

Cosmic Ray Interactions near the GZK and below

Skopelos, 2005

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Overview

introduction

simulation of air-showers

MC - hybrid

observables

long-lat

detector-simulations

input from hadronic models

relevant physical input: σ , dn/dxF ,
mult for muons, K - forward scattering

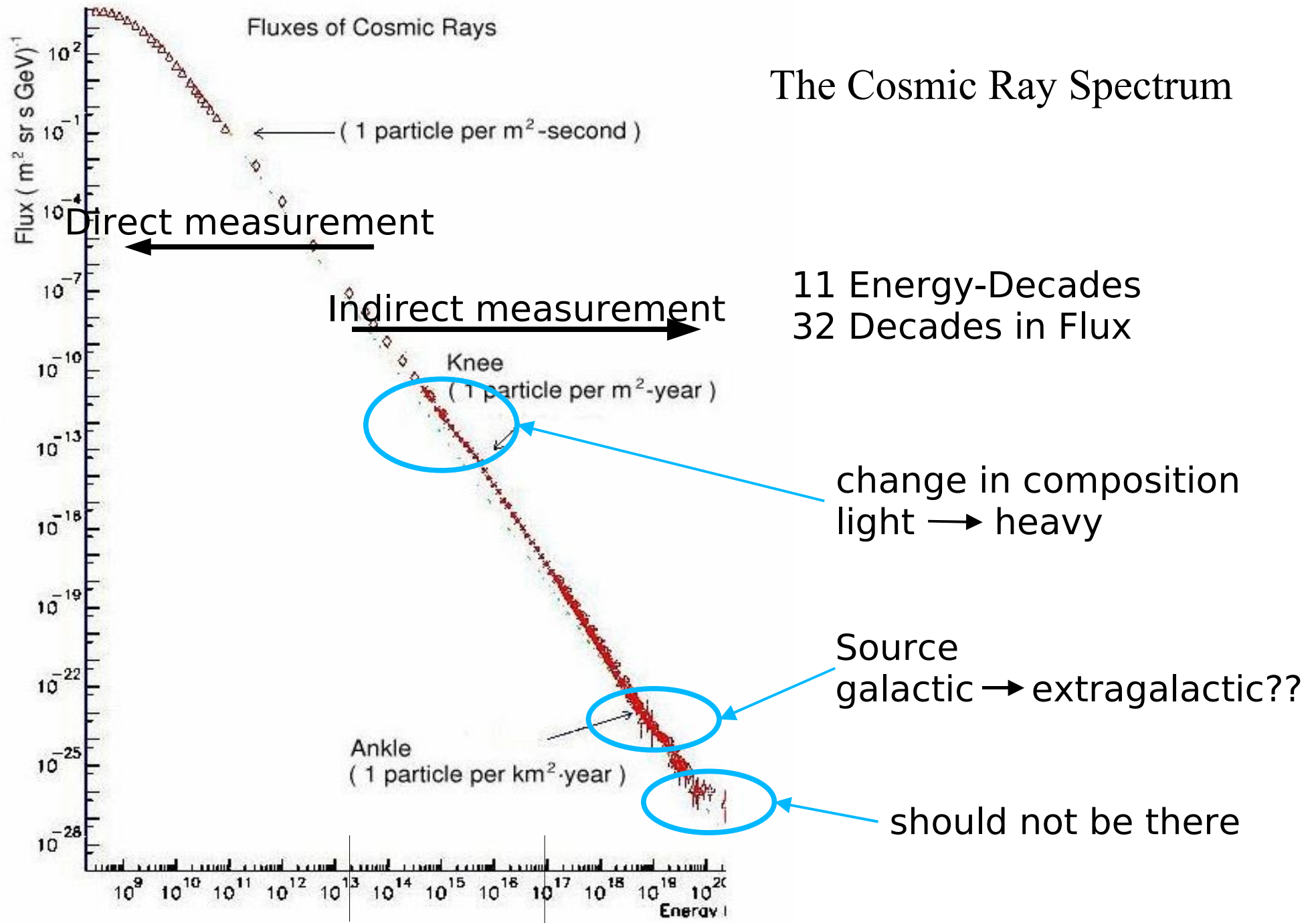
hadronic models and air shower data:

muons bundles

Kascade data

Lateral distribution functions

The Cosmic Ray Spectrum



RHIC

LHC

Status quo in ultra high energy cosmic rays

Hires and AGASA disagree on:

- normalization of cosmic ray flux
- trans-GZK events
- clustering of cosmic rays

Auger:

so far SD measurement calibrated with FD
no events above $1e20$

Detection of UHECR

Direct detection not possible:

Flux very low:

$E > 10^{20}$ eV

-> 1 particle/km²/century

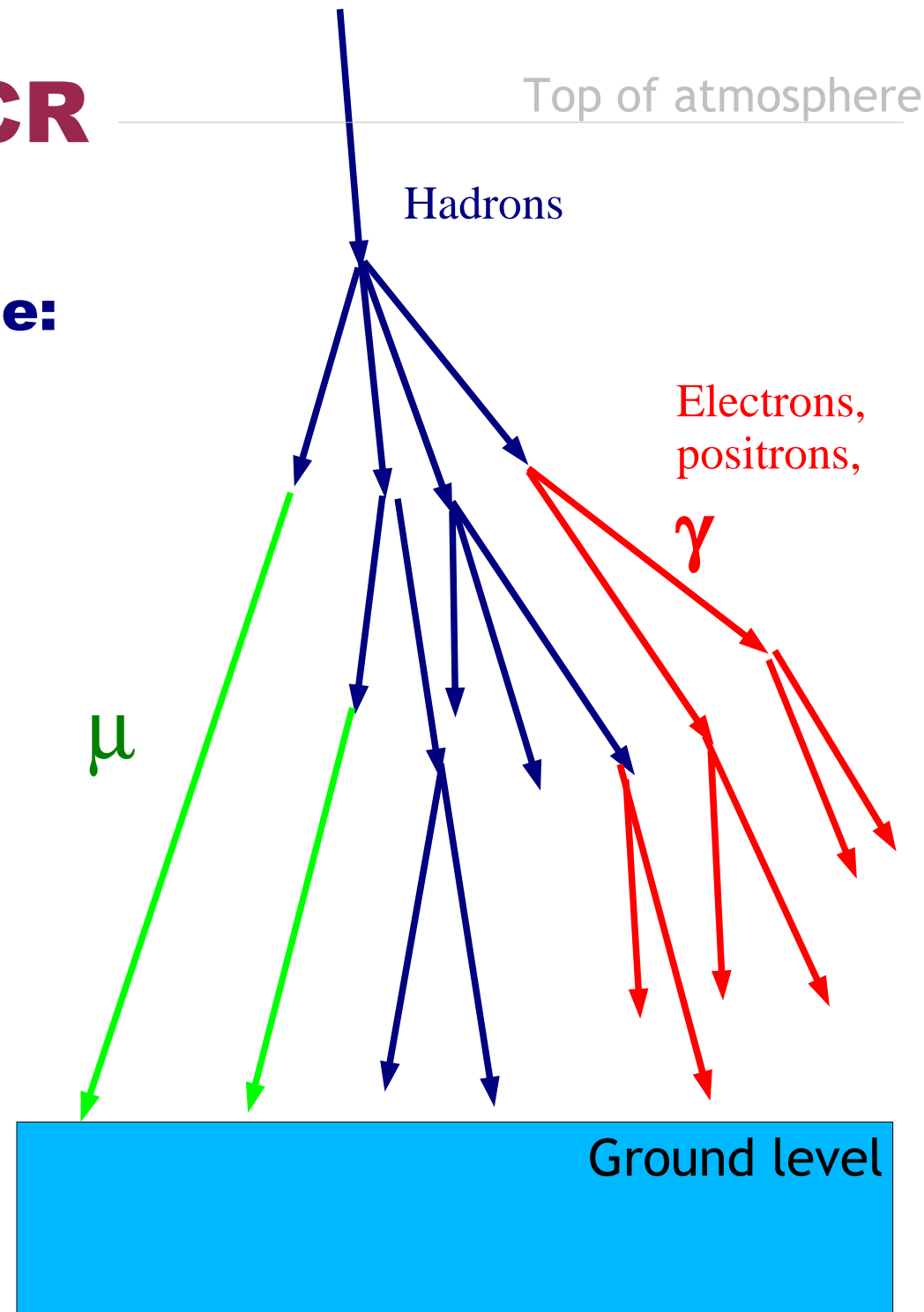
indirect detection via

AIR SHOWERS

induced by UHECRs

Reconstruct primary
from shower properties:

- Energy,
- Arrival direction,
- Particle type



Traditional Monte-Carlo method for air shower simulations

Choose some high energy and low energy hadronic model

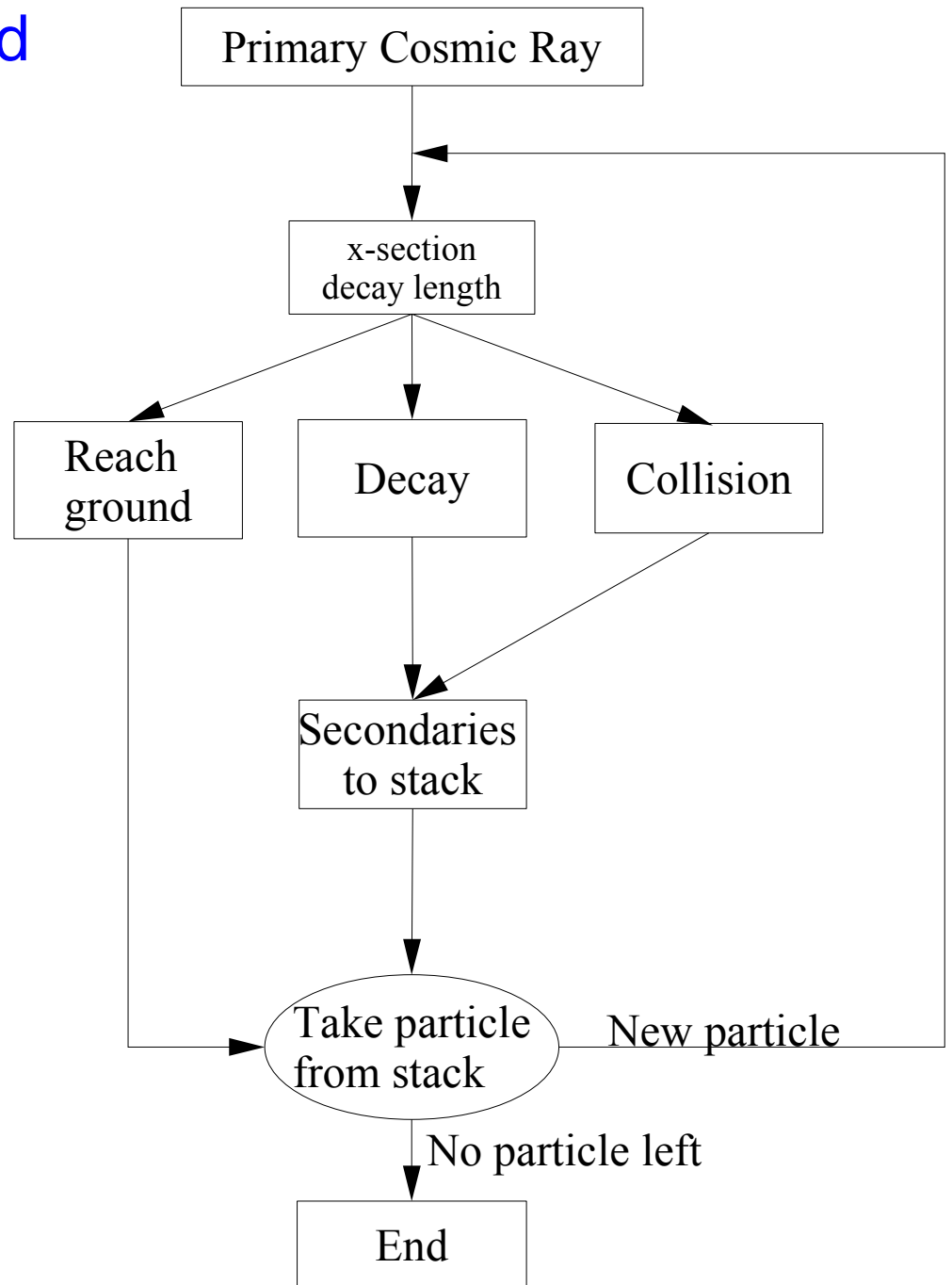
choose model for electromagnetic interactions

Interaction length dependent on thickness of material

decay length dependent on geometrical path

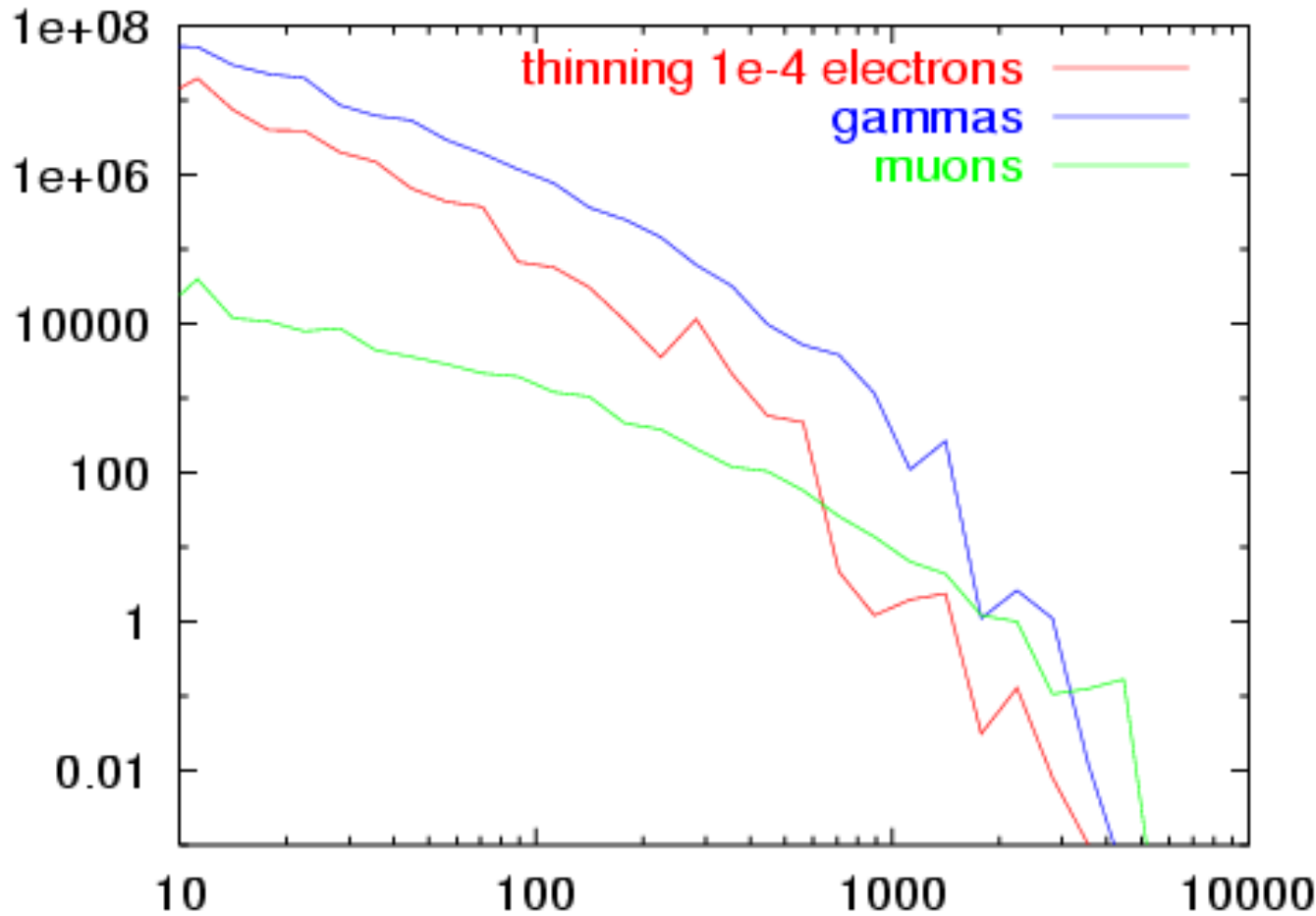
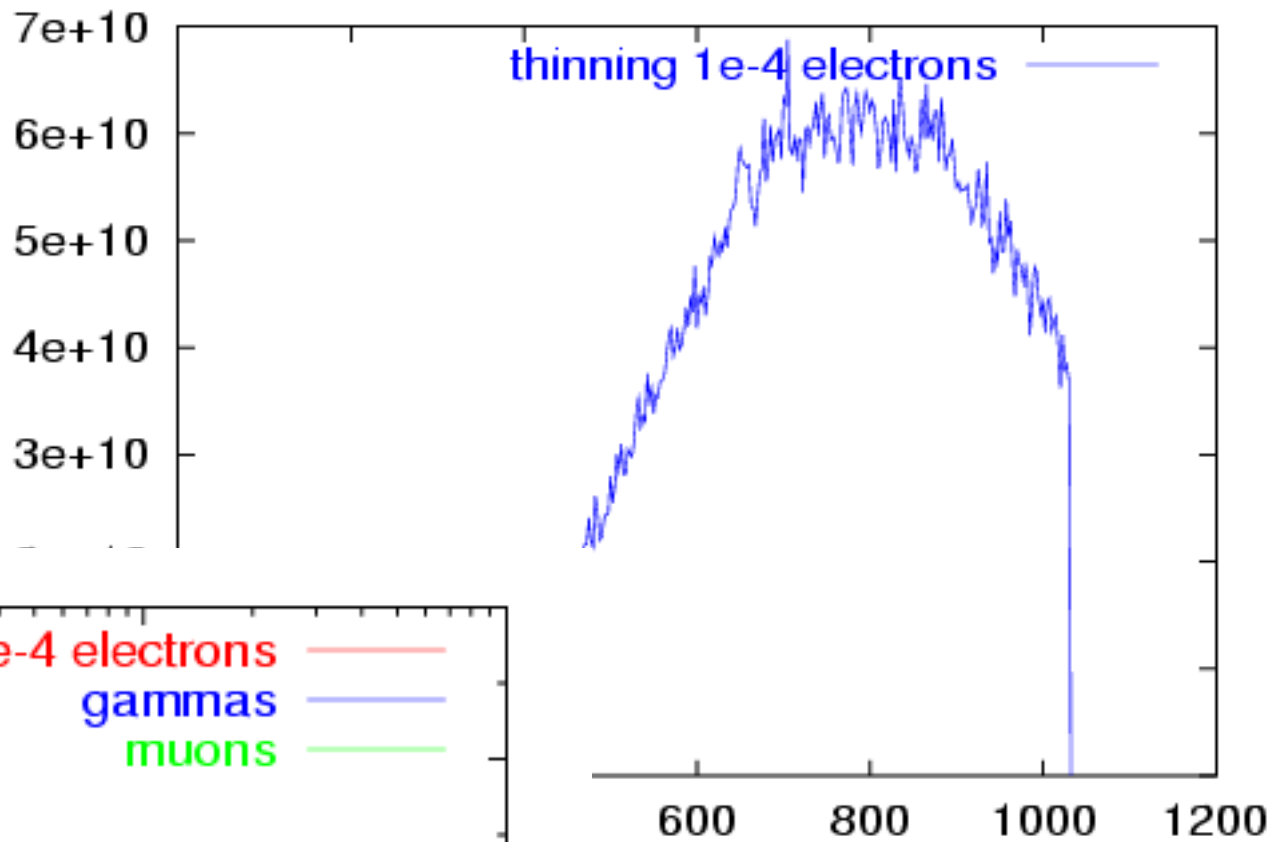
- MOCCA (Hillas),
- CORSIKA (Heck, Knapp et al),
- AIRES (Sciutto)

At 10^{10} GeV at $X_{\max} \sim 6 \times 10^9$ particles ---> 1 year to simulate



Thinning

Below E_{th}
follow only a fraction
of secondaries
to save CPU-time



need better (faster)
methods for high
energies

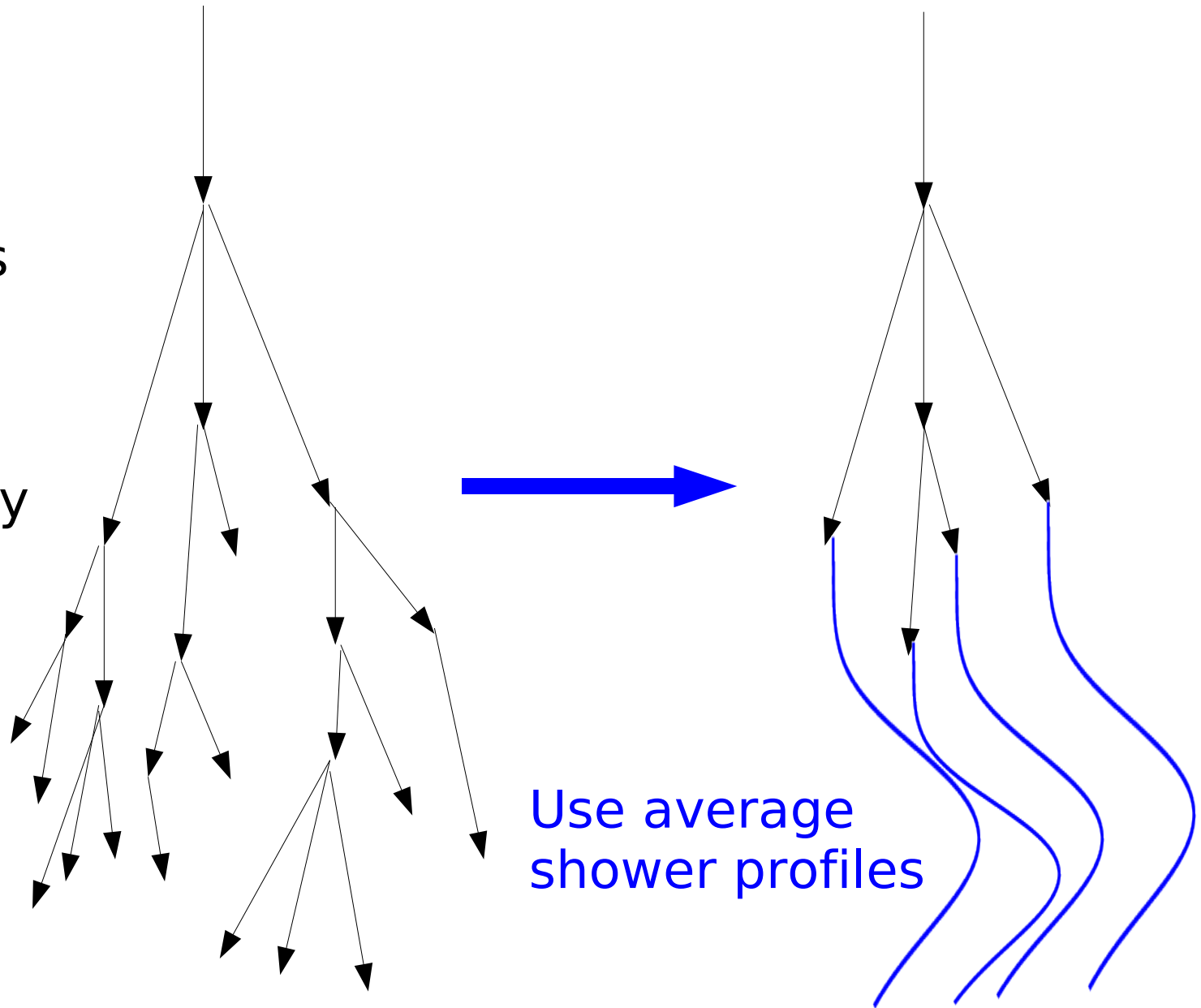
→ hybrid approach

1D Hybrid approach with Cascade Equations

neglect initial lateral spread

initial fluctuations by MC

compute shower profile numerically



Cascade equations

$$\frac{\partial h_n(E, X)}{\partial X} = -h_n(E, X) \left(\frac{1}{\lambda_n(E)} + \frac{B_n}{E X} \right) + \sum_m \int_{E_{\min}}^{E_{\max}} h_m(E', X) \left(\frac{W_{mn}(E', E)}{\lambda_m(E')} + \frac{B_m D_{mn}(E', E)}{E' X} \right) dE'$$

h_n : number of hadrons n per dE

E : Energy X : slant depth

B_n : decay constant

λ_n : interaction length

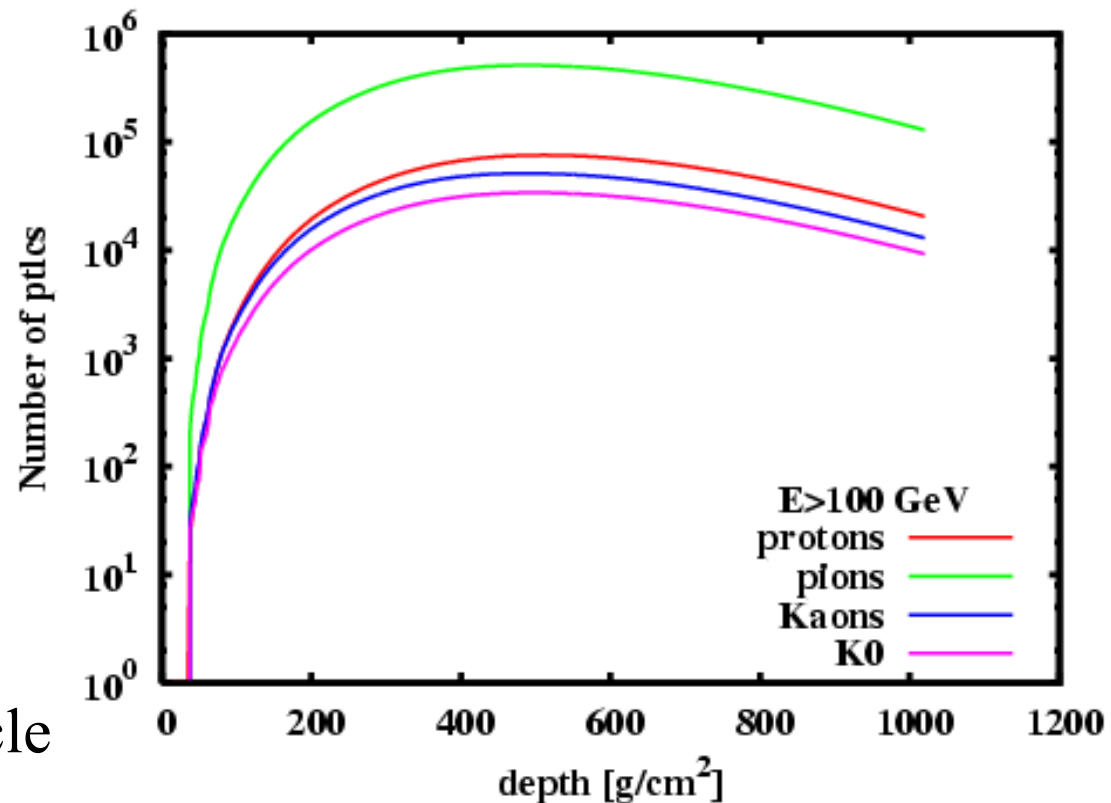
W_{mn} : collision function

D_{mn} : decay function

initial condition:

$h_n(E, X) = \delta(E - E_0)$ for a given particle

Solution:
 $h_n(E, X)$



The physics of hadronic interactions

$$\lambda_n = \frac{A m_u}{\sigma_{inel}}$$

Mean free path of particle n in the atmosphere

$$W_{mn}(E, E') = \frac{dN_{mn}}{dE'}$$

Spectrum of produced particles of type n from m+Air collision

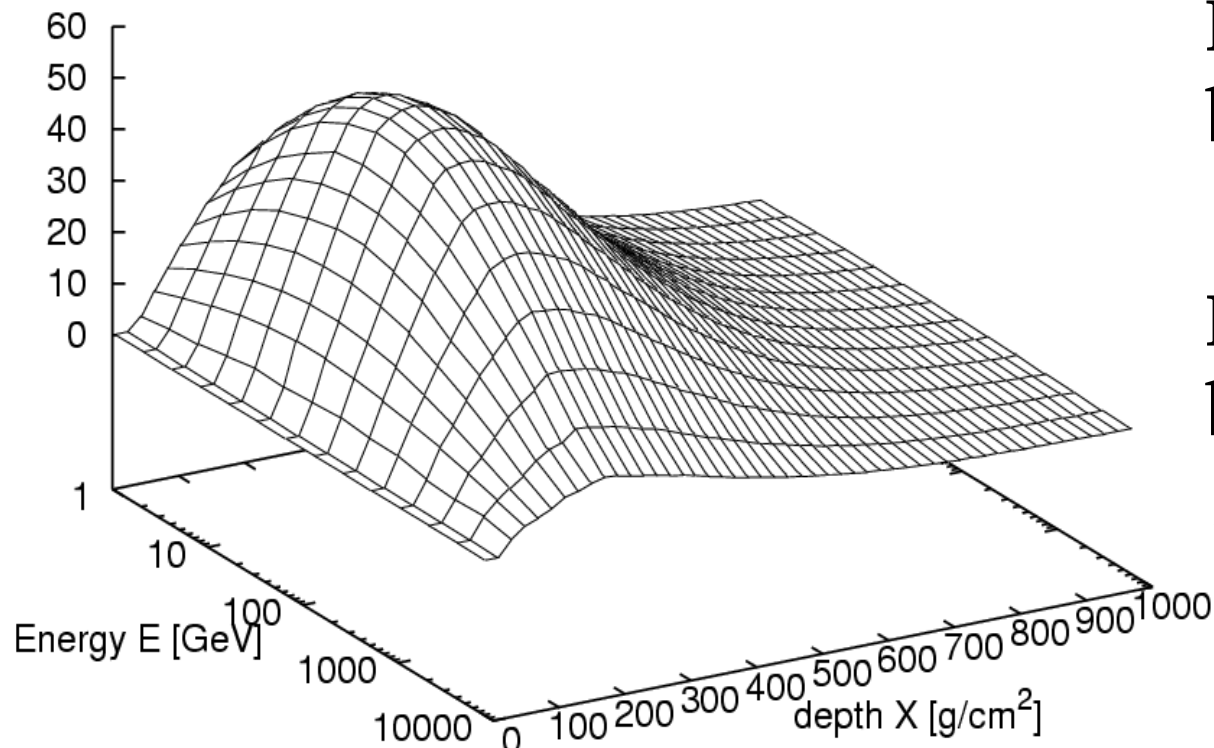
tabulated from MC

The Rest: Decays, electromagnetic interactions, muon propagation, under control (?)

Source function for low energy particles

$$\frac{\partial h_n^{source}(E, X)}{\partial X} = \sum_m \int_{E_{min}}^{E_{max}} h_m(E', X) \left(\frac{W_{mn}(E', E)}{\lambda_m(E')} + \frac{B_m D_{mn}(E', E)}{E' X} \right) dE'$$

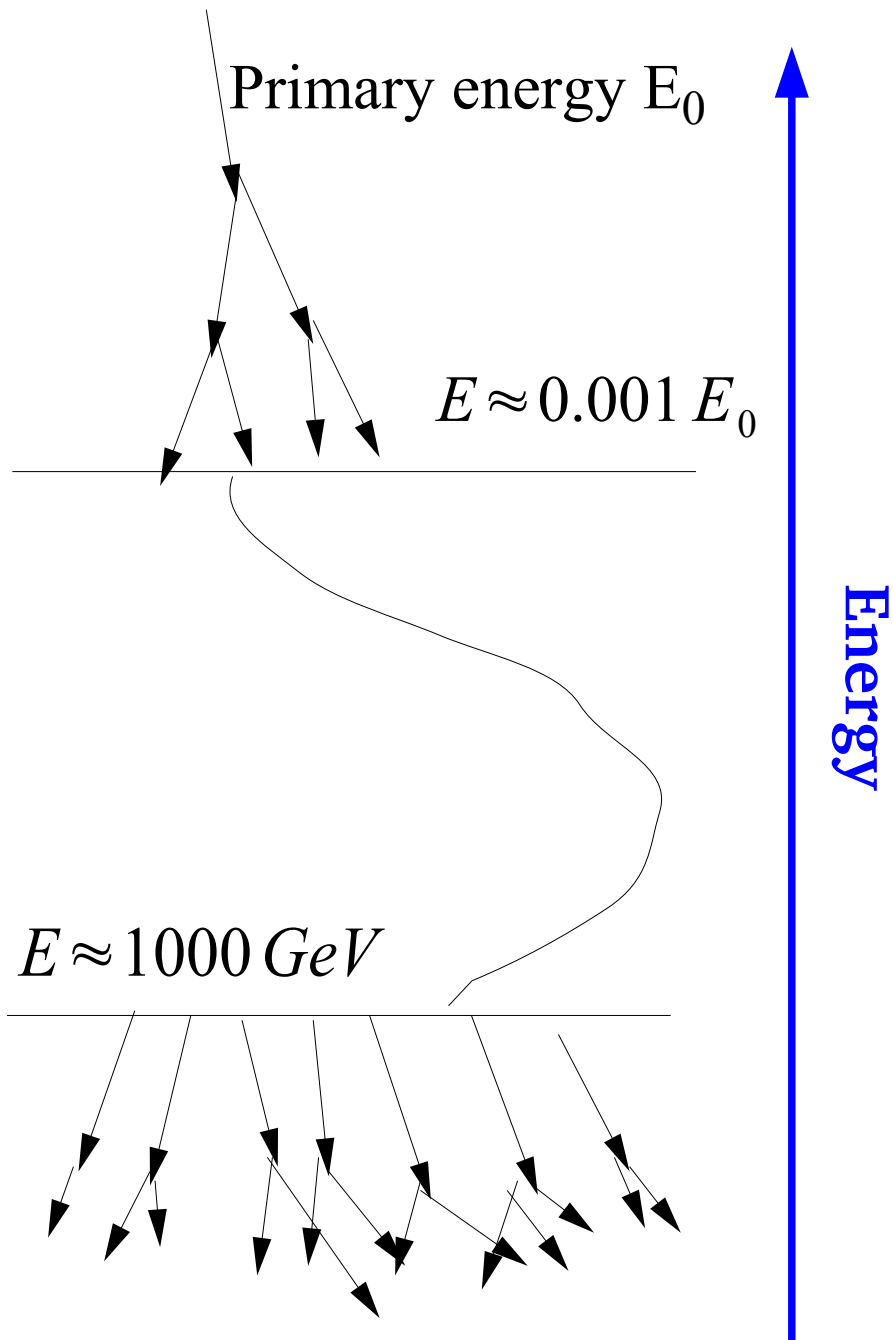
$E \frac{dh_{\pi}^{source}(E, X)}{dX}$



Production of hadrons
below $\sim 10^4$ GeV

Electrons and photons
below ~ 10 GeV

Place particles along shower axis



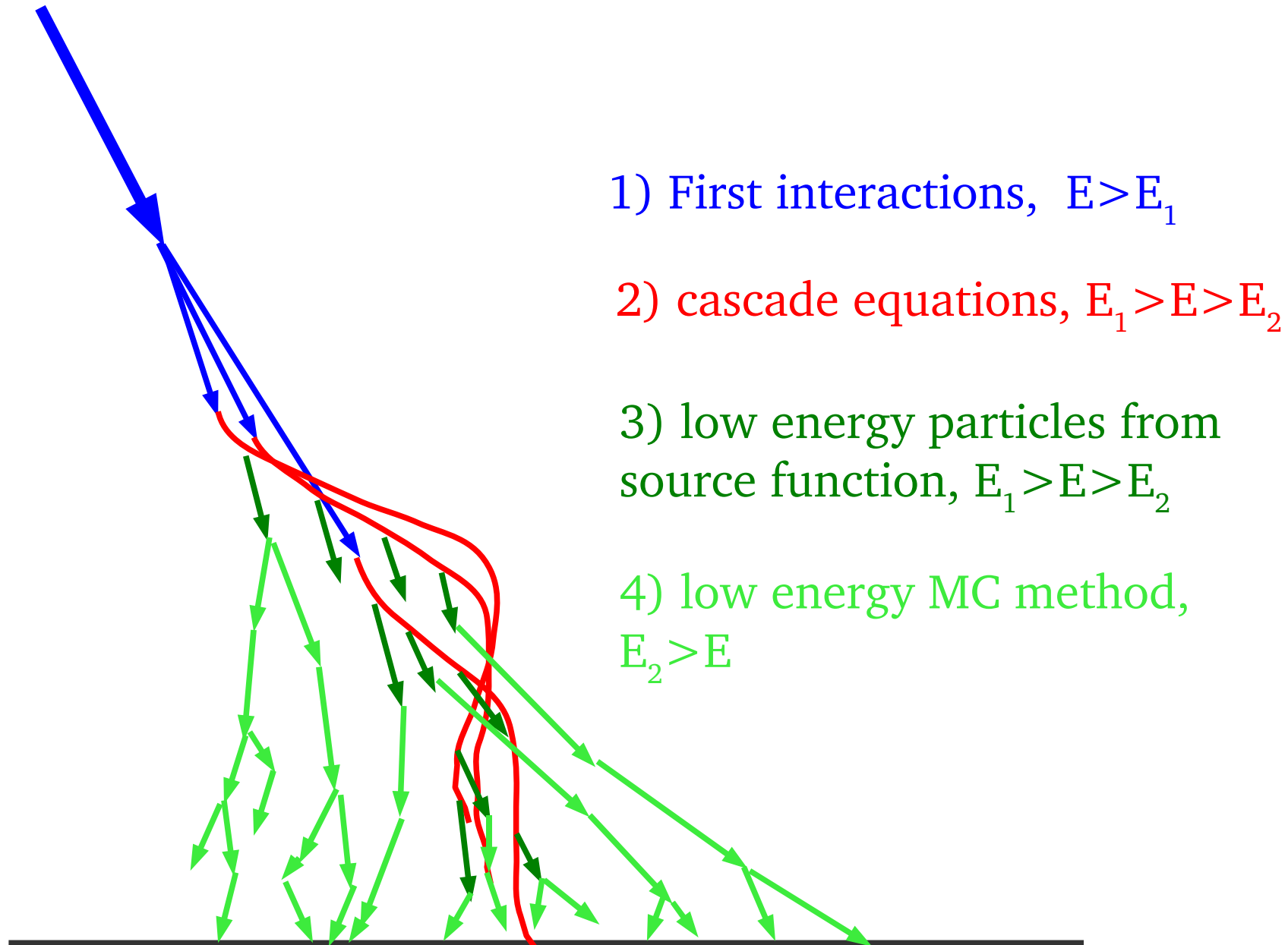
The initial part of the shower is calculated explicitly to get natural fluctuations.

Once multiplicities grow cascade equations solve the transport analytically. No thinning required !!

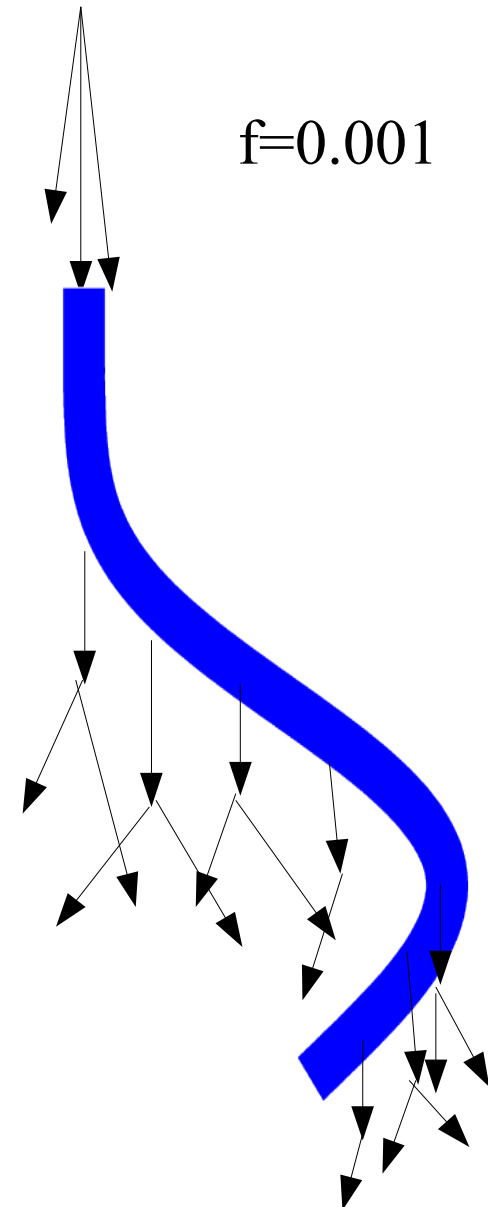
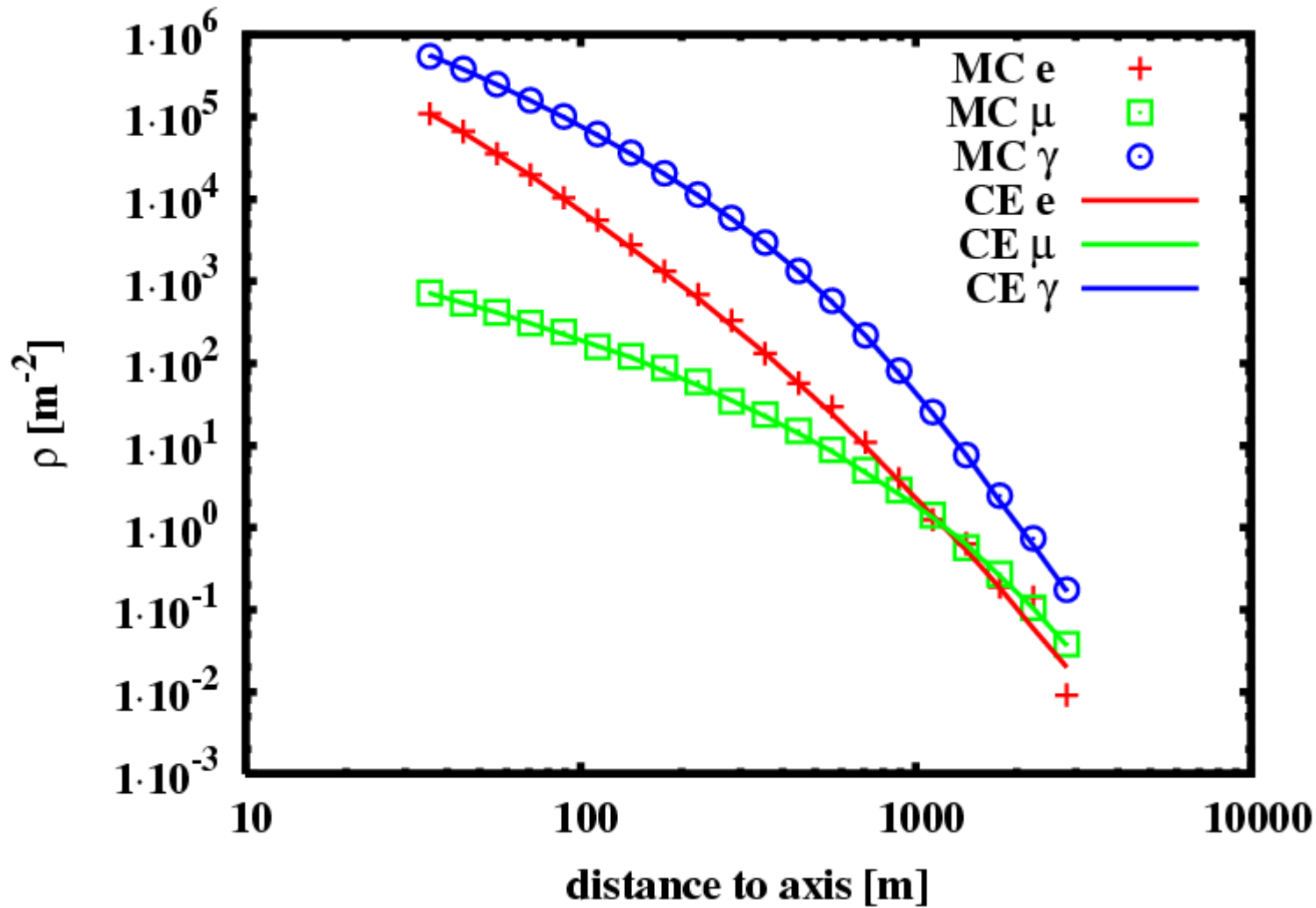
The low energy part again can be done explicitly by Monte Carlo to get the lateral spread of the shower.

40x faster than pure MC

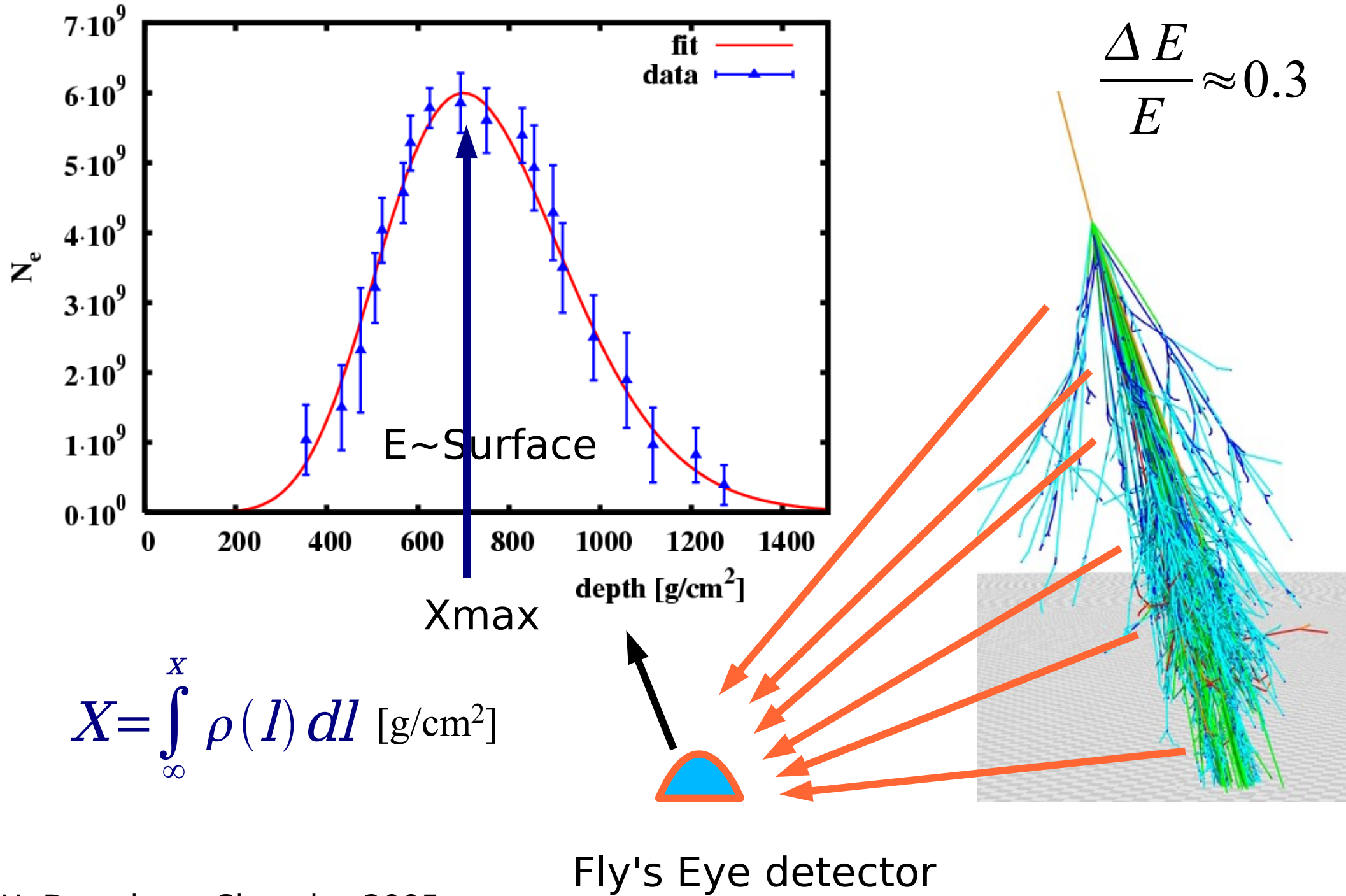
Visualization of procedure



Lateral distributions
with same high energy part



Longitudinal measurement with fluorescence light



Lateral distribution functions

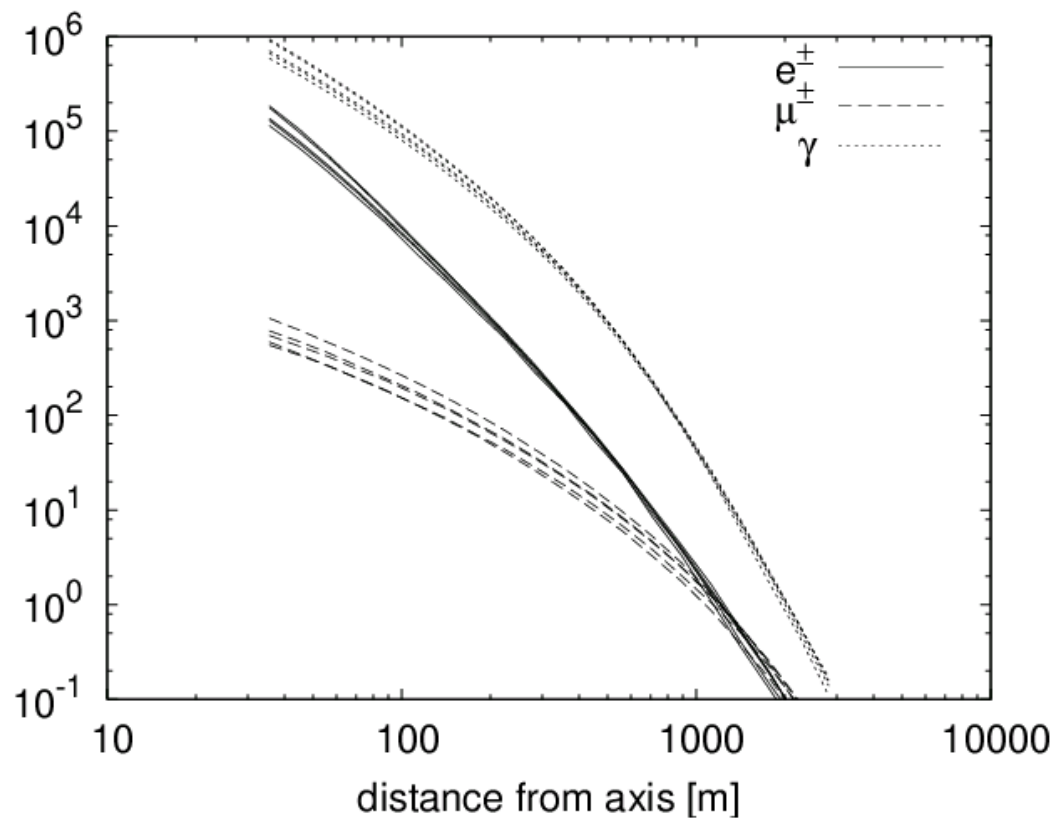
Fluctuations of the LDF of different showers cancel
at some empirical distance from the shower axis
--> Energy estimator

try to extract information about muons --> composition estimator

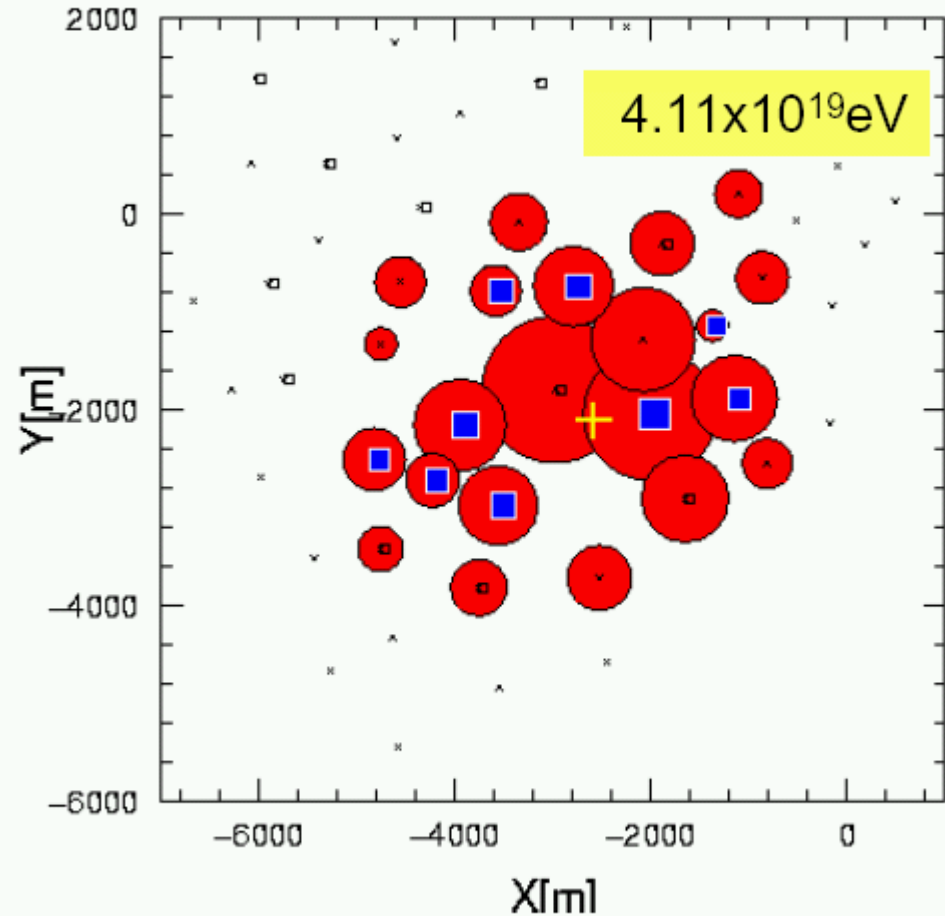
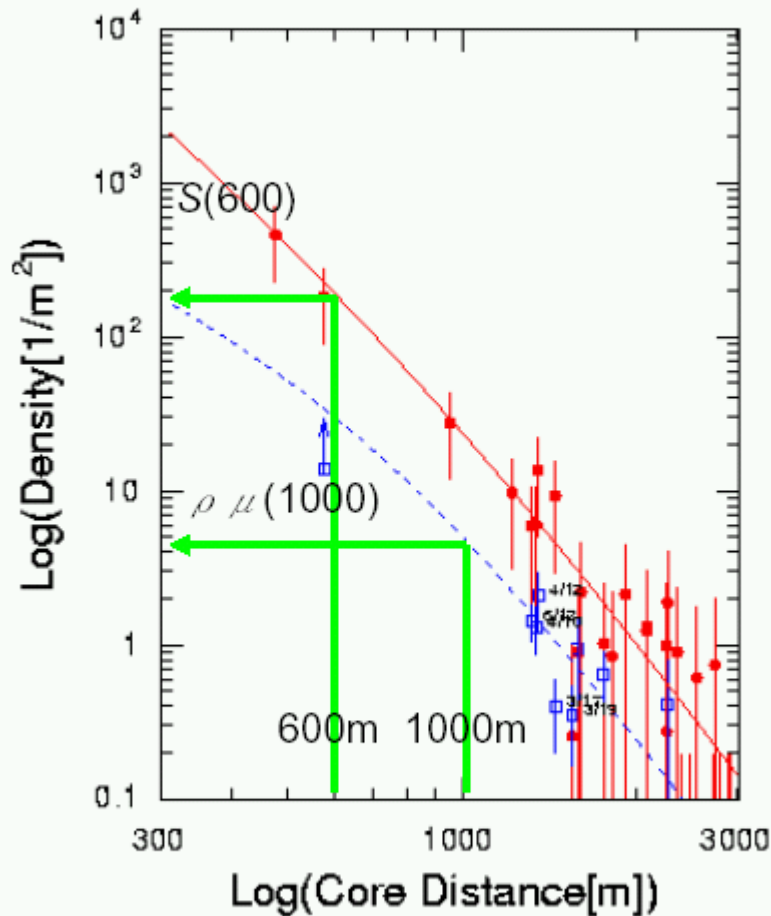


Common ground detectors:

- plastic scintillators
- water cherenkov



Event sample & observables



Energy estimator (charged particle density @600m): $S(600)$

From Kenji Shinozaki (AGASA)

Event generators commonly used in Air Shower Applications

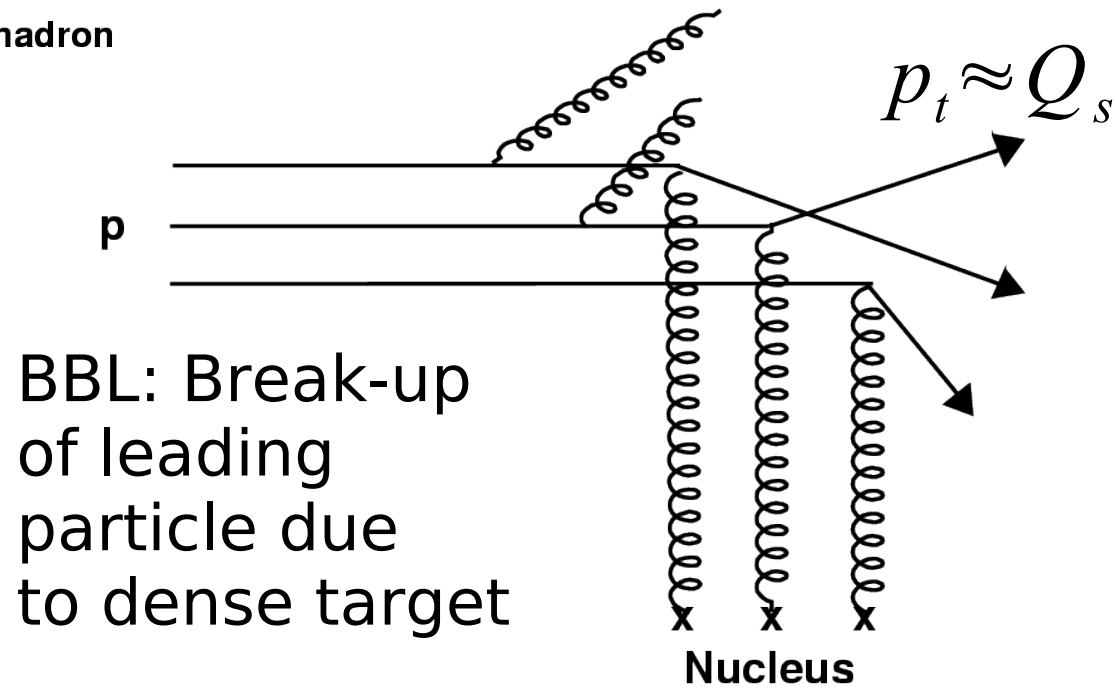
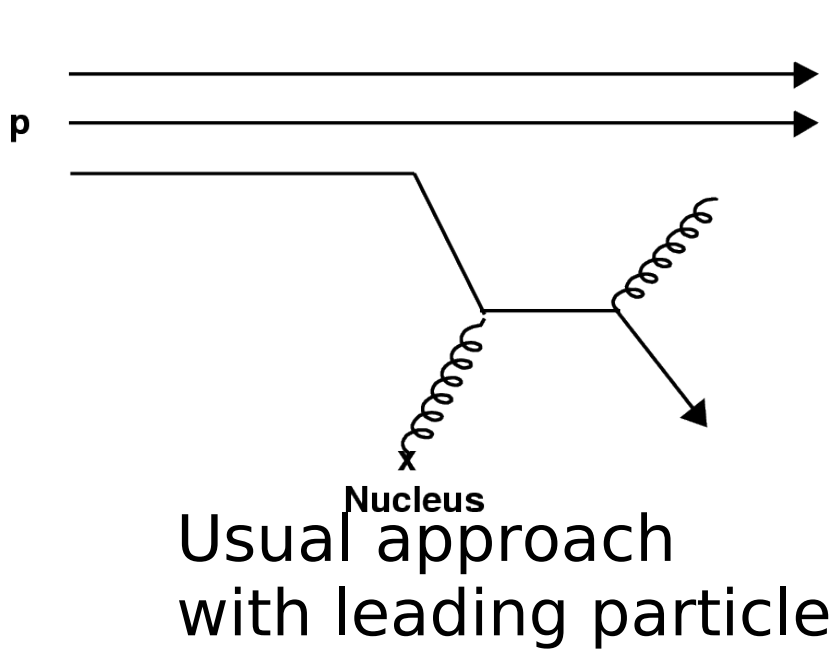
QGSjet01 (Ostapchenko, Kalmykov)
multiple soft and semihard Pomeron exchange

QGSjet II (Ostapchenko)
enhanced Pomeron diagrams

Sibyll 2.1 (Fletcher, Engel et al.)
pQCD event generator with multiple soft int.
energy dependent p_t -cutoff

BBL 1.0 (HJD, Dumitru, Strikman)
Implements parton saturation for forward
scattering

BBL Monte Carlo



Valence quarks,
gluon distribution:

$$P_i(x) = f_i(Q_s^2(x), x)$$

$$\langle p_t \rangle \approx Q_s(x)$$

valence quarks: GRV94 PDF ($xf(x)$ dominant at high x)

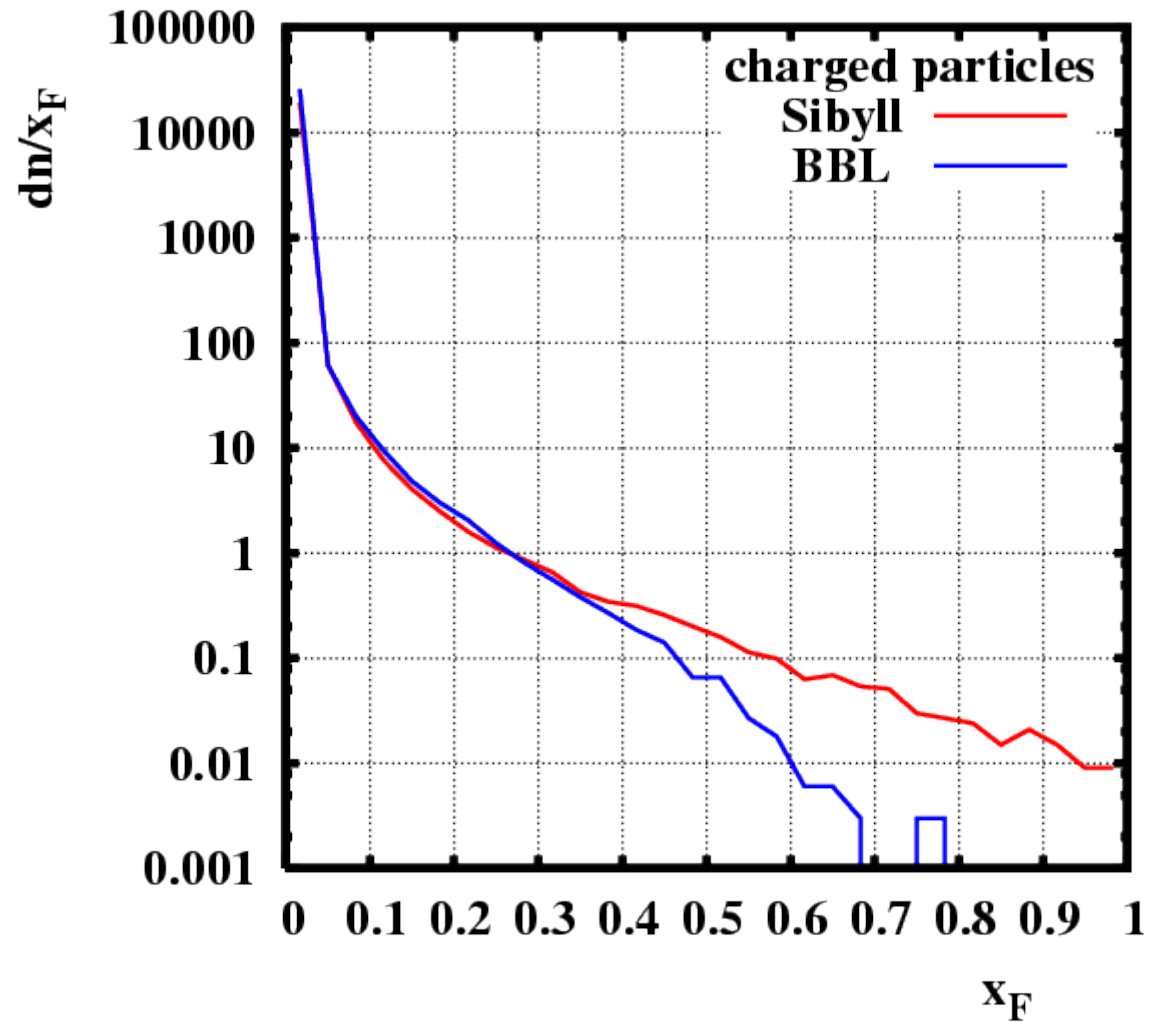
gluons: $x g(x, q_t^2) \propto \frac{1}{\alpha_s} \min(q_t^2, Q_s^2(x)) (1-x)^4$

(Kharzeev, Levin, Nardi, NPA 2004)

Suppression of forward scattering

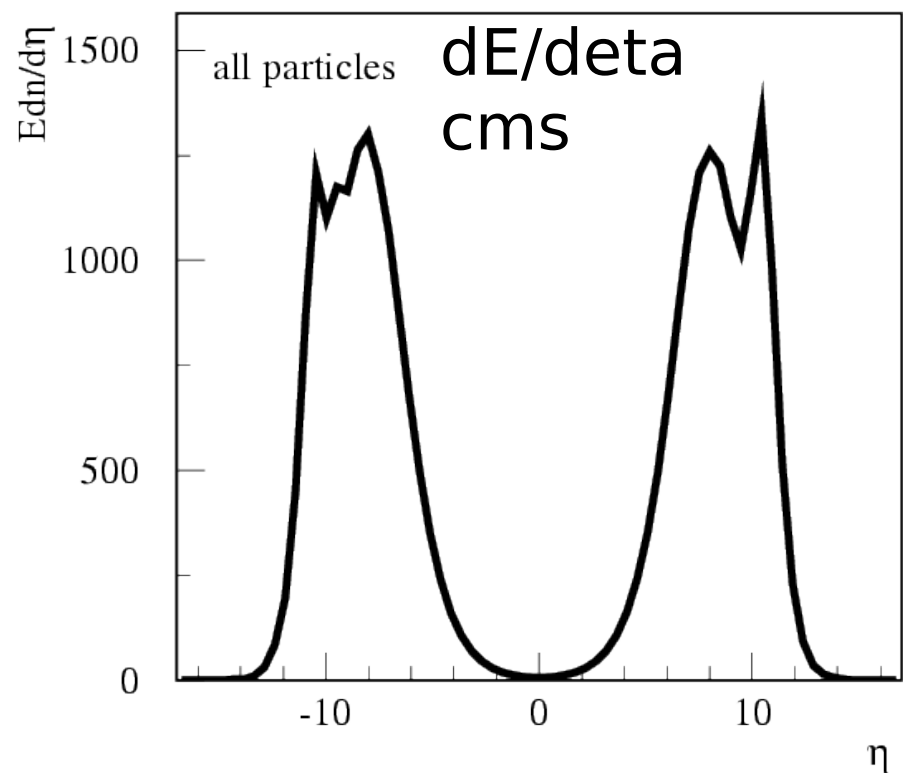
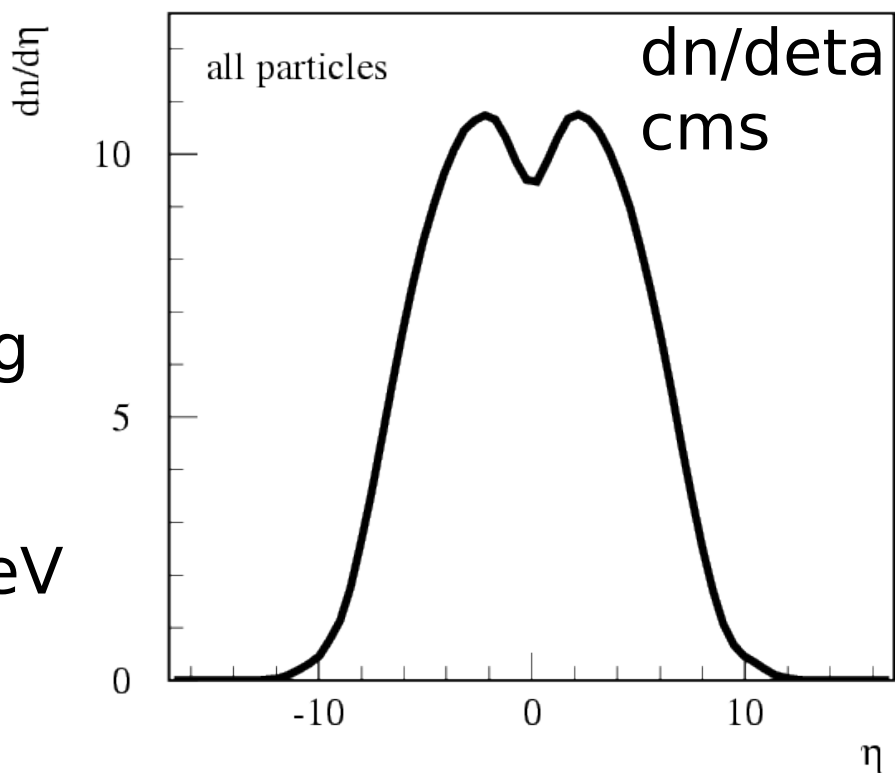
Central p-N scattering

Forward scattering suppressed due to independent fragmentation of leading quarks

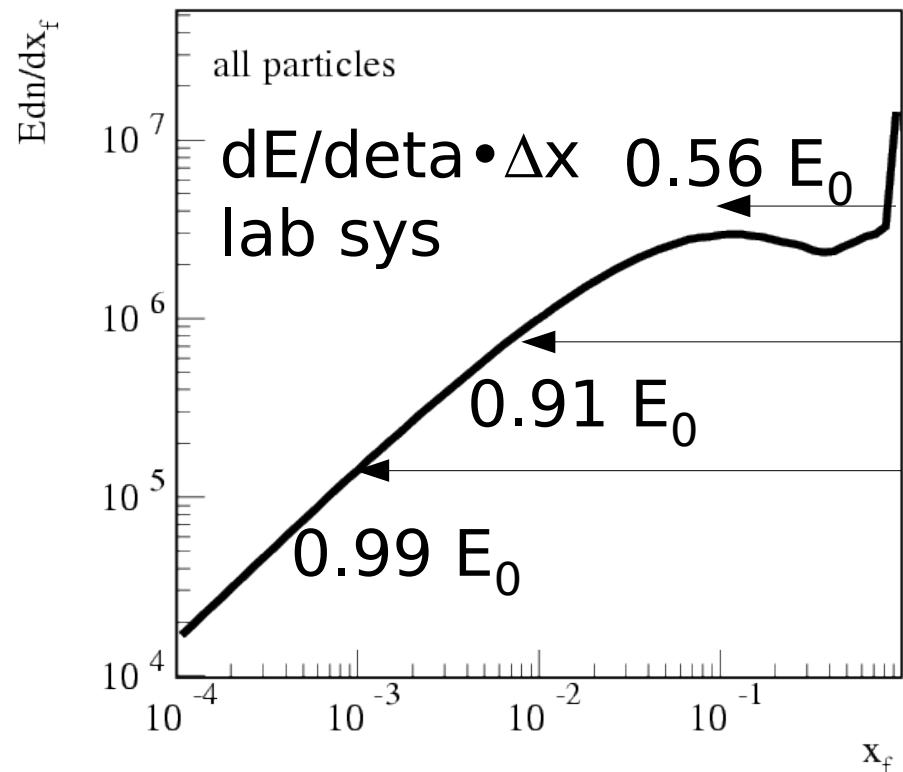
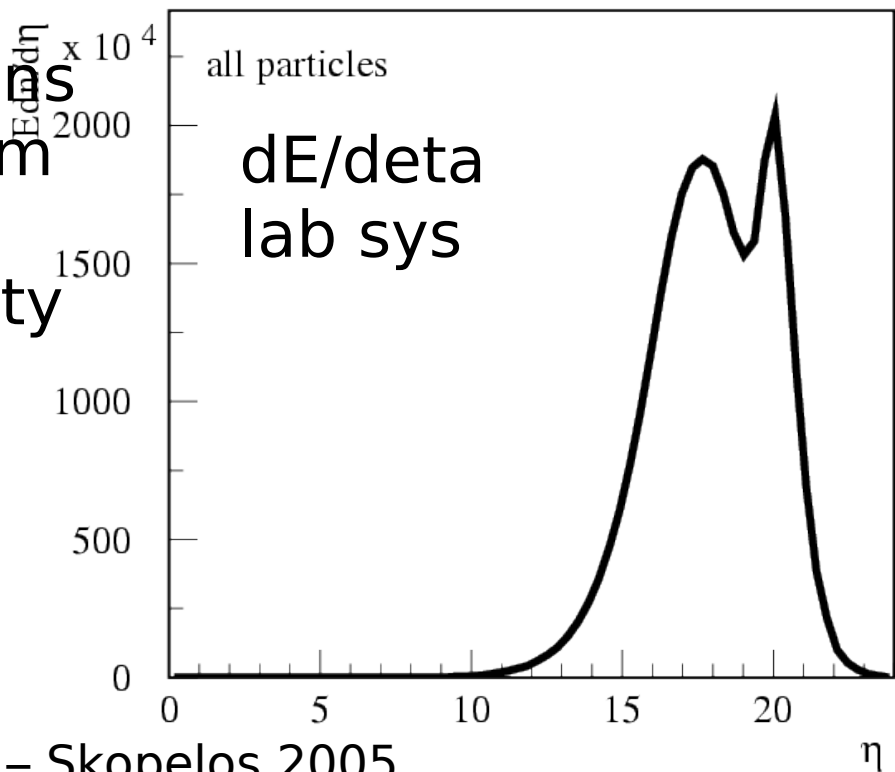


Forward scattering

for p-p
14000 GeV
cms



still, muons
come from
high
multiplicity

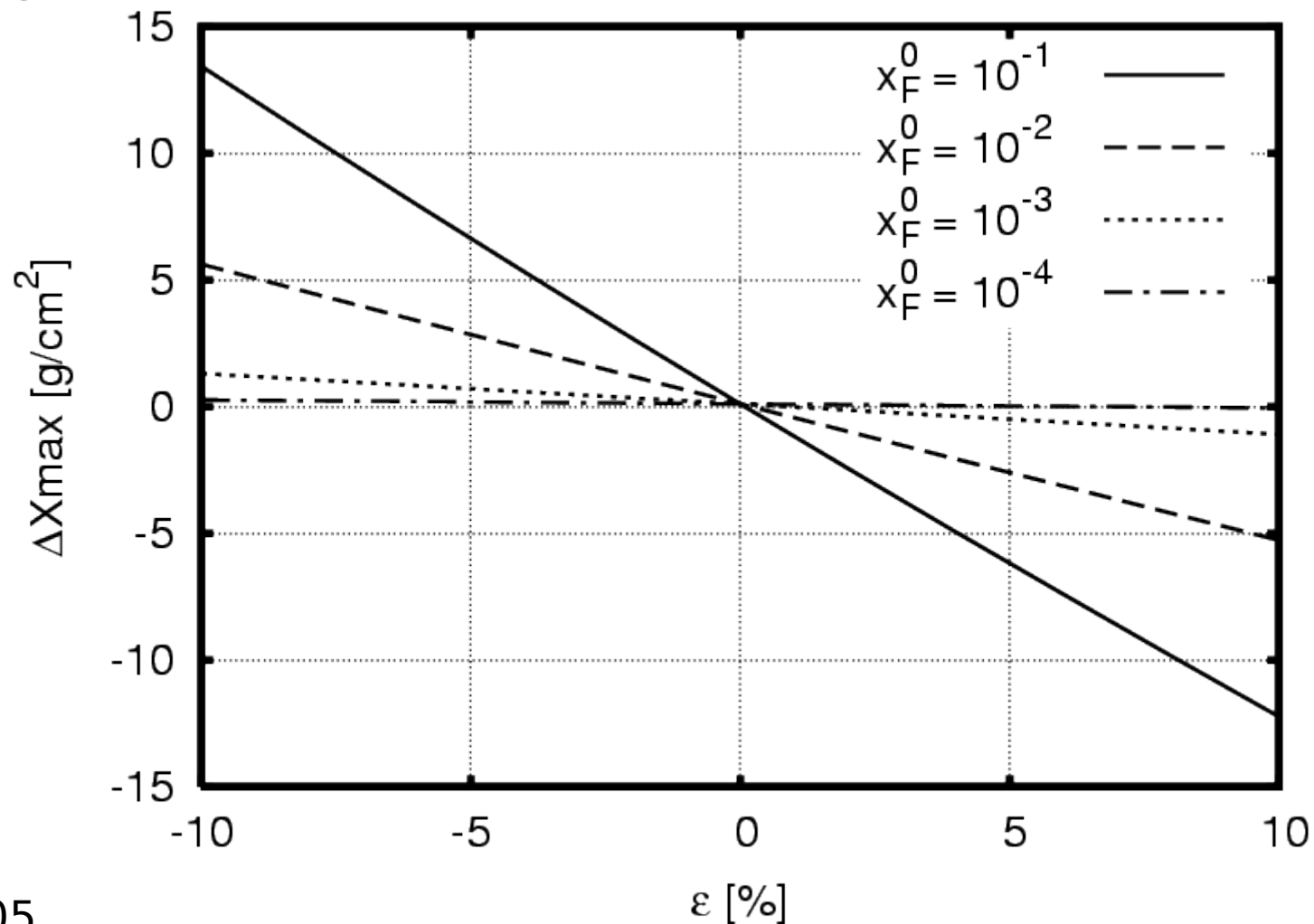


Sensitivity of X_{\max} to forward scattering

Modify dn/dx_F by factor $(1+\varepsilon)$ for $x < x_F^0$
 $(1-\varepsilon')$ for $x > x_F^0$ to conserve energy

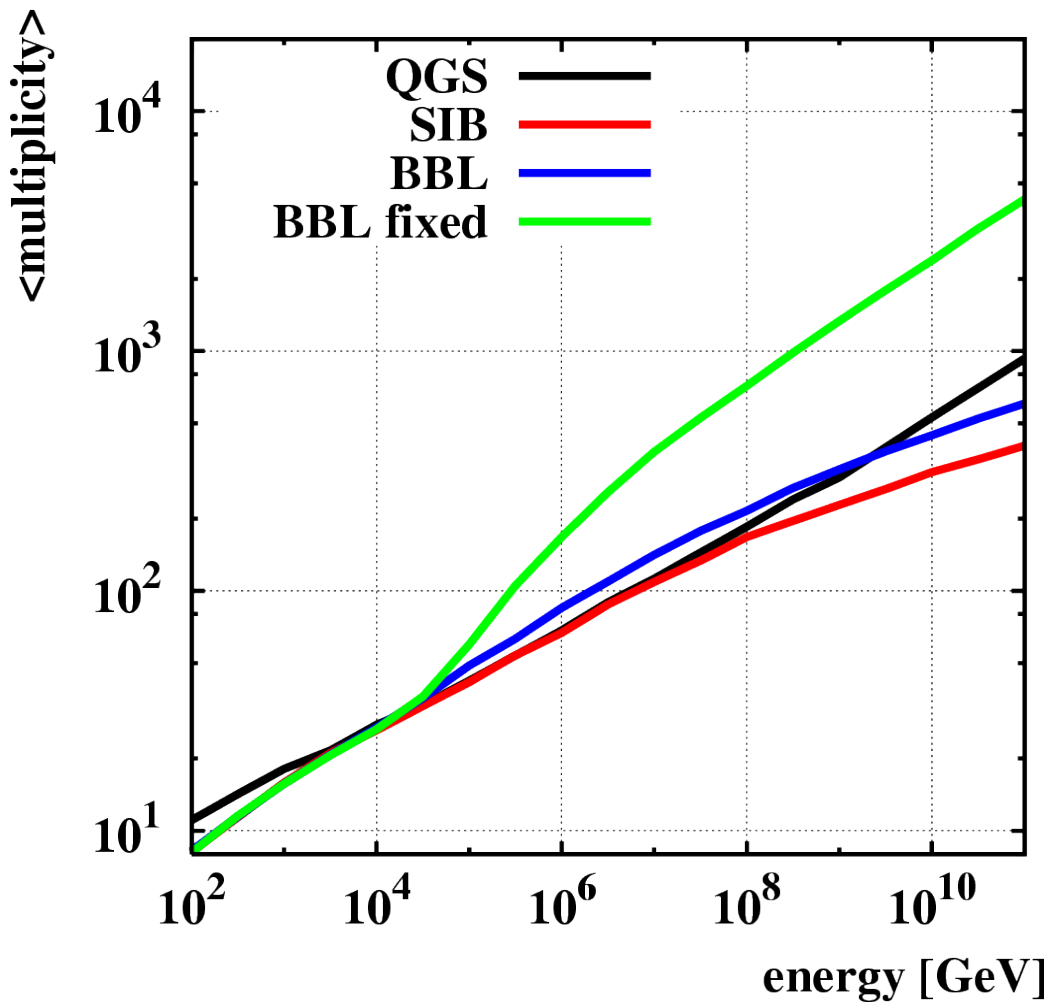
Only possible
with transport model
since dn/dx_F is
input

plot ΔX_{\max}
versus ε



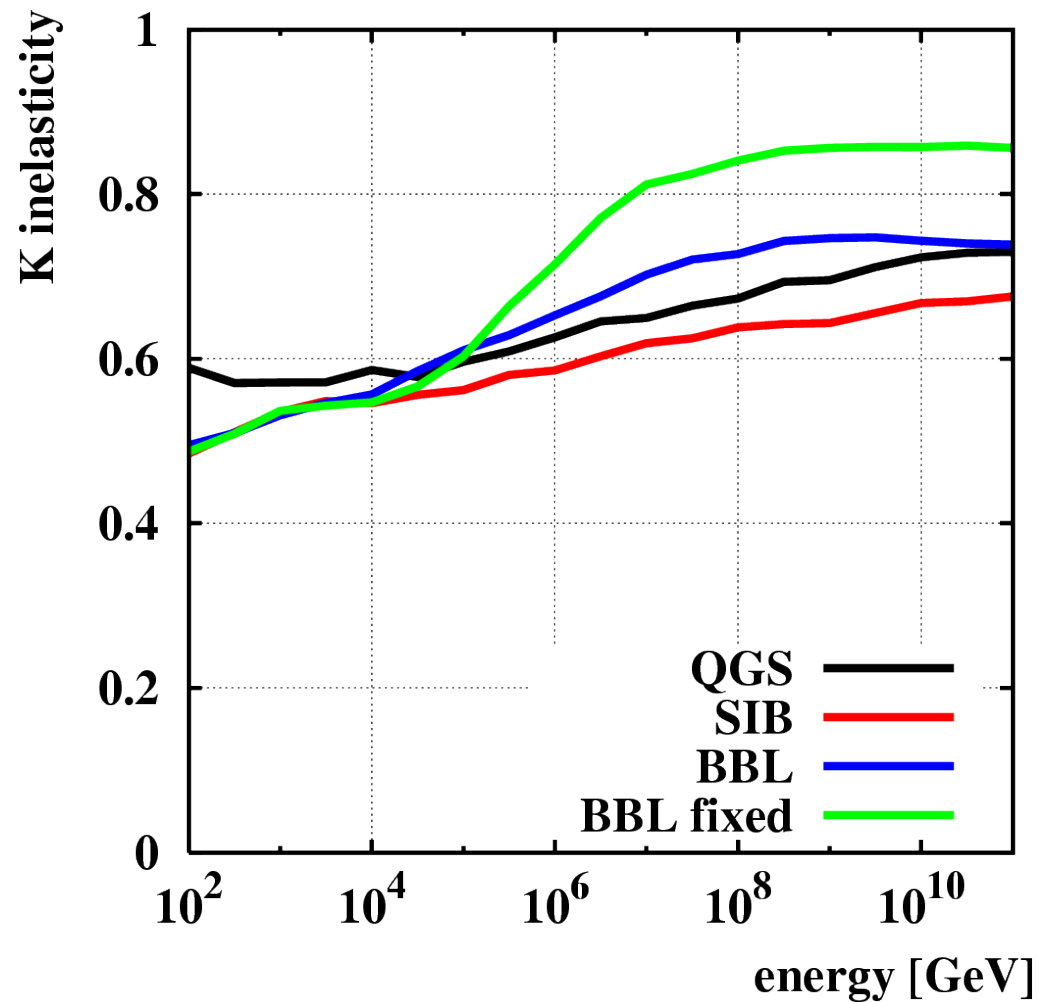
Event shape

Multiplicity



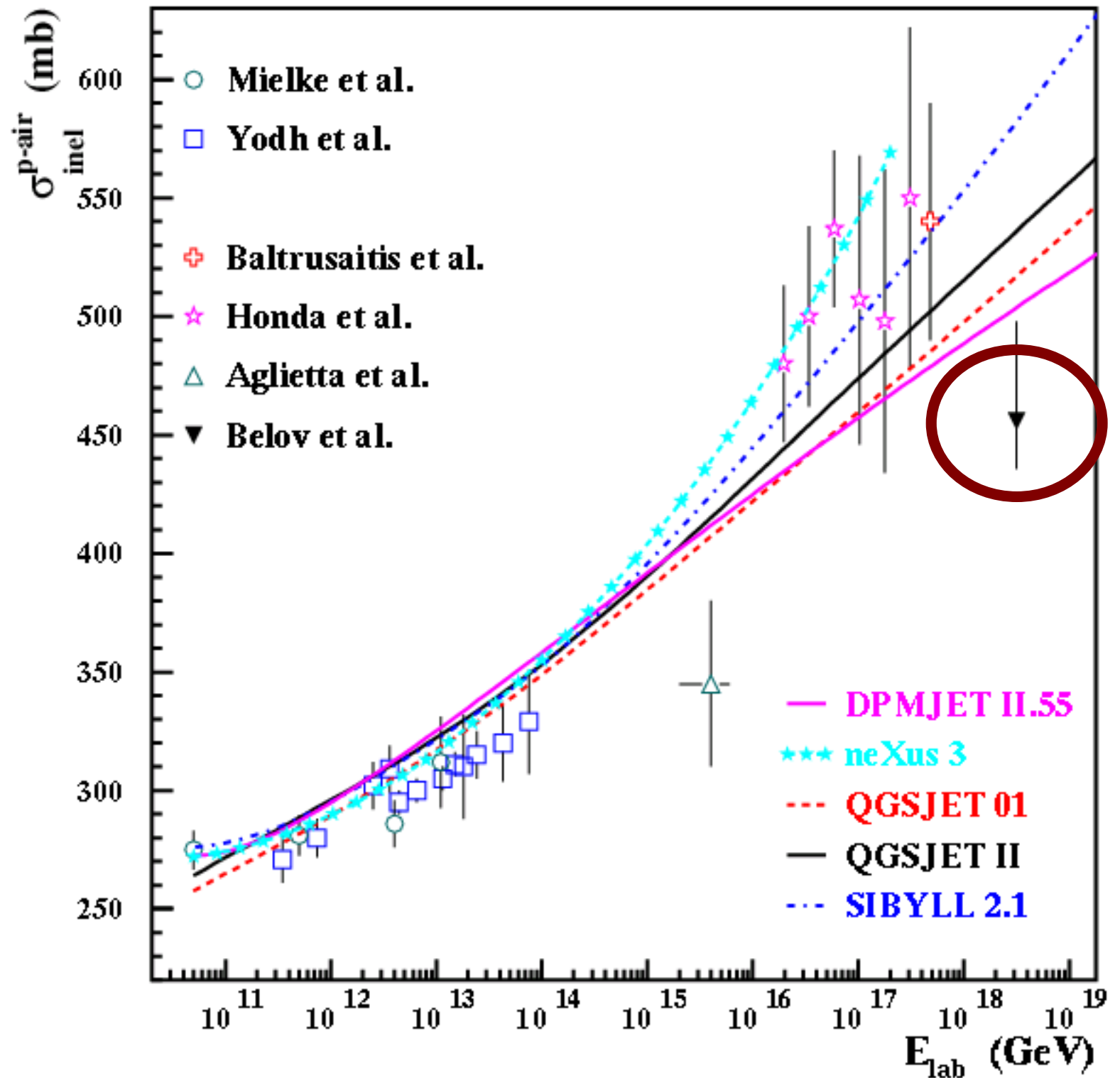
Inelasticity

$K=1 - \langle x_F \text{ of fastest particle} \rangle$



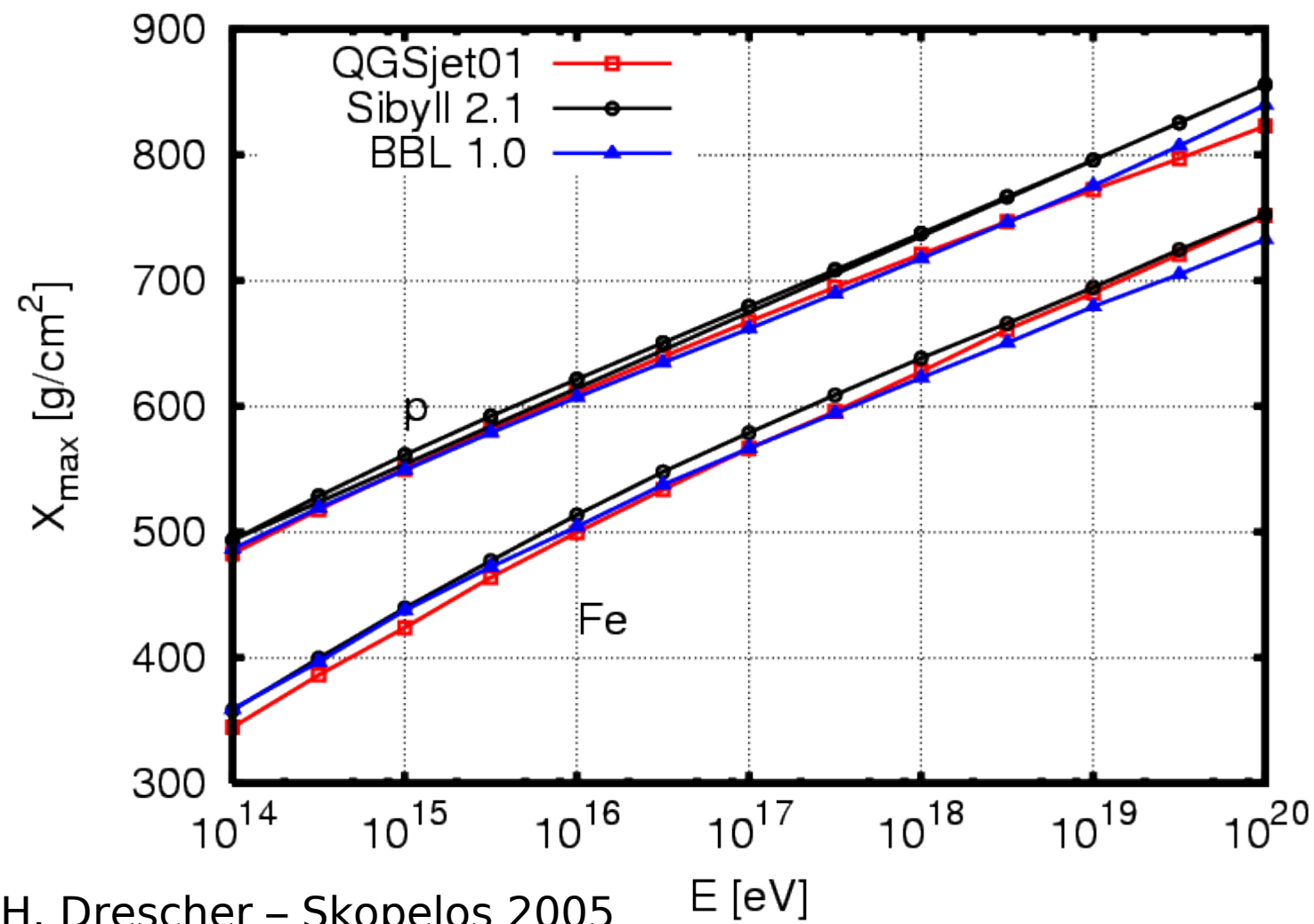
Inelastic p-Air cross section of commonly used models

Latest point from Hires (Belov) quite low



X_{\max} for proton and Iron primaries

(Iron primary per superposition)



QGSjet and BBL very similar up to 10^{19} eV for proton primary

for Iron has BBL lowest X_{\max} at highest energies

X_{\max} dominated by σ_{inel} and inelasticity K

How much of X_{\max} is due to forward scattering
how much due to inelastic cross-section ???

X_{\max} computation for the same
inelastic cross-section:

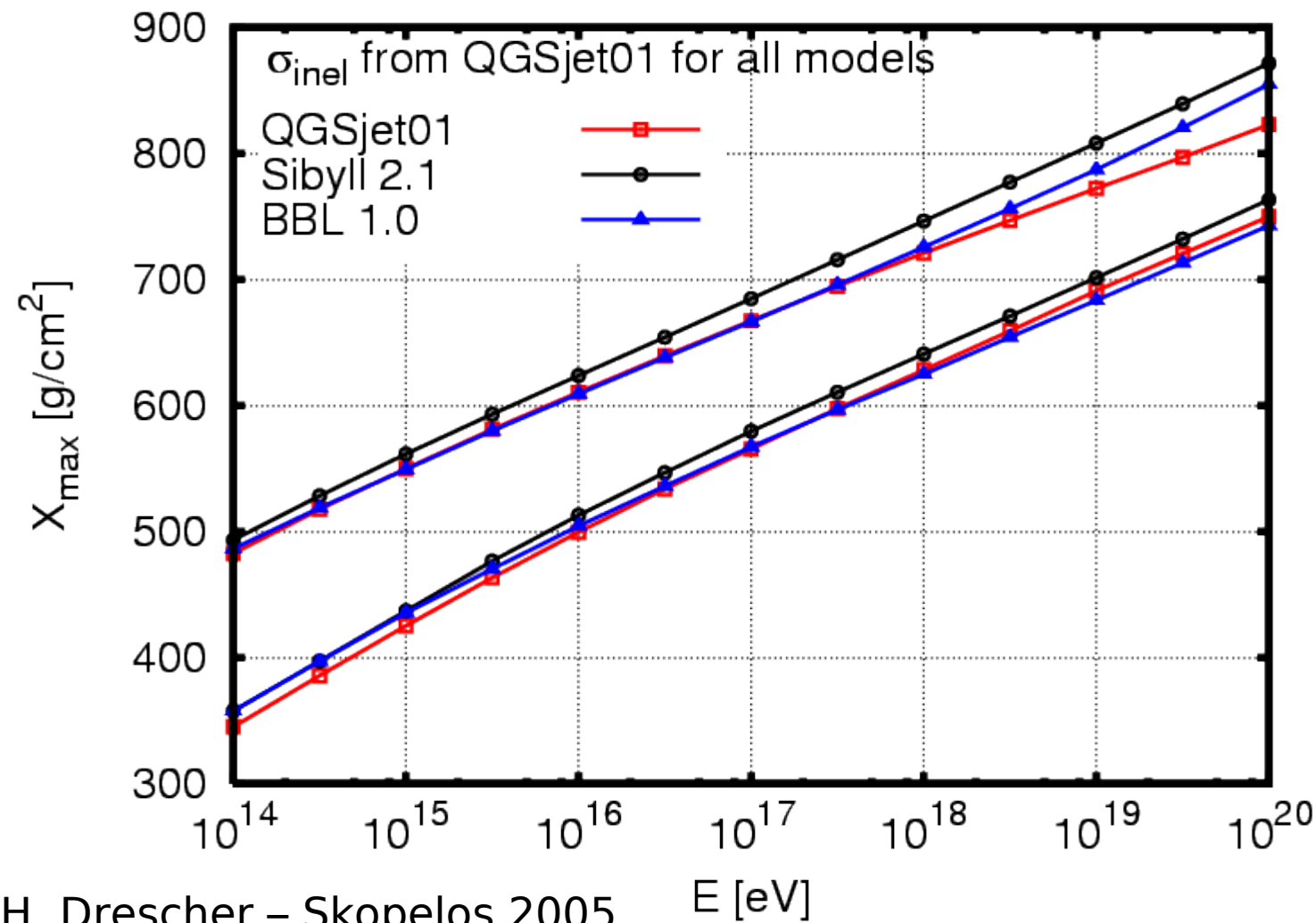
Can be done quickly with hybrid method:
fast computation of a mean shower

QGSjet σ_{inel} is lower
than that of Sibyll

Shifts SIB and BBL
 X_{\max} by 10-15 g/cm²

Differences due to
forward scattering
are even higher:
ca. 50 g/cm² at
highest energies

Models differ more
than they appear



Inelasticity to quantify forward scattering

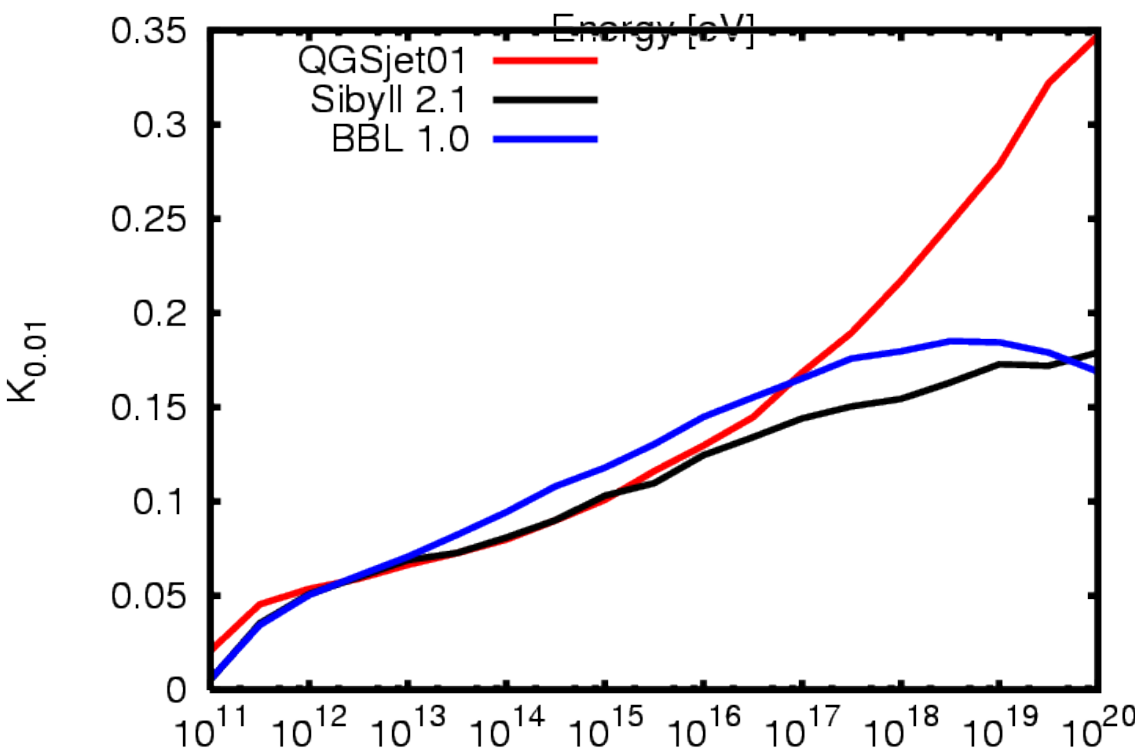
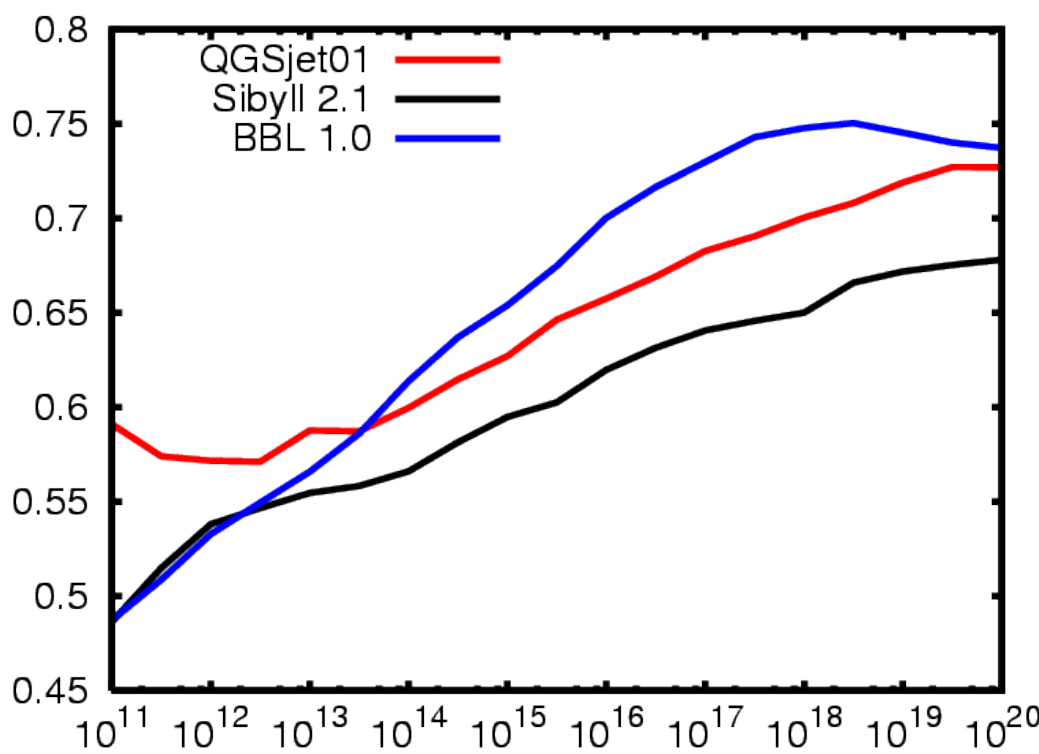
$K = 1 - \langle x_F \rangle$ of fastest particle

BBL has higher Xmax but higher K than QGSjet

$$K_{0.01} = 1 - \int_{0.01}^1 x_F \frac{dn}{dx_F} dx_F$$

QGSjet has flatter x_F distribution

need more than K to describe forward scattering



Hadronic models and experimental
cosmic ray air shower data

Muon bundles

Kascade results

Lateral Distributions (in Auger)

essential: results for muons

Muon Bundles from CosmoLEP

LEP experiments measure cosmic ray air-showers
underground

only high energy muons reach the detector

L3: -30 m 15 GeV cut-off

DELPHI: -100 m 50 GeV

ALEPH: -125 m 70 GeV cut-off

Measure multiplicity of muons in detectors

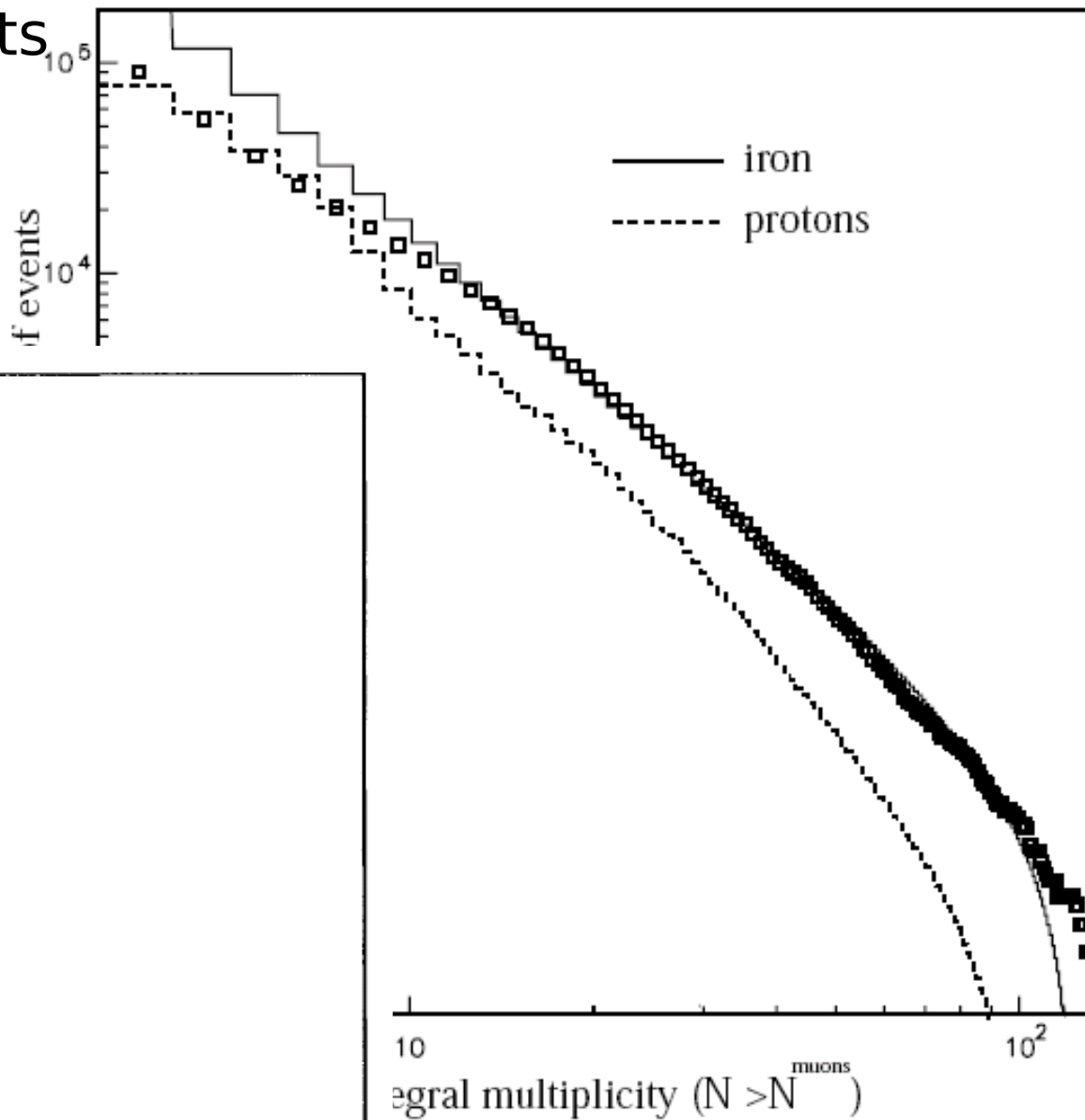
strategy: forward folding

with flux of cosmic ray spectrum from other experiments

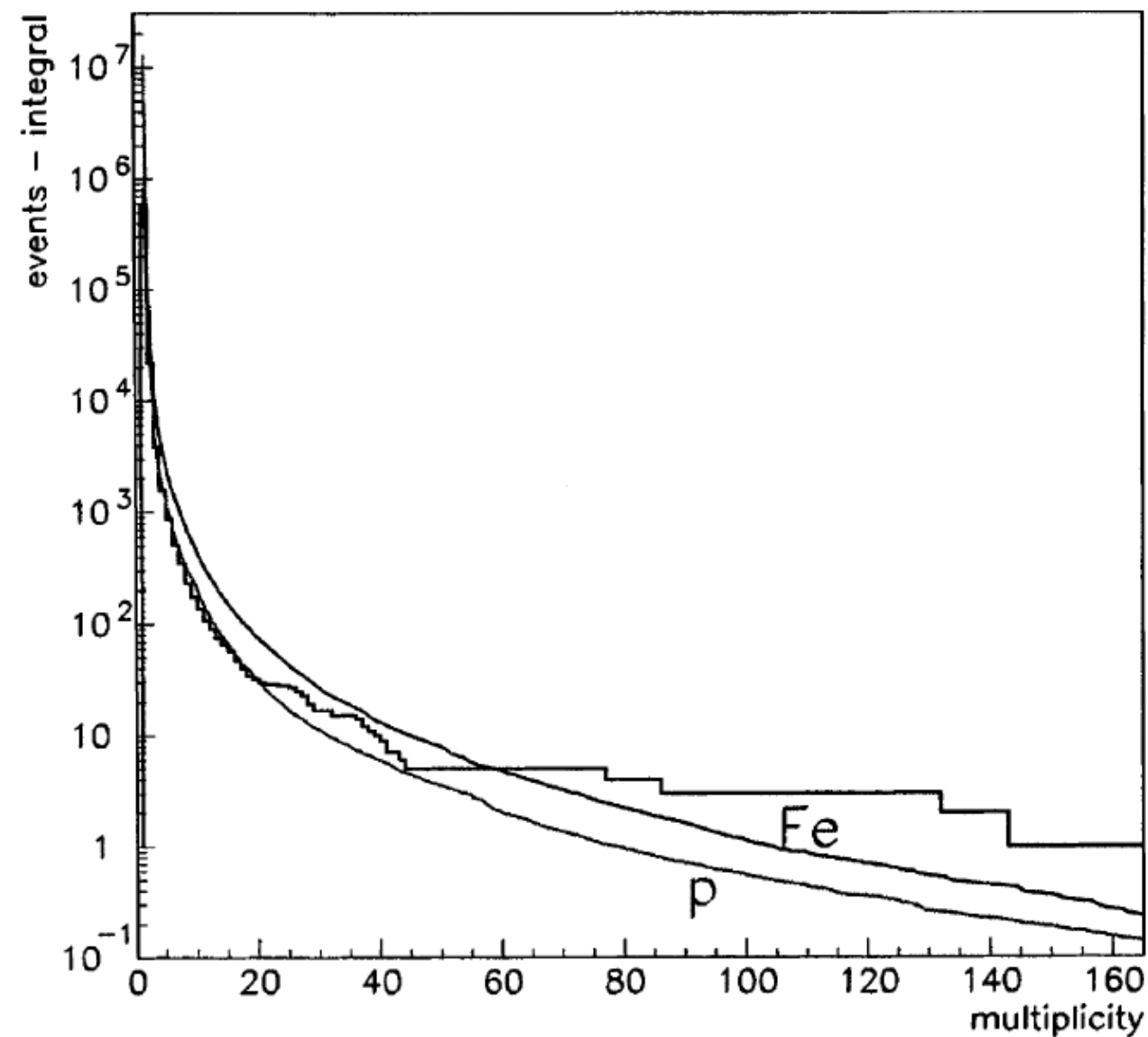
simulate detector response for extreme assumptions

pure proton and pure Iron

Delphi results



Aleph results

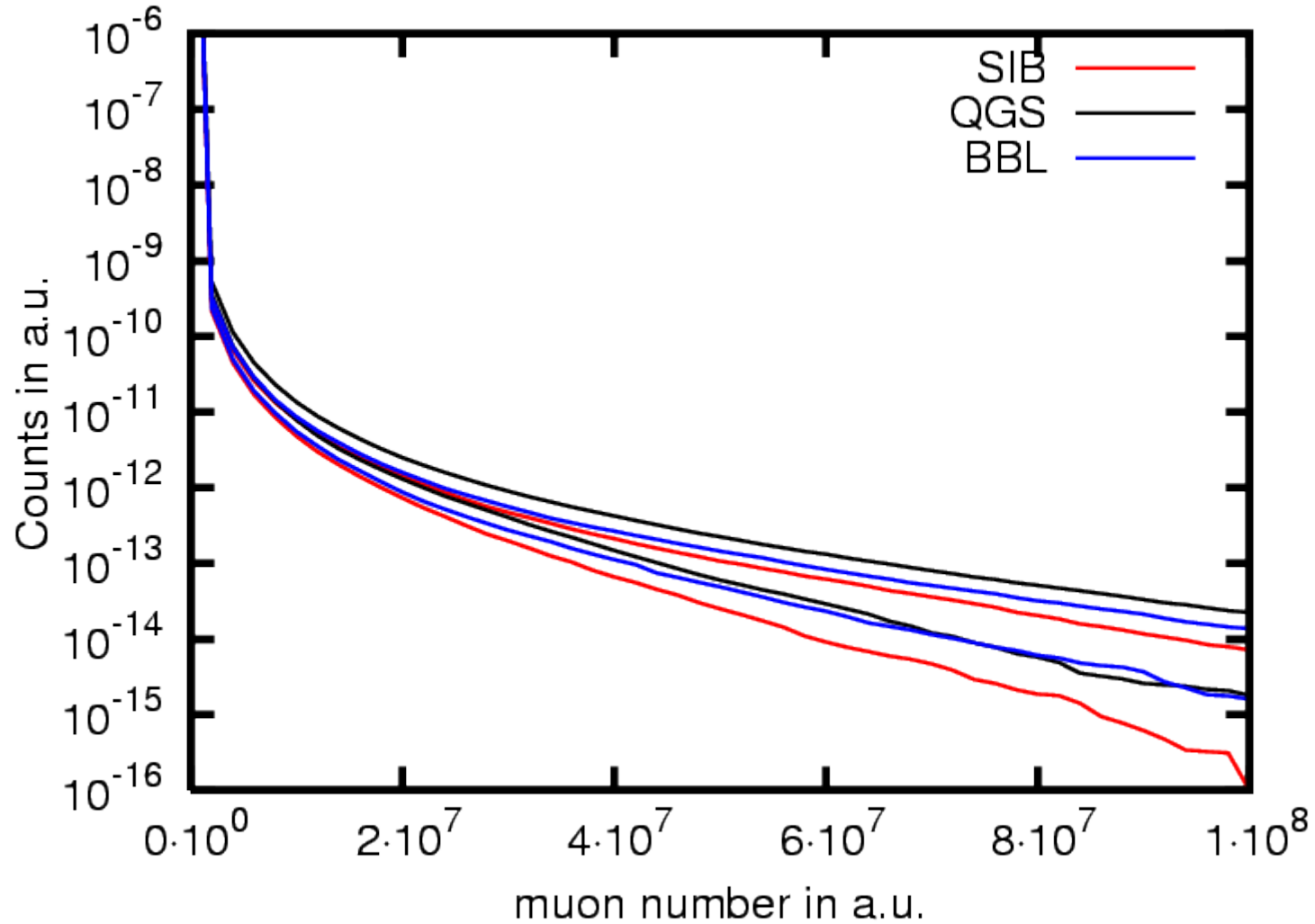


Abundance at high muon multiplicity (integrated)

Simulation results for the three models

(only vertical showers were simulated, $E_{\text{muon}} > 50 \text{ GeV}$)

Assumed
cosmic ray flux:
 $F \sim E^{-3}$



Upper line for Iron
lower line for proton

Kascade

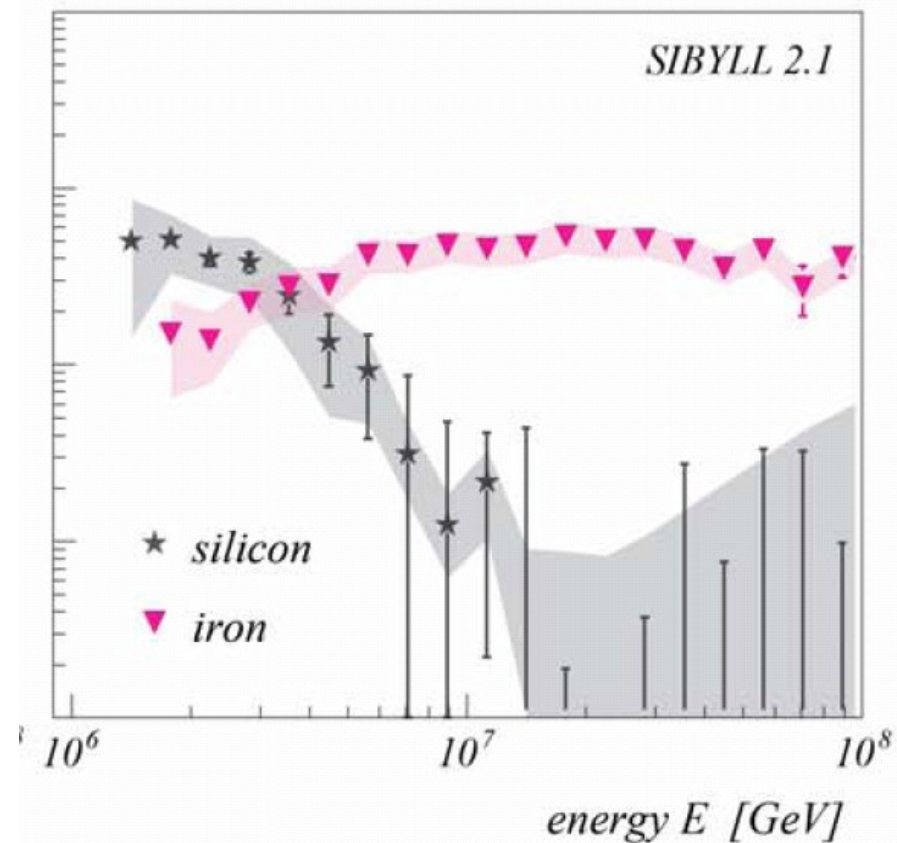
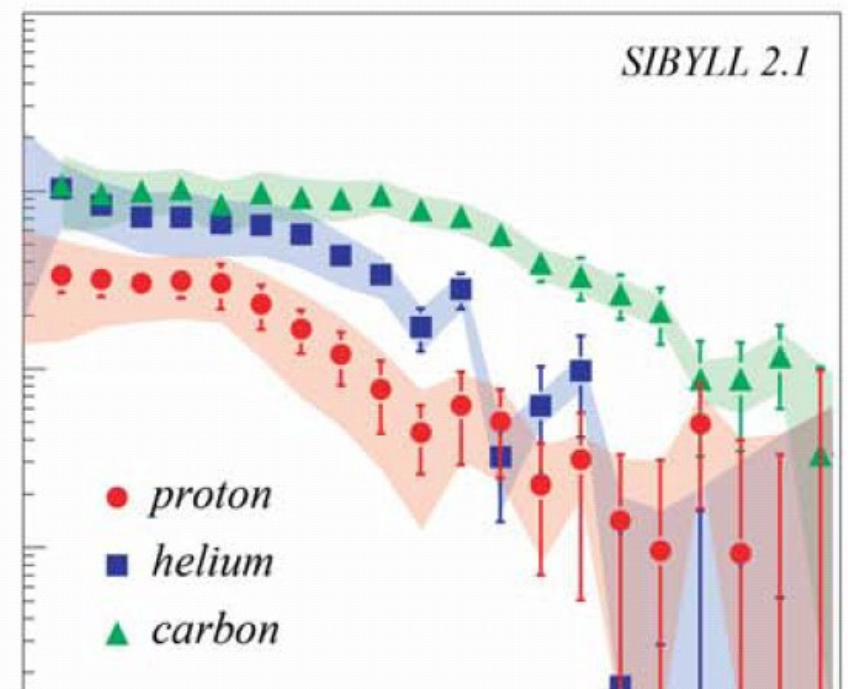
get composition around the knee (10^{15} - 10^{16} eV) by measuring electrons and muons with a ground array (Karlsruhe)

To deconvolve composition, they fit model results for different primaries to experimental data

here: goodness of fit on $N_{el} - N_{muon}$ plane

From Holger Ulrich (Kascade)

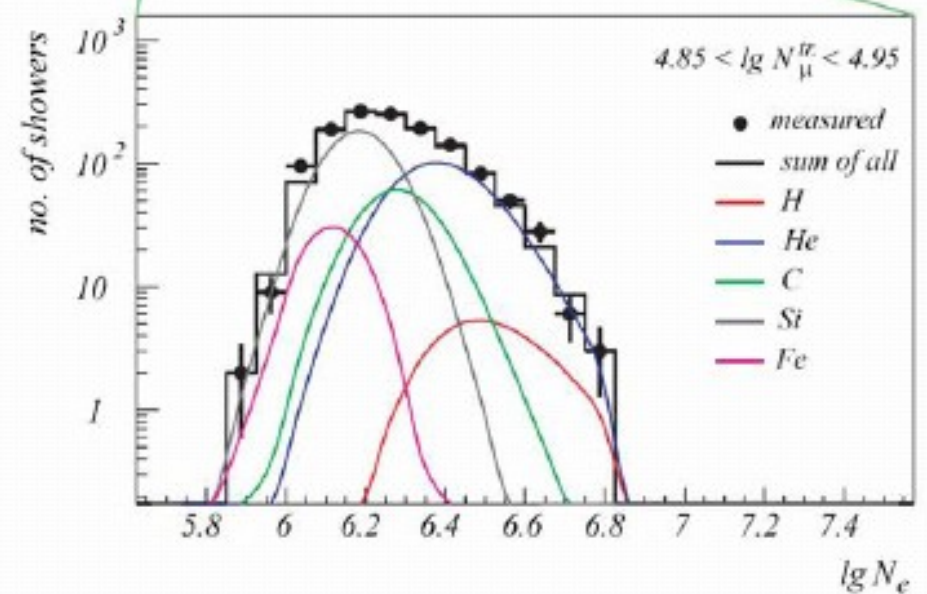
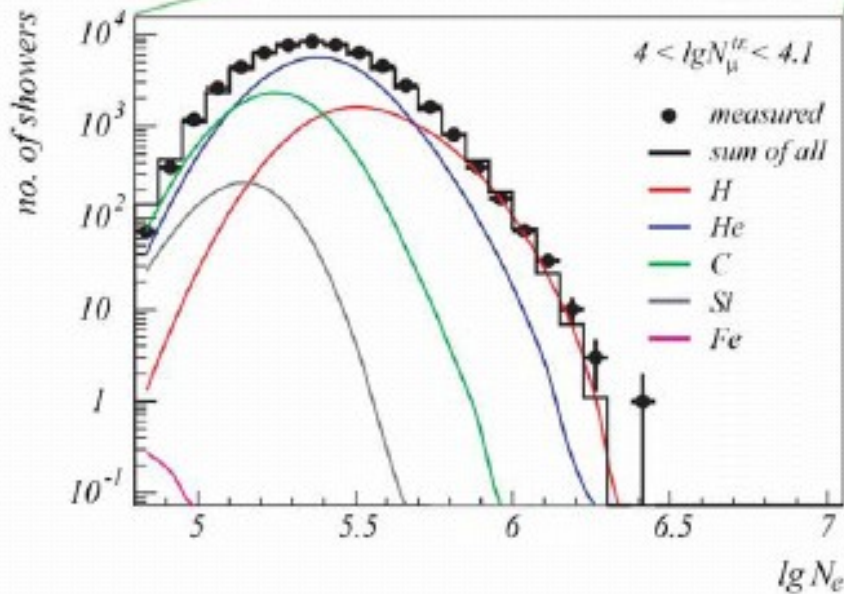
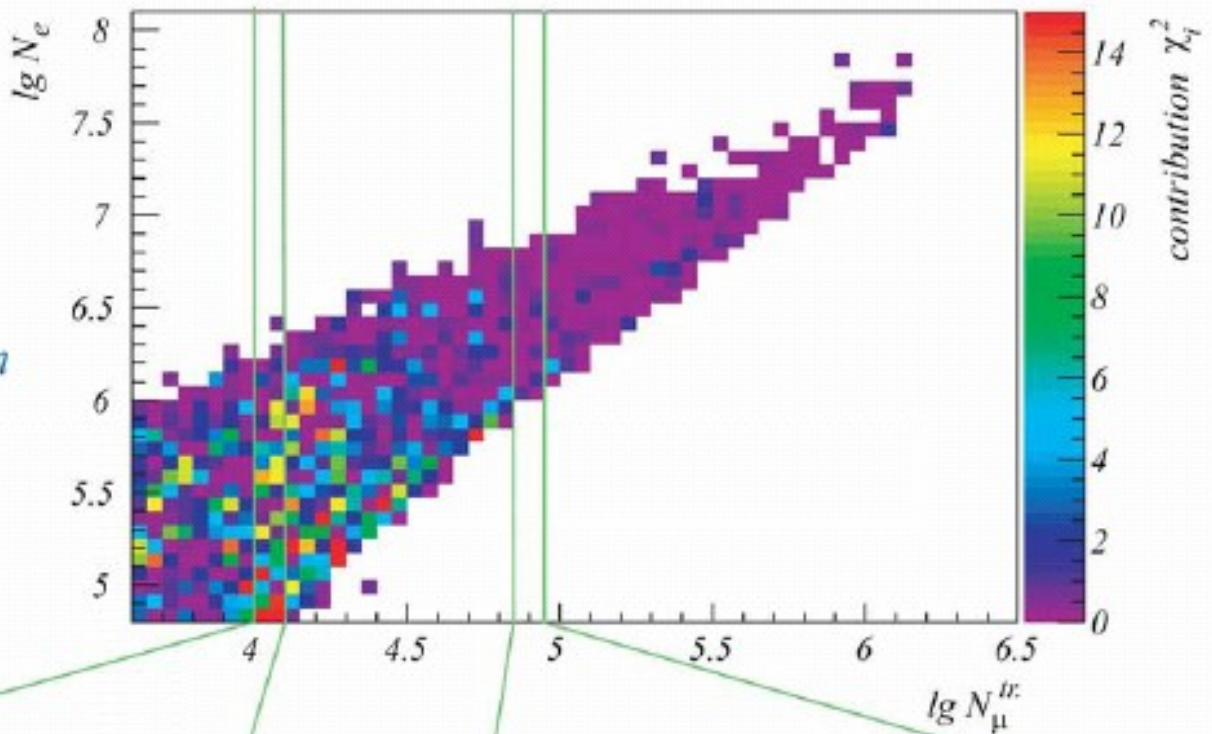
H. Drescher – Skopelos 2005



QGSJet 01 - result
Description of data

*forward folding of solution with
 calculated probabilities, calculation
 of how the data would look like*

*comparison between calculated
 and measured data: χ^2*

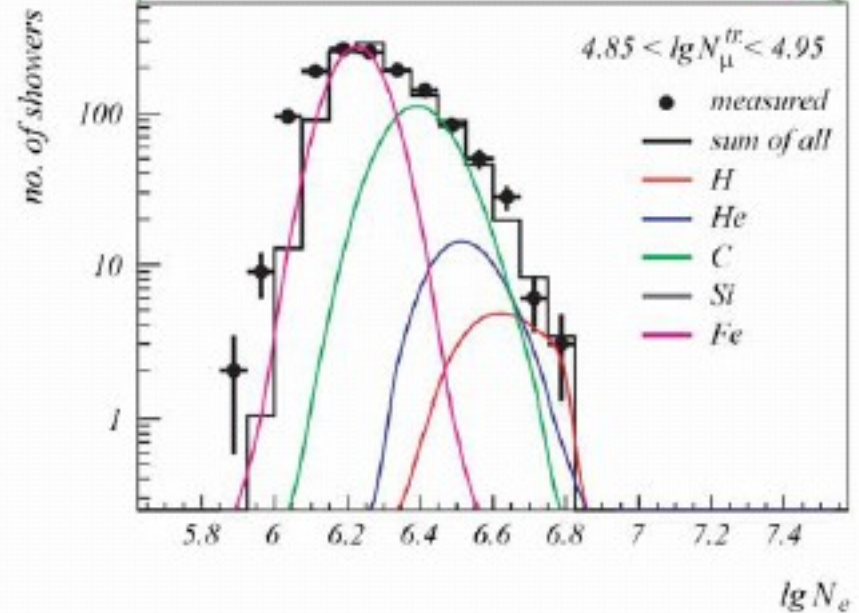
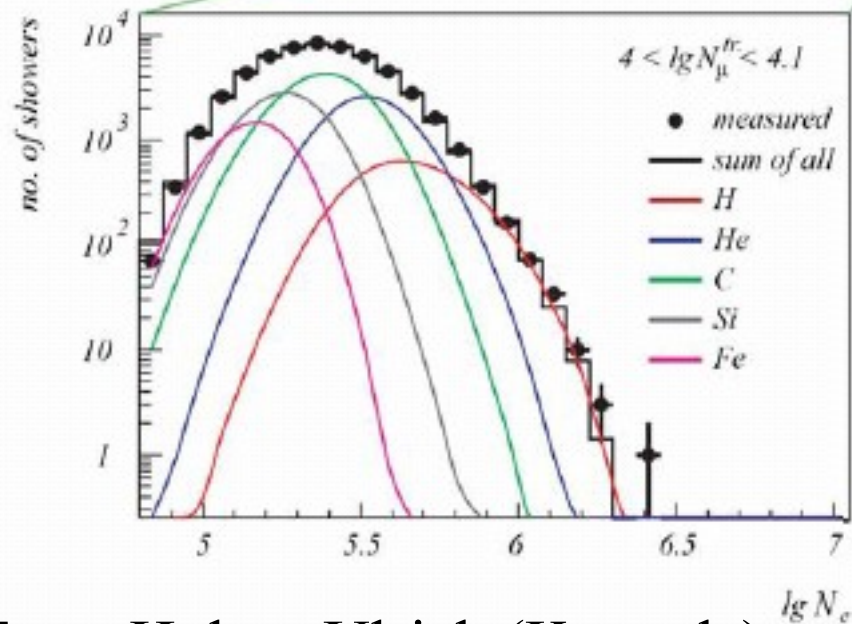
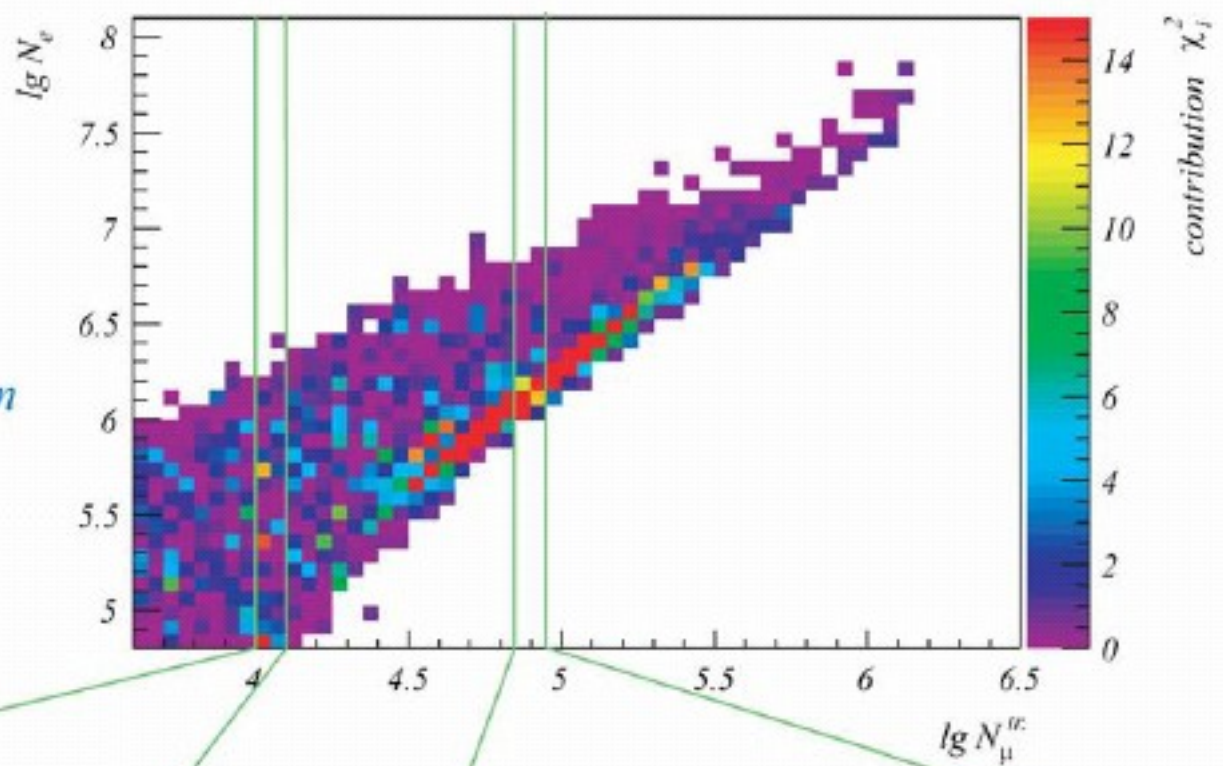


QGSJET01: electron poor

SIBYLL 2.1 - result
Description of data

*forward folding of solution with
 calculated probabilities, calculation
 of how the data would look like*

*comparison between calculated
 and measured data: χ^2*

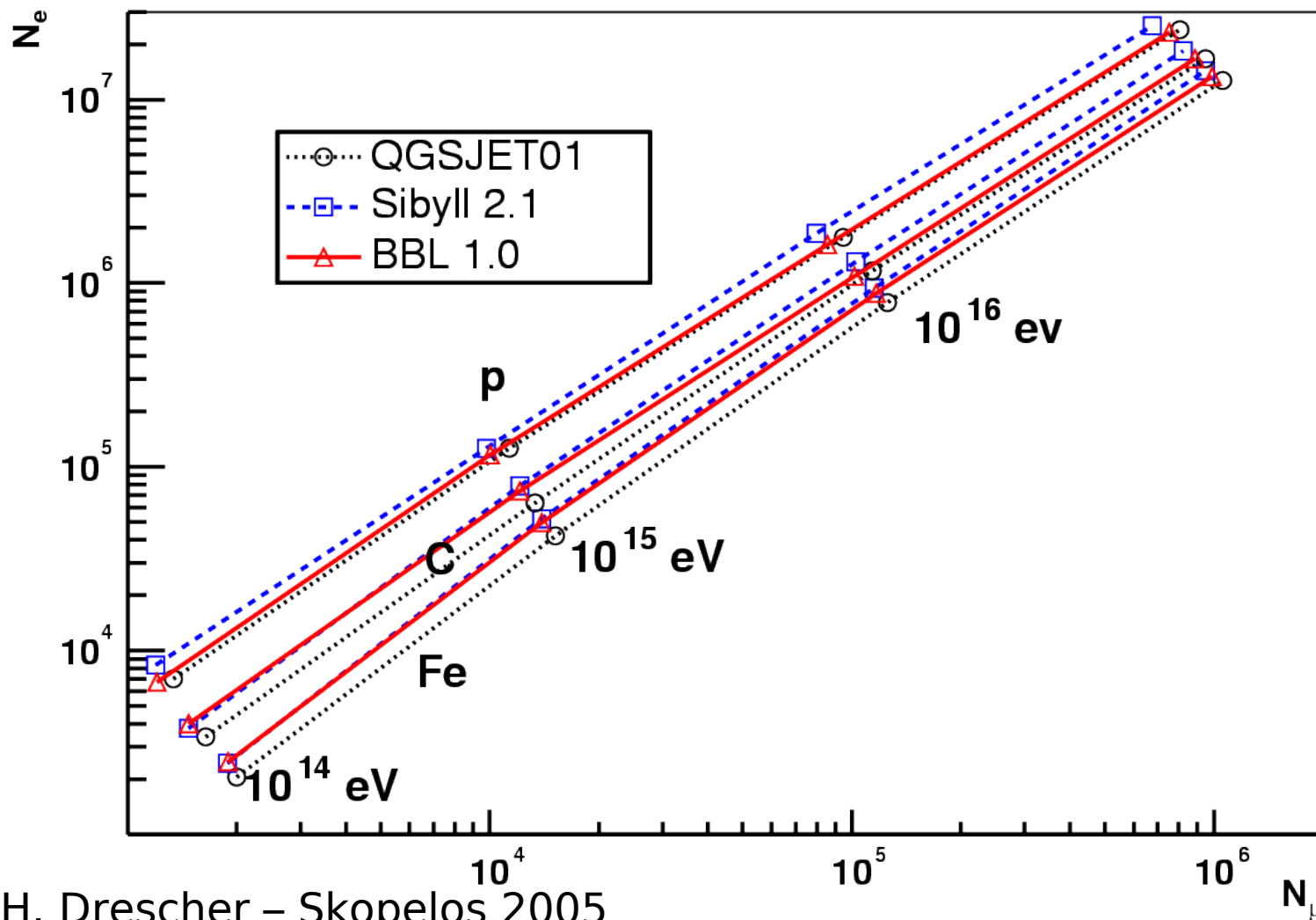


From Holger Ulrich (Kascade)

Too electron rich/muon poor

Electron and Muon numbers

BBL reproduces Sibyll at low energies (by construction) and is similar to QGSjet at higher energies.
--> same as a function of primary mass A.

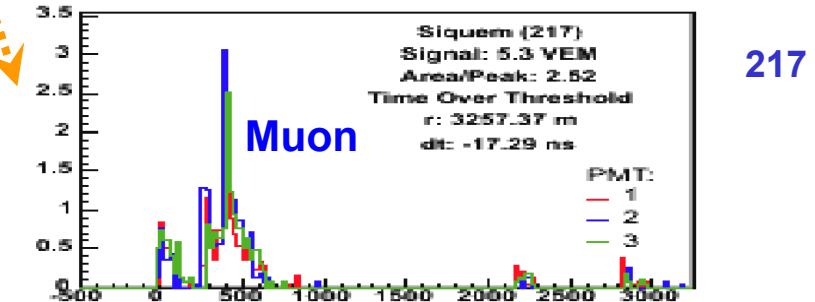
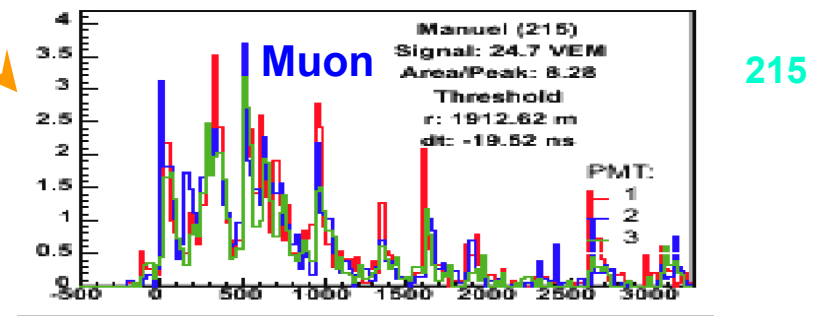
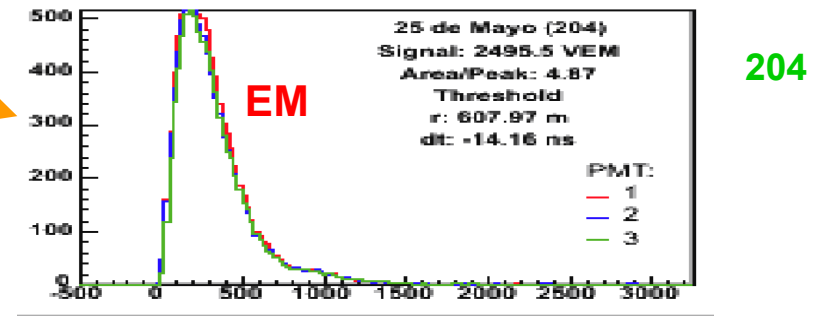
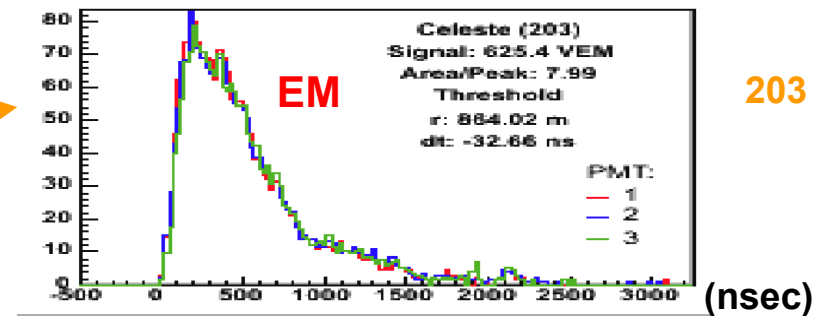
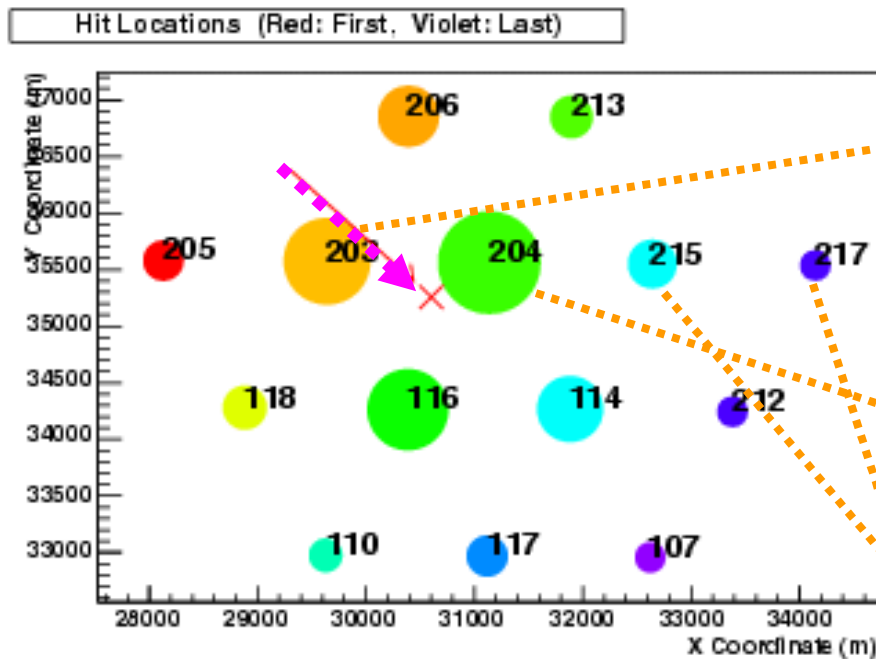


Maybe helps to resolve known problems of models at KASCADE ;-)

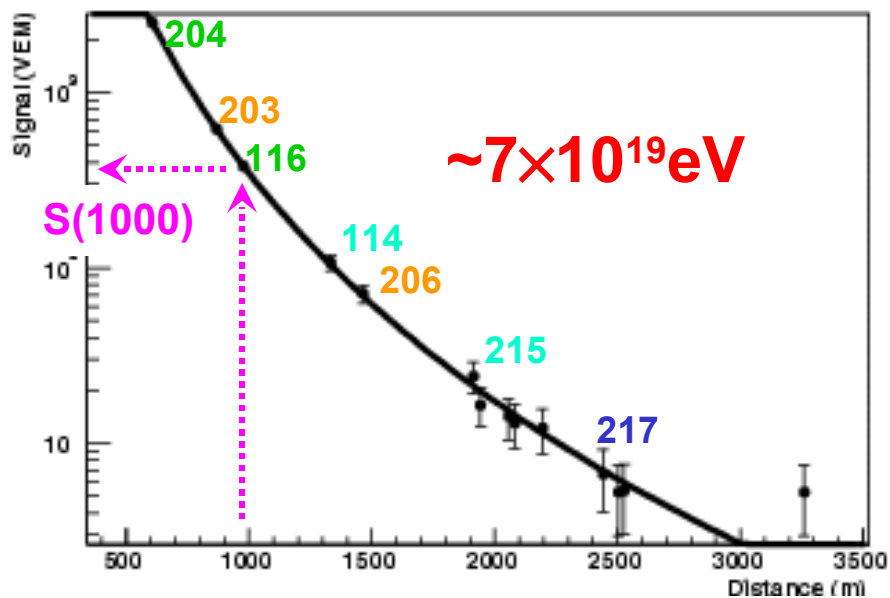
but: full implementation of N-N scattering necessary

A Typical Auger Event (Zenith~35°)

From
Katsushi Arisaka

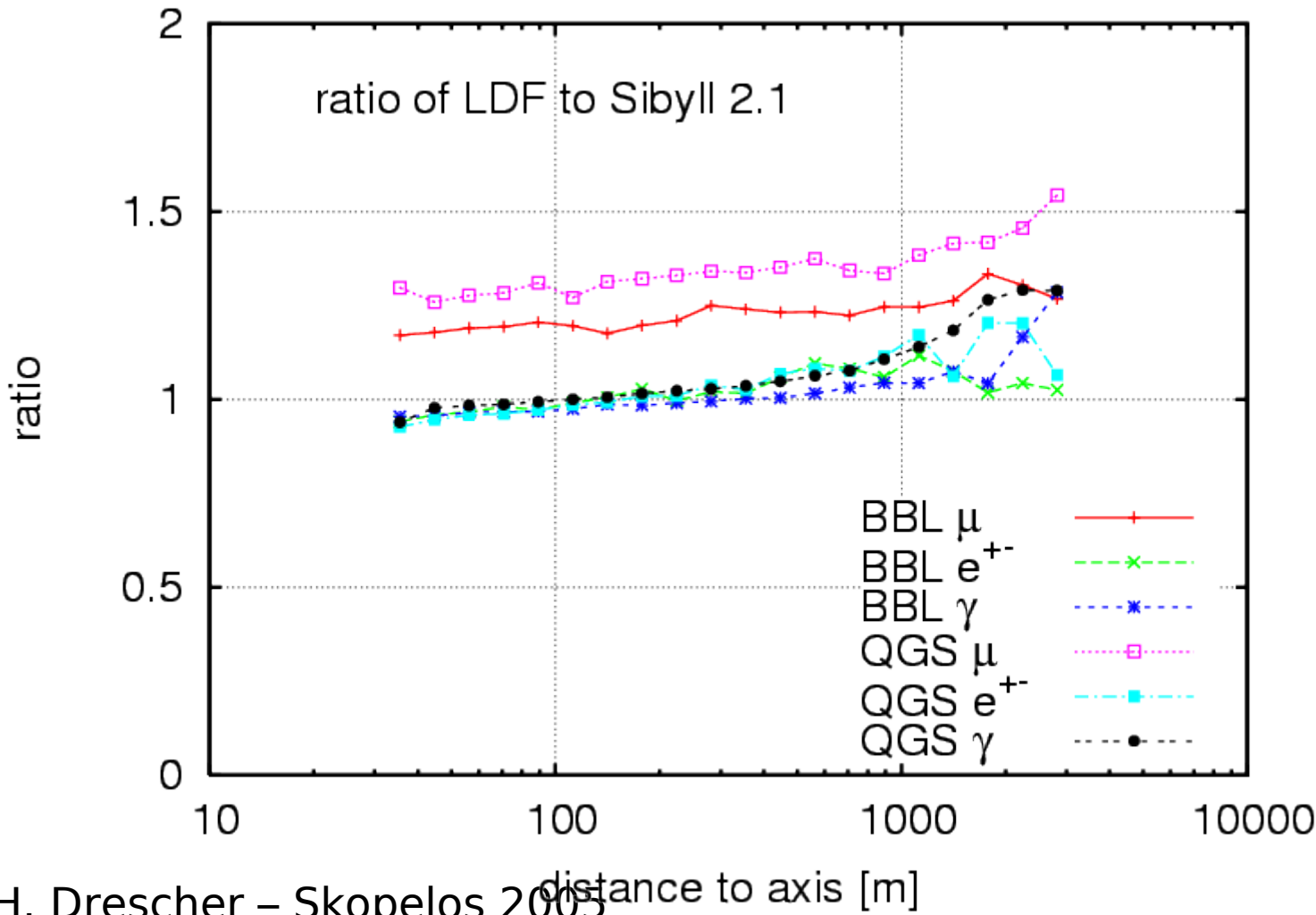


Lateral Distribution Function



Lateral Distribution Functions

Forward suppression in BBL leads to enhanced particle production at small x --> more muons



Slope for em particles quite insensitive to high energy hadronic part

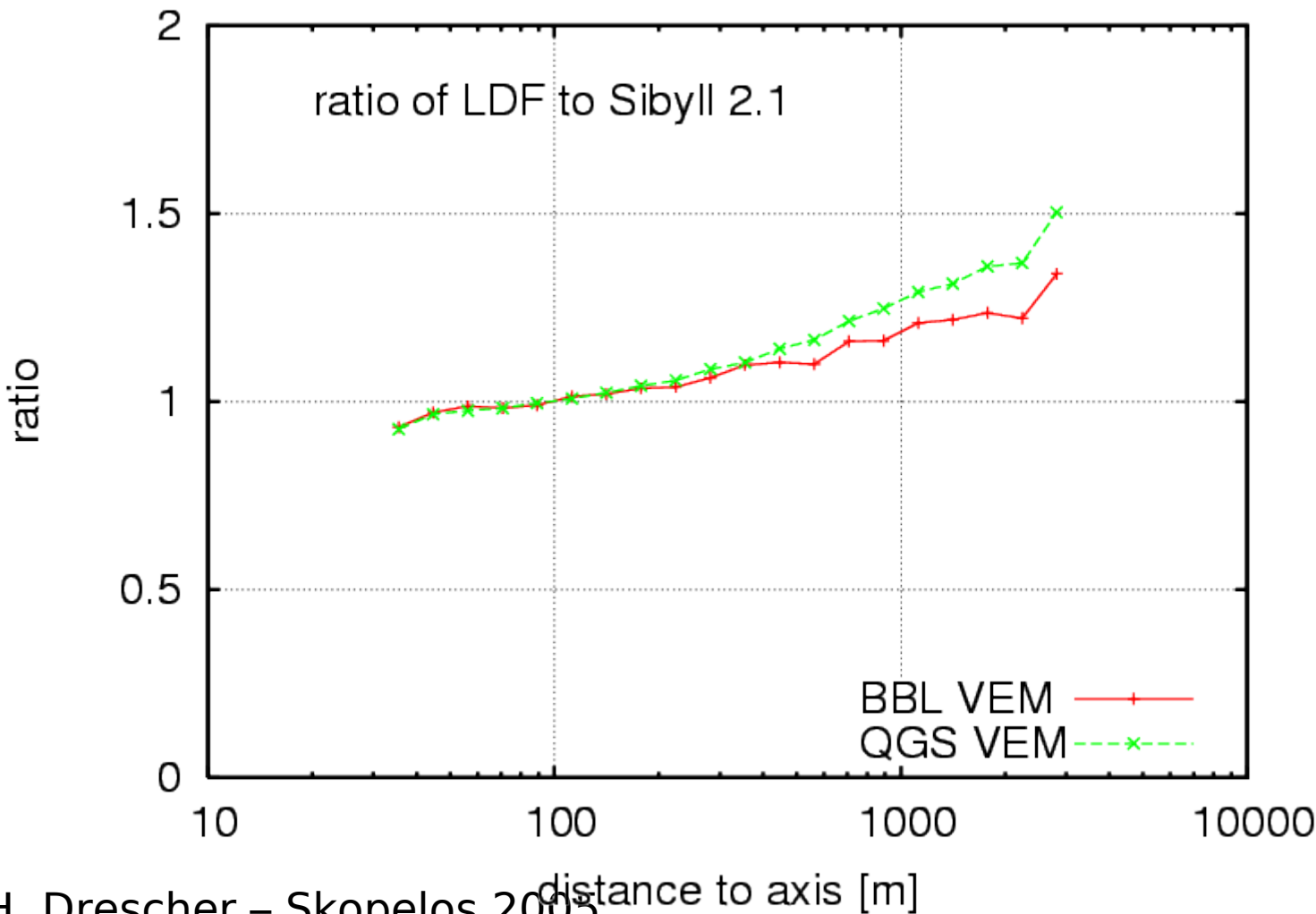
Vertical proton 10^{19} eV showers, Auger-altitude

LDF in VEM units

flatter LDF since γ dominate at low r, μ at large r

~ 25%
difference at
1000m

Energy calibration
of Auger with MC
is still 25 higher
than fluorescence!



Summary

- fast air shower simulations with hybrid approach
- Forward scattering suppression with BDL
- Air Shower data and hadronic interaction models

try to use hadronic models in air shower simulations

- for the interpretation of data
- to learn about interactions

(if possible both)