

**Space charge: PART I**

**Analytical solution of the Laplace equ. for  
a coaxial cavity**

**Space charge: PART II**

**Space charge scenarios for the ALICE TPC**

**Space charge: PART III**

**Space point distortions due to space charges  
and more ...**

# **Space charge: PART I**

**Analytical solution  
of the Laplace equ.  
for a coaxial cavity**

**Stefan Rossegger**

# What is needed ...

... to calculate the field distortions due to space charges?

- ➔ A solution of the **3D-Laplace equation** for a given charge accumulation within the TPC geometry
  - **Numerical solution?** ...
    - ... high accuracy takes time
  - **Analytical approach** by means of Green's functions
    - ... that's a nice idea ... ;-)

# Theoretical base ...

- If a fast converging Green's function is found, one just multiplies it with the charge distribution and integrates over its volume ...

$$\Phi(r, \phi, z) = \frac{1}{\epsilon_0} \int r' dr' \int d\phi' \int dz' \rho_c(r', \phi', z') \cdot G(r, \phi, z, r', \phi', z')$$

$$E(r, \phi, z) = -\nabla\Phi(r, \phi, z)$$

... which results in the potential and the electric field components

# Which Green's function?

A solution is given in [1]

[1] W. R. Smythe, **Static and dynamic electricity**, New York-Toronto-London: McGraw-Hill 1950.

$$G(r, \phi, z; r', \phi', z') = \frac{1}{\pi L} \sum_{m=0}^{\infty} \sum_{n=1}^{\infty} (2 - \delta_{m0}) \cos[m(\phi - \phi')] \sin(\beta_n z) \sin(\beta_n z') \frac{R_{mn1}(r_{<}) R_{mn2}(r_{>})}{I_m(\beta_n a) K_m(\beta_n b) - I_m(\beta_n b) K_m(\beta_n a)}.$$

$$R_{mn1}(r) = K_m(\beta_n a) I_m(\beta_n r) - I_m(\beta_n a) K_m(\beta_n r).$$

**Good news:** non-oscillating solution in r-direction, we get  $\mathbf{E}_r$

**Bad news:**  $\mathbf{E}_z$ ,  $\mathbf{E}_\phi$  show slow convergence, or are even divergent (in r'-plane)!

CLEARLY NOT SUFFICIENT for a 3D SPACE-CHARGE-DISTR. !!!

# ... more Green's functions ?

After a few months of fighting with merciless Bessel functions, two more solution were derived [CERN-OPEN-2009-003]

(with different regions of slow convergence ...)

(1)

$$G(r, \phi, z; r', \phi', z') = \frac{1}{2\pi} \sum_{m=0}^{\infty} \sum_{n=1}^{\infty} (2 - \delta_{m0}) \cos[m(\phi - \phi')] \frac{R_{mn}(r)R_{mn}(r')}{N_{mn}^2} \frac{\sinh(\beta_{mn}z_{<}) \sinh(\beta_{mn}(L - z_{>}))}{\beta_{mn} \sinh(\beta_{mn}L)},$$

$$R_{mn}(r) = Y_m(\beta_{mn}a)J_m(\beta_{mn}r) - J_m(\beta_{mn}a)Y_m(\beta_{mn}r).$$

(2)

$$G(r, \phi, z; r', \phi', z') = \frac{1}{\pi L} \sum_{m=0}^{\infty} \sum_{n=1}^{\infty} (2 - \delta_{m0}) \cos[m(\phi - \phi')] \sin(\beta_n z) \sin(\beta_n z') \frac{R_{mn1}(r_{<})R_{mn2}(r_{>})}{I_m(\beta_n a)K_m(\beta_n b) - I_m(\beta_n b)K_m(\beta_n a)}.$$

$$R_{mn1}(r) = K_m(\beta_n a)I_m(\beta_n r) - I_m(\beta_n a)K_m(\beta_n r).$$

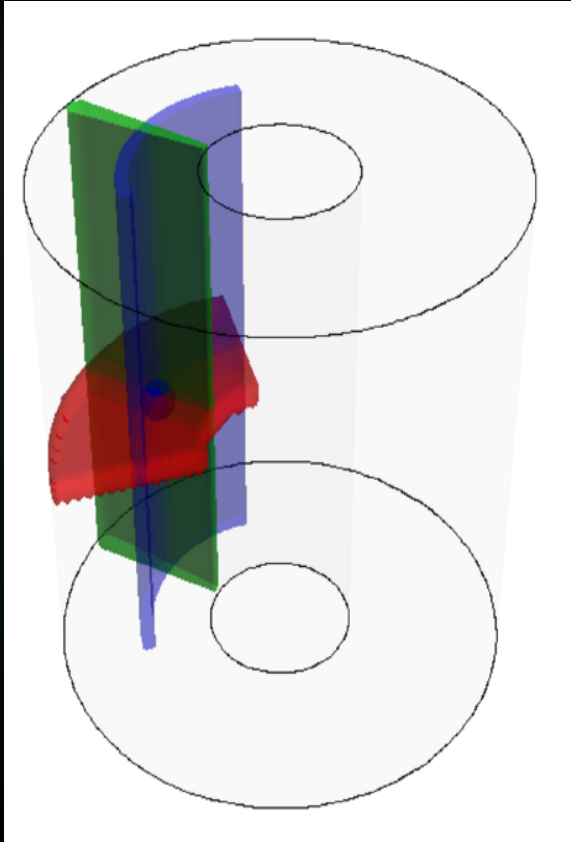
(3)

$$G(r, \phi, z; r', \phi', z') = \frac{1}{L} \sum_{k=1}^{\infty} \sum_{n=1}^{\infty} \frac{\cosh[\mu_{nk}(\pi - |\phi - \phi'|)]}{\mu_{nk} \sinh(\pi \mu_{nk})} \sin(\beta_n z) \sin(\beta_n z') \frac{R_{nk}(r)R_{nk}(r')}{N_{nk}^2}.$$

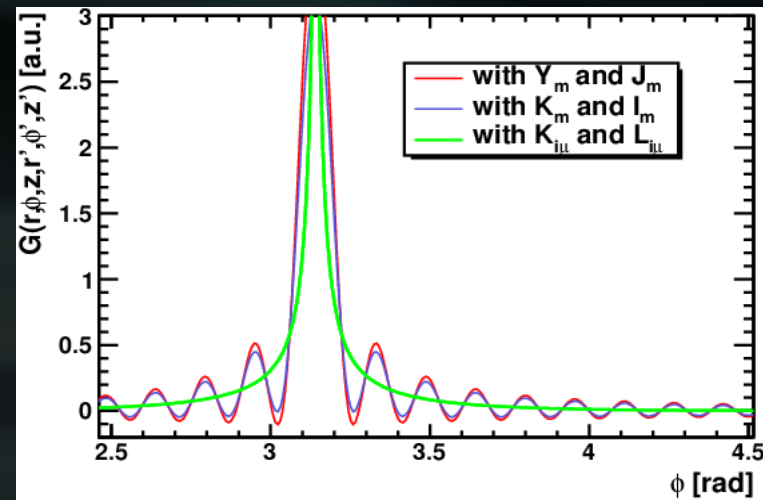
$$R_n(r) = L_{i\mu}(\beta_n a)K_{i\mu}(\beta_n r) - K_{i\mu}(\beta_n a)L_{i\mu}(\beta_n r).$$

# Convergence properties ...

Regions of slow convergence



- (1) allows fast converging  $E_z$
- (2) allows fast converging  $E_r$
- (3) allows fast converging  $E_\phi$



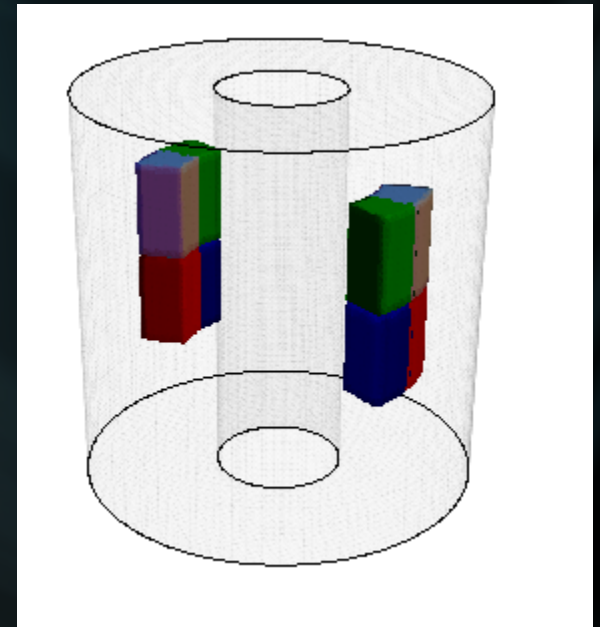
FINALLY, handling of any 3D case (not only radial symmetric ones) including correct Dirichlet boundaries in  $r$  and  $z$ -direction !!

# Evaluation ...

Pre-calculation of the fields for a const. charge within a “Piece-Of-Cake” volume (P.O.C charge) at every possible position within the volume

**THIS RESULTS IN A LOOKUP TABLE**

Approximation of the desired charge accumulation through SIMPLE SUMMATION of the (weighted) P.O.C charges ... VERY FAST!





# How fast? How accurate?

Tested for some realistic cases of:  $\rho(r,z) = (A_1 + A_2 z) / r^2$

Comparison of discretization with the analytic solutions:

- A slicing of “P.O.C” boxes of about **10 cm** size
- Summation limit for the Look-Up table generation was less than 30 terms

➡ **Accuracy** around  $10^{-7}$  for all E field components

➡ Calculation time **less than a 1 sec !**

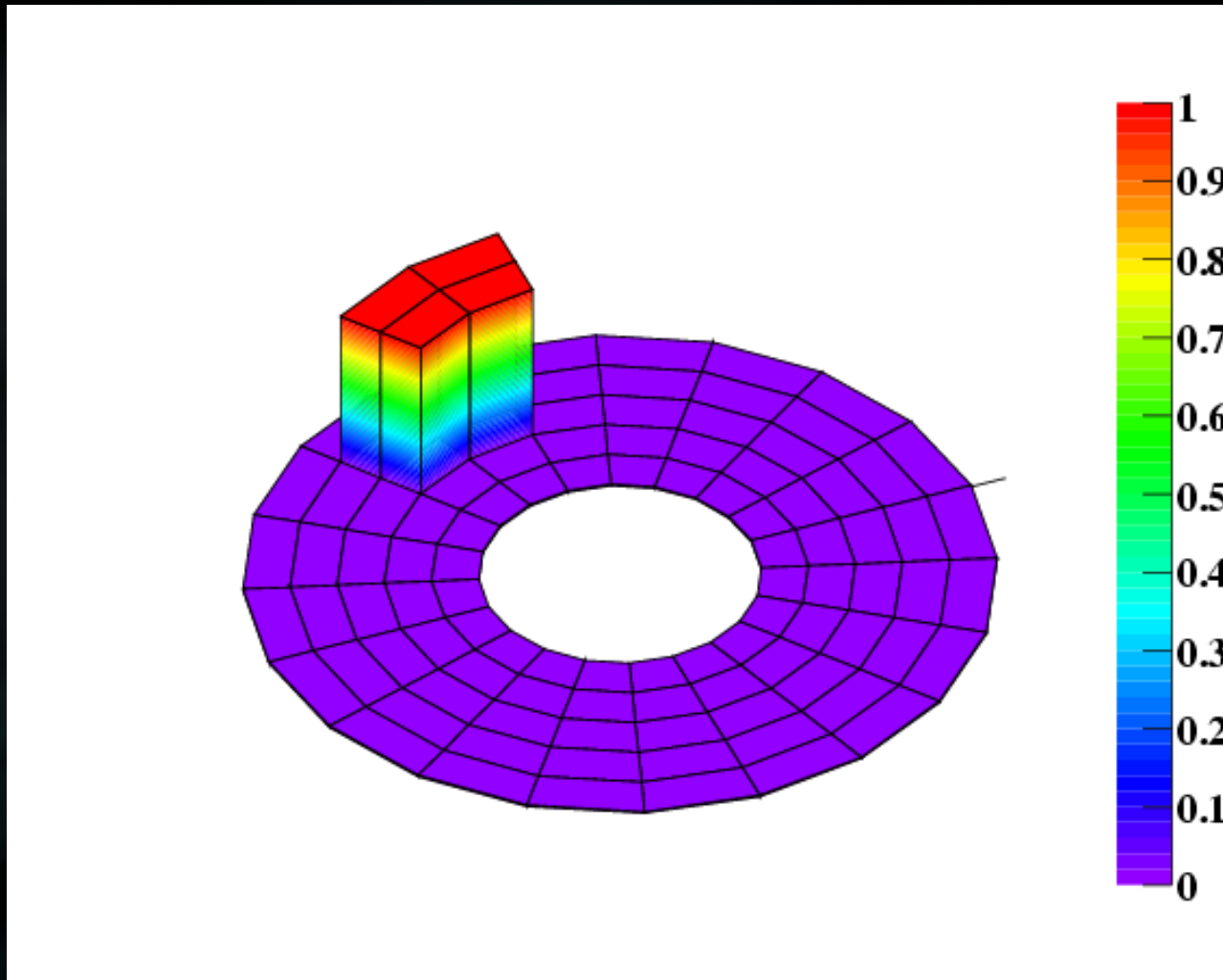
# CONCLUSION

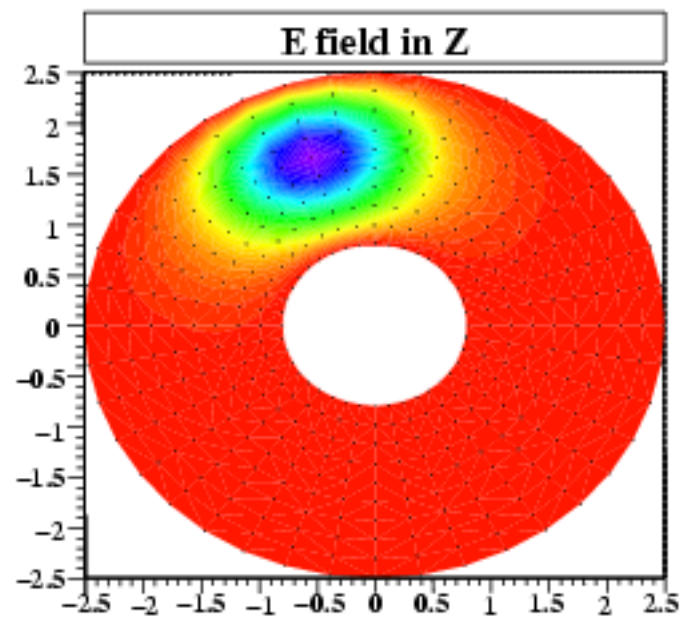
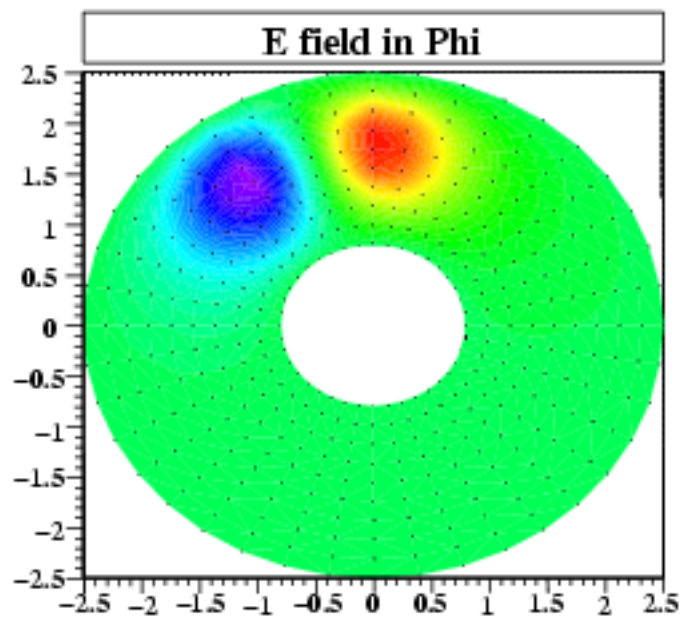
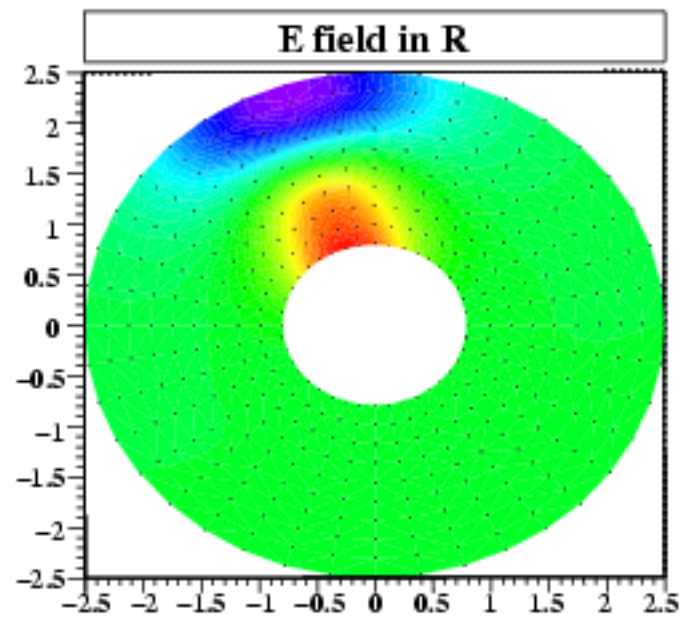
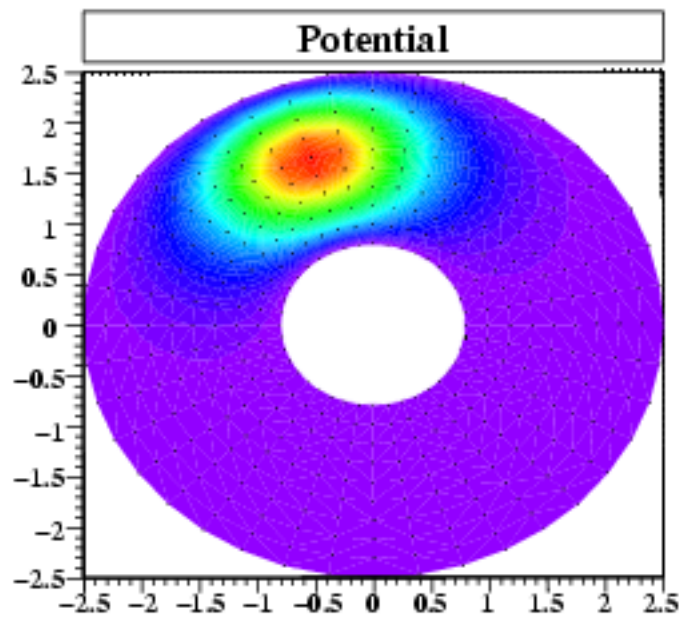
- With the presented ANALYTICAL SOLUTIONS you can calculate EVERY FIELD component for ANY DESIRED SPACE CHARGE DISTRIBUTION.
- BUT, special mathematical packages, like GSL and Alg831, are needed and the summation needs time
- With a LOOK UP TABLE (of const. 'P.O.C' charges) you can calculate a highly accurate approximation for every possible configuration within “NO TIME”

# Example 1:

SIMPLE Charge distribution

$\rho(r, \phi, z) = 1$  ... within a box

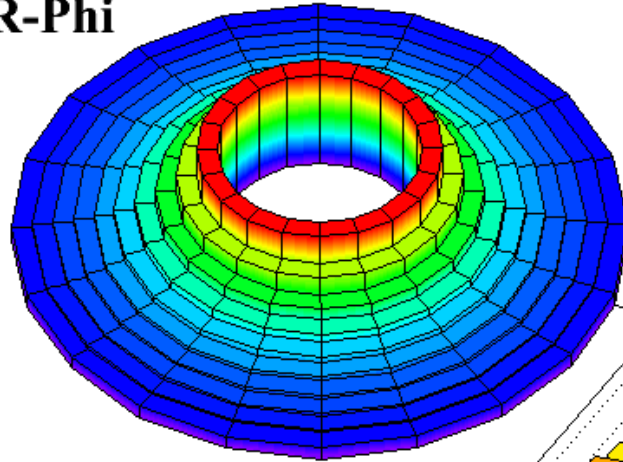




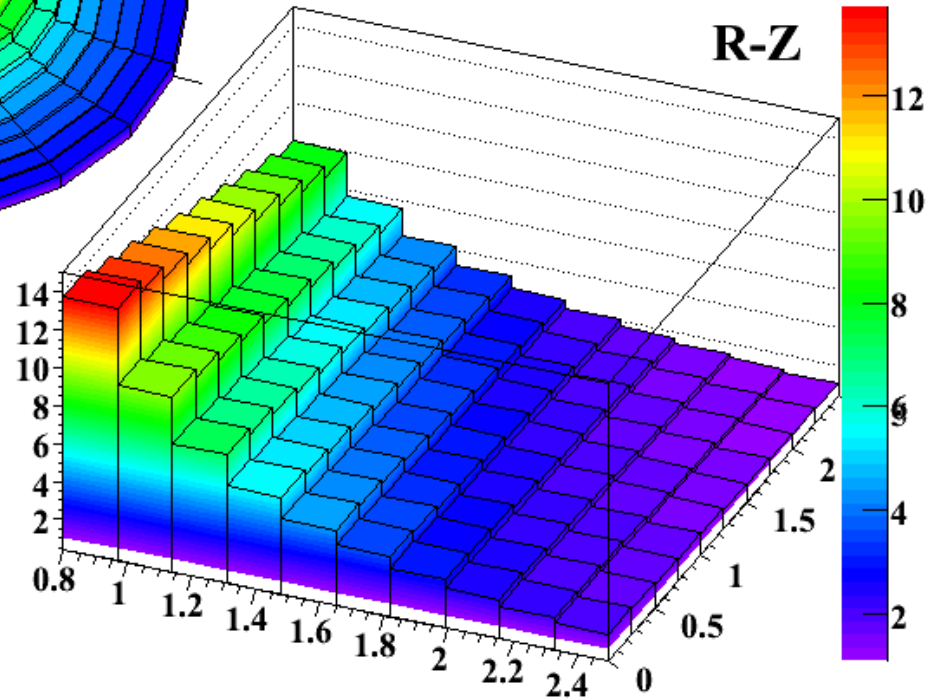
$$\rho(r,z) = 1/r^2 * k * z \dots \text{(used in STAR)}$$

# Example 2:

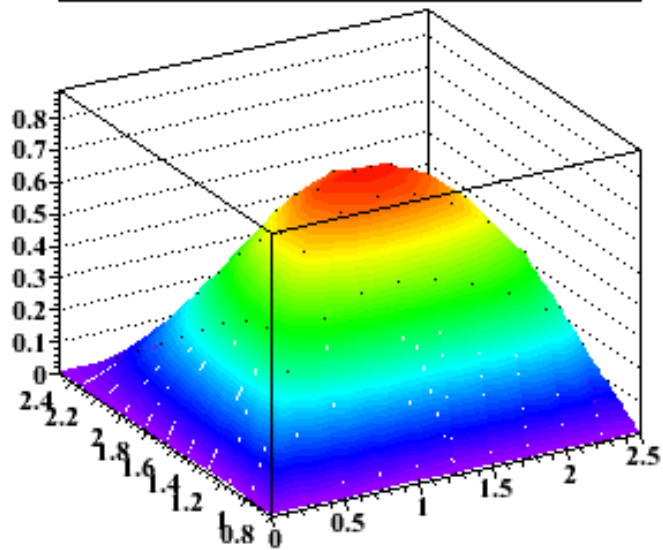
R-Phi



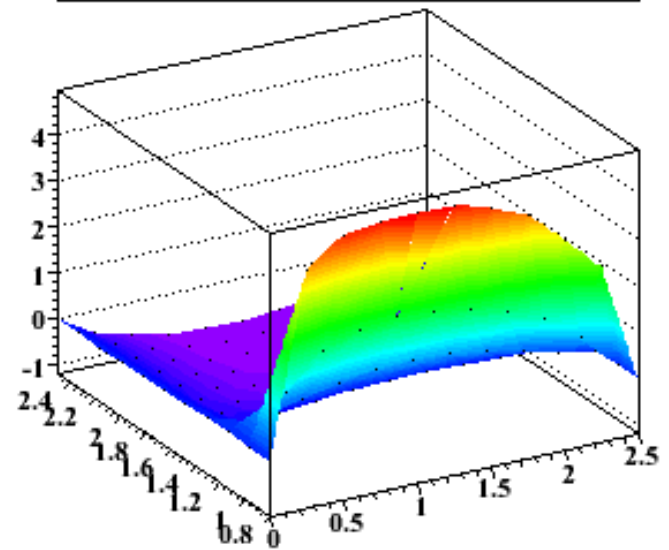
R-Z



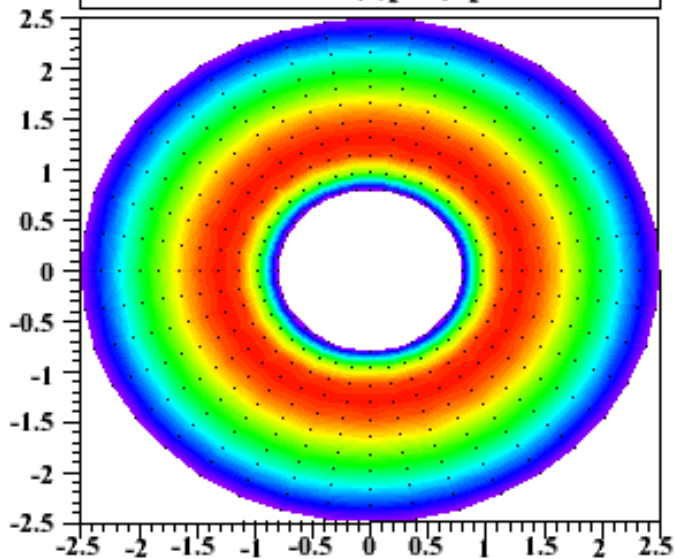
Potential: (r,z)-plane



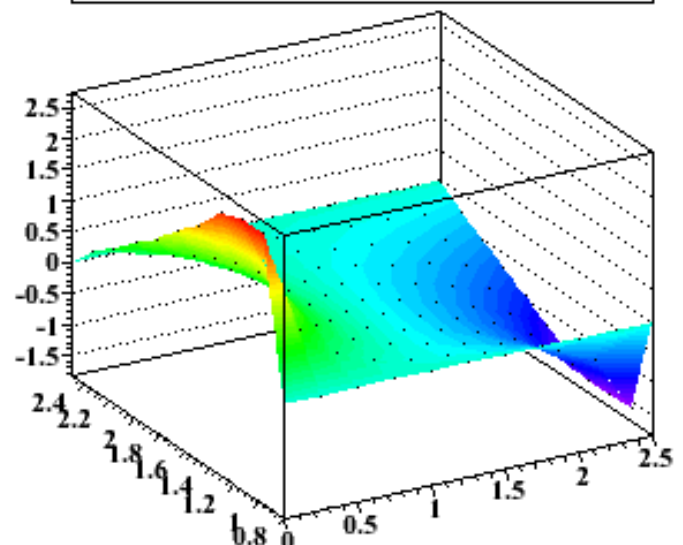
E field in R



Potential: (r,phi)-plane



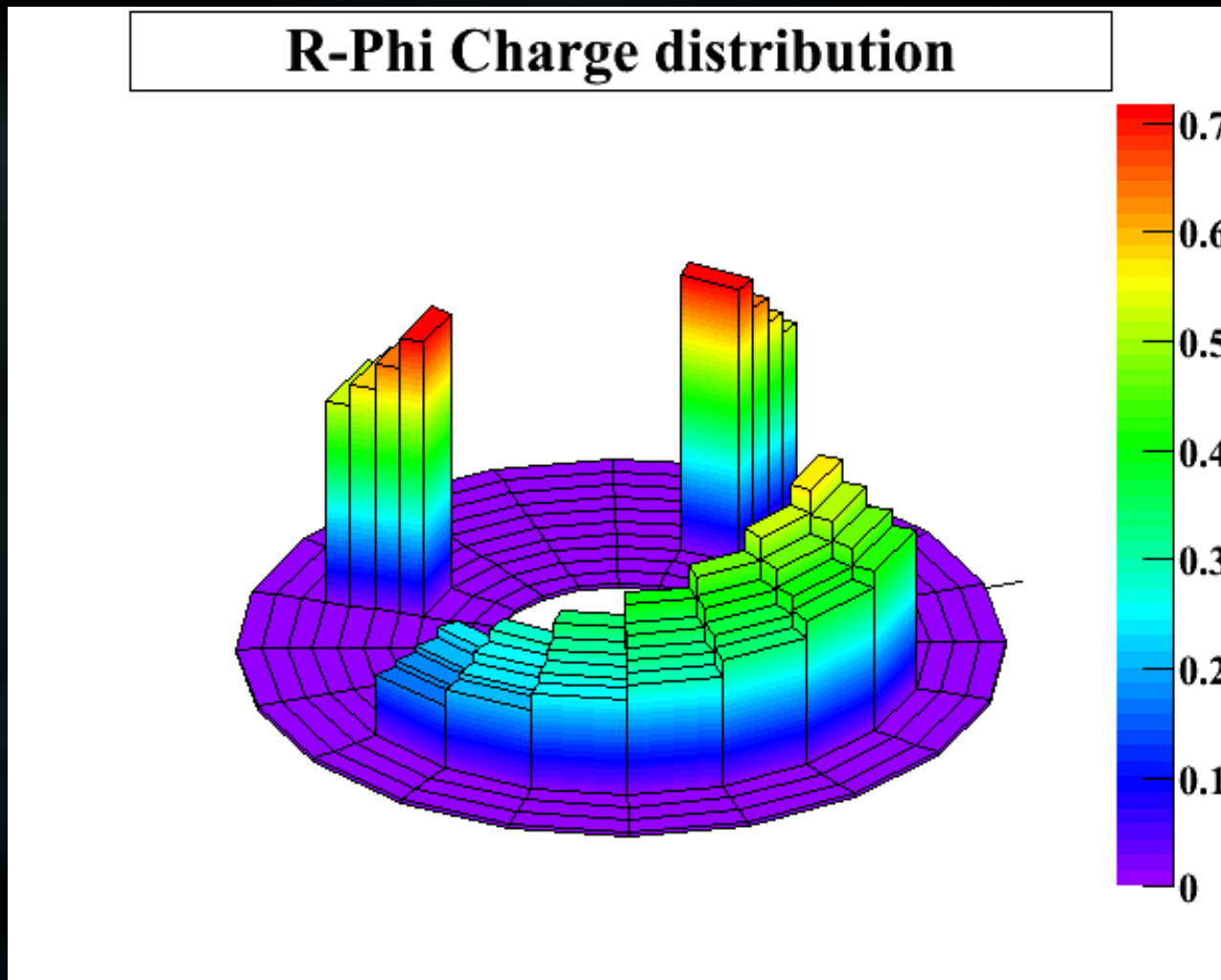
E field in Z

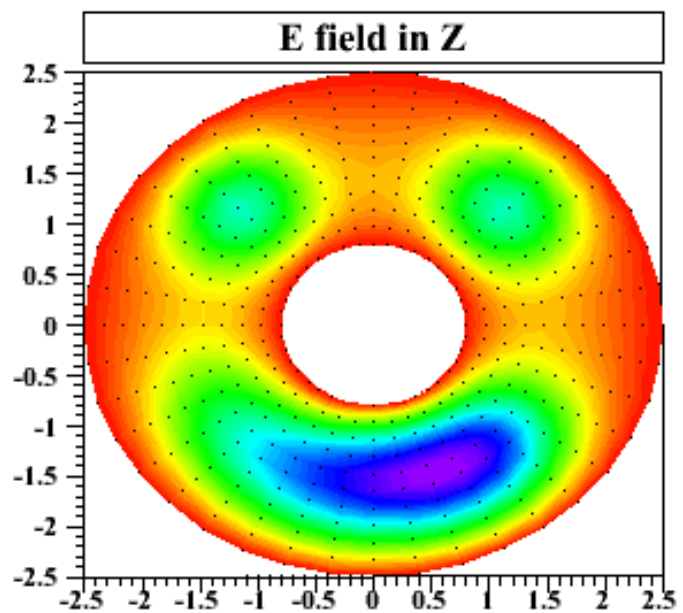
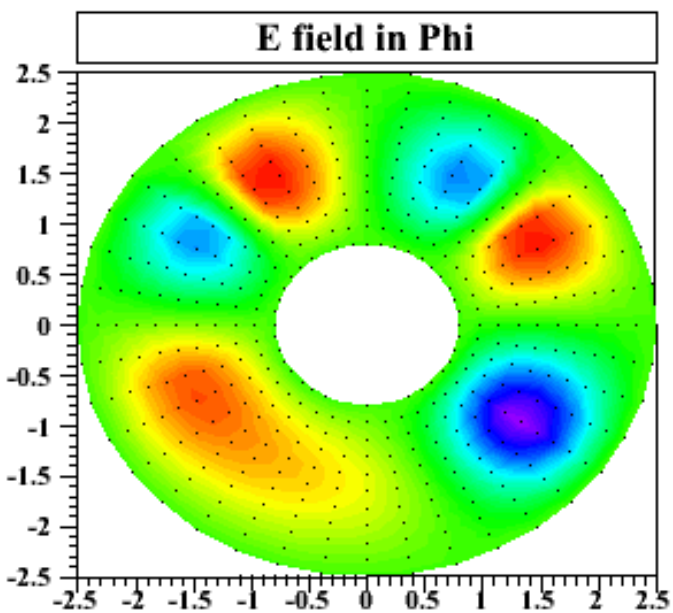
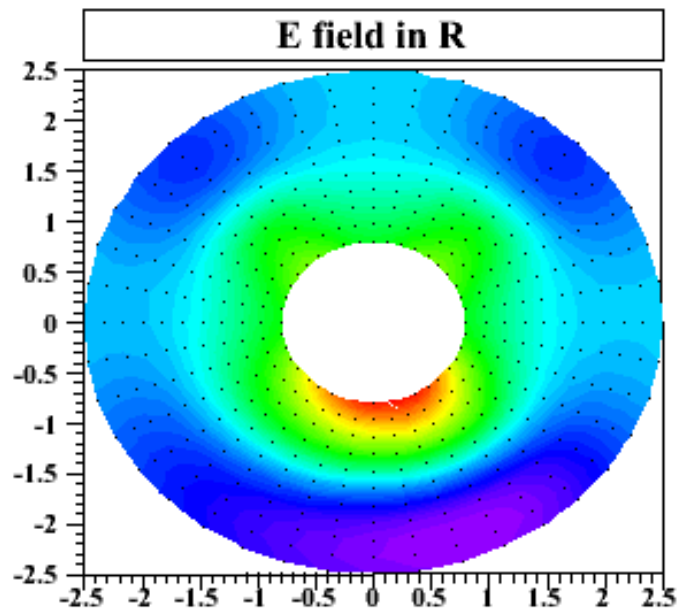
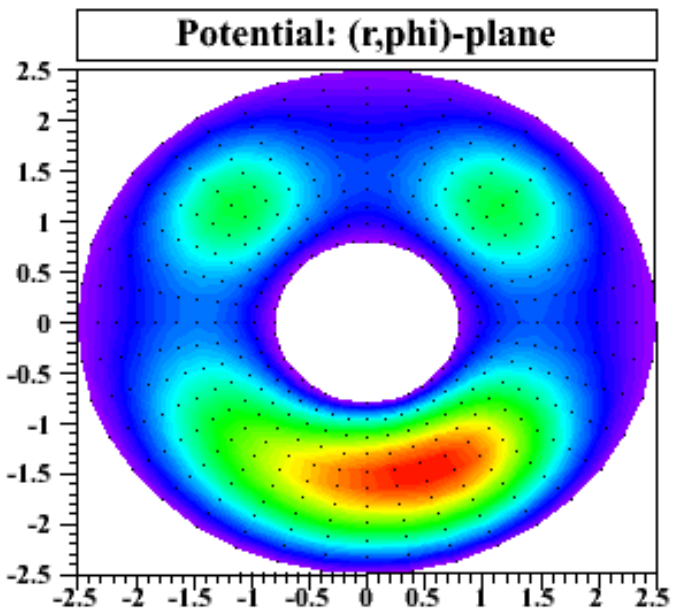


# Example 3:

FANCY Charge distribution

$\rho(r, \phi, z) = ??$  who can guess ??







# BACKUP slide:

(4)<sup>th</sup> solution: Green's function with fast convergence close to the Point Charge, in case you don't want to integrate over const. “Piece-Of-Cake” boxes ...

$$\begin{aligned}
 G(r, \phi, z, r', \phi', z') &= \\
 &= G_C(r, \phi, r', \phi', z - z') - G_C(r, \phi, r', \phi', z + z') - G_C(r, \phi, r', \phi', 2L - z - z') + \\
 &\quad + \frac{1}{2\pi} \sum_{m=0}^{\infty} \sum_{n=1}^{\infty} (2 - \delta_{m0}) \cos[m(\phi - \phi')] \frac{R_{mn}(r) R_{mn}(r')}{N_{mn}^2} \frac{g_s(\beta_{mn})}{\beta_{mn}} \quad (5.56)
 \end{aligned}$$

$$\begin{aligned}
 g_s &= g_0 - \frac{1}{2} [e^{-\beta(z > -z <)} - e^{-\beta(z > +z <)} - e^{-\beta(2L - z > -z <)}] = \\
 &= \frac{e^{-\beta(2L - z > +z <)} + e^{-\beta(2L + z > -z <)} - e^{-\beta(2L + z > +z <)} - e^{-\beta(4L - z > -z <)}}{2(1 - e^{-2\beta L})} \quad (5.54)
 \end{aligned}$$

$$\begin{aligned}
 G_C(r, \phi, r', \phi', \alpha) &= \\
 &= G_F(r, \phi, r', \phi', \alpha) + H_C(r, \phi, r', \phi', \alpha) \\
 &= \frac{1}{4\pi} \frac{1}{\sqrt{r^2 - 2rr' \cos(\phi - \phi') + r'^2 + \alpha^2}} - \dots \\
 &- \frac{1}{2\pi^2} \sum_{m=0}^{\infty} \epsilon_m \cos[m(\phi - \phi')] \int_0^{\infty} d\lambda \cos(\lambda\alpha) [I_m(\lambda r) c_m(\lambda) + K_m(\lambda r) d_m(\lambda)]
 \end{aligned}$$

$$\begin{aligned}
 \text{with } \epsilon_m &= (2 - \delta_{m0}) \\
 c_m(\lambda) &= K_m(\lambda b) R_m(a, r') / N_m \\
 d_m(\lambda) &= I_m(\lambda a) R_m(r', b) / N_m \\
 R_m(s, t) &= I_m(\lambda s) K_m(\lambda t) - I_m(\lambda t) K_m(\lambda s) \\
 N_m &= I_m(\lambda a) K_m(\lambda b) - I_m(\lambda b) K_m(\lambda a)
 \end{aligned}$$

(to be published)