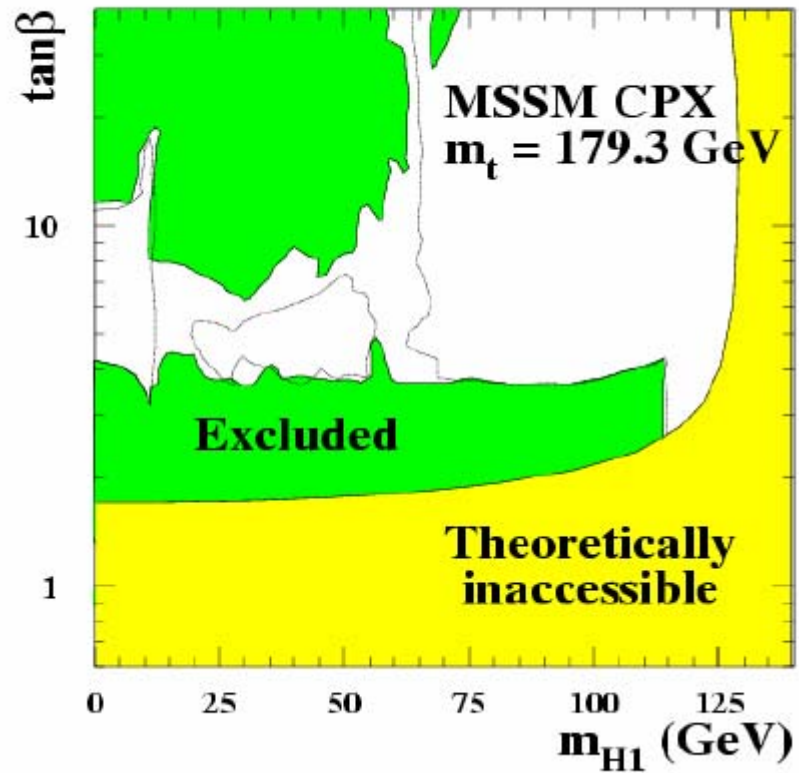
- In low energy SUSY, there are extra CP-violating phases beyond the CKM ones, associated with complex SUSY breaking parameters.

- One of the most important consequences of CP-violation is its possible impact on the explanation of the matter-antimatter asymmetry.

Electroweak baryogenesis may be realized even in the simplest SUSY extension of the SM, but demands new sources of CP-violation associated with the third generation sector and/or the gaugino-Higgsino sector.

- CP-violation in the Higgs sector appears at the loop-level, but can still have important consequences for Higgs physics

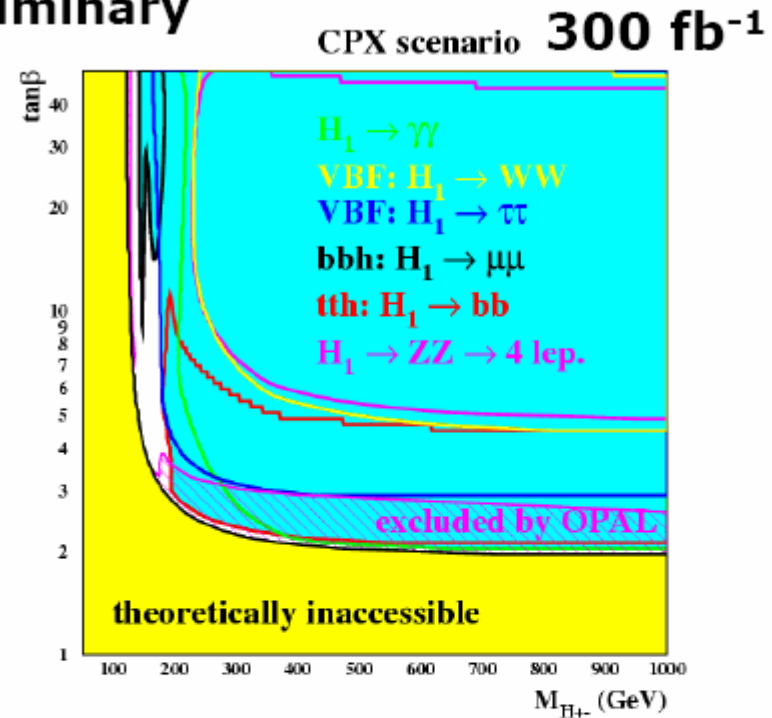
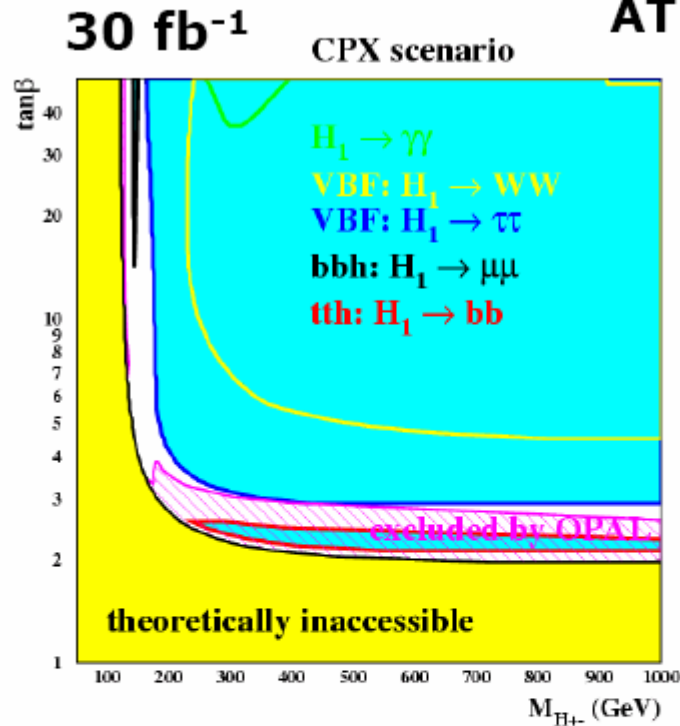
LEP



ZH_1 suppressed and/or $H_2 \rightarrow H_1H_1$ allowed

Light Higgs Boson H_1

ATLAS preliminary



- border of discovery region at low $\tan\beta$ mostly determined by availability of inputs (VBF >110 GeV, ttH and $\gamma\gamma > 70$ GeV)
- border at low $M_{H^{\pm}}$ due to decoupling of H_1 from W,Z and t
- for VBF channels: assume same efficiencies for contribution of CP even and CP odd states (needs to be checked)
- for ttH: efficiencies for CP even and odd bosons are the same

M. Schumacher, SUSY 04

Tevatron/LHC

Looking for $H_2 \rightarrow H_1 H_1$

- Standard signatures not sufficient to probe the presence of Higgs bosons decaying into lighter Higgs states.
- Lighter states have weak couplings to the weak gauge bosons, but large couplings to third generation down quarks and leptons.
- Possibility of looking for two taus and two bottoms (jets) signatures at LHC in the weak boson fusion production channel of two CP-odd like Higgs bosons. (J. Gunion et al. with 300 inverse fb at the LHC)
- Encourage the study of $gg \rightarrow H_2$, $t\bar{t}H_2$, $W/Z H_2$ and $WW/ZZ H_2$ with subsequent decay $H_2 \rightarrow H_1 H_1$, using the extra leptons from W/Z 's.

CPsuperH

- Code to compute Higgs spectrum, couplings and decay modes in the presence of CP-violation

Lee, Pilaftsis, M.C., Choi, Drees, Ellis, Lee, Wagner.'03

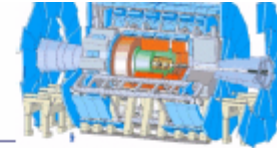
- CP-conserving case: Set phases to zero. Similar to HDECAY, but with the advantage that charged and neutral sector treated with same rate of accuracy.
- Combines calculation of masses and mixings by M.C., Ellis, Pilaftsis, Wagner. with analysis of decays by Choi, Drees, Hagiwara, Lee and Song.
- Available at

<http://theory.ph.man.ac.uk/~jslee/CPsuperH.html>

Prospects for the Search for a Doubly Charged Higgs with ATLAS

Kamal Benslama
Columbia University

Physics Motivation



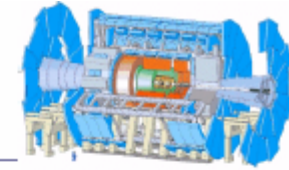
L-R symmetric model would be a natural extension of the SM

- $SU(2)_L \times SU(2)_R \times U(1)_{B-L}$
- predicts new fermions:
heavy Majorana neutrino
- predicts new gauge bosons:
 W_R
- predicts new Higgs sector

$$\Delta_R = (\Delta_R^0, \Delta_R^+, \Delta_R^{++})$$

$$\Delta_L = (\Delta_L^0, \Delta_L^+, \Delta_L^{++}) \text{ (if Lagrangian is invariant under } L \leftrightarrow R \text{ symmetry)}$$

$$\Delta_L^{++} \rightarrow l^+ l^+$$



Signal : consider $\Delta_L^{++} \rightarrow l^+ l^+$, only e or μ
 (cannot decay to quark pair because of charge conservation)

Backgrounds : $Wt\bar{t}$, $W^+W^+(QCD)$, $W^+W^+(EW)$, $WZqq$, $t\bar{t}$

$$n_{lep} = 2$$

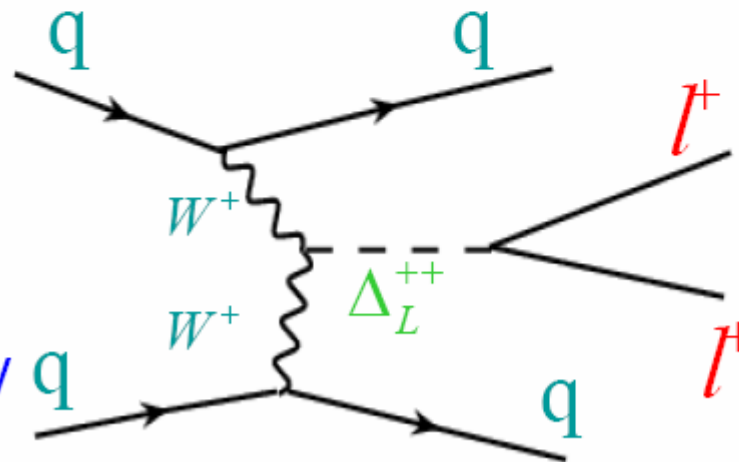
$$\Delta\phi_{ll} > 2.5$$

$$\Delta P_T^{ll} = f(M_{ll})$$

Forward jets tagging:

$E_{j1} > 200$ GeV, $E_{j2} > 200$ GeV
 well separated in eta

↖ Main background

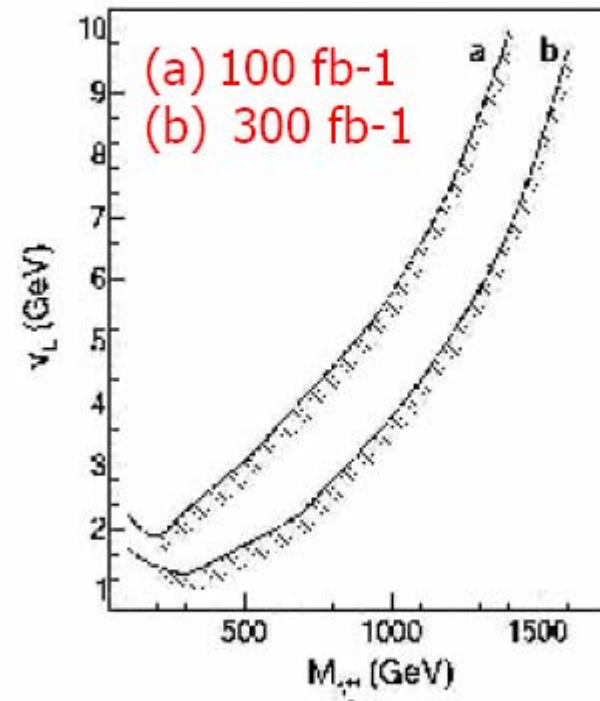


Missing transverse momentum cut

Results



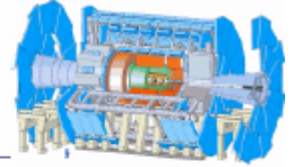
| 100 fb-1 | Δ^{++} 300 GeV | Δ^{++} 800 GeV | total backg |
|---|--------------------------|--------------------------|-------------|
| Isolated leptons | 330 (384) | 59 (69) | 133/13 |
| $ \Delta\phi_{ll} > 2.5 $ | 253 (289) | 56 (65) | 75/8.3 |
| $\Delta_{P_T}^{ll} > (\frac{M_{ll}}{2} + 50)$ | 220 (260) | 50 (59) | 37/2.5 |
| Fwd Jet tagging | 156(185) | 40 (47) | 17/1.4 |
| ptmiss | 152(180) | 34 (40) | 3.0/0.1 |



Nb of Events after each cut

Discovery reach

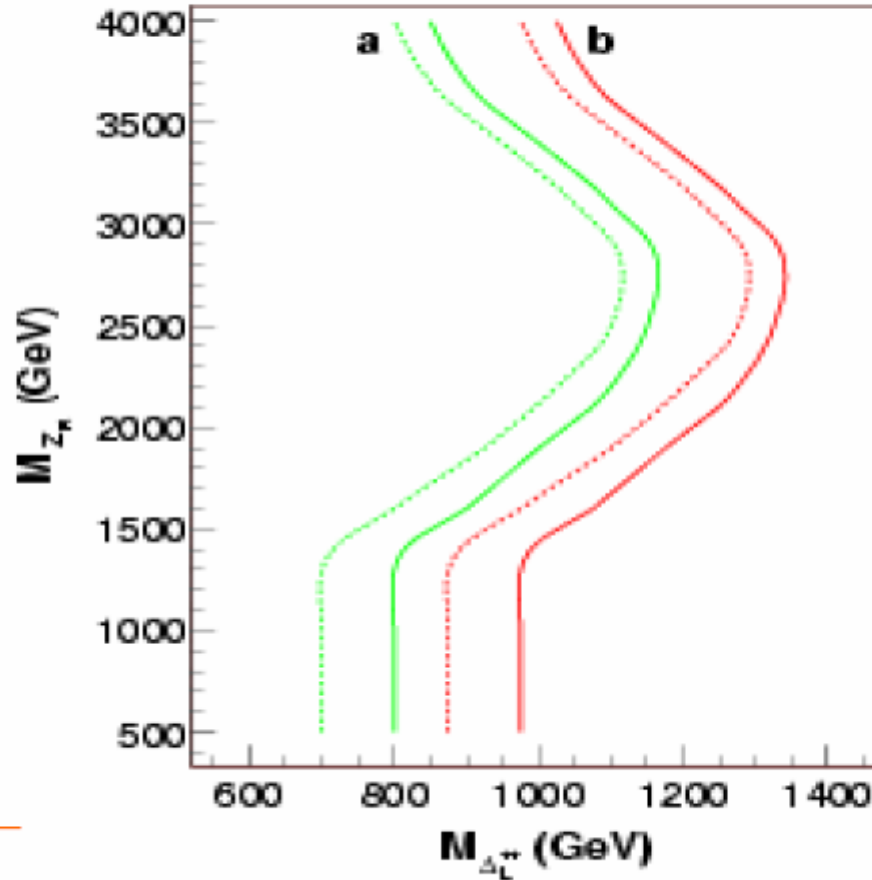
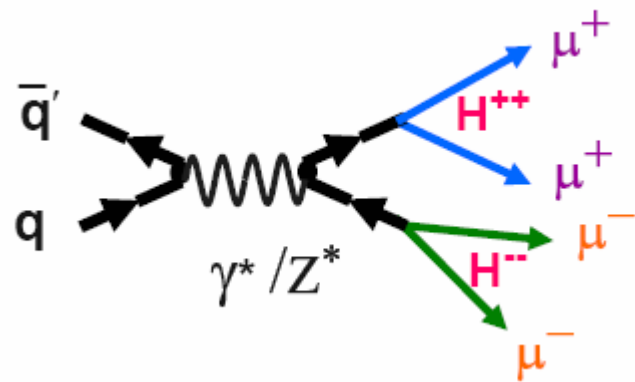
$$\Delta_L^{++} \Delta_L^{--} \rightarrow l^+ l^+ l^- l^-$$



- (a) 100 fb-1
- (b) 300 fb-1

Full = 3 leptons are observed

Dashed = 4 leptons are observed



Similar results for Δ_R^{++}

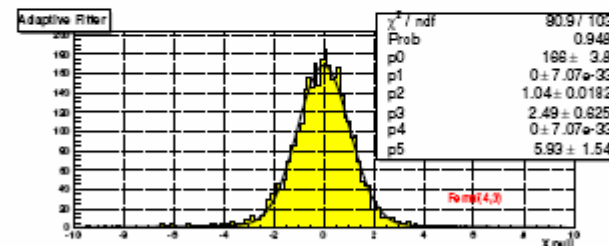
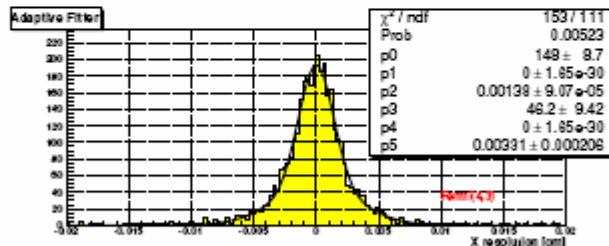
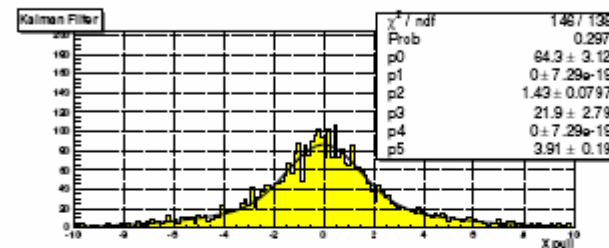
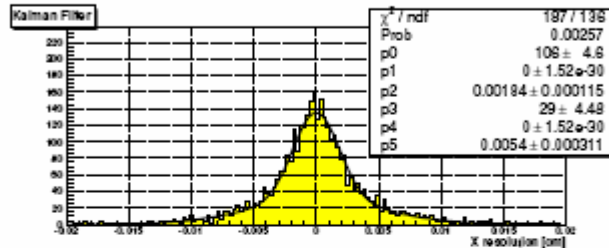
Robust Vertex Reconstruction in Heavy Flavor Events

Ariel Schwartzman
Princeton University

- Standard methods for primary vertex reconstruction:
 - Biased by the presence of secondary vertex tracks. Square cuts designed to reject outliers have the effect of reducing the vertex resolution by removing true primary vertex tracks.
- The Adaptive Fitter:
 - Robust, iterative Kalman Filter. Secondary vertex tracks are down-weighted, depending on the x^2 contribution to the vertex and the presence of all other neighbor tracks. All tracks are considered in the fit.

Adaptive algorithm down-weights SV tracks from small decay length vertices.

Adaptive Primary Vertex Performance (I)



- Multiplicity = 18.7
- Resolution = 13.8 um.
- Pull = 1.04.
- Tails = 7.9%

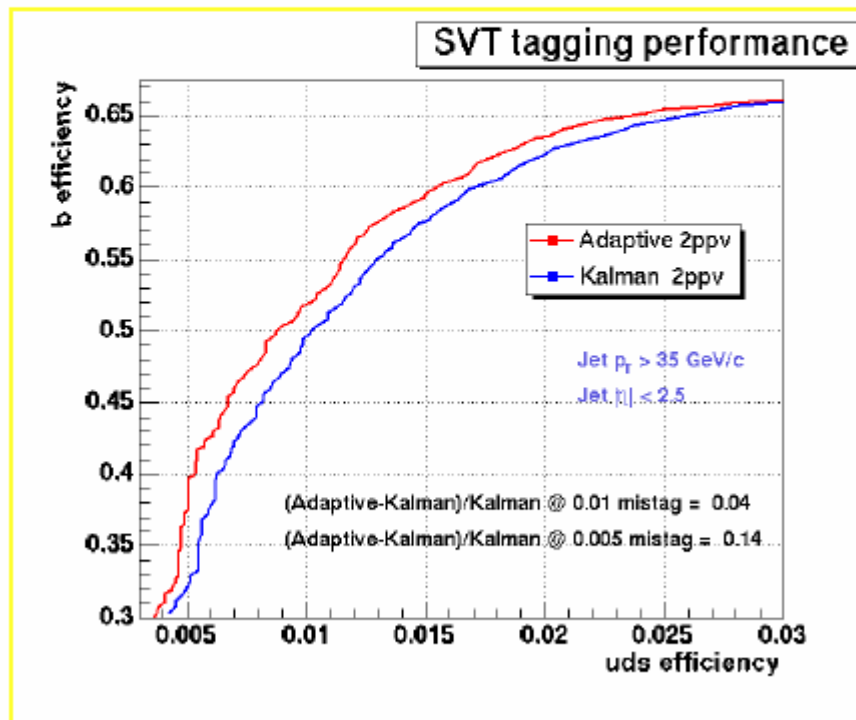
Zbb simulated events

| | | |
|-------------------------|--------|--------|
| Resolution improvement: | 18.3um | 13.8um |
| Pull improvement: | 1.43 | 1.04 |
| Tails improvement: | 48.2% | 7.9% |

$$\text{Resolution} = X(\text{reco}) - X(\text{true})$$

$$\text{Pull} = (X(\text{reco}) - X(\text{true})] / \text{Sigma}X(\text{reco})$$

Secondary Vertex Performance using Adaptive Primary Vertexing



New Adaptive PV allows a better separation between primary and secondary vertices.

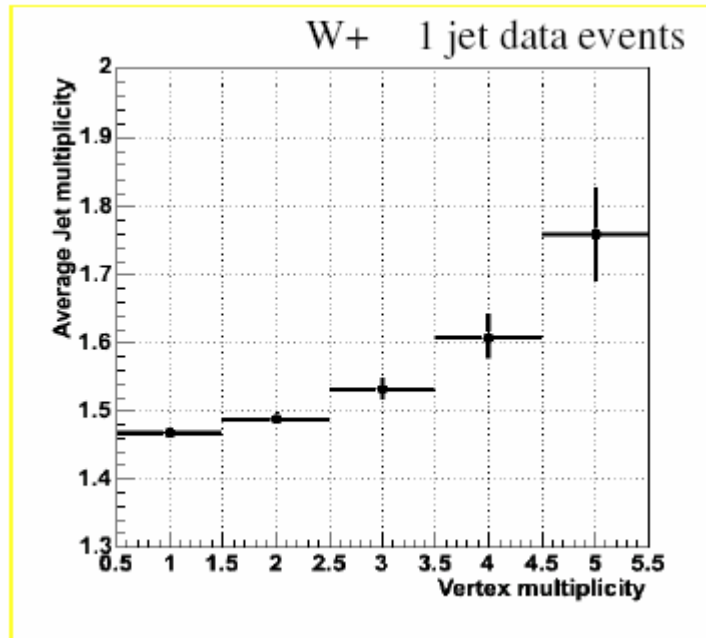
At a same mistag rate, secondary vertex b-tagging efficiency in the simulation is improved.

Further improvement expected by using SV adaptive fitting.

Improving Calorimeter Jet Reconstruction Using Tracks and Vertices

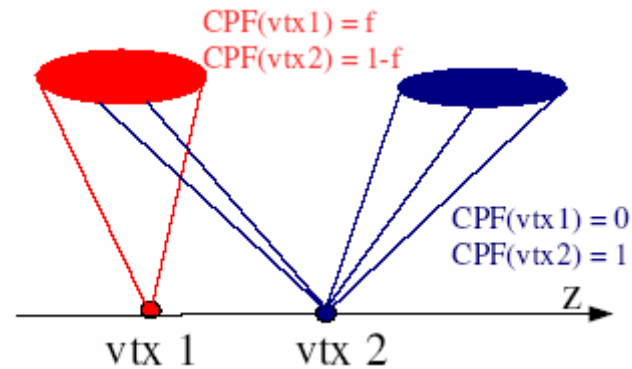
Ariel Schwartzman
Princeton University

Jet Multiplicity at High Luminosity



- Average jet multiplicity increases with number of primary vertices.
- Jet kinematics is distorted at low p_T (extra low- p_T jets)
 - Additional minimum bias interactions can give rise to extra jets.

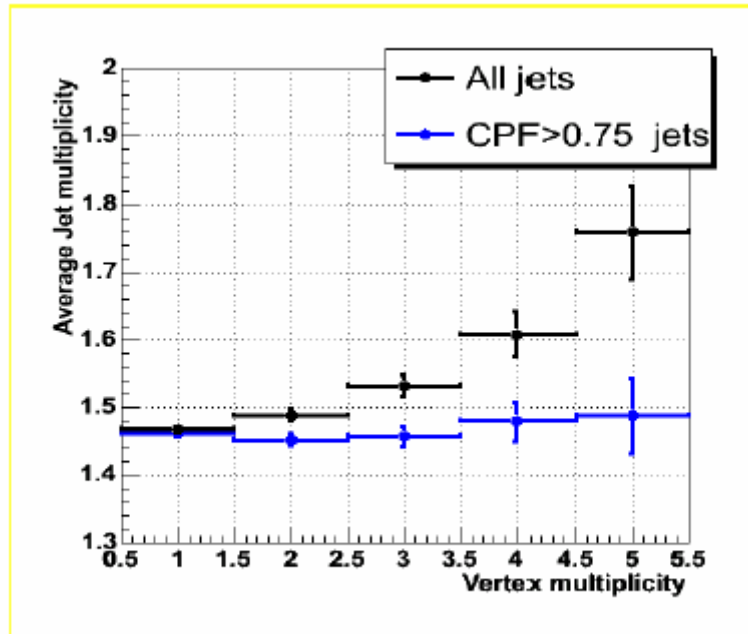
Jet-Vertex Identification



associate *jet j* to the vertex *i* for which $CPF(jet_j, vtx_i)$ is maximum.

CPF is the charged particle energy fraction of jet *j* from vertex *i*.

Jet Multiplicity at High Luminosity



The dependence of jet multiplicity on the number of interactions is significantly reduced for jets with more than 75% of its energy coming from the hard scatter vertex.

Jet-Vertex algorithm allows to identify (remove) soft jets arising from min-bias interactions.

Improving Jet Energy Resolution Using Tracks

Classify tracks:

$R(\text{vtx}) < 0.5$, $R(\text{cal}) < 0.5$: **IN** jet

$R(\text{vtx}) < 0.5$, $R(\text{cal}) > 0.5$: **Out-of-cone**

Technique: for every charged hadron (track) matched to a jet:

- *subtract* the expected energy deposited in the calorimeter.
- *add* the track momentum.

$$E_{jet} = E_{jet} + (1 - R^{cal}(E_{track}))E_{track}$$

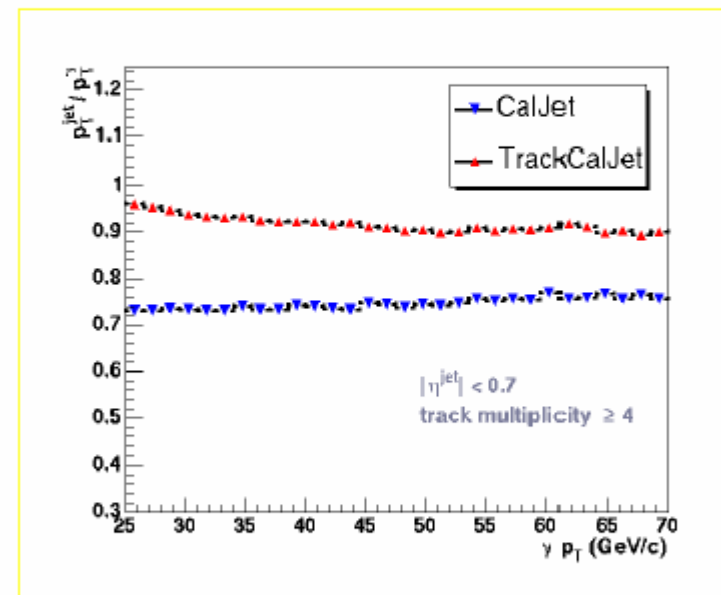
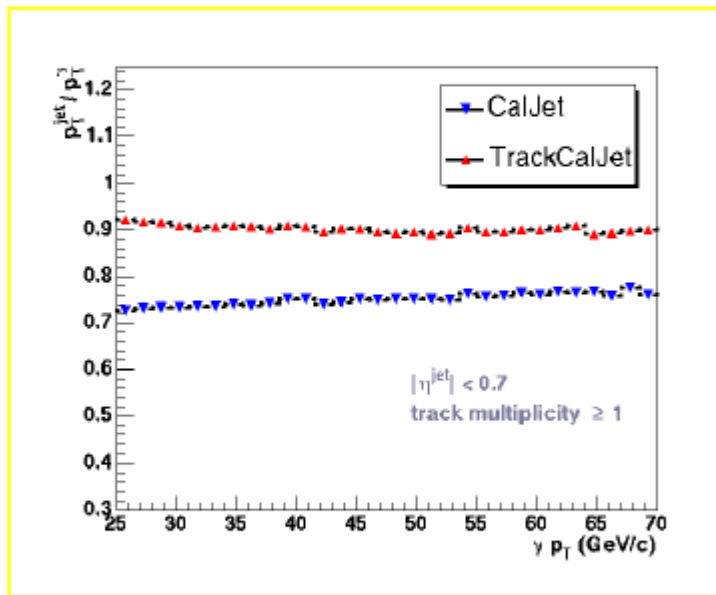
$$R^{cal}(E_{track}) = E^{cal} / E^{track}$$

is the single pion
calorimeter response.

- *Add* the energy of out-of-cone tracks.

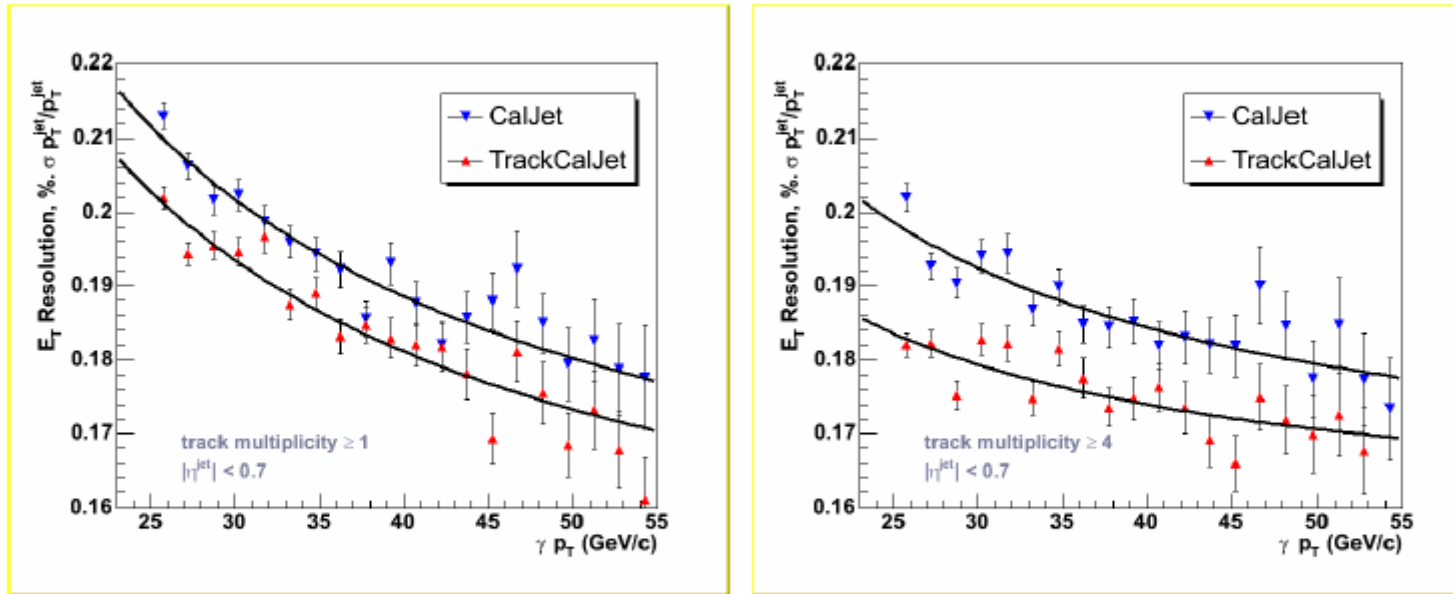
$$E_{jet} = E_{jet} + E_{track}$$

Jet Energy Scale in photon+jet Data



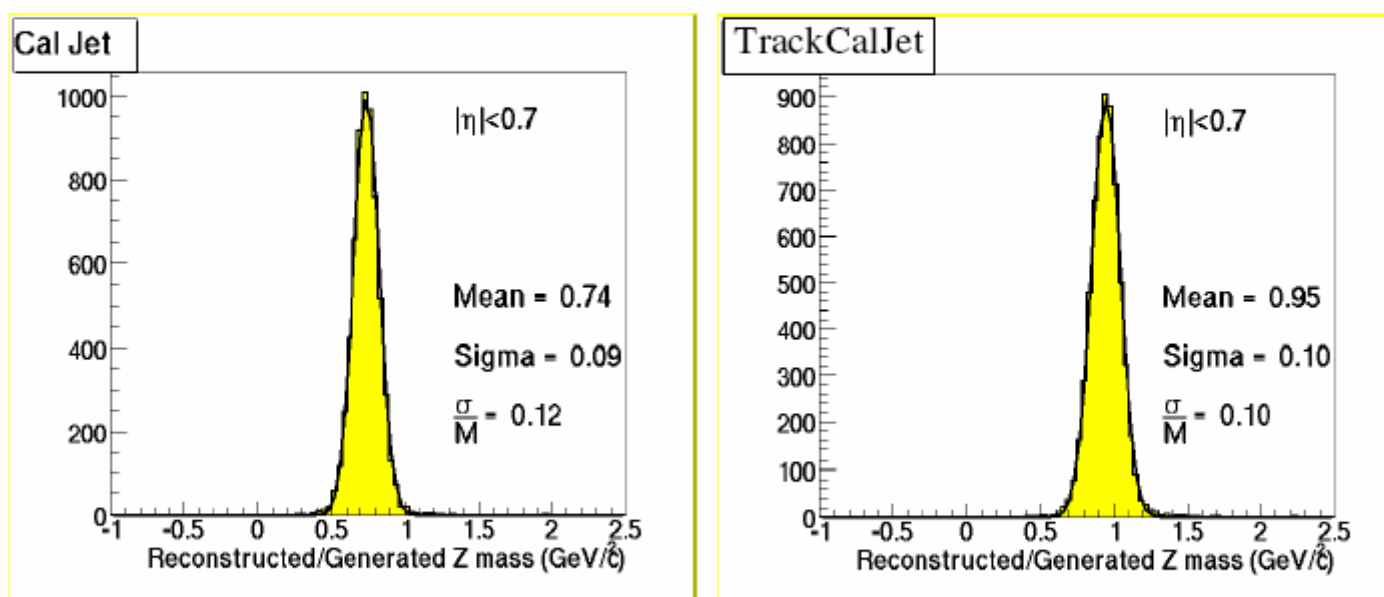
Jet energy offset is significantly improved with the use of tracks.

Jet Energy Resolution in photon+jet Data



Jet energy resolution dependence on Jet p_T for low and high track-multiplicity jets. **$\sim 10\%$ Improvement.**

Dijet Mass Resolution in Zqq events

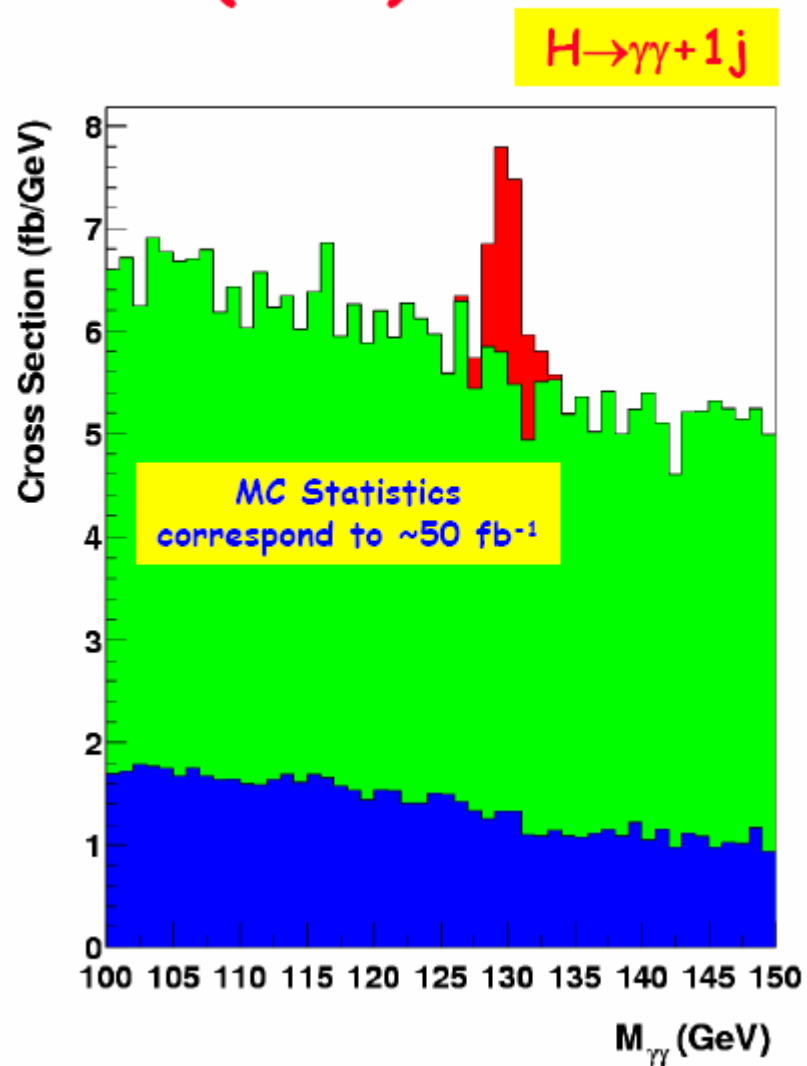
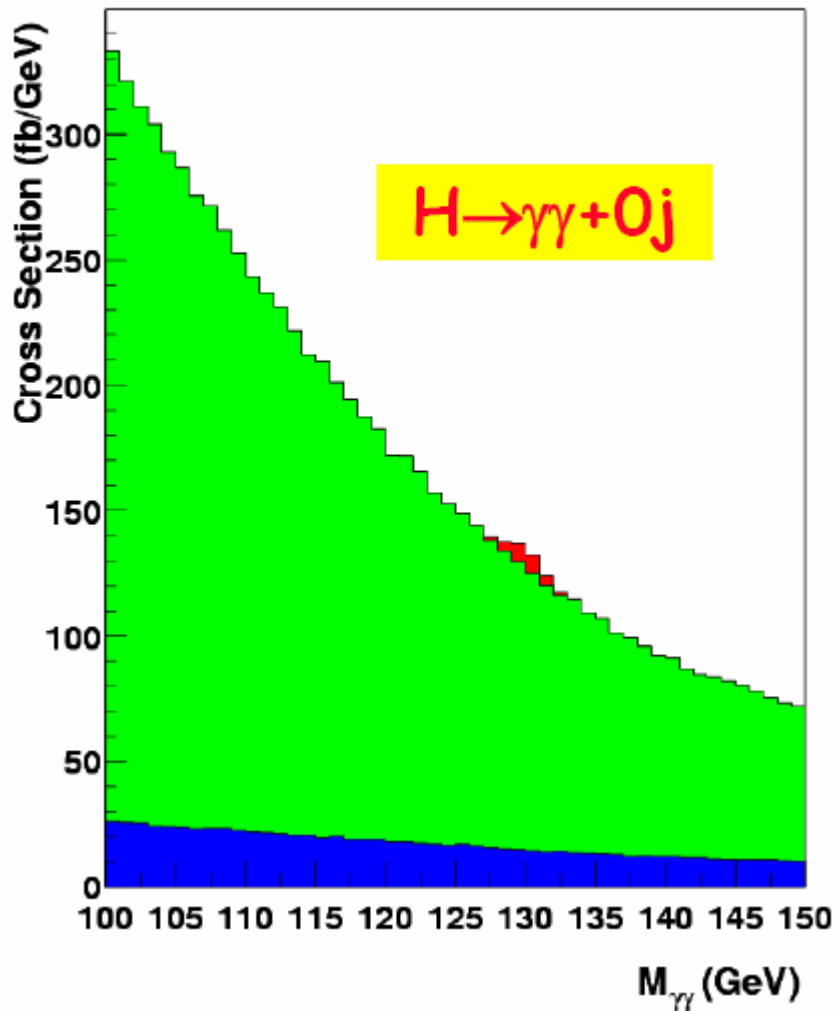


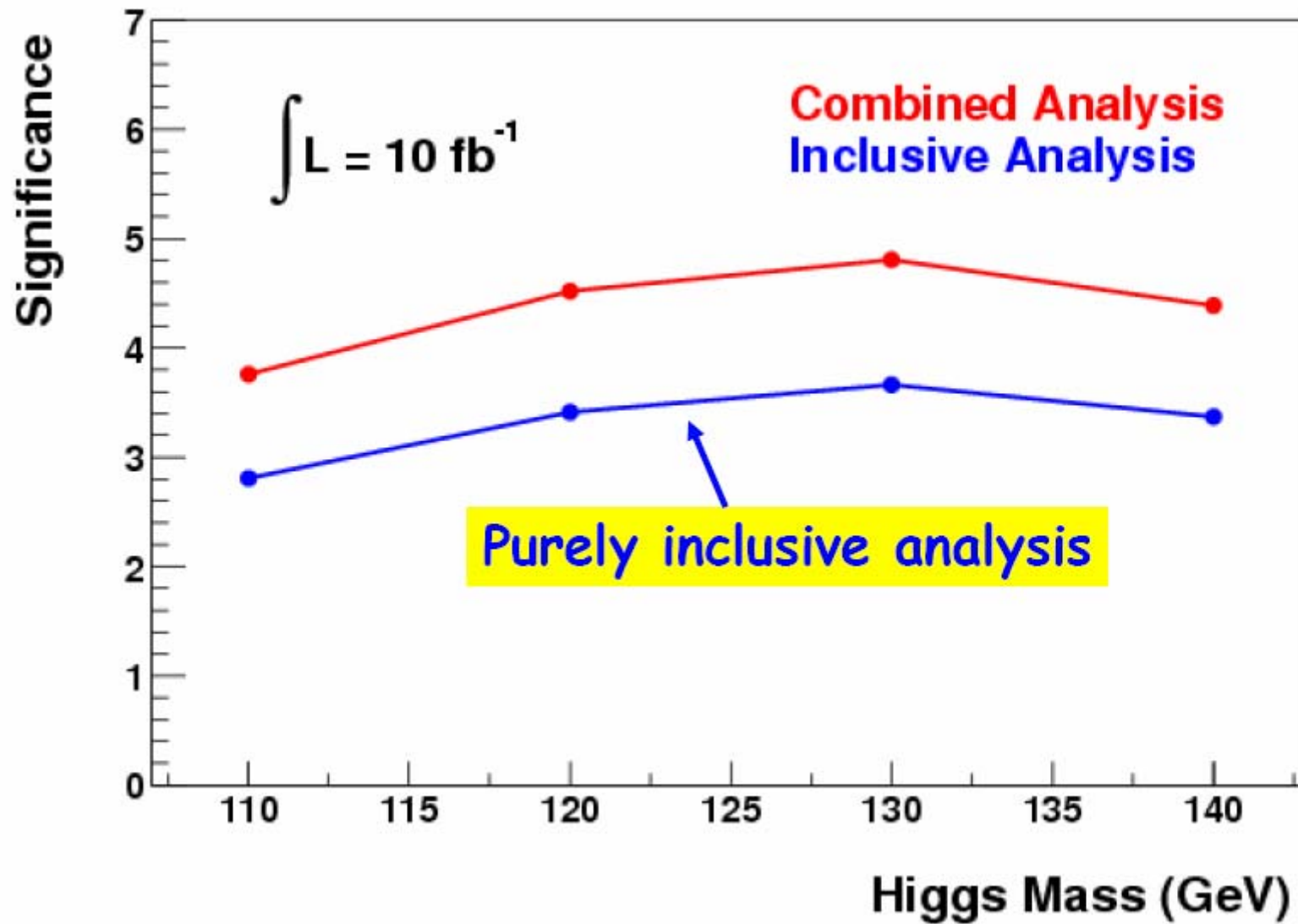
15% mass resolution improvement.

Higgs + 1jet Signatures

Yaquan Fang, Bruce Mellado, William Quayle
and Sau Lan Wu
University of Wisconsin-Madison

Optimization (cont)

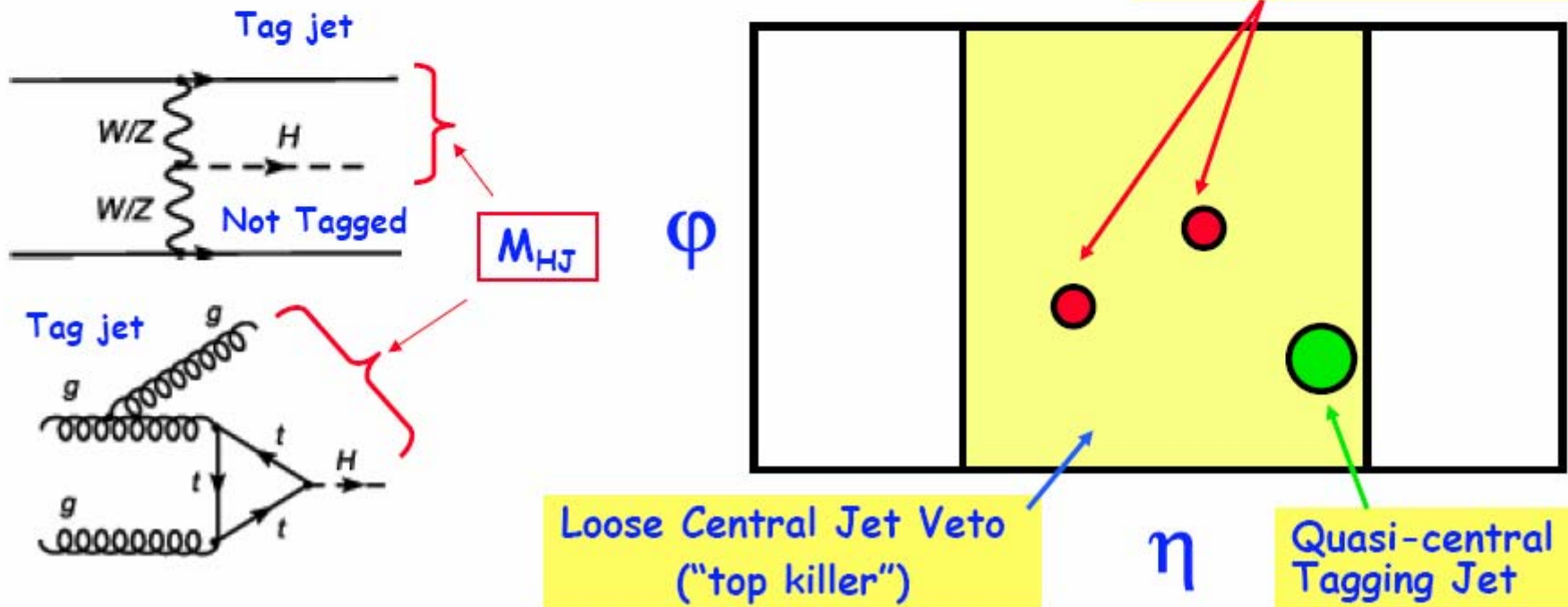


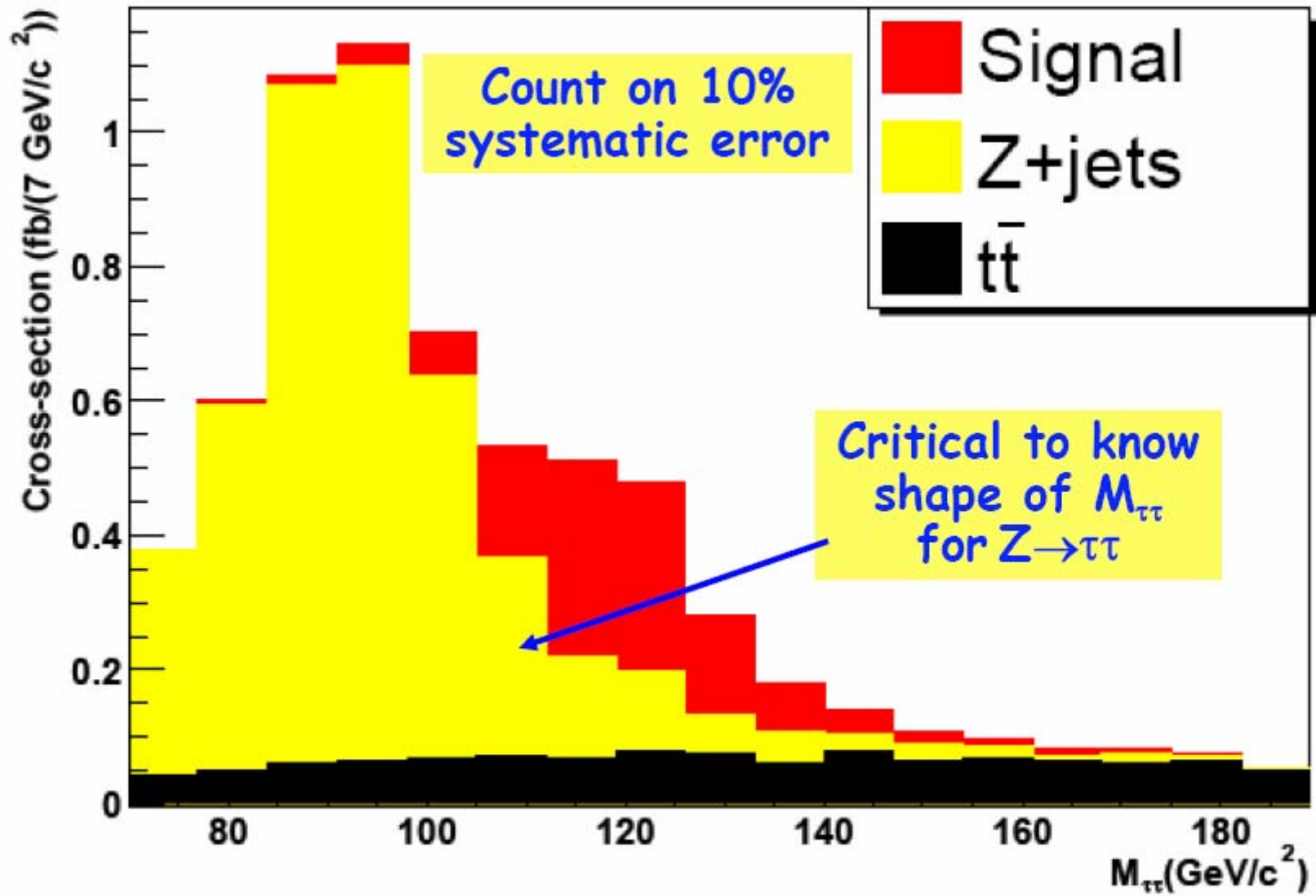


H($\rightarrow\tau\tau$)+1jet at the LHC

✚ Tag one semi-central jet, require $P_{TH} > 100$ and $M_{HJ} > 700$ GeV and a loose central jet veto ("top killer")

➤ Allow significant contribution from $gg \rightarrow h$





Signal Significance for 30 fb⁻¹

($M_H=120$ GeV 10% systematic error on background)

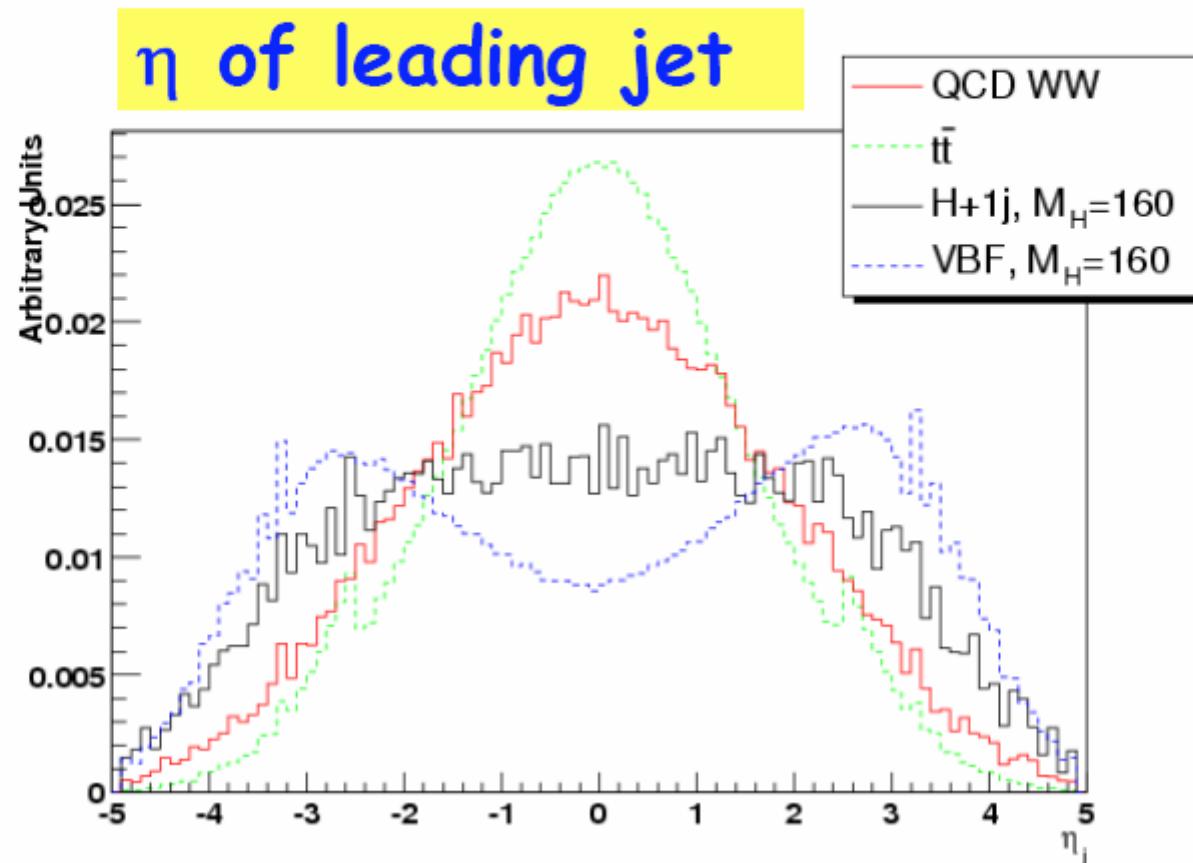
| Higgs Mass (GeV/c ²) | 110 | 120 | 130 | 140 | 150 |
|---|-----|-----|-----|-----|-----|
| Signal Significance for cut analysis (σ) | 4.3 | 5.0 | 4.8 | 3.6 | 2.1 |
| Signal Significance for NN analysis (σ) | 5.5 | 6.6 | 6.3 | 4.8 | 2.8 |

Overlap With H+2jets (VBF)

the overlap is 24%

$H \rightarrow WW^{(*)} + \geq 1 \text{ jet}$

- Use basic property that leading jet associated with Higgs tends to be more forward than in QCD backgrounds



Conclusion

We still have a lot to learn
before the LHC turns on.

If you think we're busy now,
wait until we have data from the LHC