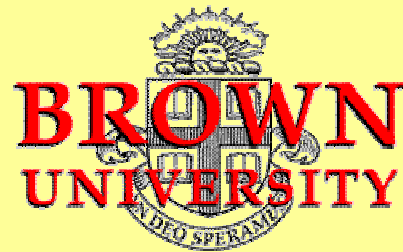


# Probing Quantum Gravity at Colliders

**Greg Landsberg**



**TeV4 LHC Workshop**

*February 3-5, 2005 at Brookhaven National Laboratory*

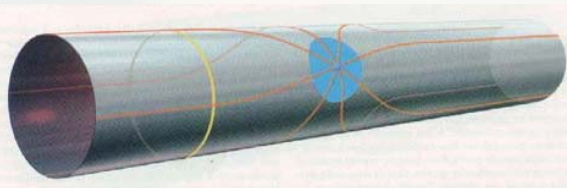
# Outline

- Realms of Quantum Gravity
- Some Lessons of the Tevatron
- Beyond the Tevatron
- Conclusions

# Models of Extra Dimensions

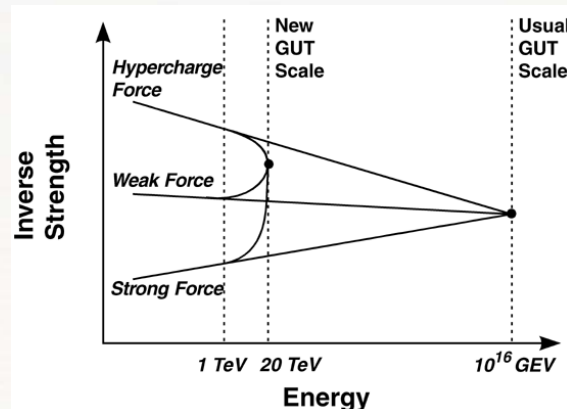
## ADD Scenario:

- Pro: “Eliminates” the hierarchy problem by stating that physics ends at a TeV scale
- Only gravity lives in the “bulk” space
- Size of ED’s ( $n=2-7$ ) between  $\sim 100 \mu\text{m}$  and  $\sim 1 \text{ fm}$
- Black holes at the LHC
- Con: Doesn’t explain how to make ED large



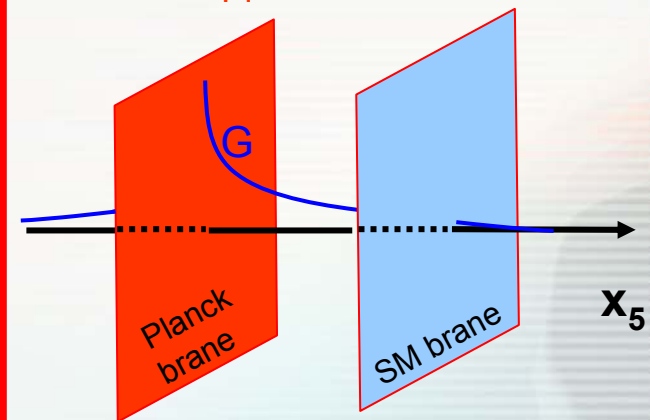
## Small ED:

- Pro: Could lower GUT scale by changing the running of couplings
- At least gauge bosons ( $g/\gamma/W/Z$ ) “live” in ED’s
- Size of ED’s  $\sim 1 \text{ TeV}^{-1}$  or  $\sim 10^{-19} \text{ m}$
- Con: Gravity is not in the picture



## RS Model:

- Pro: A rigorous solution to the hierarchy problem via localization of gravity
- Gravitons (and possibly other particles) propagate in a single ED, w/ special metric
- Higgs-like particle: radion
- Con: Size of ED as small as  $\sim 1/M_{\text{Pl}}$  or  $\sim 10^{-35} \text{ m}$

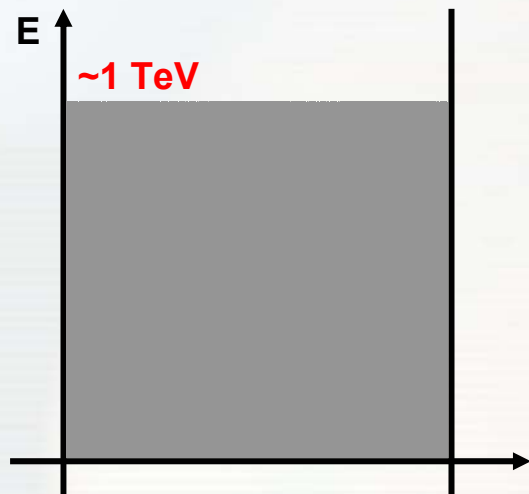




# Kaluza-Klein Spectrum

## ADD Model:

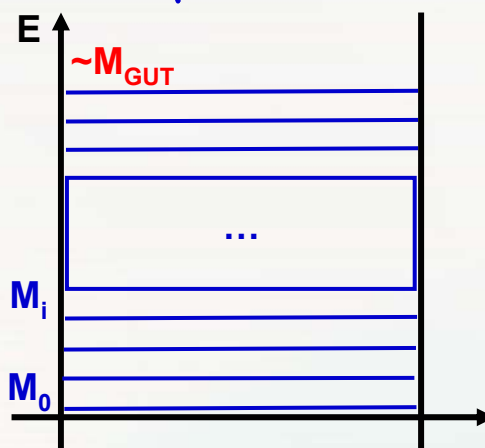
- Winding modes with energy spacing  $\sim 1/r$ , i.e. 1 meV – 100 MeV
- Can't resolve these modes – they appear as continuous spectrum
- Coupling:  $G_N$  per mode; compensated by large number of modes



## Small ED:

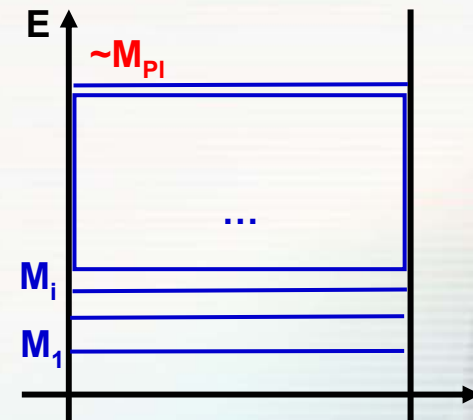
- Winding modes with nearly equal energy spacing  $\sim 1/r$ , i.e.  $\sim \text{TeV}$
- Can excite individual modes at colliders or look for indirect effects
- Coupling:  $\sim g_w$  per mode

$$M_i = \sqrt{M_0^2 + i^2/r^2}$$



## RS Model:

- “Particle in a box” with special AdS metric
- Energy eigenvalues are given by zeroes of Bessel function  $J_1$
- Light modes might be accessible at colliders
- Coupling:  $G_N$  for zero mode;  $1/\Lambda_\pi^2$  for the others

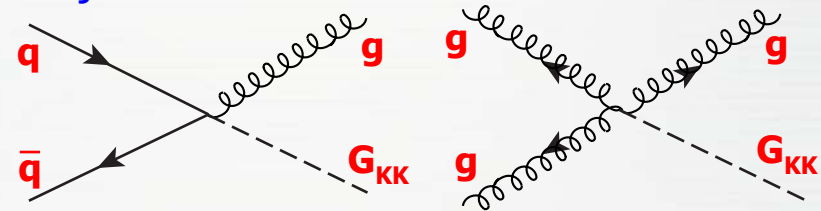


$$M_0 = 0; M_i = M_1 x_i/x_1 \approx M_1, 1.83M_1, 2.66M_1, 3.48M_1, 4.30M_1, \dots$$

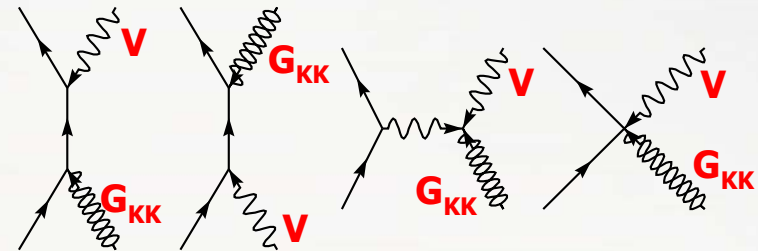
# Collider Signatures for ED

- Kaluza-Klein gravitons couple to the energy-momentum tensor, and therefore contribute to most of the SM processes
- For Feynman rules for  $G_{KK}$  see:
  - [Han, Lykken, Zhang, PRD **59**, 105006 (1999)]
  - [Giudice, Rattazzi, Wells, NP **B544**, 3 (1999)]
- Since graviton can propagate in the bulk, energy and momentum are not conserved in the  $G_{KK}$  emission from the point of view of our 3+1 space-time
- Depending on whether the  $G_{KK}$  leaves our world or remains virtual, the collider signatures include single photons/Z/jets with missing  $E_T$  or fermion/vector boson pair production
- Graviton emission: direct sensitivity to the fundamental Planck scale  $M_D$
- Virtual effects: sensitive to the ultraviolet cutoff  $M_S$ , expected to be  $\sim M_D$  (and likely  $< M_D$ )
- The two processes are complementary

## Real Graviton Emission (Large ED) Monojets at hadron colliders

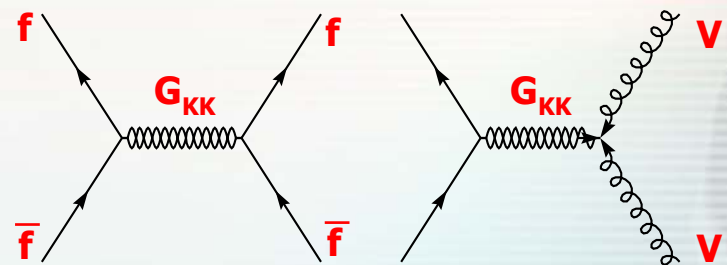


## Single VB at hadron or $e^+e^-$ colliders



## Virtual Effects (all ED)

### Fermion or VB pairs at hadron or $e^+e^-$ colliders



# L'EPilogue (Large ED)

Experiment	$e^+e^- \rightarrow \gamma G$					$e^+e^- \rightarrow ZG$					Color coding
	n=2	n=3	n=4	n=5	n=6	n=2	n=3	n=4	n=5	n=6	
ALEPH	1.28	0.97	0.78	0.66	0.57	0.35	0.22	0.17	0.14	0.12	≤184 GeV
DELPHI	1.38	<b>1.02</b>	0.84	<b>0.68</b>	0.58	<del> </del>	<del> </del>	<del> </del>	<del> </del>	<del> </del>	≤189 GeV
L3	1.02	0.81	0.67	0.58	0.51	0.60	0.38	0.29	<b>0.24</b>	<b>0.21</b>	>200 GeV
OPAL	1.09	0.86	0.71	0.61	0.53	<del> </del>	<del> </del>	<del> </del>	<del> </del>	<del> </del>	$\lambda=-1$ $\lambda=+1$

## Virtual Graviton Exchange

Experiment	$e^+e^-$	$\mu^+\mu^-$	$\tau^+\tau^-$	$qq$	$ff$	$\gamma\gamma$	$WW$	$ZZ$	Combined
ALEPH	1.04 0.81	0.65 0.67	0.60 0.62	0.53/0.57 0.46/0.46 (bb)	1.05 0.84	0.81 0.82	<del> </del>	<del> </del>	0.75/1.00 (<189)
DELPHI	<del> </del>	0.59 0.73	0.56 0.65	<del> </del>	0.60 0.76	0.83 0.91	<del> </del>	<del> </del>	0.60/0.76 (ff) (<202)
L3	0.98 1.06	0.56 0.69	0.58 0.54	0.49 0.49	0.84 1.00	0.99 0.84	0.68 0.79	<del> </del>	1.1/1.0 (<202)
OPAL	1.15 1.00	0.62 0.66	<del> </del>	<del> </del>	0.62 0.66	0.89 0.83	<del> </del>	0.63 0.74	1.17/1.03 (<209)

LEP Combined: 1.2/1.1

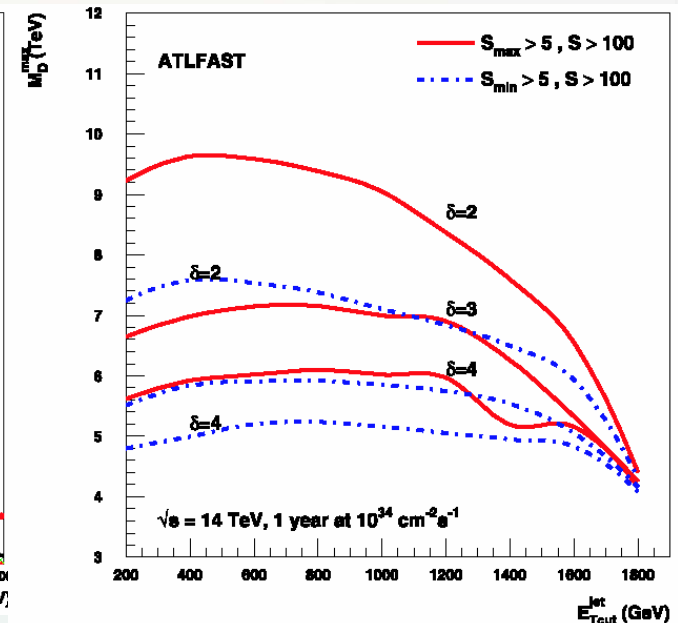
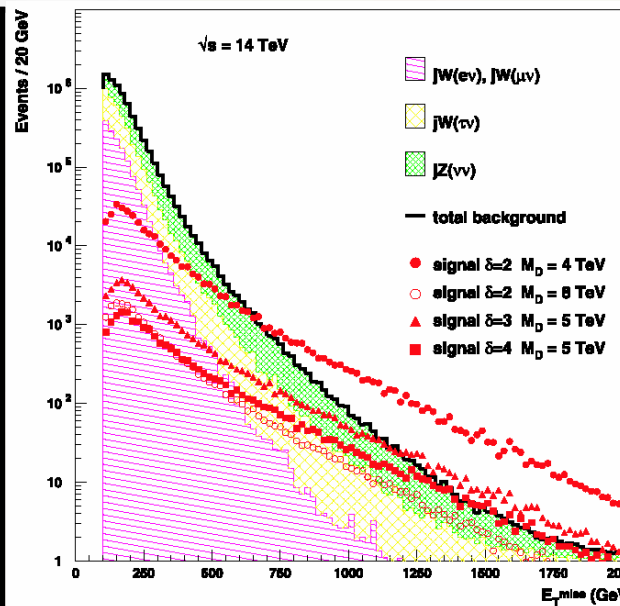
# Tevatron Lessons

- Sensitivity achieved in Run II has superseded combined sensitivity of LEP and continues to increase
- Planck scale up to  $\approx 2$  TeV will be probed in the next couple years before the start of the LHC
- What are some of the lessons that we have learned?

# Lesson 1: Monojets are Tough!

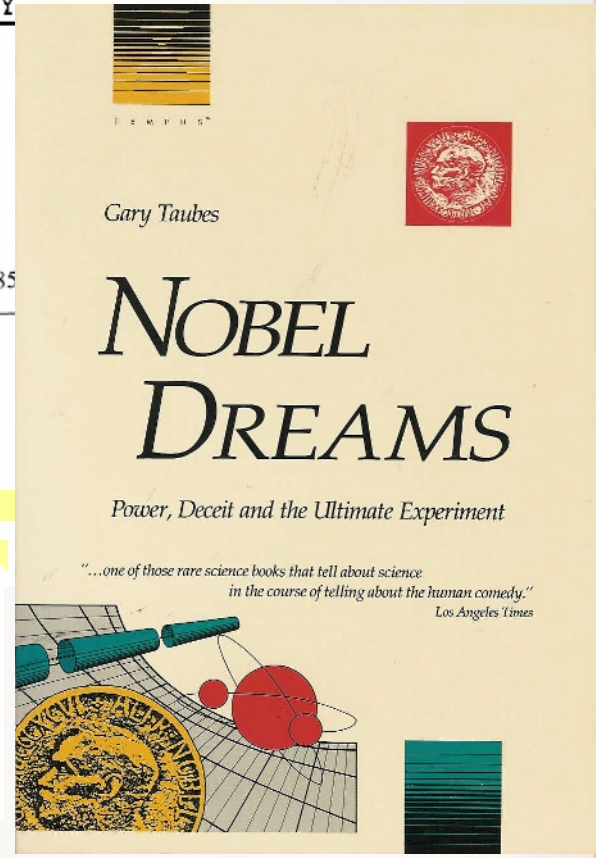
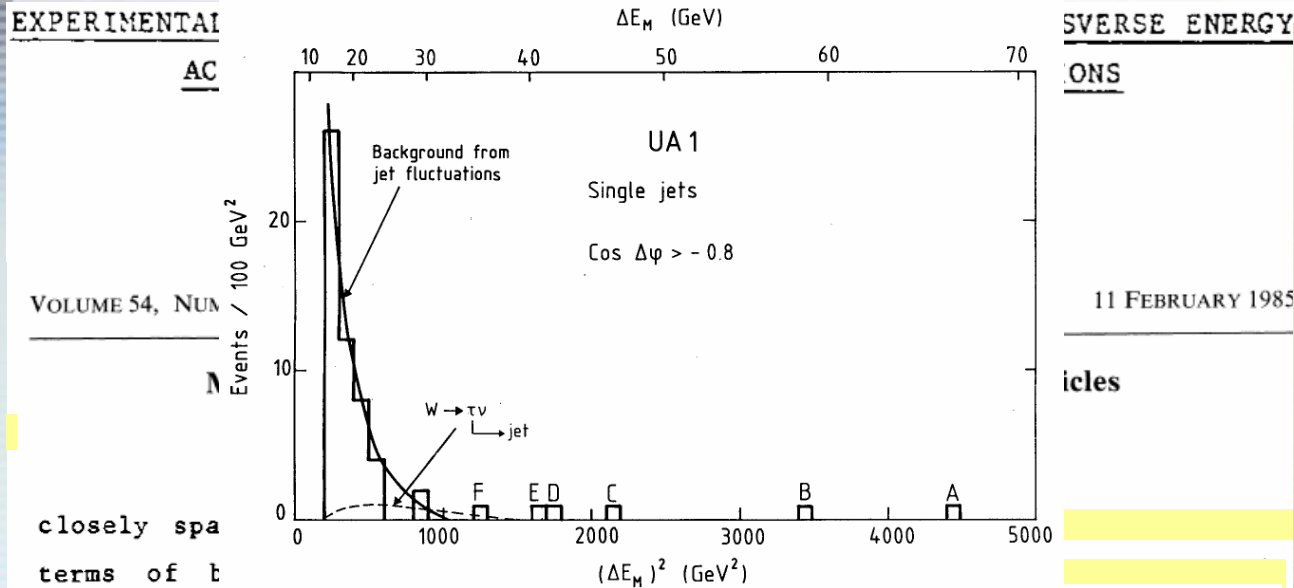
- Production of single jets is a classical signature for large ED, and one of the first studied phenomenologically [Giudice, Rattazzi, Wells, Nucl. Phys. **B544**, 3 (1999)]; ATLAS studies [Vacavant/Hinchliffe, J. Phys. G: Nucl. Part. Phys. **27**, 1839 (2001)]
- Nevertheless, it was one of the last signatures actually probed experimentally by DØ [PRL **90**, 251802 (2003)] and CDF [PRL **92**, 121802 (2004)]
- Challenges are **understanding and suppression of instrumental backgrounds and non-Gaussian tails** in  $ME_T$ 
  - These challenges are not unique to the Tevatron and I expect them to be quite important at the LHC as well (or so history taught us!)

n	$M_D$ reach Run I	$M_D$ reach Run II	$M_D$ reach LHC 100 fb <sup>-1</sup>
2	1100 GeV	1400 GeV	8.5 TeV
3	950 GeV	1150 GeV	6.8 TeV
4	850 GeV	1000 GeV	5.8 TeV
5	700 GeV	900 GeV	5.0 TeV





# Monojets: Tainted History

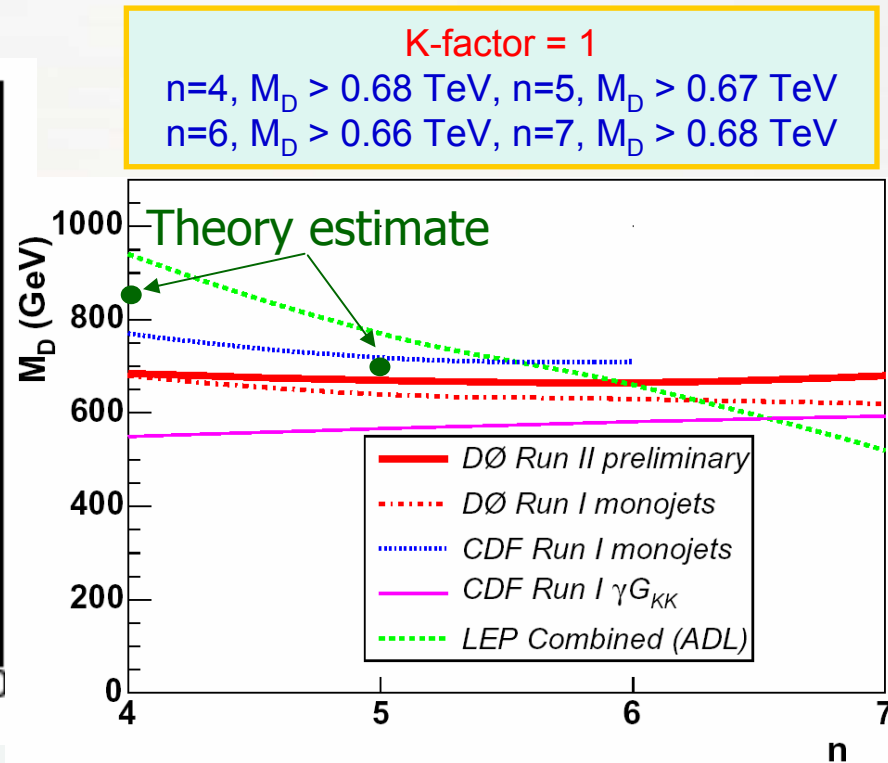
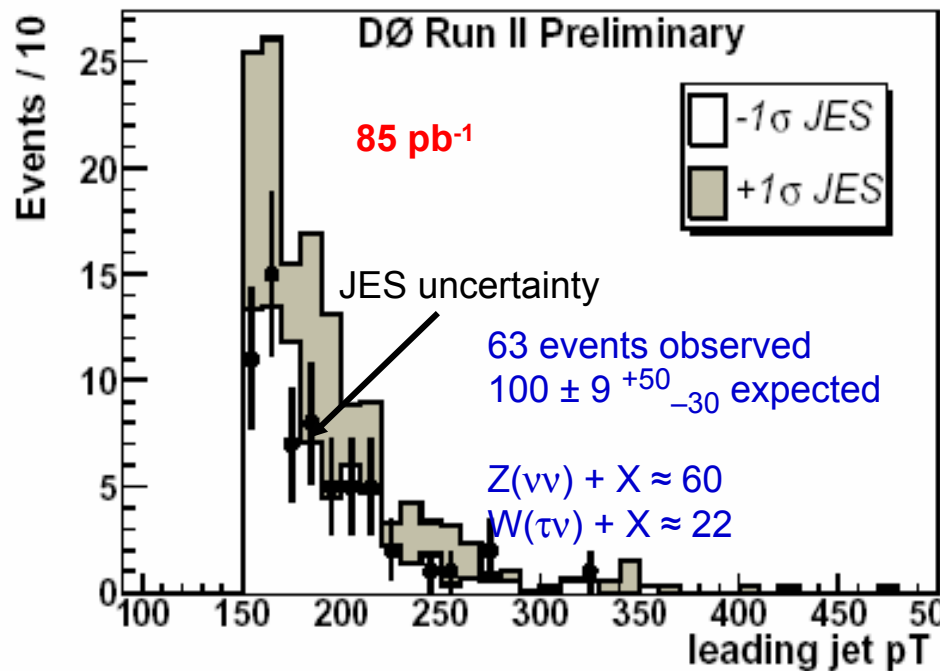


The [redacted] at the CERN  $p\bar{p}$  collider has caused ripples of excitement throughout the particle physics world, since they [redacted]

- The signature was deemed doomed and nearly forgotten
- It took many years for successful monojet analyses at a hadron collider to be completed (CDF/DØ)

# Monojets and Jet Energy Scale

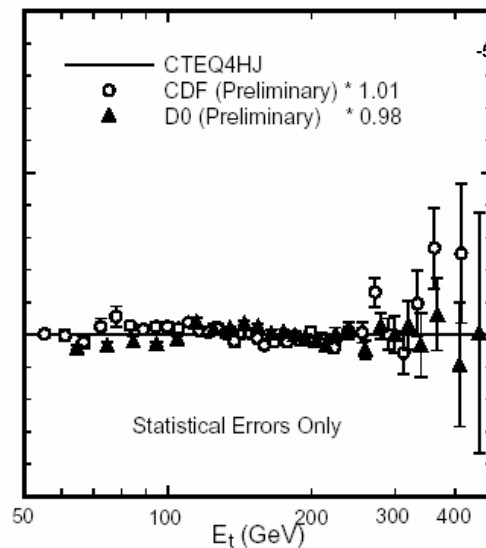
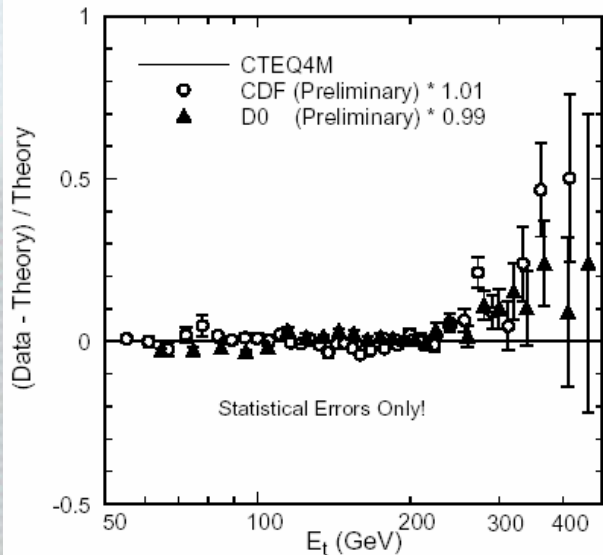
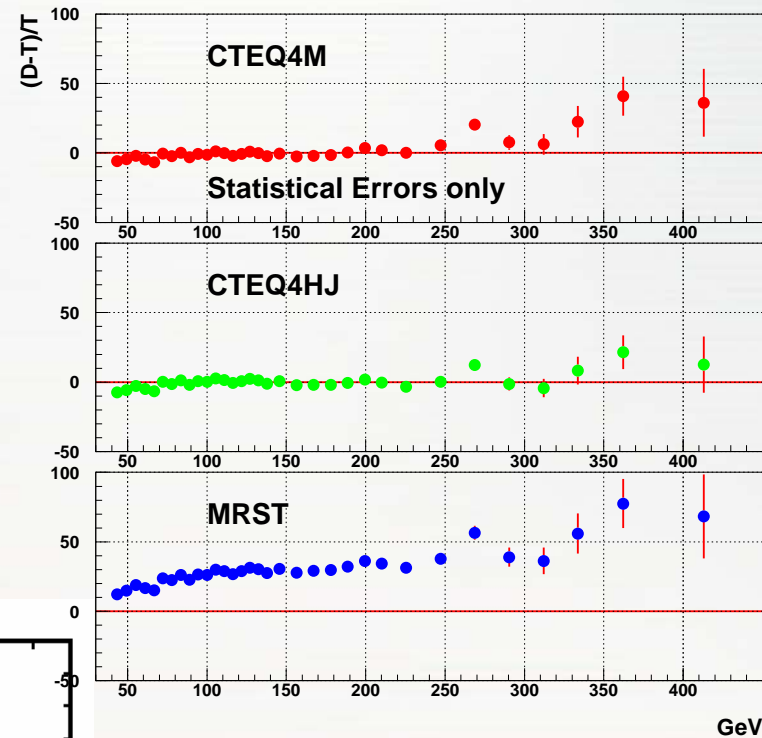
- Tevatron Run II: major systematics comes from the **jet energy scale**
- For the DØ Run II case it will be reduced soon by a factor of 2-3, but will still be important, if not dominant!
- **Expect this uncertainty to be a major source of systematics at the LHC**
- Getting **JES uncertainty down to a few per cent level at the LHC would take time**



# Lesson 2: Tricks of the PDF's

- PDF effects could play nasty tricks in searches for new physics
  - Cf. apparent excess of high- $p_T$  jets observed by CDF in Run I
  - Later explained by modification of gluon part of PDF (still poorly constrained!)
  - Will have impact on ED searches at the LHC

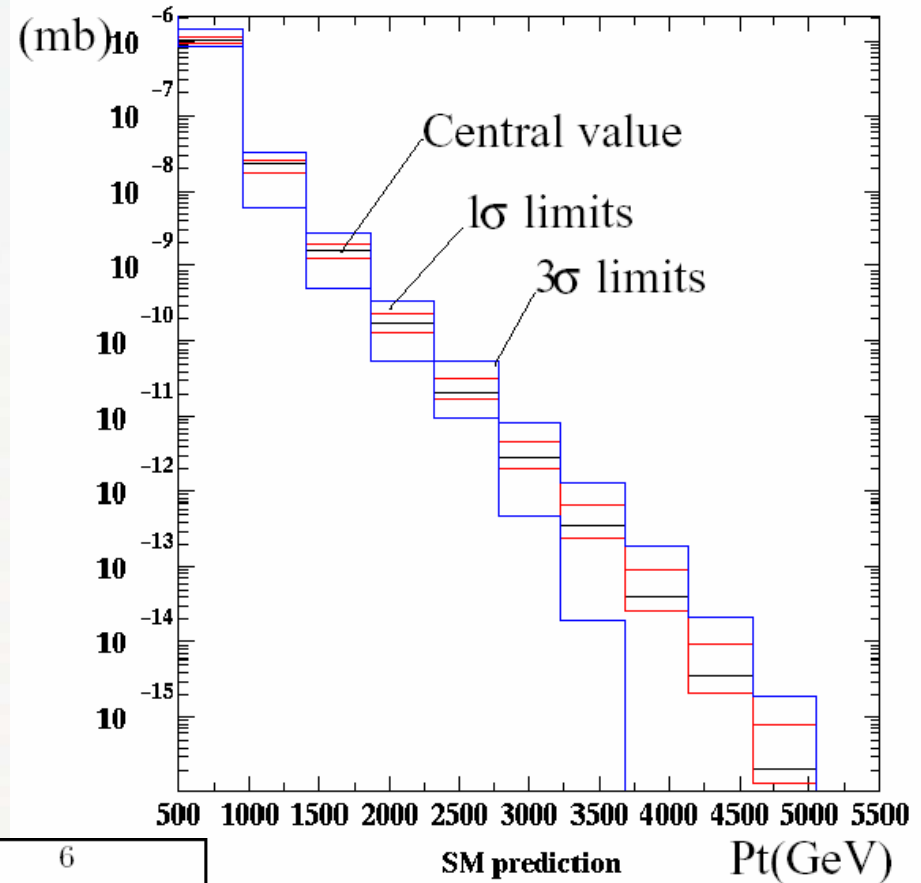
Inclusive Jet Cross Section (CDF Preliminary)



# PDF Effects at the LHC

- CTEQ6M error set provides a great tool for studies of the PDF uncertainty effects
- Example: **PDF effects in dijet production**
  - Offers direct sensitivity to ED models via virtual exchange
  - Accounting for the **PDF uncertainty results in significantly reduced sensitivity to ED effects**
  - Expect comparable effect for monojet production

[S. Ferrag, hep-ph/0407303]



	2 extra-dimensions	4 extra-dimensions	6 extra-dimensions
Theoretically	5 TeV	5 TeV	5 TeV
including PDF uncertainties	< 2 TeV	< 3 TeV	< 4 TeV

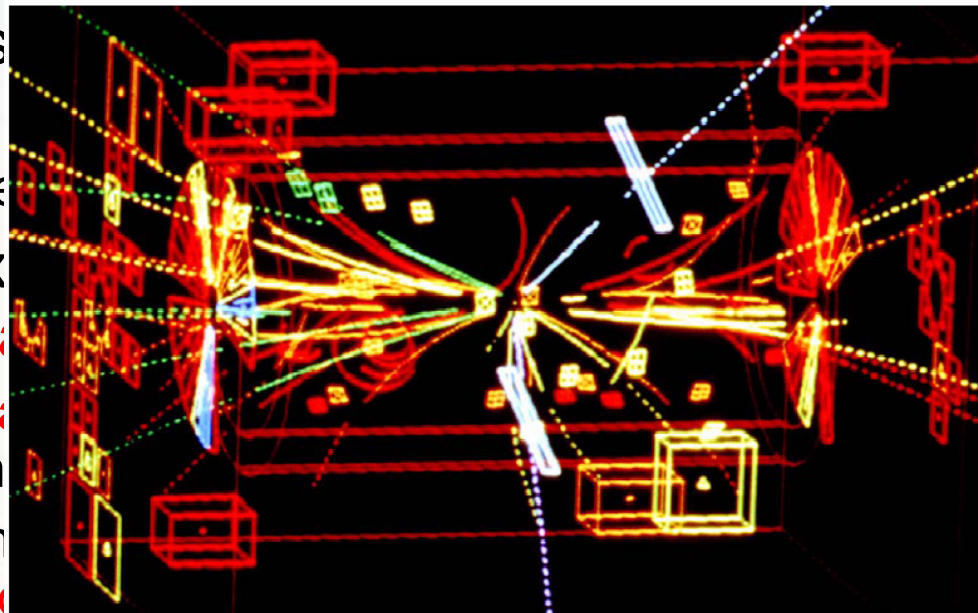


# Lesson 3: Drell-Yan and Diphotons

- Not all the Tevatron lessons are “negative”!
- **Drell-Yan and diphoton production** remain one of the best channels for “quick” discovery of ED effects at the LHC
  - Offer sensitivity to all three types of extra dimensions as well as other types of new physics (TC,  $Z'$ , ...)
- Drell-Yan **historically has been a fruitful channel for discoveries** ( $J/\Psi$ ,  $Y$ ,  $Z$ )
  - So maybe diphotons for the Higgs!

- Well-understood and NLO (diphoton)
  - Small dependence on  $\alpha_s$

- Tevatron experiments **performed for**
  - Have a “statistical” measurement
  - Have been **collaborative**



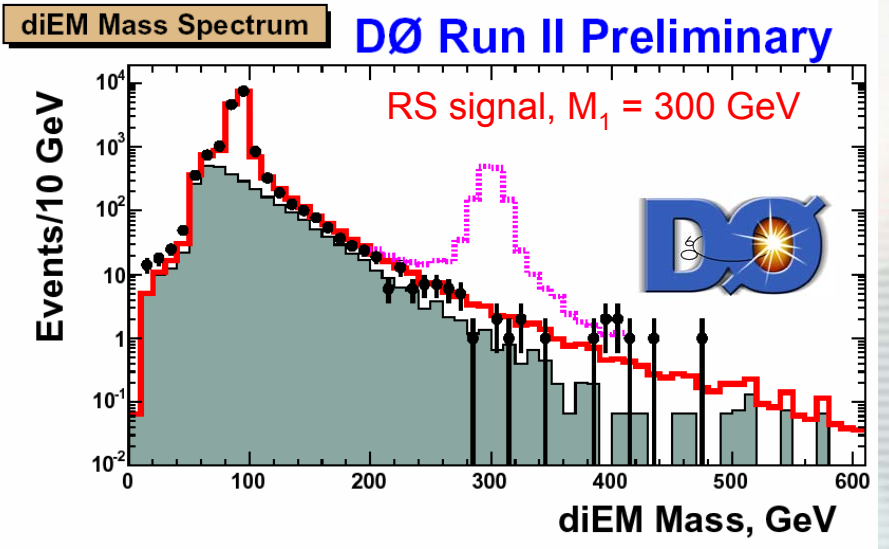
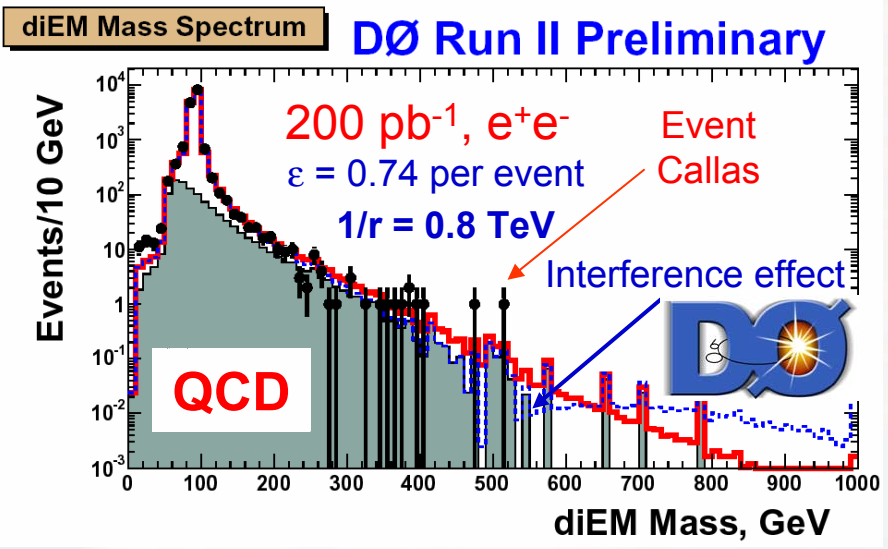
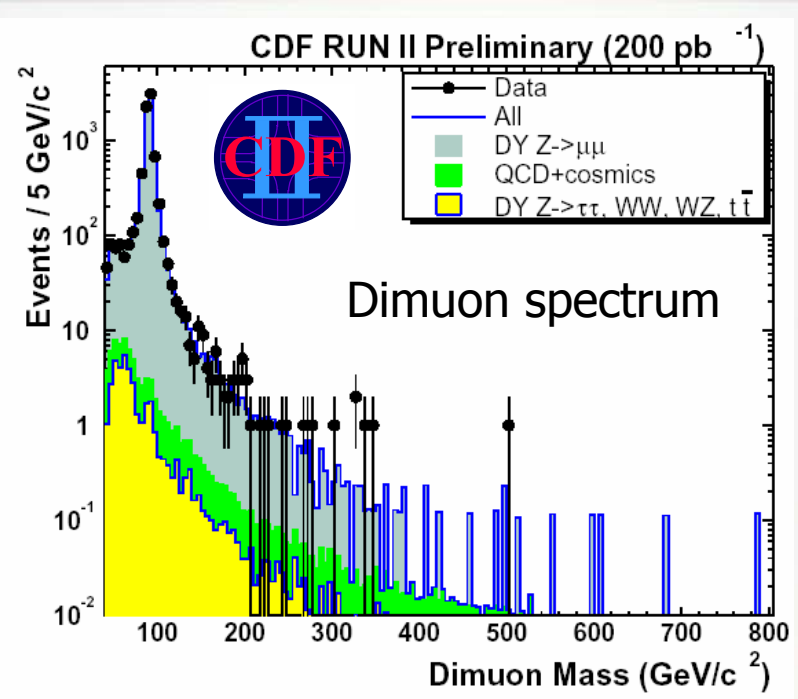
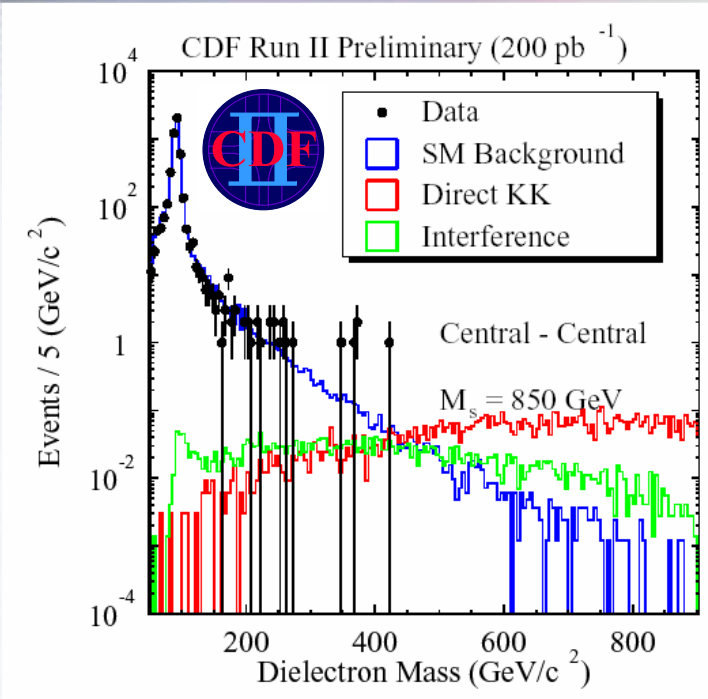
to NNLO (DY)

can be detector:

for in situ

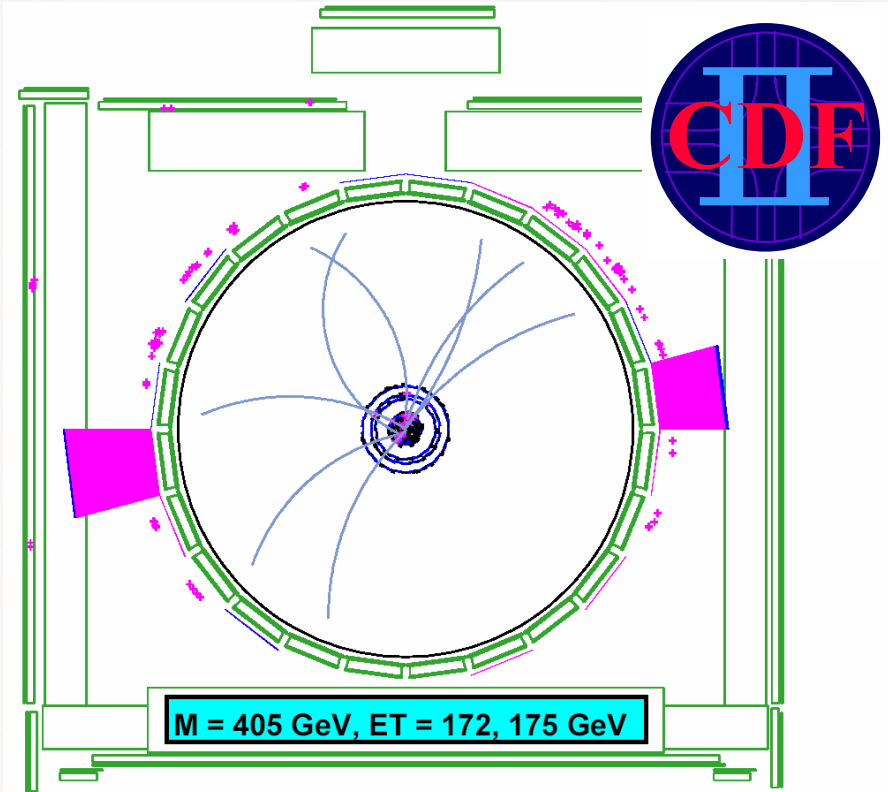
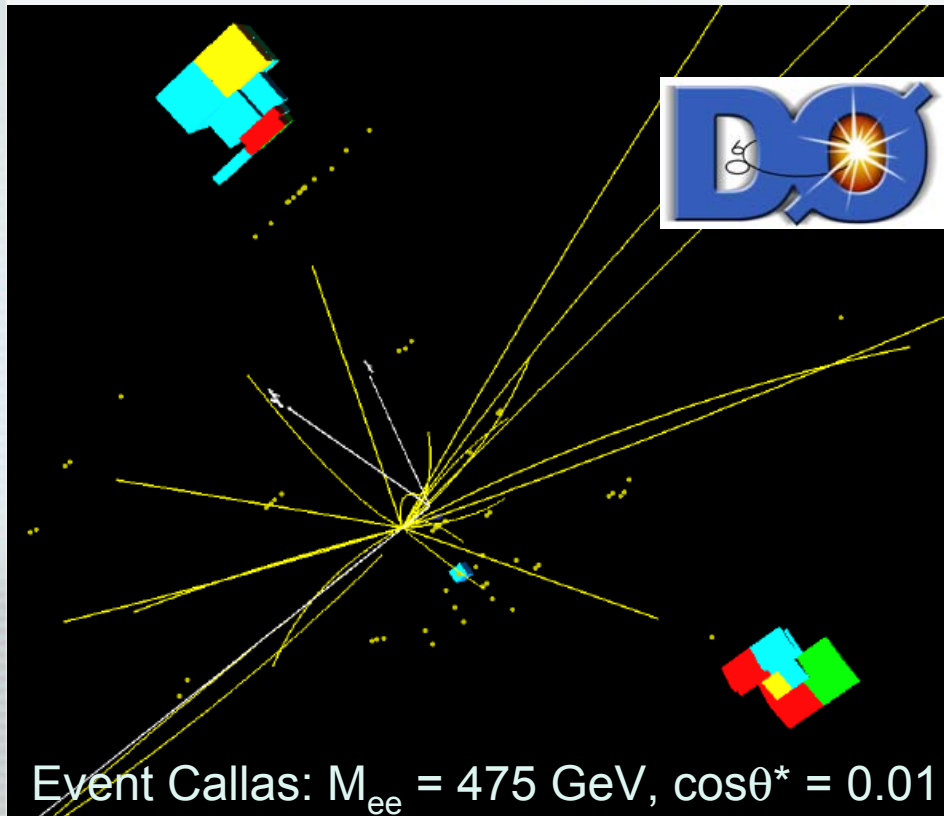
Tevatron  
sensitivity to date

# Invariant Mass Spectra



# Spectacular Candidate Events

- Two highest-mass candidate events seen by DØ ( $e^+e^-$ ) and CDF ( $\gamma\gamma$ )



# Beyond the Tevatron

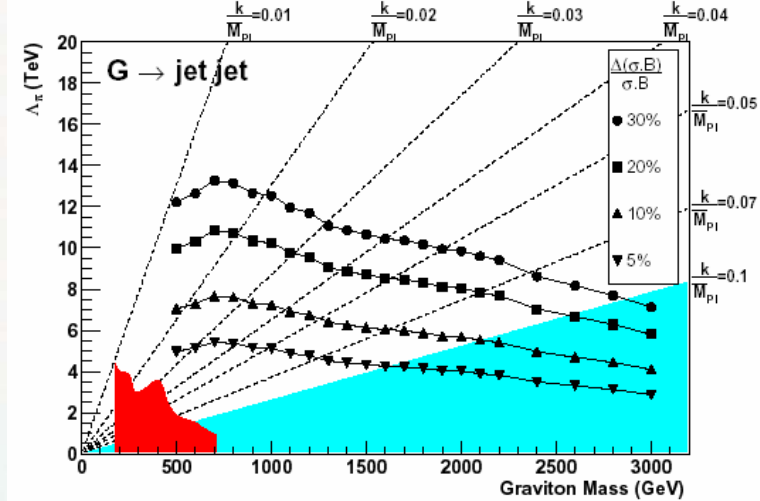
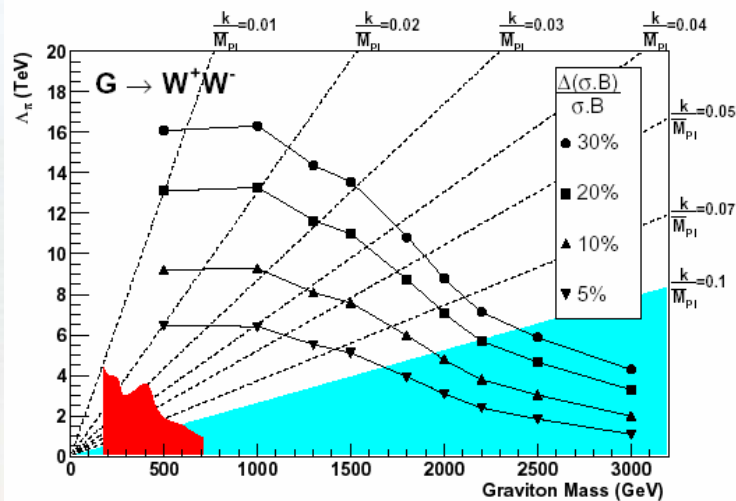
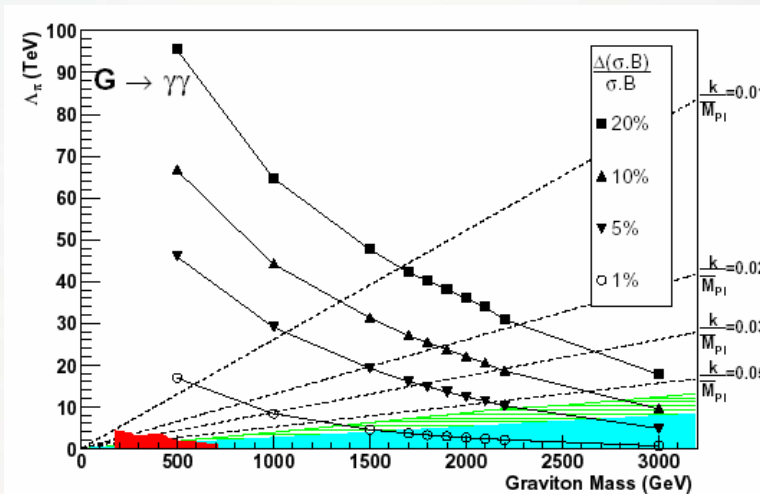
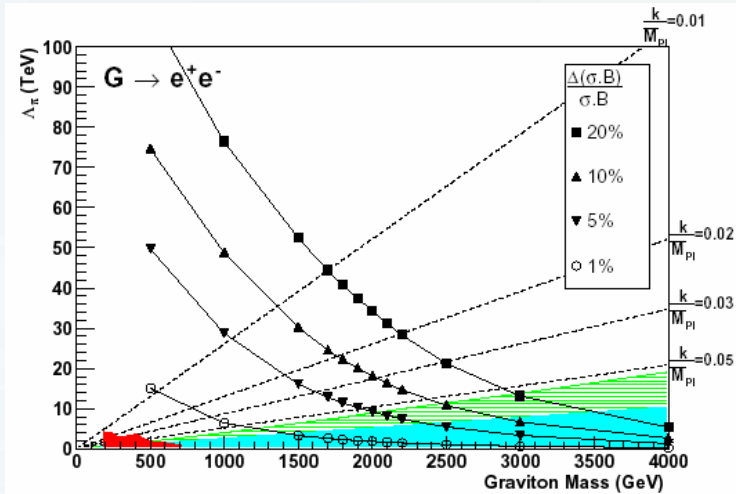
- Independent of whether the first hints of quantum gravity are seen at the Tevatron or not, LHC is the machine to study this physics
  - Sensitivity to  $M_{Pl}$  is about factor of five higher
  - Some measurements are completely unique to LHC
  - Qualitatively new physics may open up (e.g. black holes)
  - Signals for quantum gravity (if found) can be established firmly by measuring couplings, energy dependence, spin of the graviton, etc.
  - Most of the allowed model space in various scenarios can be ruled out if no indications for the signal is seen
  - Detailed studies of quantum properties of gravity may be possible



# Example: Graviton Couplings

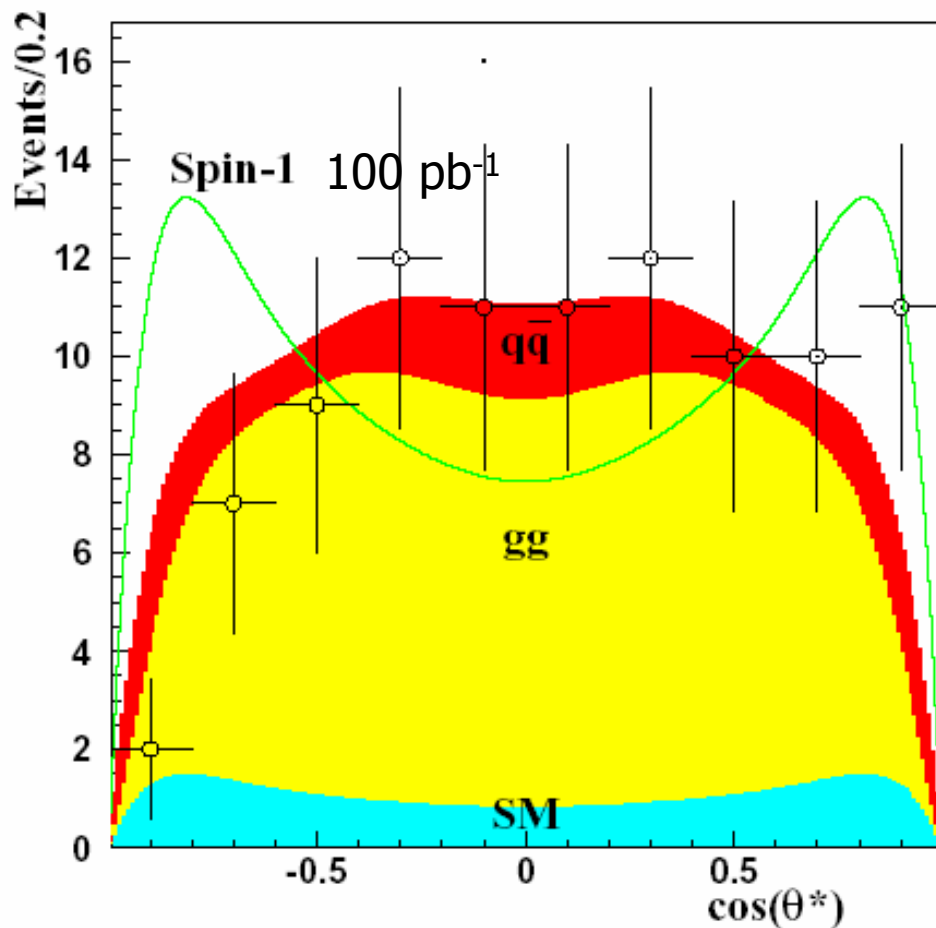
- **Universality of couplings** is one of the fundamental property of gravity
- Here is what LHC can do in the **RS scenario**:

[Allanach et al, JHEP 0212, 039 (2002)]



# Example: Graviton Spin

Shown: Randall-Sundrum scenario with  $M_G = 1.5$  TeV  
 Error bars correspond to  $100 \text{ pb}^{-1}$



[Allanach et al., JHEP **0009**, 019 (2000)]

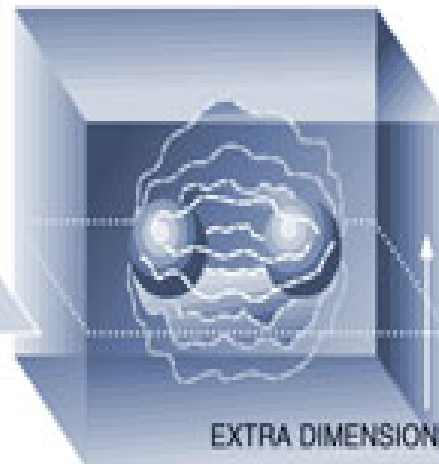
- Establishing the **spin-2 nature of the graviton** propagator is a unique measurement possible at the LHC
  - Applies both to the Randall-Sundrum and large ED effects

## Black Holes on Demand

The New York Times  
ON THE WEB

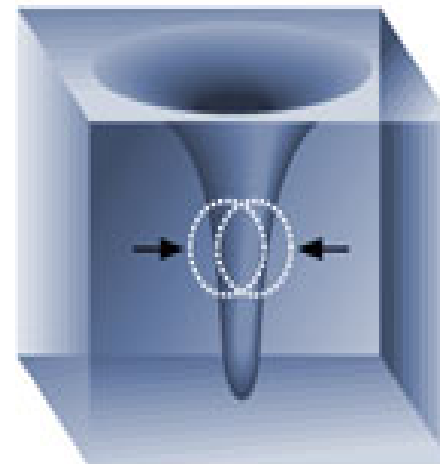
Scientists are exploring the possibility of producing miniature black holes on demand by smashing particles together. Their plans hinge on the theory that the universe contains more than the three dimensions of everyday life. Here's the idea:

*Particles collide in three dimensional space, shown below as a flat plane.*

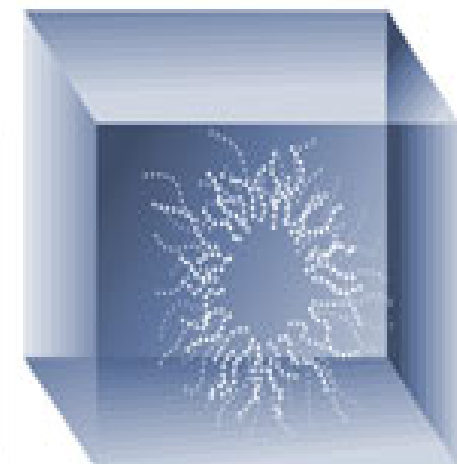


As the particles approach in a particle accelerator, their gravitational attraction increases steadily.

When the particles are extremely close, they may enter space with more dimensions, shown above as a cube.



The extra dimensions would allow gravity to increase more rapidly so a black hole can form.



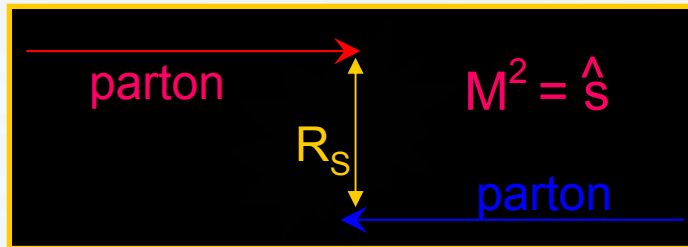
Such a black hole would immediately evaporate, sending out a unique pattern of radiation.

NYT, 9/11/01

# Black Hole Production...

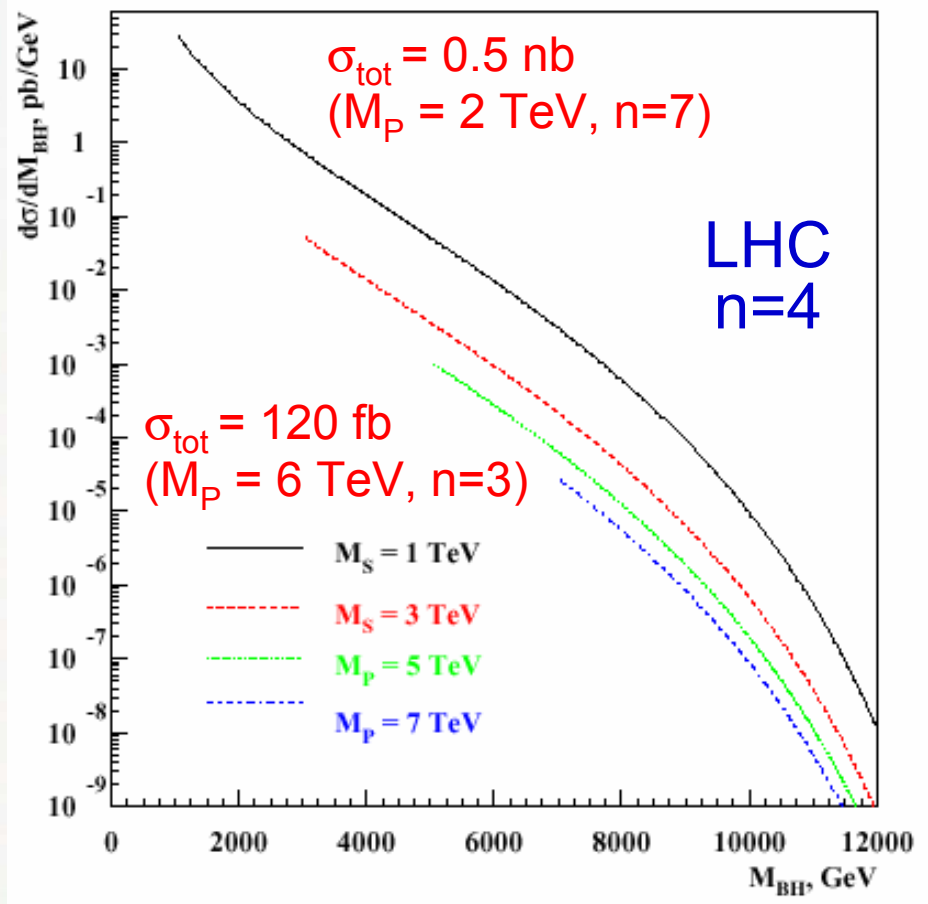
- Main idea: when the c.o.m. energy reaches the fundamental Planck scale, a BH is formed; cross section is given by the black disk approximation:

$$\sigma \sim \pi R_S^2 \sim 1 \text{ TeV}^{-2} \sim 10^{-38} \text{ m}^2 \sim 100 \text{ pb}$$



- While we do not have quantum picture of gravity for  $M_{\text{BH}} \sim M_{\text{Pl}}$ , certain expectations can be used to predict behavior of near-critical BH (e.g., “democratic” couplings)
- Quantum corrections can be minimized by looking at the proper regime (cf. Planck vs. Rayleigh-Jeans behavior in the infrared)

[Dimopoulos, GL, PRL 87, 161602 (2001)]





# ... and Decay

- Once produced, BH evaporates via black-body Hawking radiation
- Hawking temperature:

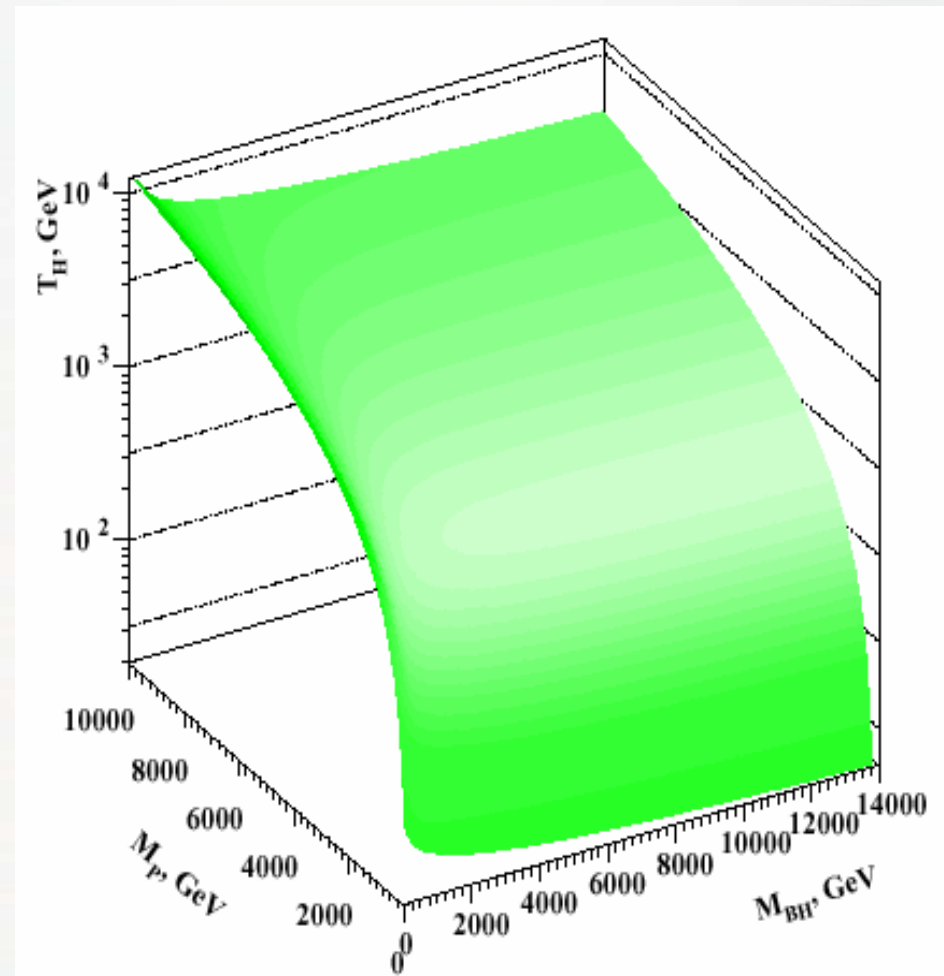
$$R_S T_H = \frac{n+1}{4\pi}$$

(in natural units  $\hbar = c = k = 1$ )

- BH radiates mainly on the brane [Emparan/Horowitz/Myers, hep-th/0003118]
- Democratic couplings to  $\sim 120$  SM d.o.f. yield probability of Hawking evaporation into jets,  $\gamma$ ,  $l^\pm$ , and  $\nu$   $\sim 75\%$ ,  $2\%$ ,  $10\%$ , and  $5\%$  respectively
- Averaging over the BB spectrum gives average multiplicity of decay products:

$$\langle N \rangle \approx \frac{M_{BH}}{2T_H}$$

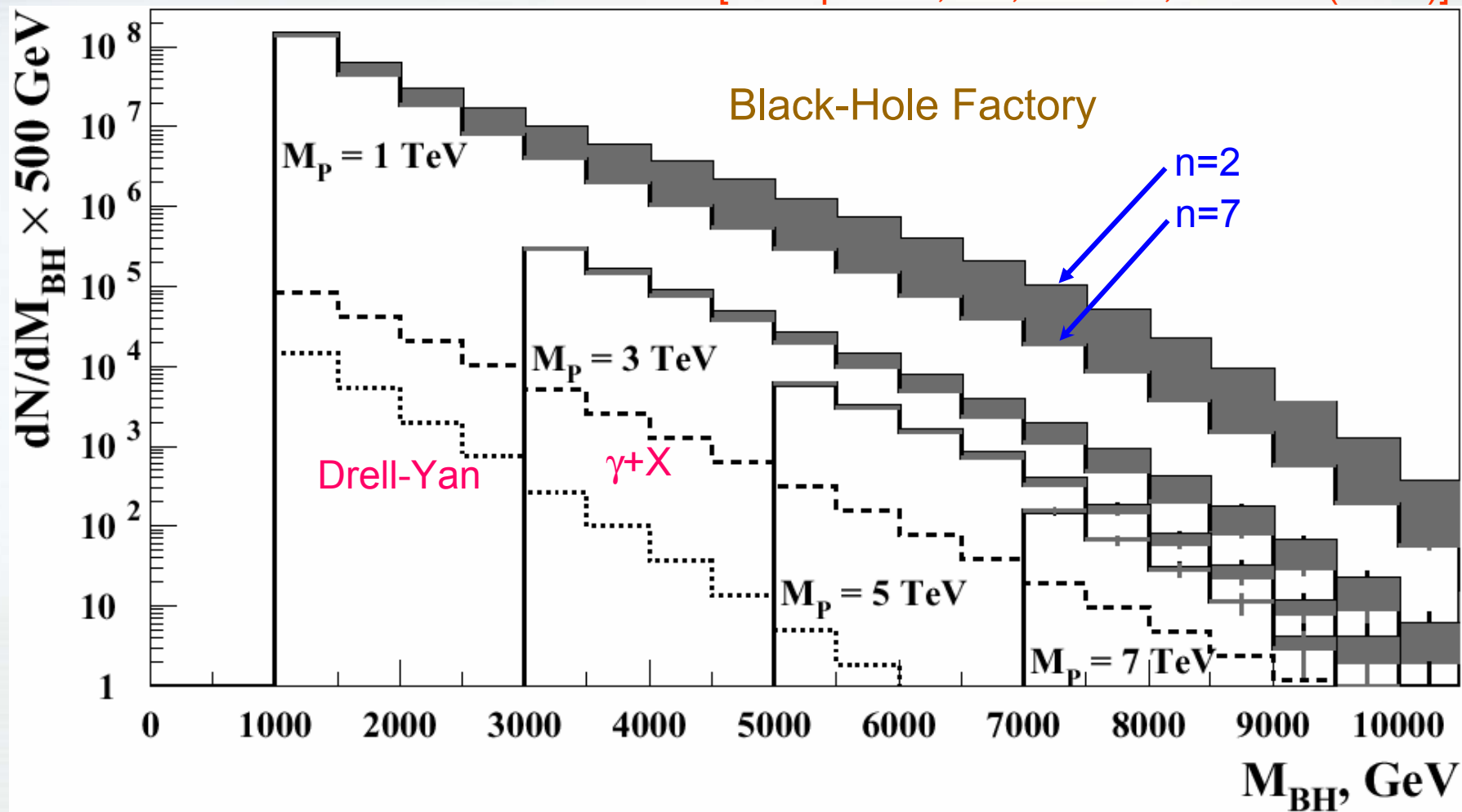
[Dimopoulos, GL, PRL 87, 161602 (2001)]



• Stefan's law: lifetime  $\tau \sim 10^{-26}$  s

# Black Hole Factory

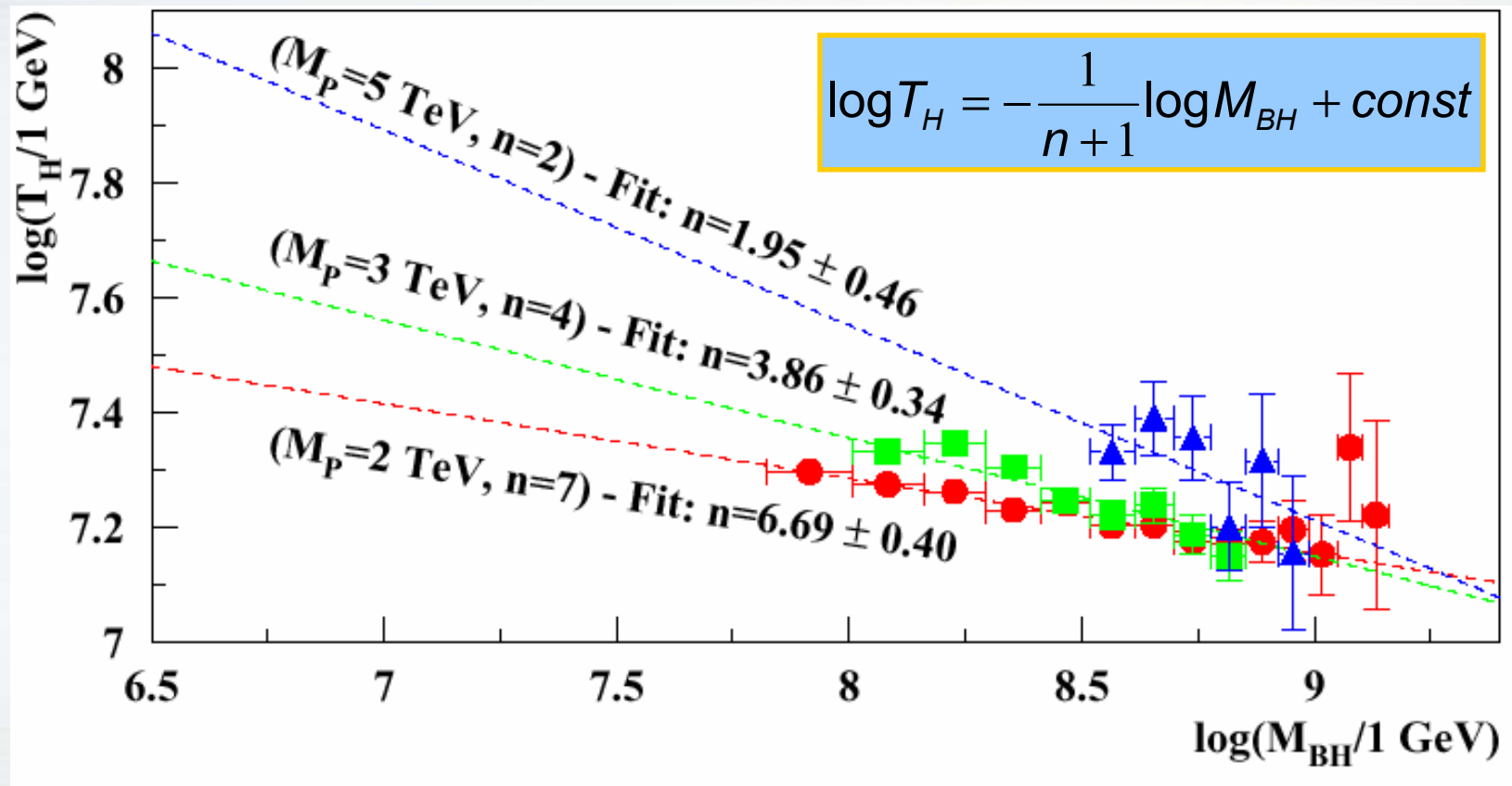
[Dimopoulos, GL, PRL 87, 161602 (2001)]



Spectrum of BH produced at the LHC with subsequent decay into final states tagged with an electron or a photon

# Space Texture at the LHC

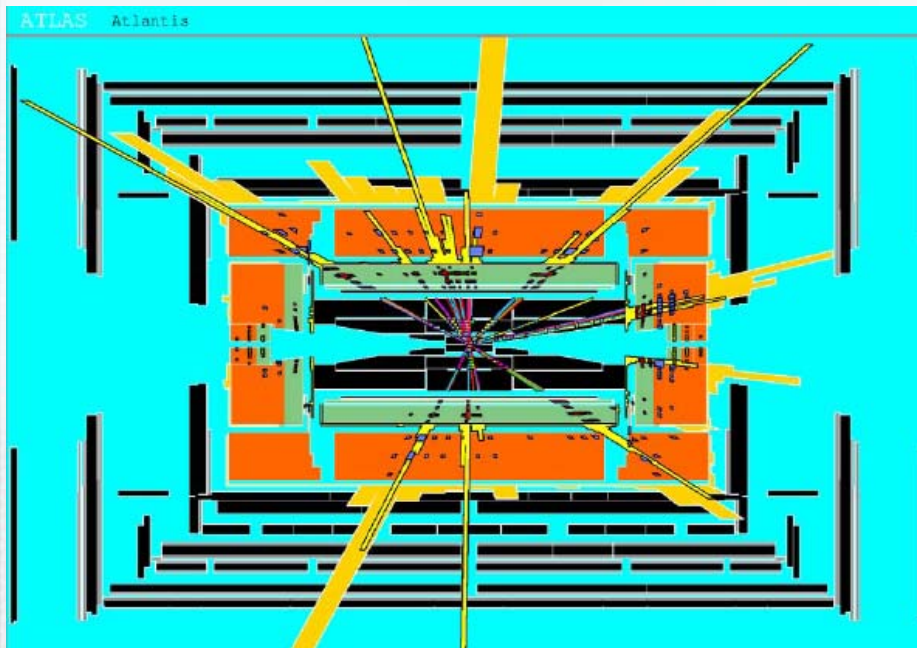
[Dimopoulos, GL, PRL 87, 161602 (2001)]



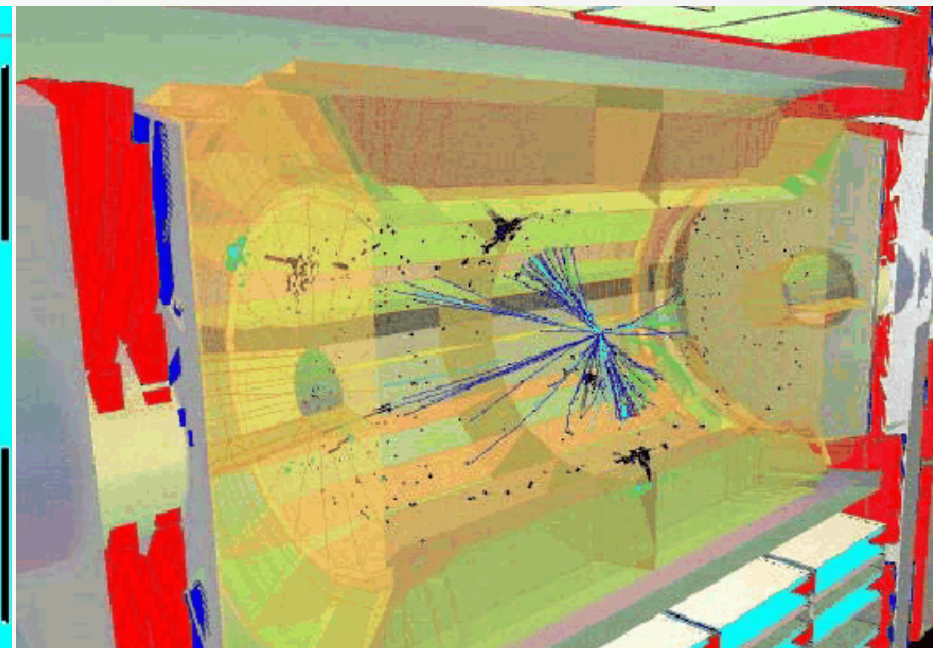
- Relationship between  $\log T_H$  and  $\log M_{BH}$  allows to find the number of ED
- This result is independent of their sizes and shape!
- This approach differs drastically from other collider signatures and would constitute a “smoking cannon” signature for low-scale (quantum) gravity

# Black Hole Events

- **First studies** have been already initiated by ATLAS and CMS
  - ATLAS – **CHARYBDIS** HERWIG-based generator with more elaborated decay model [Harris/Richardson/Webber, JHEP **08**, 033 (2003)]
  - CMS – **TRUENOIR** [Dimopoulos/GL, Snowmass-2001-P321]
- While **realistic simulations** will be needed, the BH's are really hard to miss, so I'd **wait until anything like that is seen** first!



Simulated black hole event in the ATLAS detector [from ATLAS-Japan Group]

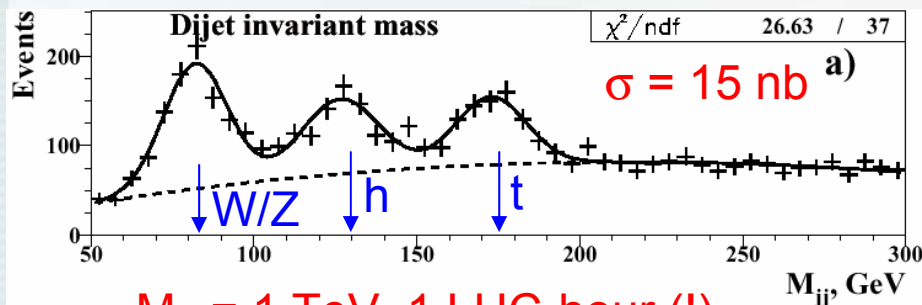


Simulated black hole event in the CMS detector [A. de Roeck & S. Wynnhoff]

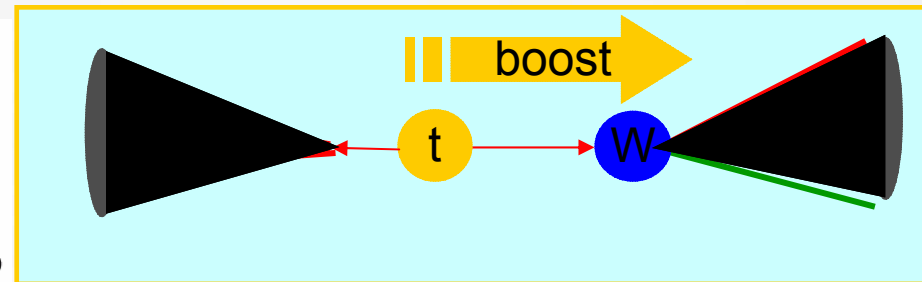


# More Fun with Black Holes

- While **we might not know about existence of new particles** with masses  $\sim 100$  GeV, **gravity certainly does**
- Opens up **an alternative way of searching for new particles**: in BH decay rather than in prompt production
- Example: **Higgs with the mass of 130 GeV** decays predominantly into a  $b\bar{b}$ -pair
- Approach: tag BH events with leptons or photons, and look at the dijet invariant mass; does not even require b-tagging!
- Use a typical LHC detector response to obtain realistic results



$M_p = 1$  TeV, 1 LHC-hour (!)



[GL, PRL **88**, 181801 (2002)]

- **Higgs observation in the black hole decays** is possible at the LHC as early as in the first day of running even with the incomplete and poorly calibrated detectors!
- For  $M_p = 1, 2, 3,$  and  $4$  TeV one needs 1 day, 1 week, 1 month, or 1 year of running to find a  $5\sigma$  signal
- Main conclusion is **confirmed by the CMS-specific analysis [Akchurin et al. Fermilab-FN-0752 (2004)]**
- Same applies to most of new particles with the mass  $\sim 100$  GeV

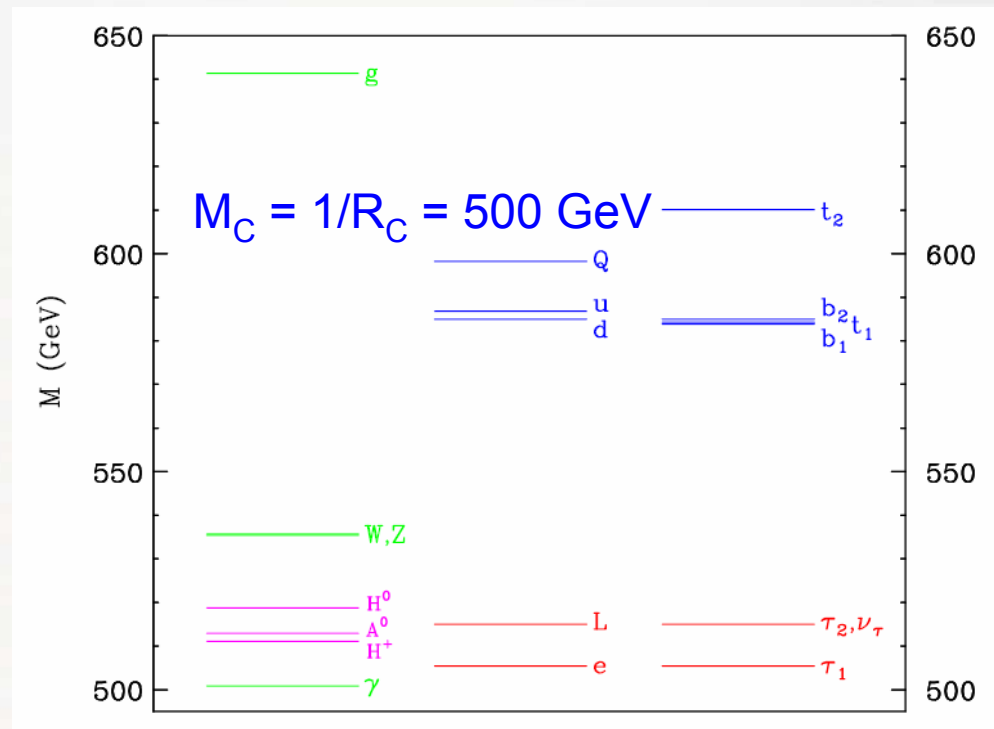
# Universal Extra Dimensions

- The most “democratic” ED model: *all* the SM fields are free to propagate in extra dimension(s) with the size  $R_c = 1/M_c \sim 1 \text{ TeV}^{-1}$  [Appelquist, Cheng, Dobrescu, PRD **64**, 035002 (2001)]
  - Instead of chiral doublets and singlets, model contains vector-like quarks and leptons
  - Gravity is not included in this model
- The number of universal extra dimensions is not fixed:
  - it’s feasible that there is just one (MUED)
  - the case of two extra dimensions is theoretically attractive, as it breaks down to the chiral Standard Model and has nice features, such as guaranteed proton stability, etc.
- Every particle acquires KK modes with the masses  $M_n^2 = M_0^2 + M_c^2$ ,  $n = 0, 1, 2, \dots$
- Kaluza-Klein number ( $n$ ) is conserved at the tree level, i.e.  $n_1 \pm n_2 \pm n_3 \pm \dots = 0$ ; consequently, the lightest KK mode is stable (and serves an excellent dark matter candidate [Cheng, Feng, Matchev, PRL **89**, 211301 (2002)])
- Hence, KK-excitations are produced in pairs, similar to SUSY particles
- Consequently, current limits (dominated by precision electroweak measurements, particularly T-parameter) are sufficiently low ( $M_c \sim 300 \text{ GeV}$  for one ED and of the same order, albeit more model-dependent for  $>1$  ED)

# UED Phenomenology

- Naively, one would expect large clusters of **nearly degenerate states** with the mass around  $1/R_C$ ,  $2/R_C$ , ...
- Cheng, Matchev, Schmaltz: **not true, as radiative corrections tend to be large** (up to 30%); thus the KK excitation mass spectrum resembles that of SUSY!
- Minimal UED model with a single **extra dimension**, compactified on an  $S_1/Z_2$  orbifold
  - **Odd fields** do not have 0 modes, so we identify them w/ “**wrong**” **chiralities** that **vanish in the SM**

- $Q, L (q, l)$  are  $SU(2)$  doublets (**singlets**) and contain both chiralities

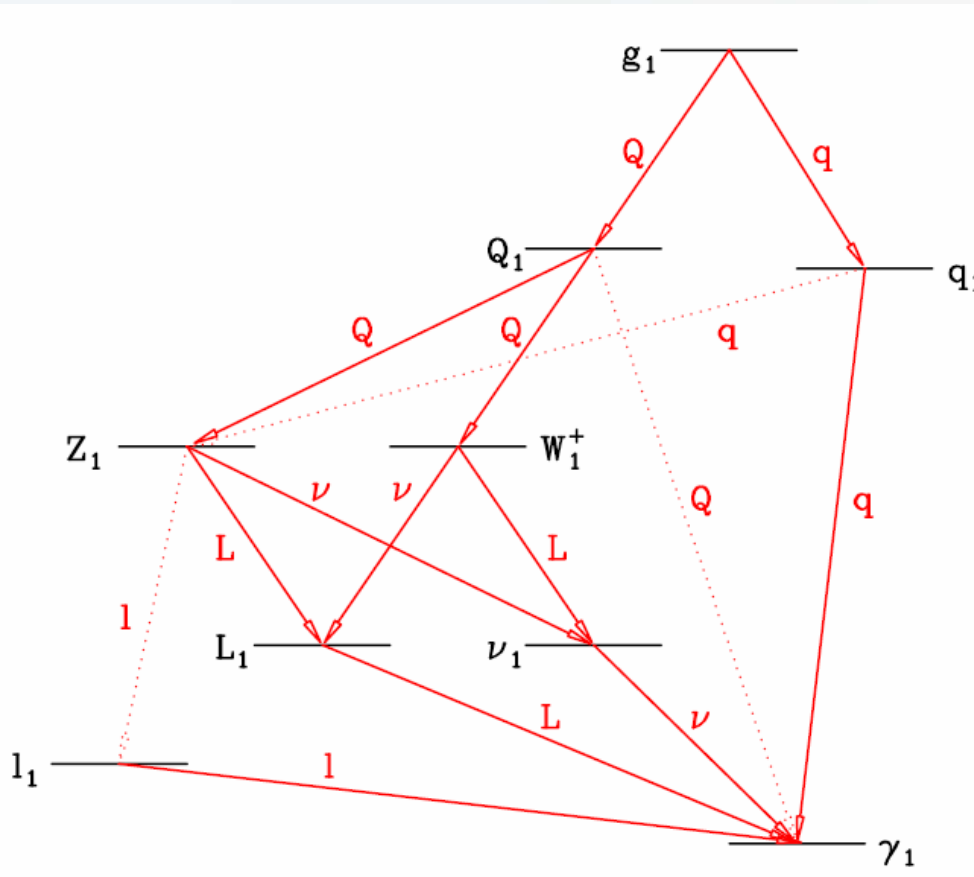


[Cheng, Matchev, Schmaltz, PRD **66**, 056006 (2002)]

# SUSY Without SUSY!

- First level KK-states spectroscopy

[Cheng, Matchev, Schmaltz, PRD **66**, 056006 (2002)]



Decay:

$$B(g_1 \rightarrow Q_1 Q) \sim 50\%$$

$$B(g_1 \rightarrow q_1 q) \sim 50\%$$

$$B(q_1 \rightarrow q \gamma_1) \sim 100\%$$

$$B(t_1 \rightarrow W_1 b, H_1^+ b) \sim$$

$$B(Q_1 \rightarrow Q Z_1 : W_1 : \gamma_1) \sim 33\% : 65\% : 2\%$$

$$B(W_1 \rightarrow \nu L_1 : \nu_1 L) = 1/6 : 1/6 \text{ (per flavor)}$$

$$B(Z_1 \rightarrow \nu \nu_1 : L L_1) \sim 1/6 : 1/6 \text{ (per flavor)}$$

$$B(L_1 \rightarrow \gamma_1 L) \sim 100\%$$

$$B(\nu_1 \rightarrow \gamma_1 \nu) \sim 100\%$$

$$B(H_1^\pm \rightarrow \gamma \gamma_1, H^{\pm*} \gamma_1) \sim 100\%$$

Production:

$$q_1 q_1 + X \rightarrow ME_T + \text{jets} (\sim \sigma_{\text{had}}/4); \text{ but: } \\ \text{low } ME_T$$

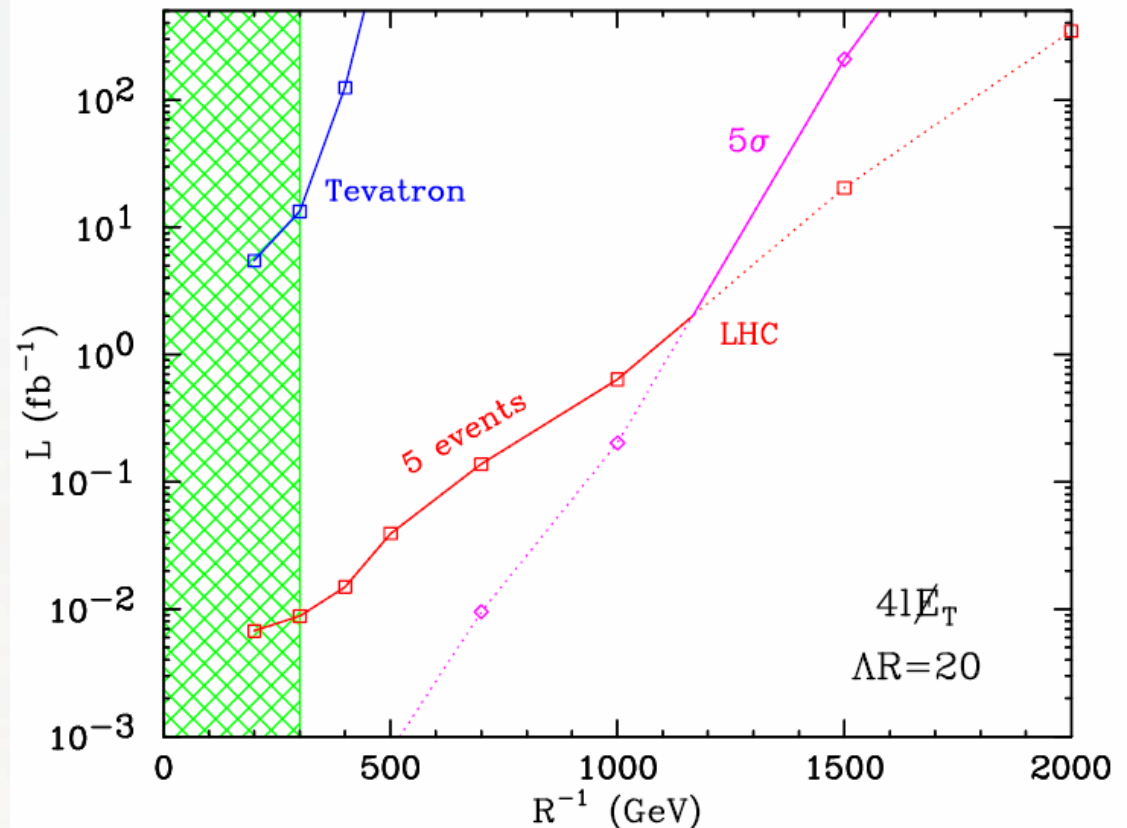
$$Q_1 Q_1 + X \rightarrow V_1 V'_1 + \text{jets} \rightarrow 2-4 \ell + ME_T \\ (\sim \sigma_{\text{had}}/4)$$



# Sensitivity in the Four-Lepton Mode

- Only the gold-plated 4-leptons +  $ME_T$  mode has been considered in the original paper
- Much more promising channels:
  - dileptons + jets +  $ME_T$  + X (x9 cross section)
  - trileptons + jets +  $ME_T$  + X (x5 cross section)
- Detailed simulations is required: would love to see this in a MC
- Has not been studied in details neither at the Tevatron, nor at the LHC!

[Cheng, Matchev, Schmaltz, PRD **66**, 056006 (2002)]



L is per experiment (single experiment)

# Conclusions

- The **Tevatron** provides an important input and insight for future searches for quantum gravity effects at **the LHC**
  - Early understanding of instrumental backgrounds
  - Jet energy scale determination
  - Optimization of  $ME_T$  resolution
- However, **some inputs could only come from the LHC itself**:
  - Real detector performance
  - Accelerator background conditions
  - Constraining gluon PDF's
- **Discoveries are unlikely to happen “on day one”** – be prepared for a long, possibly bumpy ride, but a thrilling one!
- **Nevertheless**, an order of magnitude **higher energy** at the LHC **may open doors for qualitatively new physics** (e.g. BH!) and lead to **surprises very early on**
- So, **let's build the damn thing and have fun!**