

On study of angle correlations in $\Phi \rightarrow ZZ \rightarrow 2e2\mu$ at the CMS

Michał Bluj

Soltan's Institute for Nuclear Studies, Warsaw

Outline

- Model
- Methodology
- MC Samples
- Selection
- Results

The results presented here are preliminary and can't be distributed before they get approved by the CMS collaboration.

Model

- An effective model of ΦZZ coupling with scalar ($g^{\mu\nu}$) and pseudoscalar ($\epsilon^{\mu\nu\rho\sigma} k_{1\rho} k_{2\sigma}$) terms (A.Skjold, P.Osland Phys. Lett. B329, 305 (1994), implemented in Pythia):

$$C_{\Phi ZZ} \sim m_Z^2 g^{\mu\nu} + \tan \xi \cdot \epsilon^{\mu\nu\rho\sigma} k_{1\rho} k_{2\sigma}$$

where $k_1 = (q_1 + q_2)$, $k_2 = (q_3 + q_4)$, $q_{i=1\dots 4}$ momenta of Z^0 s and leptons;

$\tan \xi$ describes deviation from SM (scalar $\xi = 0$, pseudoscalar $\xi = \pm\pi/2$, CP -violation $\xi \neq 0, \pm\pi/2$).

- Differential cross-section:

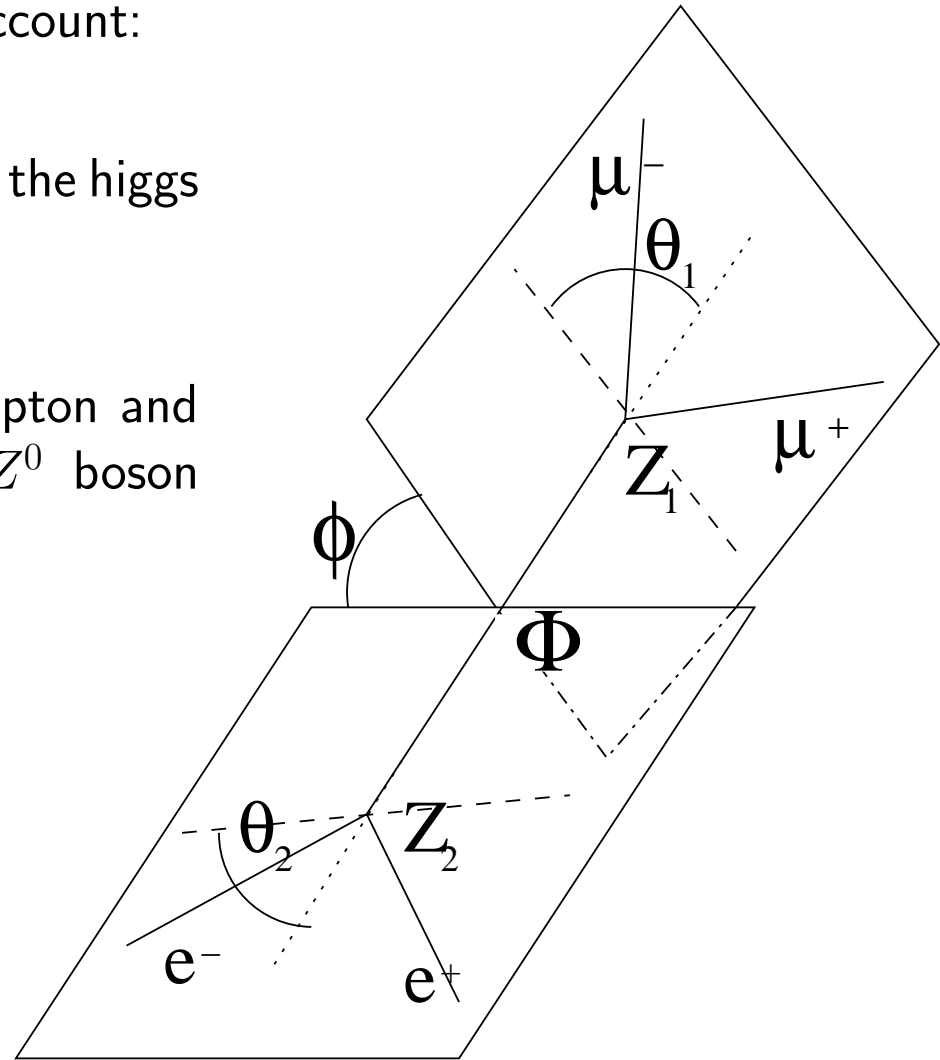
$$d\sigma(\xi) \sim \mathcal{H} + \tan \xi \cdot \mathcal{V} + \tan^2 \xi \cdot \mathcal{A}$$

where:

- scalar: \mathcal{H}
- mixing term: \mathcal{V}
- pseudoscalar: \mathcal{A}

Methodology: definition of observables

- Angular distributions distinguish between states with different ξ 's
- Two distributions taken in to account:
 - Plane angle φ
between decay planes of Z^0 s in the higgs rest frame
 - Polar angle $\theta_{1,2}$
between negatively charged lepton and the direction of Z^0 in the Z^0 boson rest frame



Methodology: maximization of likelihood function

- Definition of likelihood function

$$\mathcal{L}(\xi, R) \equiv 2 \sum_{x_i \in \text{data}} \log \mathcal{Q}(\xi, R; x_i)$$

where

$$\mathcal{Q}(\xi, R; x_i) \equiv R \cdot \mathcal{P}_S(\xi; x_i) + (1 - R) \cdot \mathcal{P}_B(x_i)$$

$\{x_i\}$ – data event,

R – fraction of signal in data sample (1st parameter of fit),

\mathcal{P}_B and $\mathcal{P}_S(\xi)$ – probability density functions for background and signal:

$$\mathcal{P}_B \equiv \mathcal{P}_B^M \cdot \mathcal{P}_B^\varphi \cdot \mathcal{P}_B^{\cos \theta_1} \cdot \mathcal{P}_B^{\cos \theta_2}$$

$$\mathcal{P}_S(\xi) \equiv \mathcal{P}_S^M \cdot (\mathcal{P}_S^\varphi \cdot \mathcal{P}_S^{\cos \theta_1} \cdot \mathcal{P}_S^{\cos \theta_2})(\xi)$$

where

\mathcal{P}^M , \mathcal{P}^φ , $\mathcal{P}^{\cos \theta_{1,2}}$ – probability density functions for $m_{4\ell}$, φ and $\cos \theta_{1,2}$ respectively obtained by MC simulation.

Methodology: determination of ξ

- Definition of signal part of Q -function according to expression for $d\sigma(\xi)$:

$$\mathcal{P}_S^\varphi \cdot \mathcal{P}_S^{\cos\theta_{1,2}}(\xi) \equiv (\mathcal{H} + \tan\xi \cdot \mathcal{V} + \tan^2\xi \cdot \mathcal{A}) / (1 + \tan\xi + \tan^2\xi)$$

– $\mathcal{H} \equiv \mathcal{P}_H^\varphi \cdot \mathcal{P}_H^{\cos\theta_{1,2}}$ and $\mathcal{A} \equiv \mathcal{P}_A^\varphi \cdot \mathcal{P}_A^{\cos\theta_{1,2}}$ – probabilities for scalar (H) and pseudoscalar (A)

– \mathcal{V} is normalized angle distribution for mixing term (V), **not probability!** (isn't positive) thus it can't be simulated separately

- Determination of \mathcal{V} from probability for $\xi = \pi/4$

$$\mathcal{J} \equiv \mathcal{P}_S^\varphi \cdot \mathcal{P}_S^{\cos\theta_{1,2}}(\xi = \pi/4) = (\mathcal{H} + \mathcal{V} + \mathcal{A})/3$$

$$\implies \mathcal{V} = 3\mathcal{J} - \mathcal{H} - \mathcal{A}$$

and finally:

$$\mathcal{P}_S^\varphi \cdot \mathcal{P}_S^{\cos\theta_{1,2}}(\xi) \equiv [(1 - \tan\xi) \cdot \mathcal{H} + \tan\xi \cdot 3\mathcal{J} + (\tan^2\xi - \tan\xi) \cdot \mathcal{A}] / (1 + \tan\xi + \tan^2\xi)$$

MC samples for “golden” channel $\Phi \rightarrow ZZ \rightarrow 2e2\mu$

- Samples generated with detector acceptance preselection:
 - $2e$ with $p_t > 5$ GeV & $\eta < 2.7$
 - 2μ with $p_t > 3$ GeV & $\eta < 2.5$
- Signal $\Phi \rightarrow ZZ \rightarrow 2e2\mu$:
 - Samples generated for three masses above ZZ threshold $m_\Phi = 200, 300, 400$ GeV
 - Samples used for construction of \mathcal{P} (scalar, pseudoscalar and $\xi = \pi/4$) – **10k. evts.**
 - For CP -violating case samples
 - $\tan \xi = 0.1, 0.4, 4$ – **5k. evts.**
 - $\tan \xi = -0.1, -0.4, -4$ – **1k. evts.** (limited statistic)
- Background:
 - $ZZ \rightarrow 2e2\mu$ – **20k. evts.**
 - $t\bar{t} \rightarrow 2e2\mu$ – **48k. evts.**
 - $Zb\bar{b} \rightarrow 2e2\mu$ – **4k. evts.**
- All samples generated with low-lumi pile-up (~ 3.5 “soft” evts./”hard” evt.).
- Full simulation & reconstruction of the CMS detector were used.

Selection

- CMS selection for Higgs boson in $\Phi \rightarrow ZZ \rightarrow 2e2\mu$ channel (by D.Futyan, D.Giordano)
 - Di-electron or di-muon trigger (“ $Z \rightarrow 2\ell$ -trigger”)
 - Reconstructed two lepton pairs e^+e^- and $\mu^+\mu^-$
 - All 4 leptons originate at one vertex
 - Isolation at tracker
 - Kinematic cuts
 - a) Cuts on p_t s of leptons ($p_{t1}, p_{t2}, p_{t3}, p_{t4}$)
 - b) Cuts on mass of Z candidates:
 - * Symmetric window around better Z candidate (Δm_{Z1})
 - * Asymmetric mass window in mass of worse Z candidate ($m_{Z2}^{min} \div m_{Z2}^{max}$)
 - c) Mass window in four lepton mass (mass of higgs candidate) ($m_H^{min} \div m_H^{max}$)
- Cuts optimization (for each higgs mass) performed automatically to maximize significance:

$$S = \sqrt{2\ln Q}, \text{ where } Q = \left(1 + \frac{N_S}{N_B}\right)^{N_S+N_B} e^{-N_S}$$

Selected cross-sections

	$m_\Phi = 200 \text{ GeV}$				$m_\Phi = 300 \text{ GeV}$				$m_\Phi = 400 \text{ GeV}$			
	sig.	ZZ	$t\bar{t}$	$Zb\bar{b}$	sig.	ZZ	$t\bar{t}$	$Zb\bar{b}$	sig.	ZZ	$t\bar{t}$	$Zb\bar{b}$
σ_{tot} (pb)	17.86	21.2	886	525	9.41	21.2	886	525	8.71	21.2	886	525
$\sigma_{tot} \cdot \epsilon \cdot BR$	7.65	11.81	817.52	116.38	5.08	11.81	817.52	116.38	4.45	11.81	817.52	116.38
rec. $2\mu 2e$	5.46	6.71	173.02	32.77	3.74	6.71	173.02	32.77	3.35	6.71	173.02	32.77
Z mass	3.89	3.74	0.09	<0.02	2.69	2.17	0.14	0.05	2.46	1.59	0.10	<0.02
Φ mass	3.34	0.58	<0.03	<0.02	2.10	0.23	<0.03	<0.02	2.02	0.16	<0.03	<0.02
S/B		~ 5.3				~ 7.5				~ 9.6		

- All cross-section, but σ_{tot} in fb.
- ϵ stands for detector acceptance efficiency
- Signal cross-section & BR assumed to be independent from value of ξ ; SM cross-section & BR are assumed.
- Only $q\bar{q} \rightarrow ZZ$ contribution for ZZ cross-section taken in account; Contribution of $gg \rightarrow ZZ$ ($\sim 20\%$ of $\sigma_{q\bar{q} \rightarrow ZZ}$) not included.
- All background types, but ZZ negligible after selection for $m_\Phi > 2m_Z$
 \implies not taken in account in further analysis.

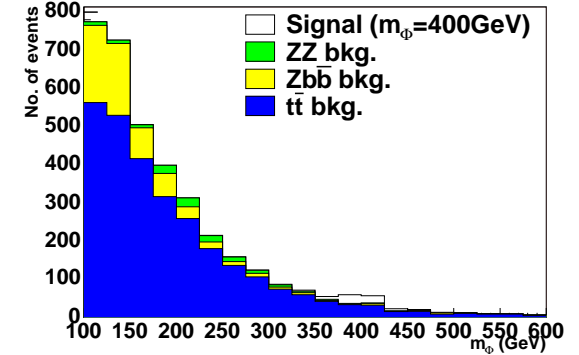
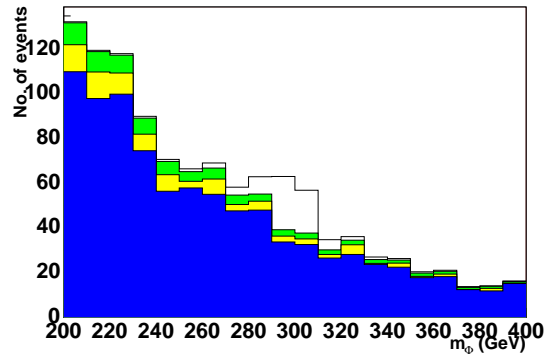
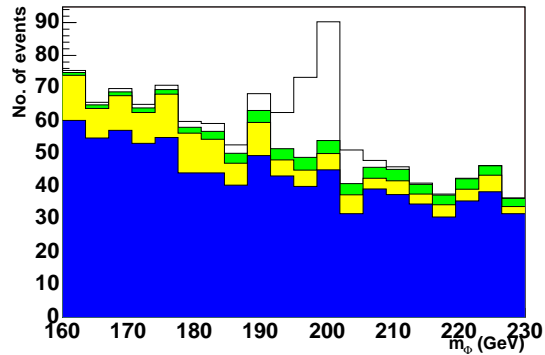
Selection: mass of 4 leptons

$m_\Phi = 200 \text{ GeV}$

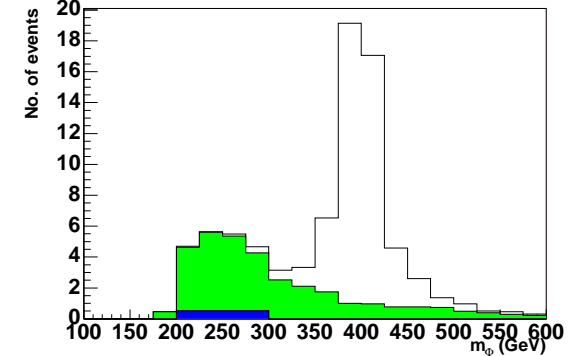
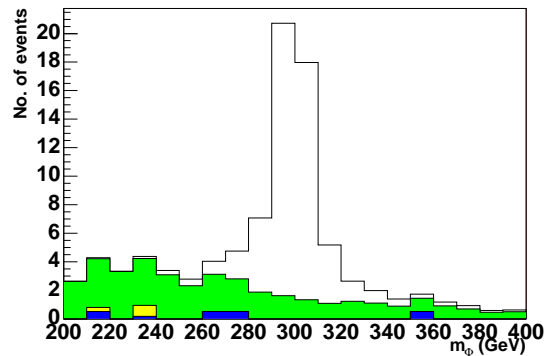
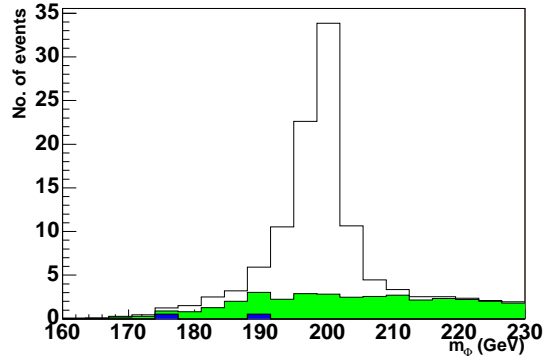
$m_\Phi = 300 \text{ GeV}$

$m_\Phi = 400 \text{ GeV}$

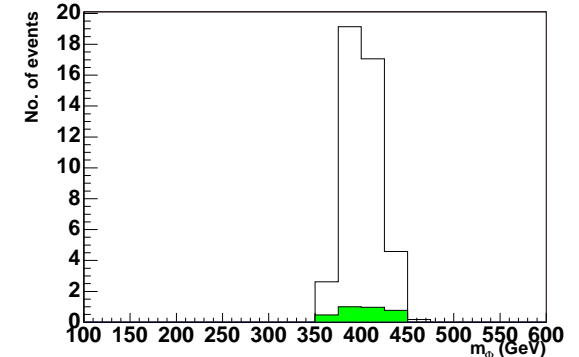
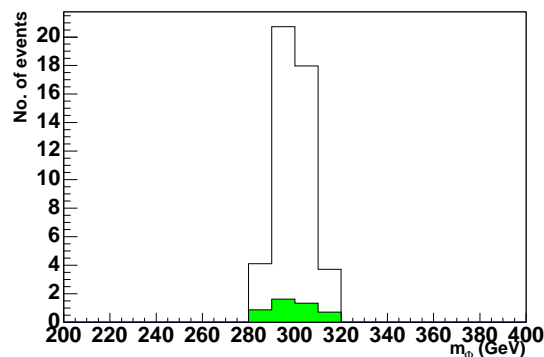
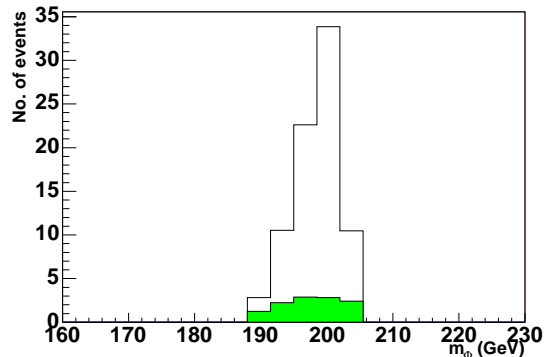
rec. $2\mu 2e$



Z mass cut



Φ mass cut



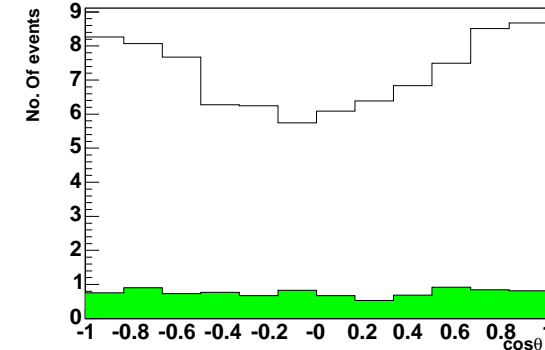
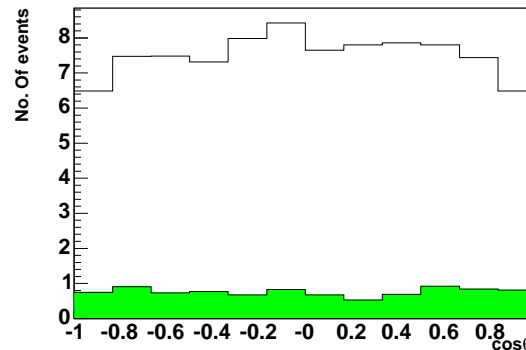
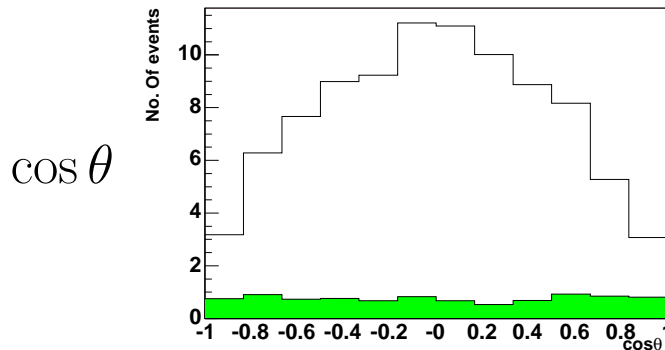
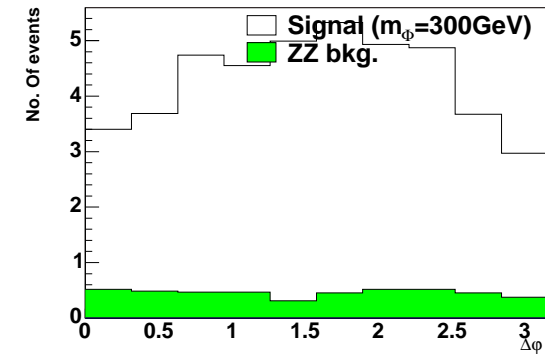
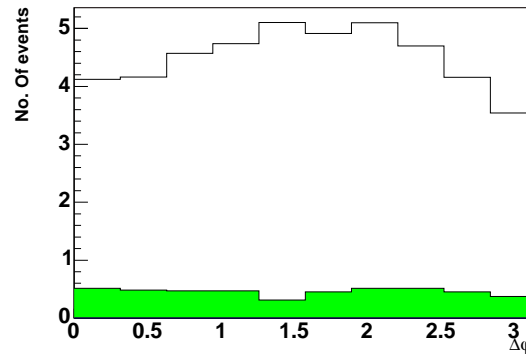
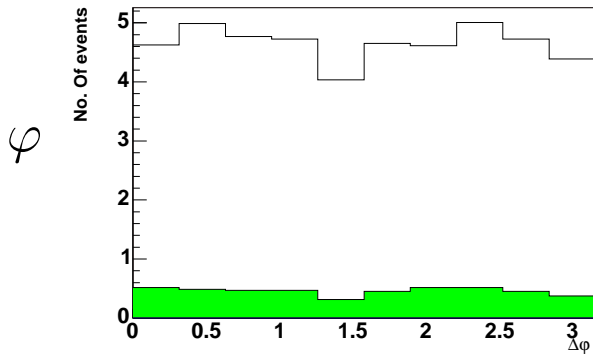
- Plots normalized to 20/fb (1 year of LHC at low lumi)

Reconstructed angle distributions

scalar

$\xi = \pi/4$

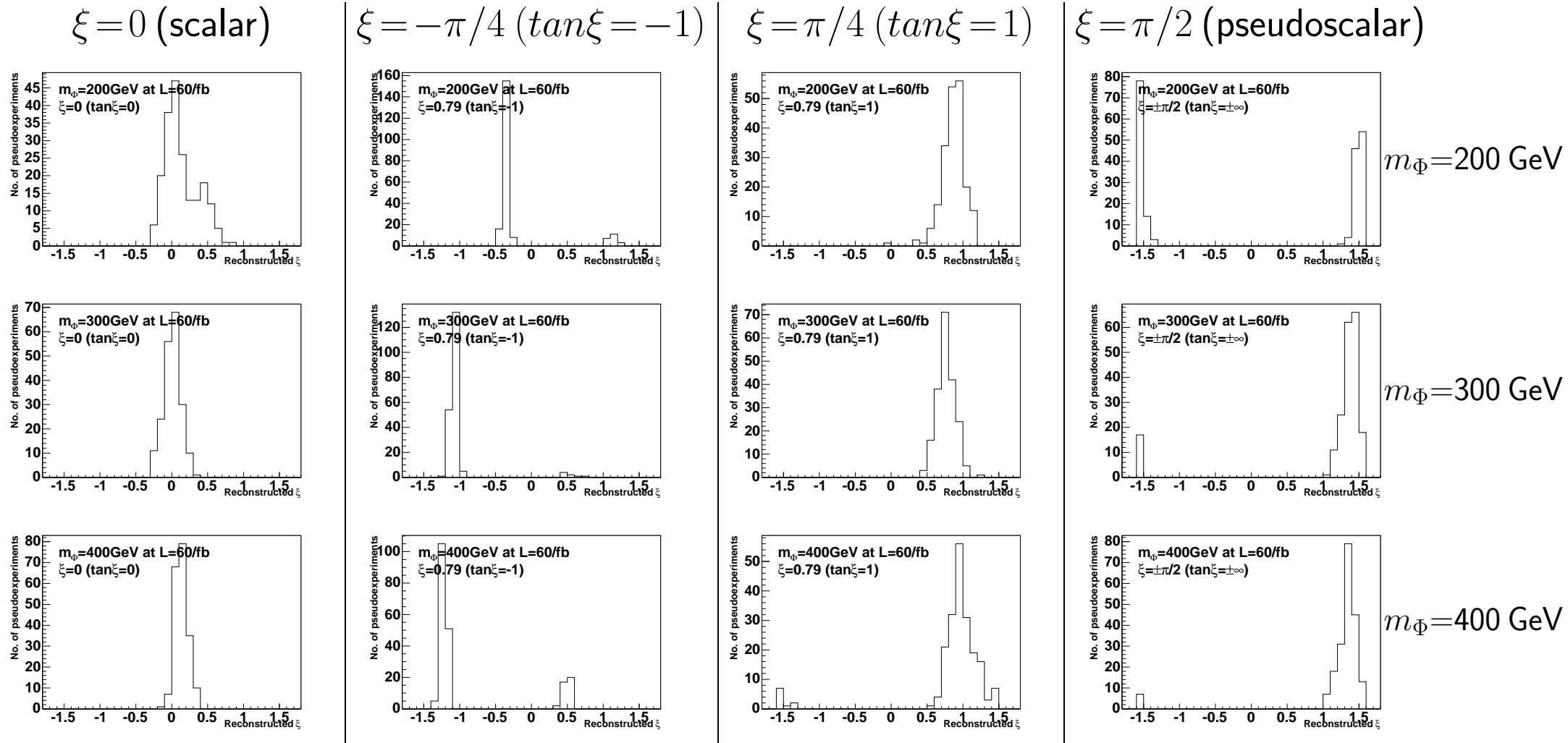
pseudoscalar



- Reconstructed angle distributions for $m_{\Phi}=300$ GeV
- Plots normalized to 20/fb (1 year of LHC at low lumi)
- Histograms for $\cos \theta$ contain sum of distributions for both θ_1 and θ_2
- Angle distributions not very smooth \implies bigger MC samples needed

Results: reconstructed ξ

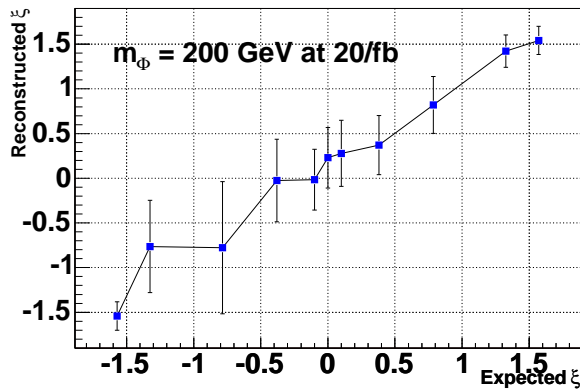
Distribution of reconstructed parameter ξ for 200 pseudoexperiments
(Results for $\mathcal{L}=60/\text{fb}$ – 3 years of LHC at low lumi.)



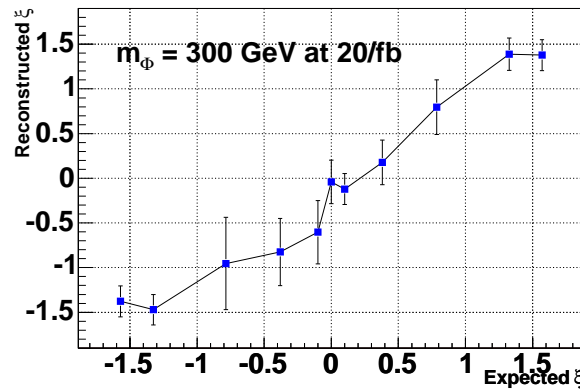
Results: ξ_{rec} in function ξ_{true}

- Mean and standard deviation of distribution of parameter ξ are taken as an estimator of ξ and $\Delta\xi$ respectively.
- Results for negative values of ξ with limited statistic.
- Results for integrated luminosities $\mathcal{L}=20, 60/\text{fb}$.

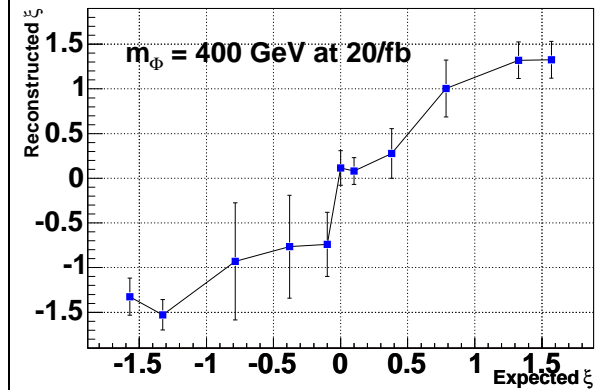
$m_\Phi = 200 \text{ GeV}$



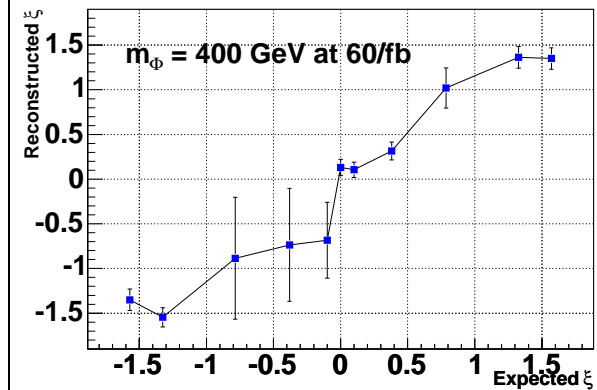
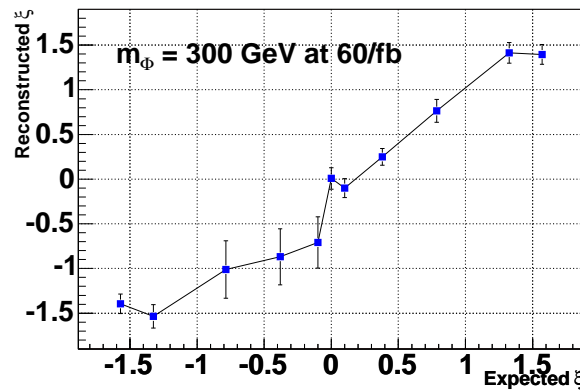
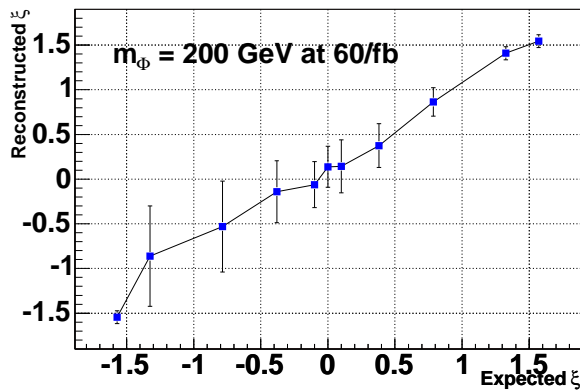
$m_\Phi = 300 \text{ GeV}$



$m_\Phi = 400 \text{ GeV}$



20/fb



60/fb

Summary

- Measurement of CP -violation in $\Phi \rightarrow ZZ \rightarrow 2e2\mu$ with the CMS is feasible
 - Sensitivity for distinguishing between CP -conserving and CP -violating cases is ~ 0.25 for 20/fb and ~ 0.15 for 60/fb.
- Quite good accuracy and precision of determination of ξ parameter:
 - Distinguishing between scalar and pseudoscalar with $S \equiv \frac{|\xi_A^{rec} - \xi_H^{rec}|}{\sqrt{(\Delta\xi_A)^2 + (\Delta\xi_H)^2}}$ bigger than ~ 3.5 for 20/fb and ~ 5.8 for 60/fb.
 - Precision of measurement of ξ in order of 0.3 for 20/fb and 0.2 for 60/fb.
- To improve results smoother angle distributions for reference probabilities \mathcal{P} are needed
 - \implies bigger MC samples for scalar, pseudoscalar and $\xi = \pi/4$.
- All $\Phi \rightarrow ZZ \rightarrow 4\ell$ channels should be used.

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